

Nuclear Physics for Beyond the Standard Model Neutrino Physics

outline

1. Neutrino interaction physics - introduction
2. Charged-Current Quasi-Elastic (CCQE) interaction
3. Neutrino baryonic resonance interaction
4. Neutrino shallow- and deep-inelastic scatterings
5. Conclusion

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King's College London

Genova HEP seminar, University of Genova, May 3, 2023

All anomalies in particle physics are Strong interaction

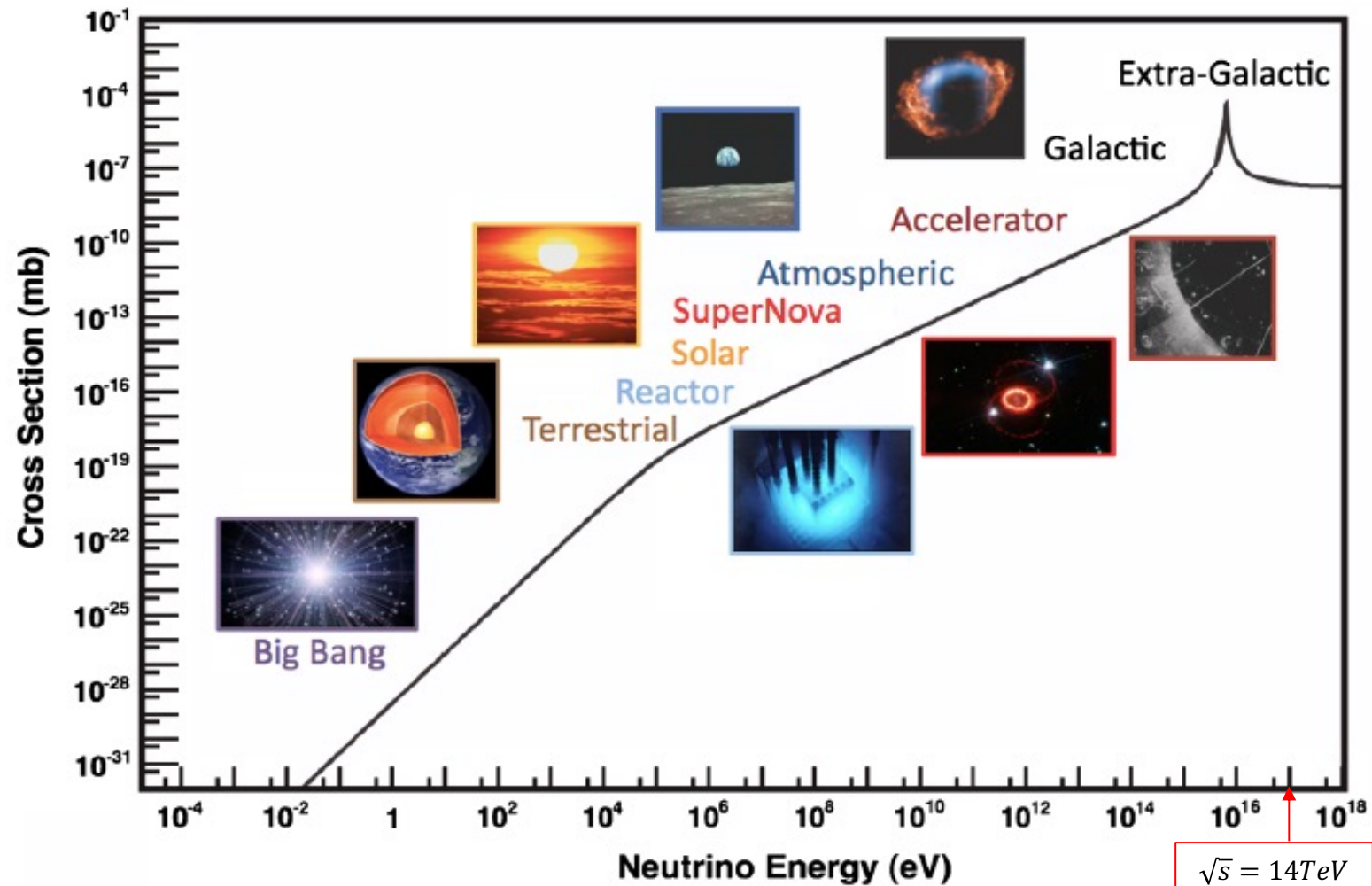
QCD

Hadron
physics

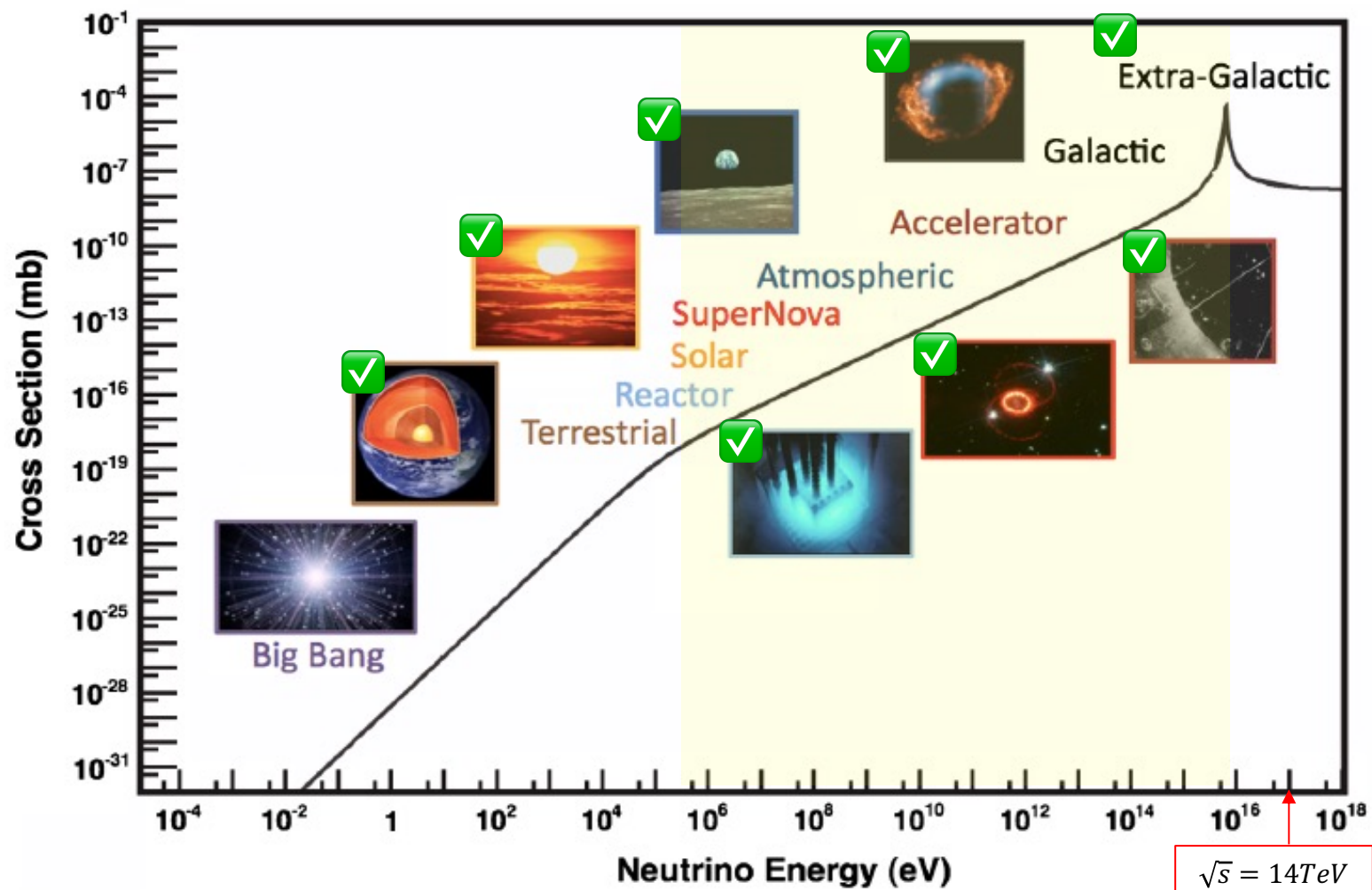
Nuclear
physics

- 1. Neutrino interaction physics - introduction**
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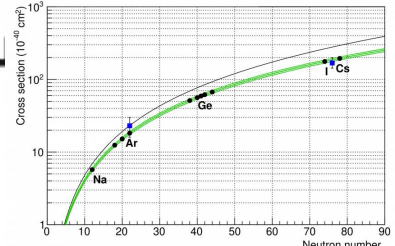
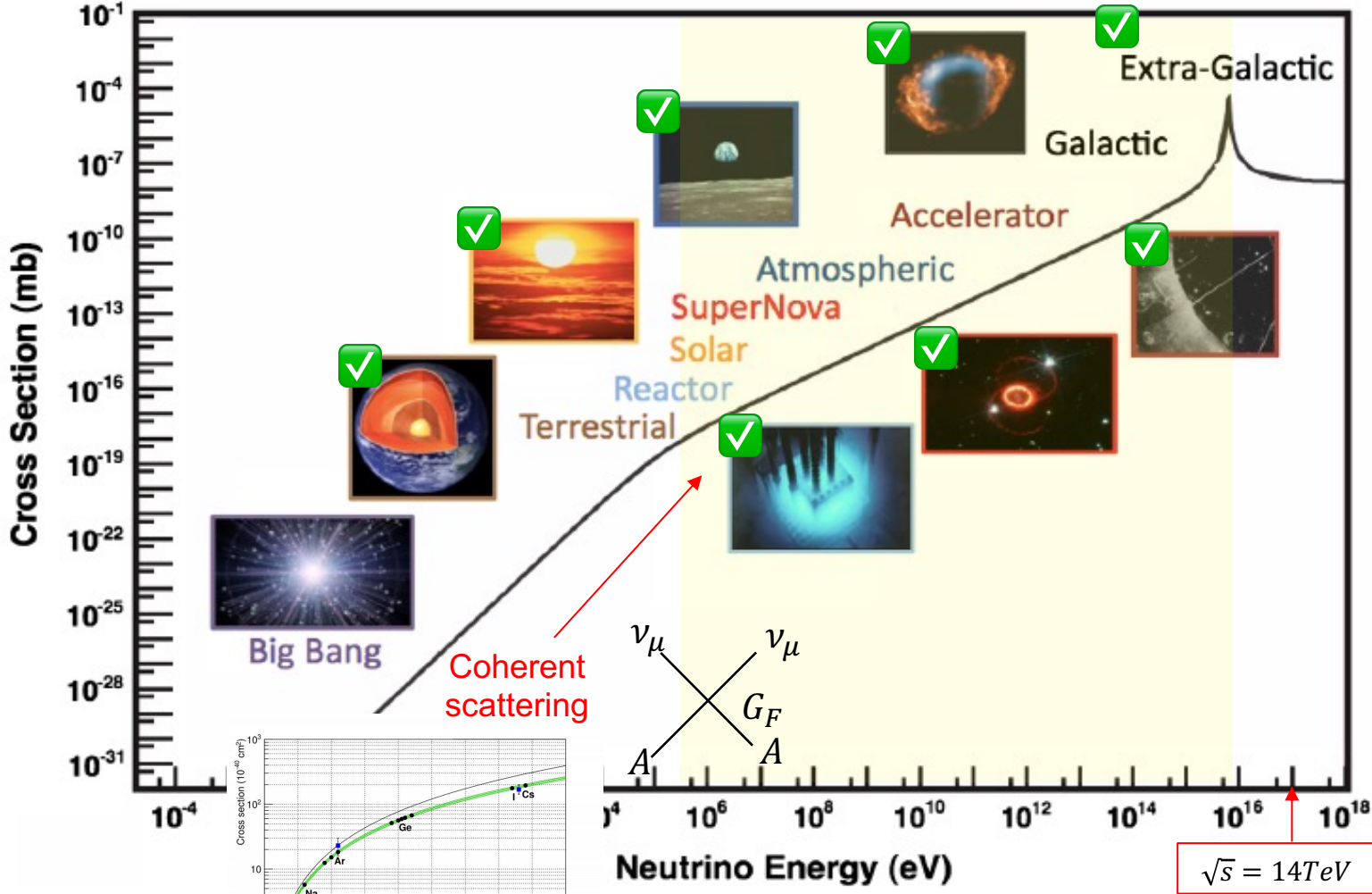
1. From eV to EeV: Neutrino cross sections across energy scales



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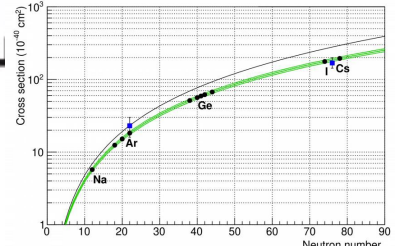
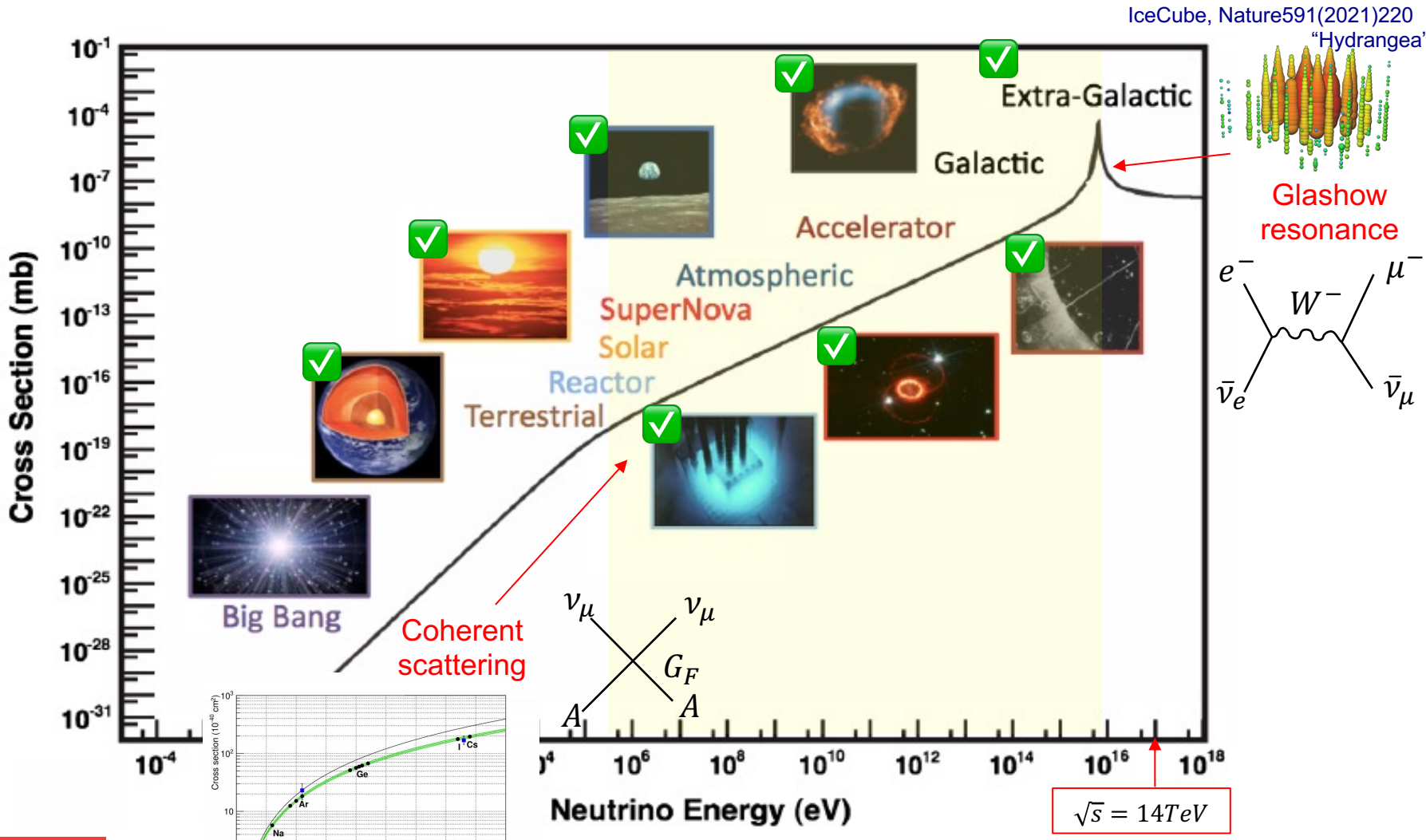
COHERENT, Science357(2017)1123
PRL126(2021)012002

katori@fnal.gov

2023/05/03



1. From eV to EeV: Neutrino cross sections across energy scales



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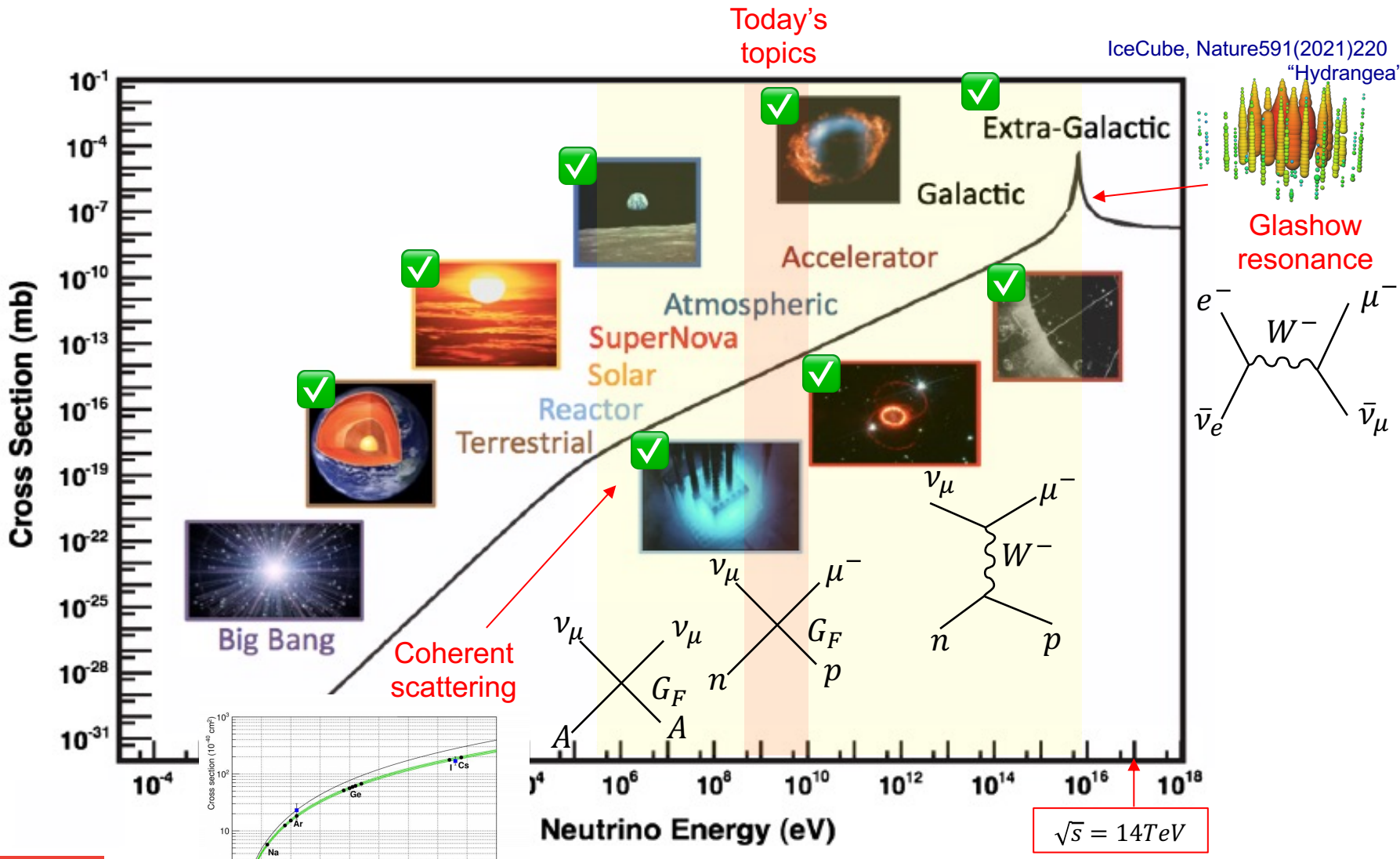
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2023/05/03

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1. From eV to EeV: Neutrino cross sections across energy scales



1. PDG: Neutrino Cross Section Measurements

PDG has a summary of neutrino cross-section data since 2012!

Focus of this talk is around a few GeV

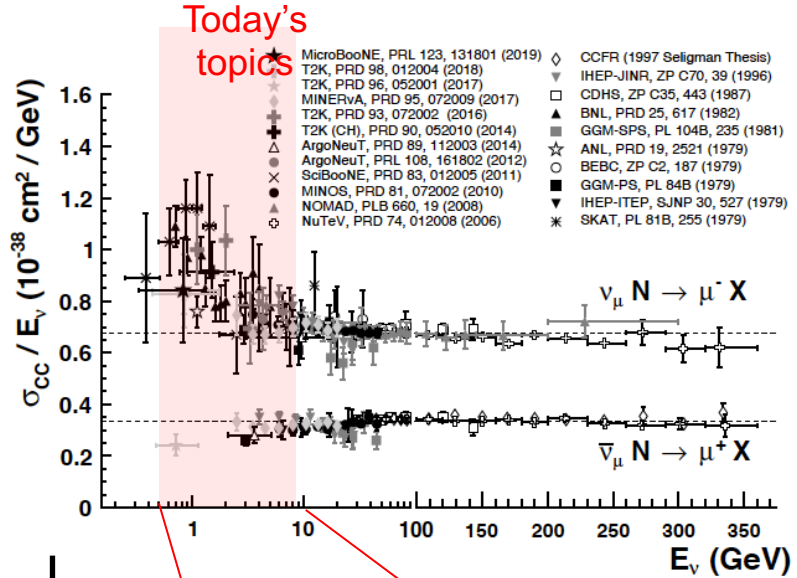
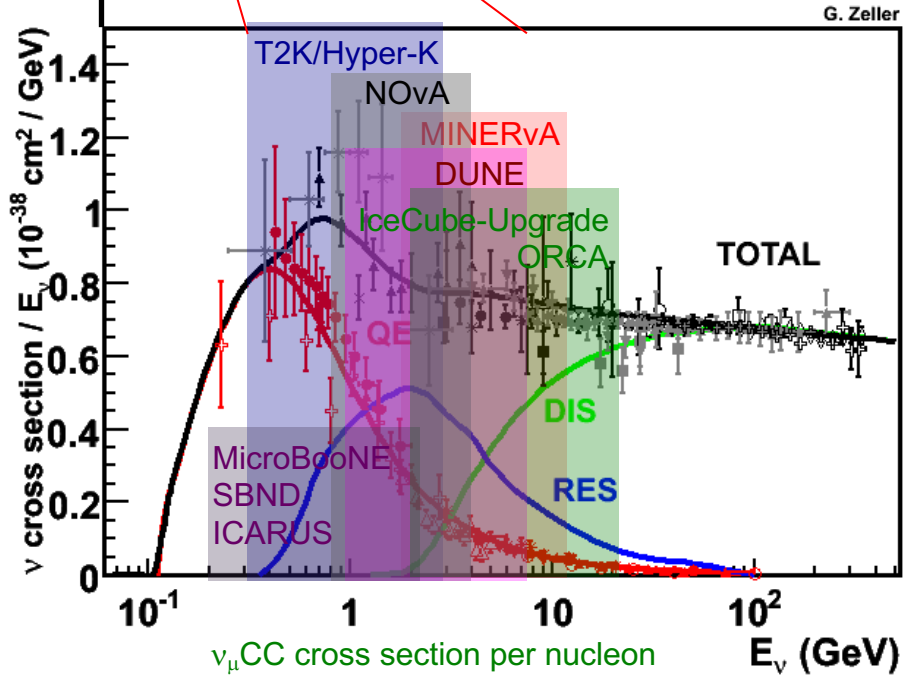


Table 52.2: Published measurements of neutrino and antineutrino CC inclusive cross sections from modern accelerator-based neutrino experiments.

experiment	measurement	target
ArgoNeuT	ν_μ [6, 7], $\bar{\nu}_\mu$ [7]	Ar
MicroBooNE	ν_μ [8, 26], ν_e [22]	Ar
MINERvA	ν_μ [9–11, 16, 17, 27], $\bar{\nu}_\mu$ [27], $\bar{\nu}_\mu/\nu_\mu$ [28]	CH, C/CH, Fe/CH, Pb/CH
MINOS	ν_μ [29], $\bar{\nu}_\mu$ [29]	Fe
NINJA	ν_μ [12], $\bar{\nu}_\mu$ [12]	H ₂ O
NOMAD	ν_μ [30]	C
SciBooNE	ν_μ [31]	CH
T2K	ν_μ [13, 14, 32–34], ν_e [23–25], $\bar{\nu}_\mu/\nu_\mu$ [15]	CH, H ₂ O, Fe

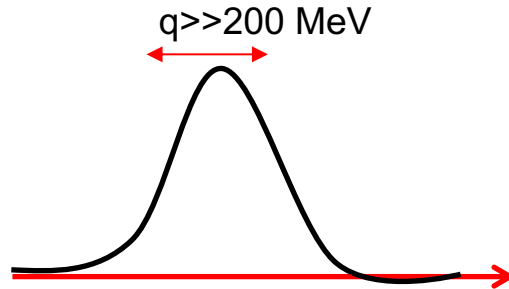


1. Neutrino interaction physics around 1-10 GeV

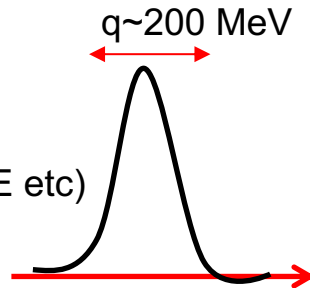
Size of wave packet \sim momentum transfer (\sim energy)

$$\hbar c = 197 \text{ MeV} \cdot \text{fm} \rightarrow 200 \text{ MeV} \sim 1 \text{ fm (size of nucleon)}$$

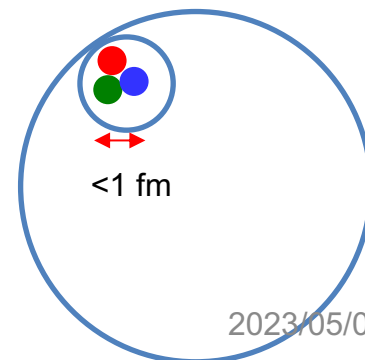
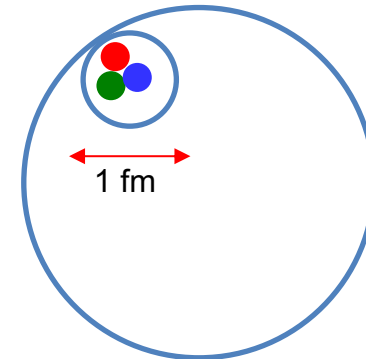
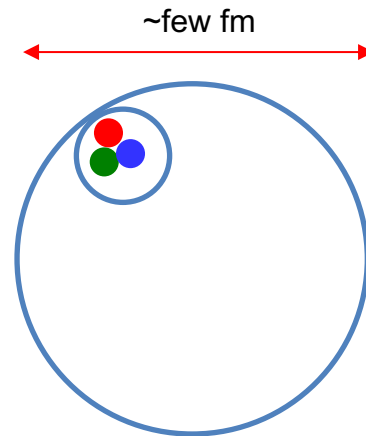
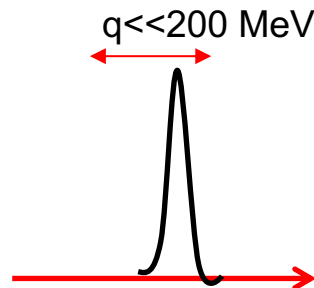
$\ll 1$ GeV neutrino beam
(solar neutrinos, etc)



~ 1 GeV neutrino beam
(T2K, NOvA, HyperK, DUNE etc)



$\gg 1$ GeV neutrino beam
(LHC, astrophysical)

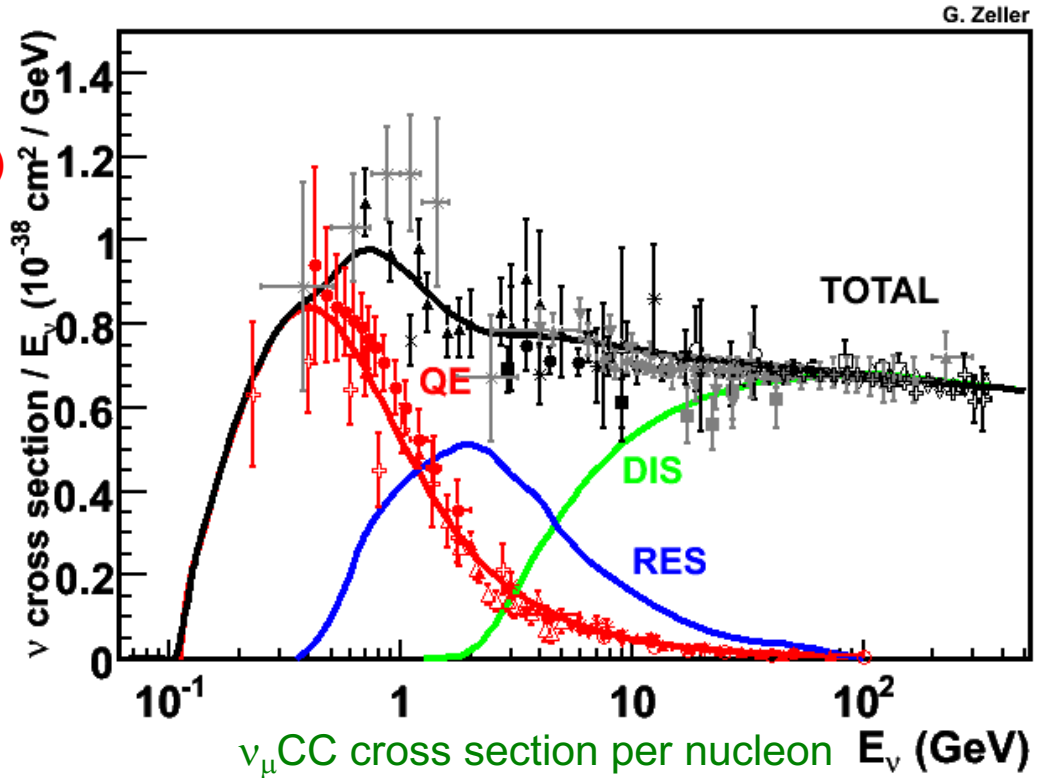
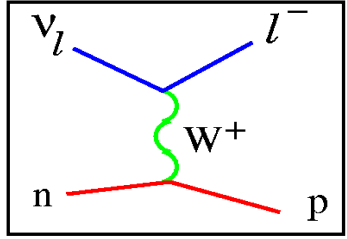


1. Neutrino interaction physics around 1-10 GeV

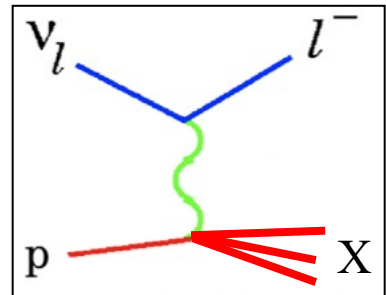
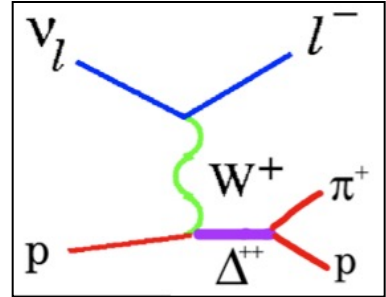
Neutrino interaction physics around 1-10 GeV

- degree of freedom change from nucleus → nucleon → parton
- There is no cut off (they all interfere)

Quasi Elastic (QE)



baryonic RESonance



Deep Inelastic Scattering (DIS)

1. Next goal of high energy physics

Establish Neutrino Standard Model (ν SM)

- SM + 3 active massive neutrinos

Unknown parameters of ν SM

1. Dirac CP phase
 2. θ_{23} ($\theta_{23}=40^\circ$ and 50° are same for $\sin 2\theta_{23}$, but not for $\sin\theta_{23}$)
 3. normal mass ordering $m_1 < m_2 < m_3$ or inverted mass ordering $m_3 < m_1 < m_2$
 4. Dirac or Majorana
 5. Majorana phases (x2)
 6. Absolute neutrino mass
- } not relevant to neutrino oscillation experiment

We need higher precision neutrino experiments around 1-10 GeV.

Low energy beam (~1 GeV)

- shorter baseline (lower flux reduction)
- lower neutrino production
- lower interaction rate
- kinematic energy reconstruction

High energy beam (~few GeV)

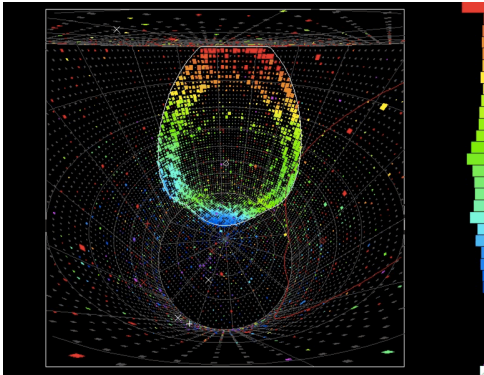
- longer baseline (higher flux reduction)
- higher neutrino production
- higher interaction rate
- calorimetric energy reconstruction

$$P_{\mu \rightarrow e}(L/E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L(km)}{E(GeV)} \right)$$

1. Next goal of high energy physics

Kinematics energy reconstruction

- problem: it assume 2-body neutrino interaction with single nucleon



$$E_{\nu}^{QE} = \frac{ME_{\nu} - 0.5m_{\mu}^2}{M - E_{\mu} + p_{\mu}\cos\theta}$$

Low energy beam (~1 GeV)

- shorter baseline (lower flux reduction)
- lower neutrino production
- lower interaction rate
- kinematic energy reconstruction

Calorimetric energy reconstruction

- problem: you need to measure energy deposit from all outgoing particles



$$E_{\nu}^{Cal} = E_{\mu} + \sum_{i=1}^{all} E_{had}^i$$

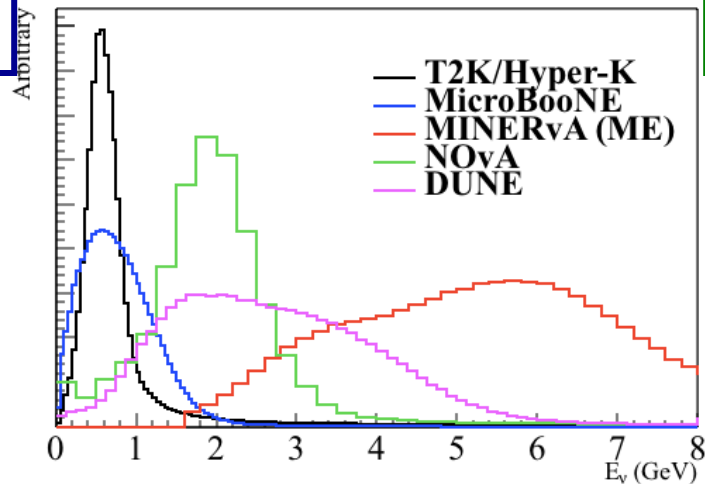
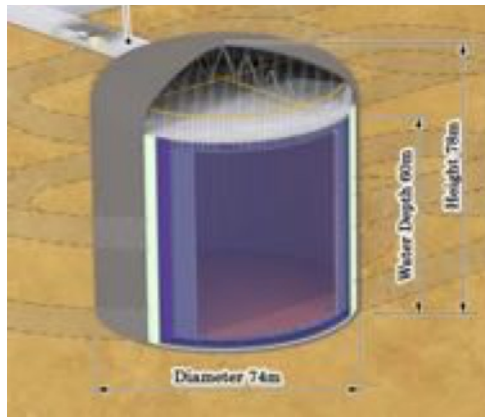
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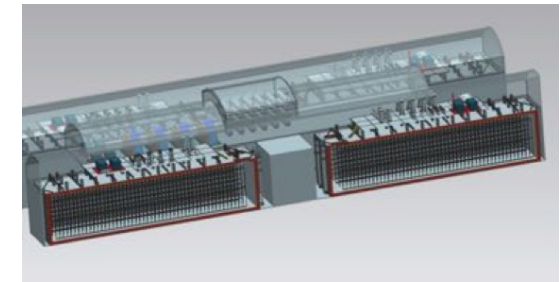
Hyper-Kamiokande (Japan)

- Water target
- 300 km baseline
- Narrow band 0.6 GeV
- Low spatial resolution
- High time resolution



DUNE (USA)

- Argon target
- 1300 km baseline
- wide band 1-4 GeV
- High spatial resolution
- Low time resolution



Low energy beam (~1 GeV)

- shorter baseline (lower flux reduction)
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High energy beam (~few GeV)

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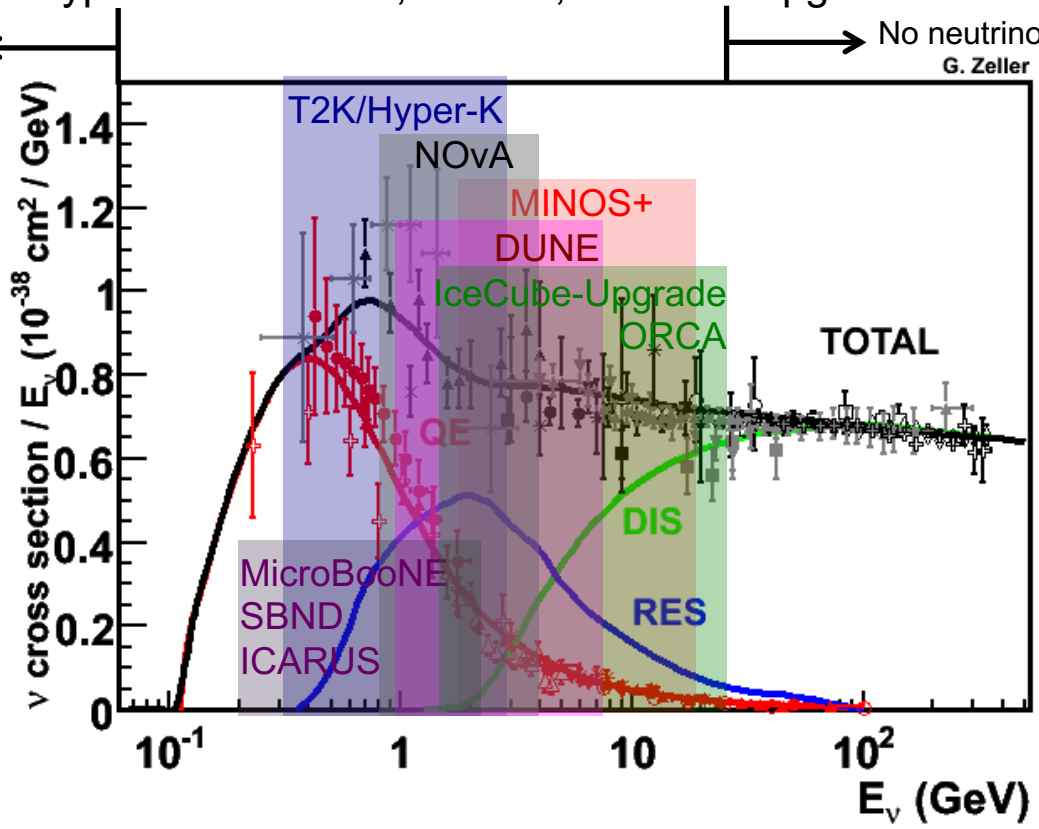
1. Next generation neutrino oscillation experiments

Current and future neutrino oscillation experiments

- J-PARC: T2K, Hyper-Kamiokande,
- Fermilab: MicroBooNE/SBND/ICARUS, MINOS+, NOvA, DUNE
- Atmospheric: Hyper-Kamiokande, ORCA, IceCube-Upgrade

Reactor and Solar neutrino oscillation

No neutrino oscillation
G. Zeller



Neutrino oscillation length reaches ~12700km (diameter of Earth)



$$P_{\mu \rightarrow e}(L/E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L(km)}{E(GeV)} \right)$$

1. Neutrino cross-section formula

Cross-section

- product of Leptonic and Hadronic tensor

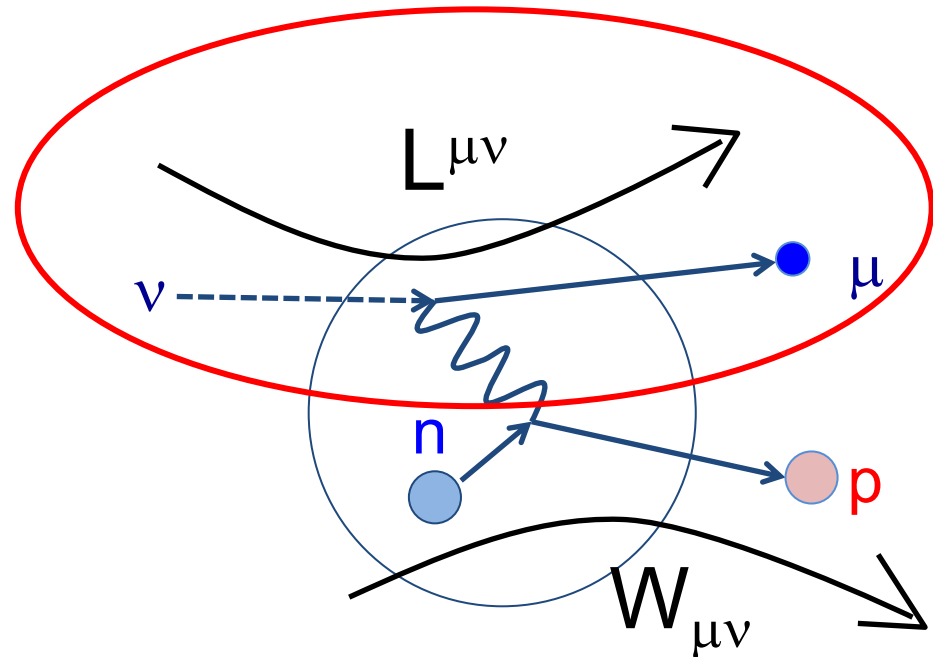
$$d\sigma \sim L^{\mu\nu}W_{\mu\nu}$$

Leptonic tensor

→ the Standard Model (easy)

Hadronic tensor

→ nuclear physics (hard)



1. Neutrino cross-section formula

Cross-section

- product of Leptonic and Hadronic tensor

$$d\sigma \sim L^{\mu\nu}W_{\mu\nu}$$

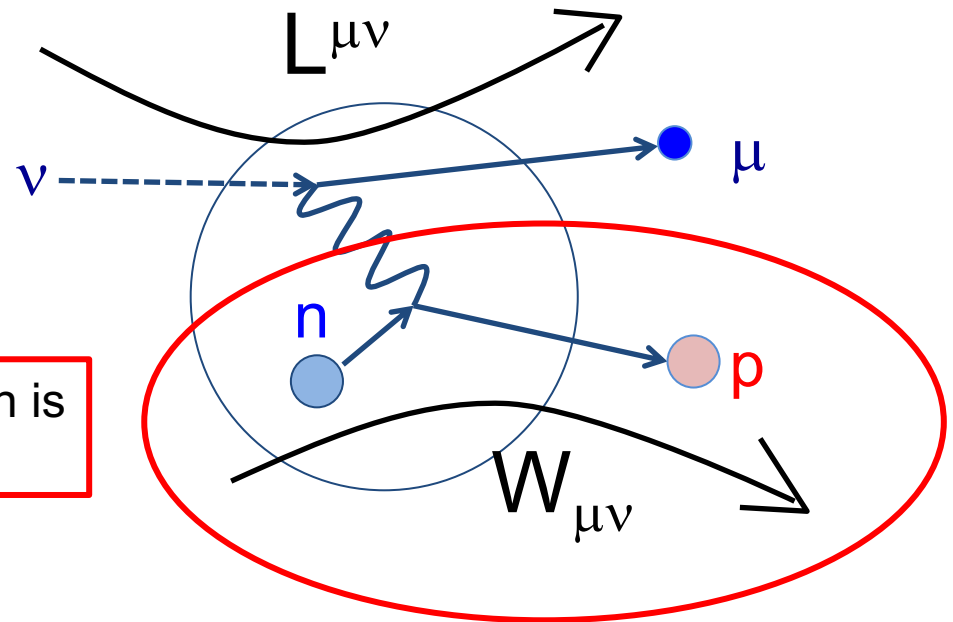
Leptonic tensor

→ the Standard Model (easy)

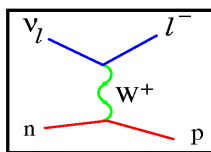
Hadronic tensor

→ nuclear physics (hard)

All complication of neutrino cross-section is how to model the hadronic tensor part



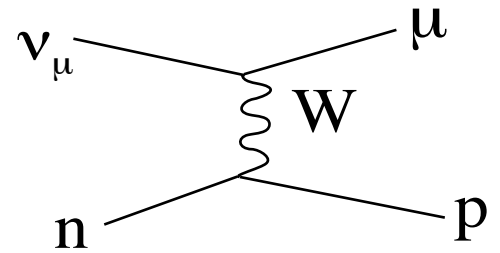
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2. Charged Current Quasi-Elastic scattering (CCQE)

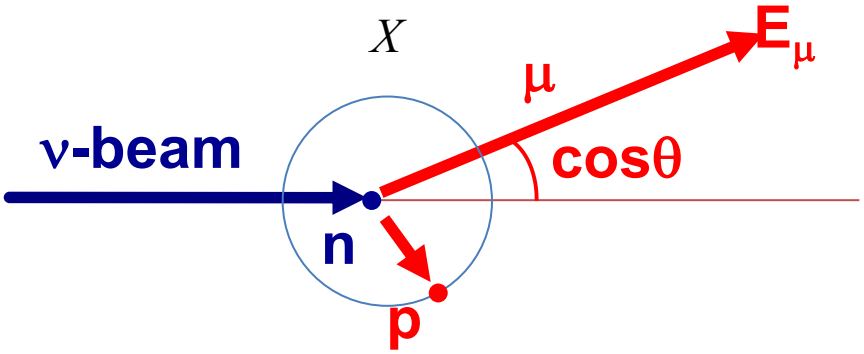
The simplest and the most abundant interaction around ~1 GeV.

$$\nu_\mu + n \rightarrow p + \mu^- \quad (\nu_\mu + X \rightarrow X' + \mu^-)$$



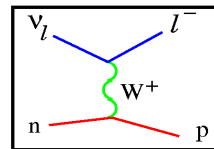
Neutrino energy is reconstructed from the observed lepton kinematics
 “QE assumption”

1. assuming neutron at rest
2. assuming interaction is CCQE



$$E_\nu^{QE} = \frac{ME_\nu - 0.5m_\mu^2}{M - E_\mu + p_\mu \cos\theta}$$

CCQE is the single most important channel of neutrino oscillation physics
 T2K, NOvA, microBoonE, Hyper-Kamiokande...etc



2. CCQE puzzle

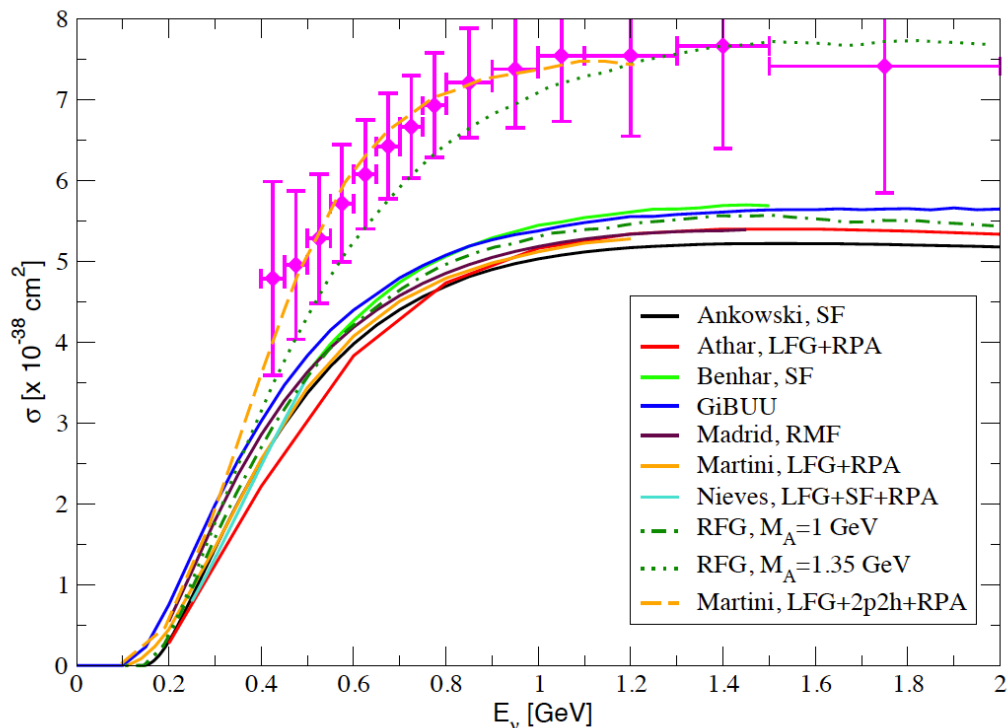
Simplest channel, but both shape and normalization disagree

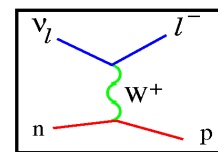
1. low Q2 suppression → Low forward efficiency? (detector?)
2. high Q2 enhancement → Axial mass > 1.0 GeV? (physics?)
3. large normalization → Beam simulation is wrong? (flux?)

CCQE interaction on nuclear targets are precisely measured by electron scattering

- Lepton universality = precise prediction for neutrino CCQE cross-section...?

MiniBooNE ν_μ CCQE cross section on ^{12}C target (per nucleon)





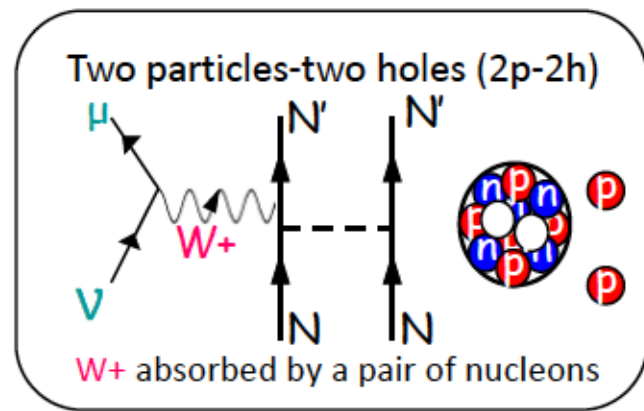
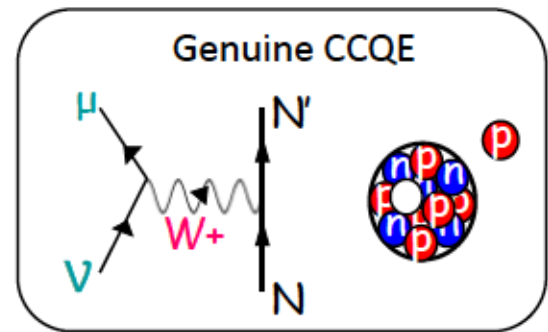
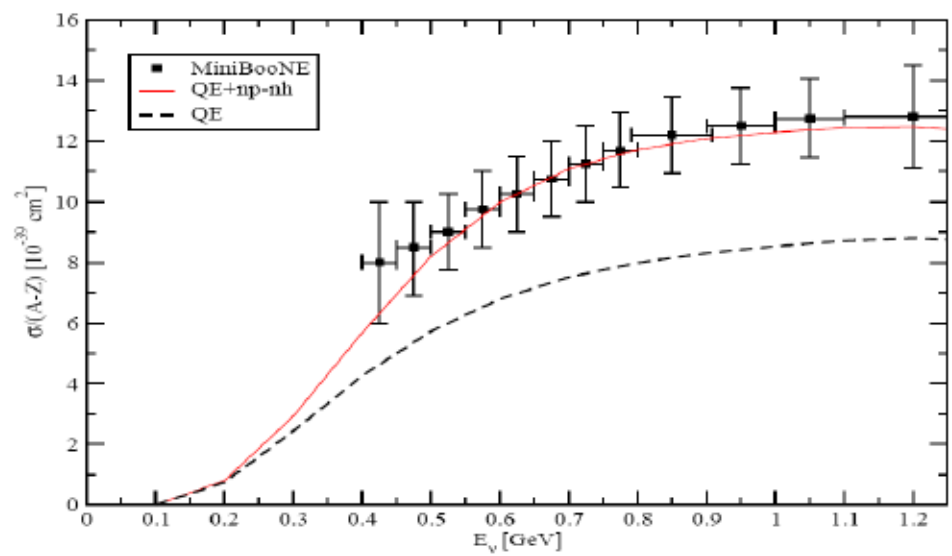
2. Solution of CCQE puzzle

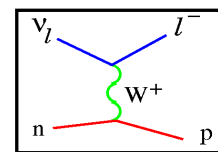
Presence of 2-body current

- CCQE is identified from single outgoing charged lepton events
- Significant fraction of events are not from 2-body neutrino-nucleon interactions
- Martini et al showed 2p-2h effect can add up ~30% more cross section

An explanation of this puzzle

Inclusion of the multinucleon emission channel (np-nh)





2. Models using 2p-2h

Flux-averaged differential cross-sections allow nuclear theorists to compare their models with data without implementing them in generators

Martini et al – Lyon 2p2ph model

Nieves et al – Valencia 2p2h model

SuSAv2 – Superscaling+MEC

Giusti et al – Relativistic Green’s function

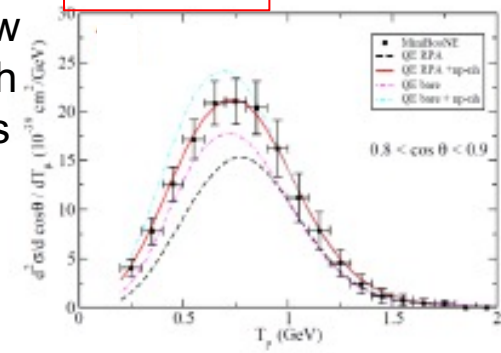
Butkevich et al – RDWIA+MEC

Lovato et al – GFMC

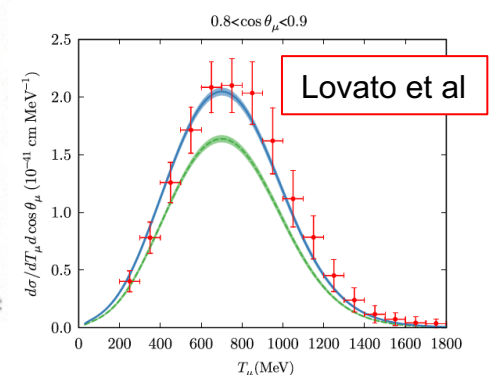
Jachowicz et al – CRPA+MEC

All models can fit with data, are they all correct models?

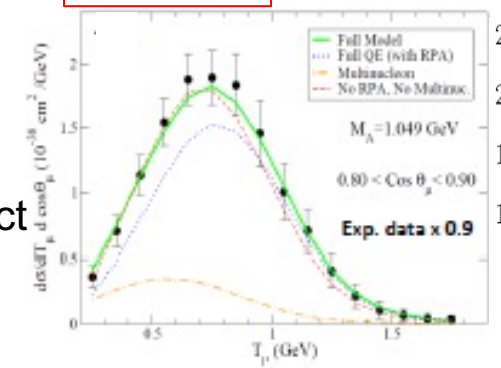
Martini et al



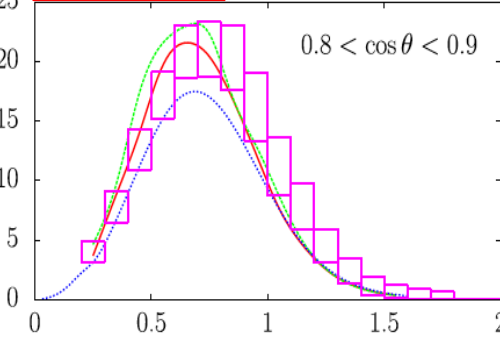
Lovato et al



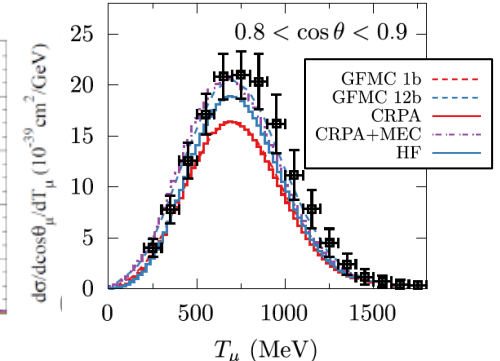
Nieves et al



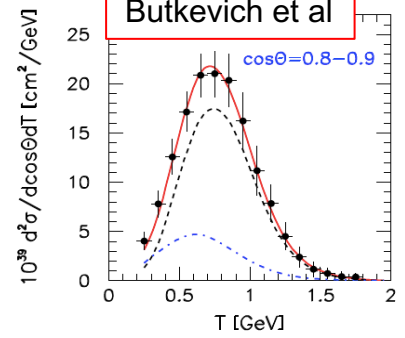
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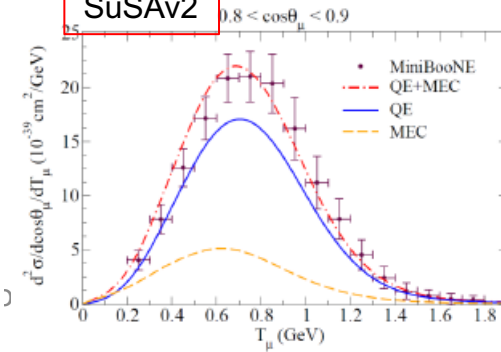
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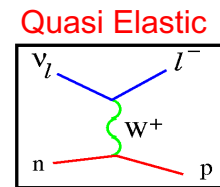


Butkevich et al



SuSAv2





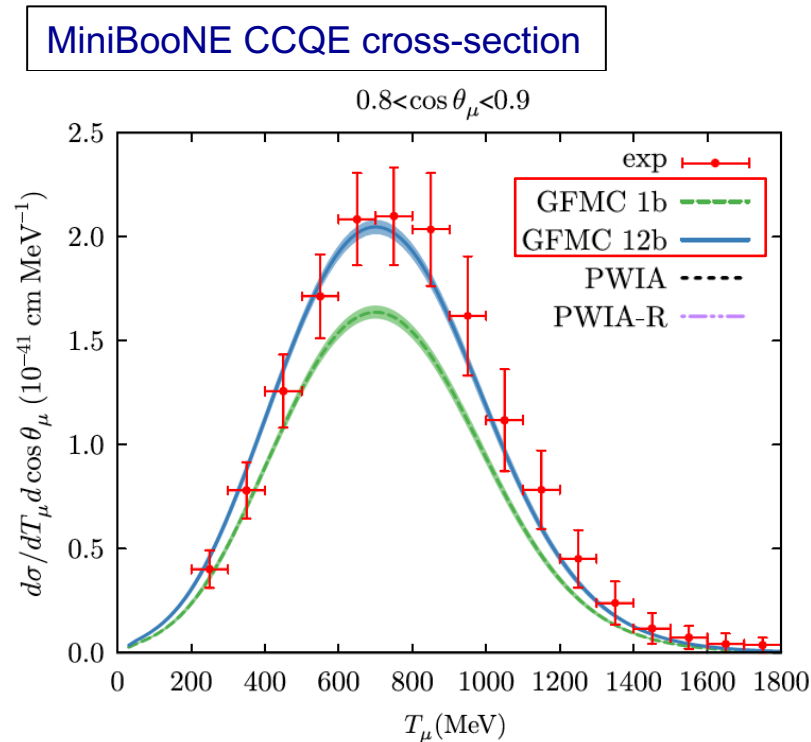
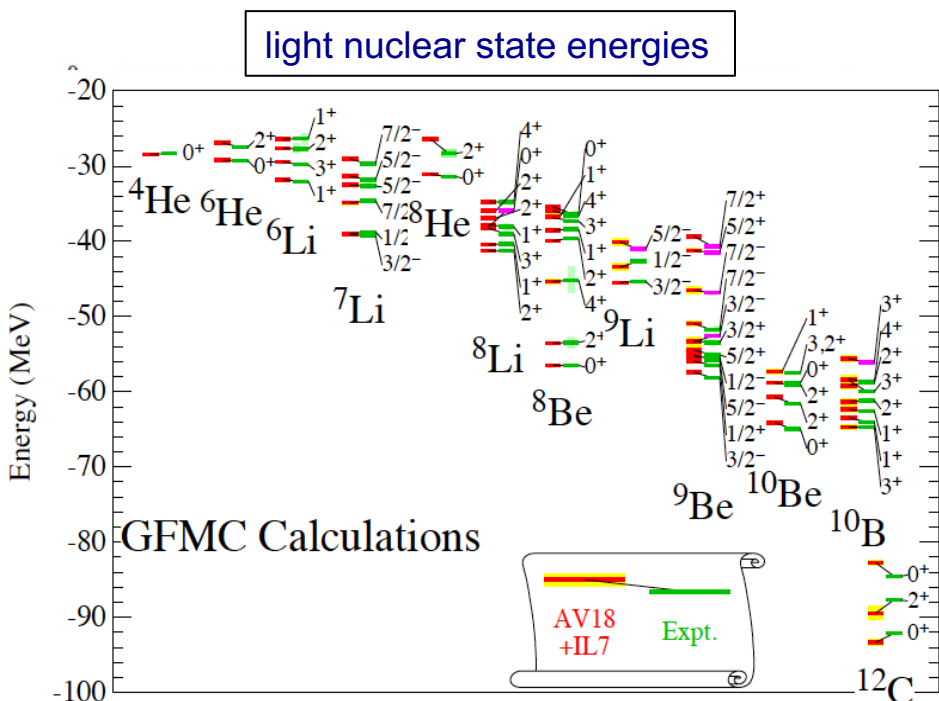
2. Nucleon correlations in neutrino physics

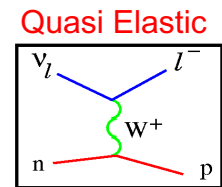
Ab-initio calculation

- Quantum Monte Carlo (QMC)
- Predicts energy levels of all light nuclei
- Consistent result with phenomenological models
- Ground state includes correct nucleon correlations

$$|\Psi_V\rangle = \mathcal{S} \prod_{i < j}^A \left[1 + \boxed{U_{ij}} + \sum_{k \neq i, j}^A \boxed{\tilde{U}_{ijk}^{TNI}} \right] |\Psi_J\rangle$$

2N potential (Av18)
3N potential (IL7)





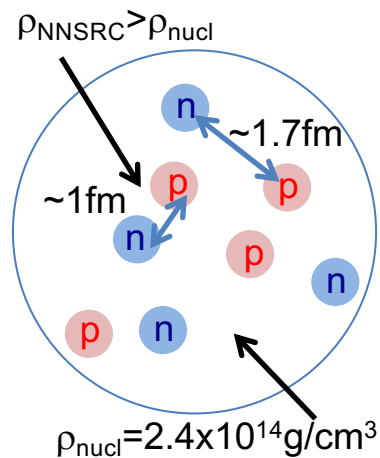
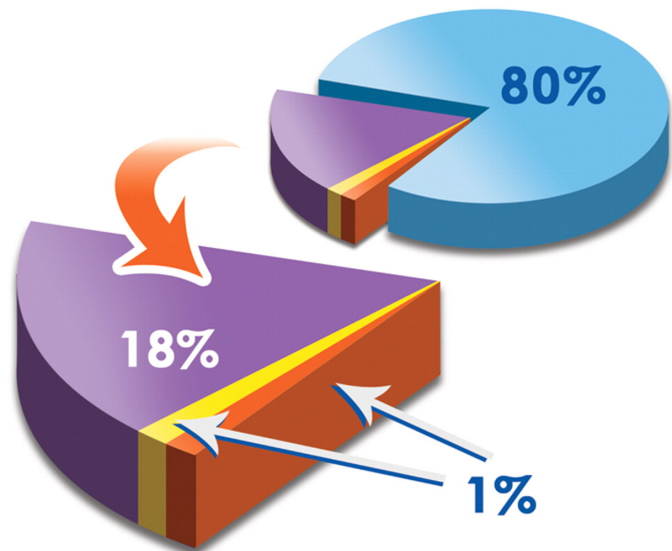
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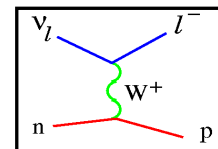
$$|\Psi_V\rangle = \mathcal{S} \prod_{i < j} \left[1 + \boxed{U_{ij}} + \sum_{k \neq i, j} \boxed{\tilde{U}_{ijk}^{TNI}} \right] |\Psi_J\rangle$$

2N potential (Av18)
3N potential (IL7)



Physics of nucleon correlation

- neutrino interaction
- $0\nu\beta\beta$
- Direct WIMP detection
- EMC effect
- etc

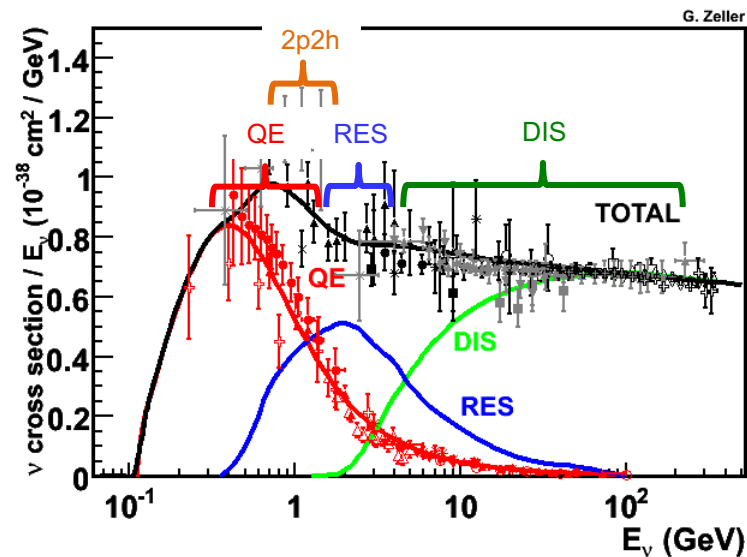
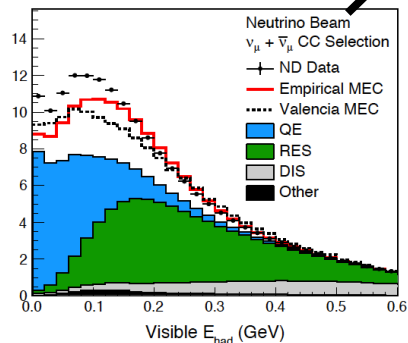
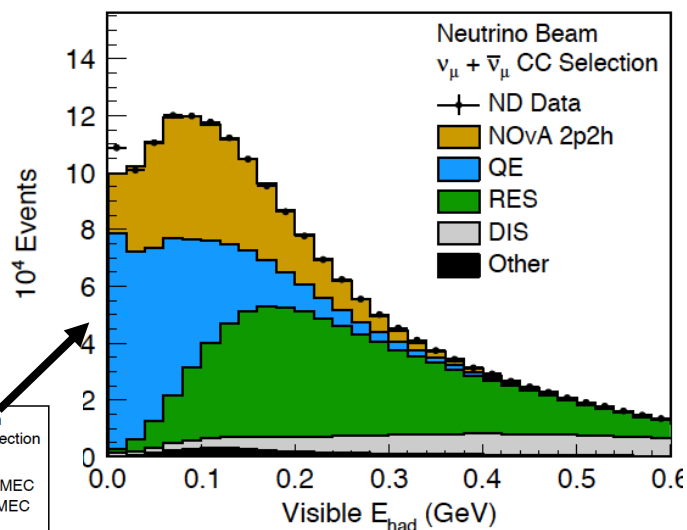


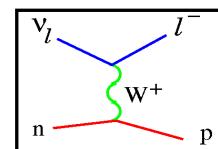
2. Nucleon correlations in neutrino physics

2-particle 2-hole (2p2h) effect

- Essential to describe data
- The biggest topic in nuxsec community (T2K, NOvA, MINERvA, MicroBooNE, etc)
- 2p2h models in generators don't describe data without heavy tuning
- High resolution detector (LArTPC, emulsion, etc) can find what is going on?

NOvA near detector data-MC comparison after fit

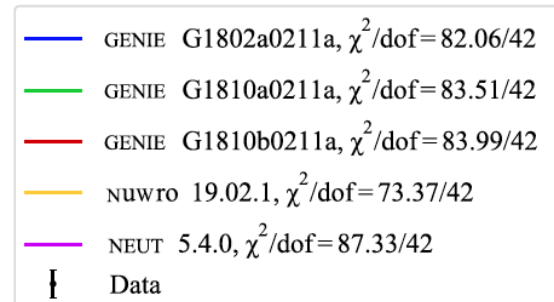
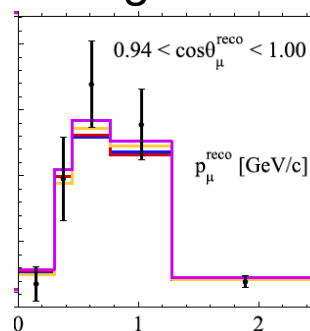




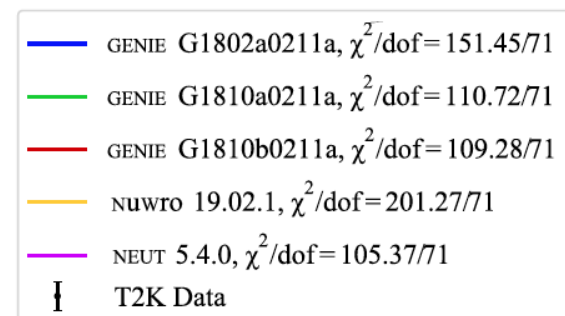
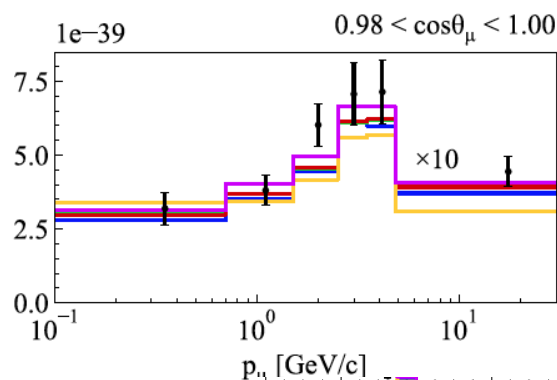
2. Data tension

Data tension – external: T2K vs. MINERvA vs. MicroBooNE
- Different kinematic coverage, different target

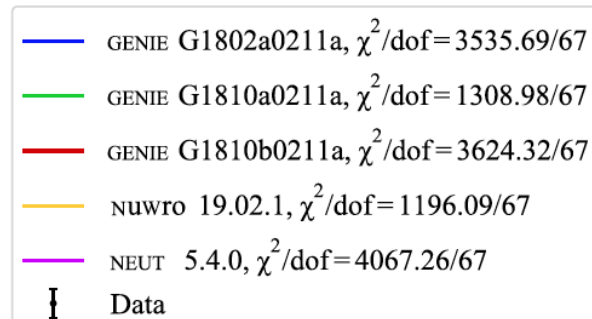
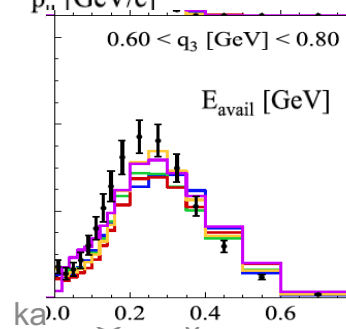
MicroBooNE CC inclusive double differential cross-section



T2K CC inclusive double differential cross-section



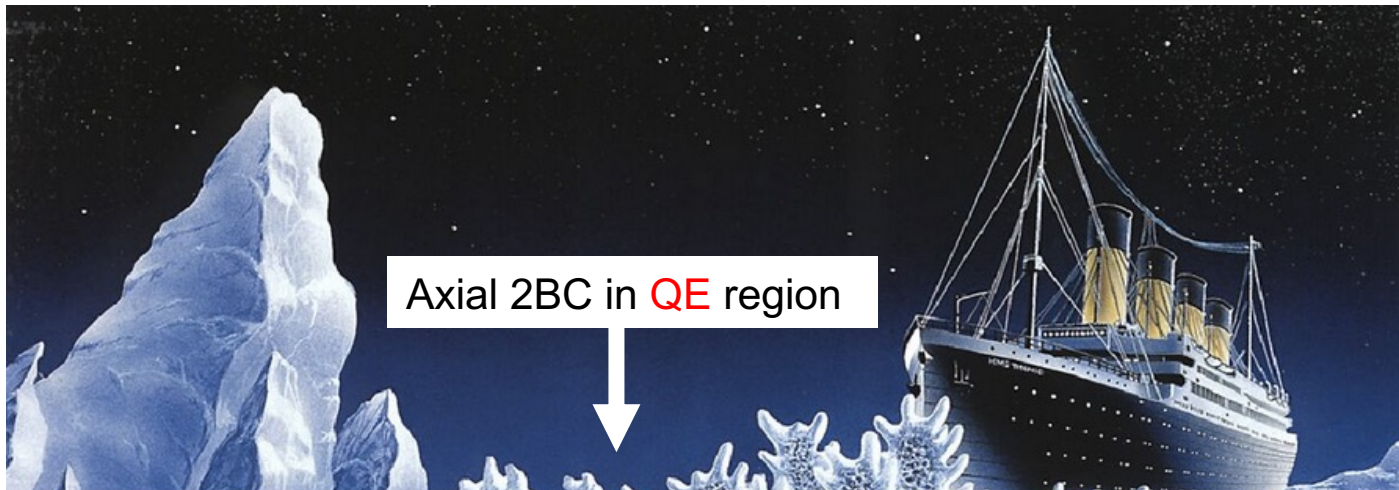
MINERvA CC inclusive double differential cross-section



1. Neutrino interaction physics - introduction
2. Charged-Current Quasi-Elastic (CCQE) interaction
- 3. Neutrino baryonic resonance interaction**
4. Neutrino shallow- and deep-inelastic scatterings
5. Conclusions

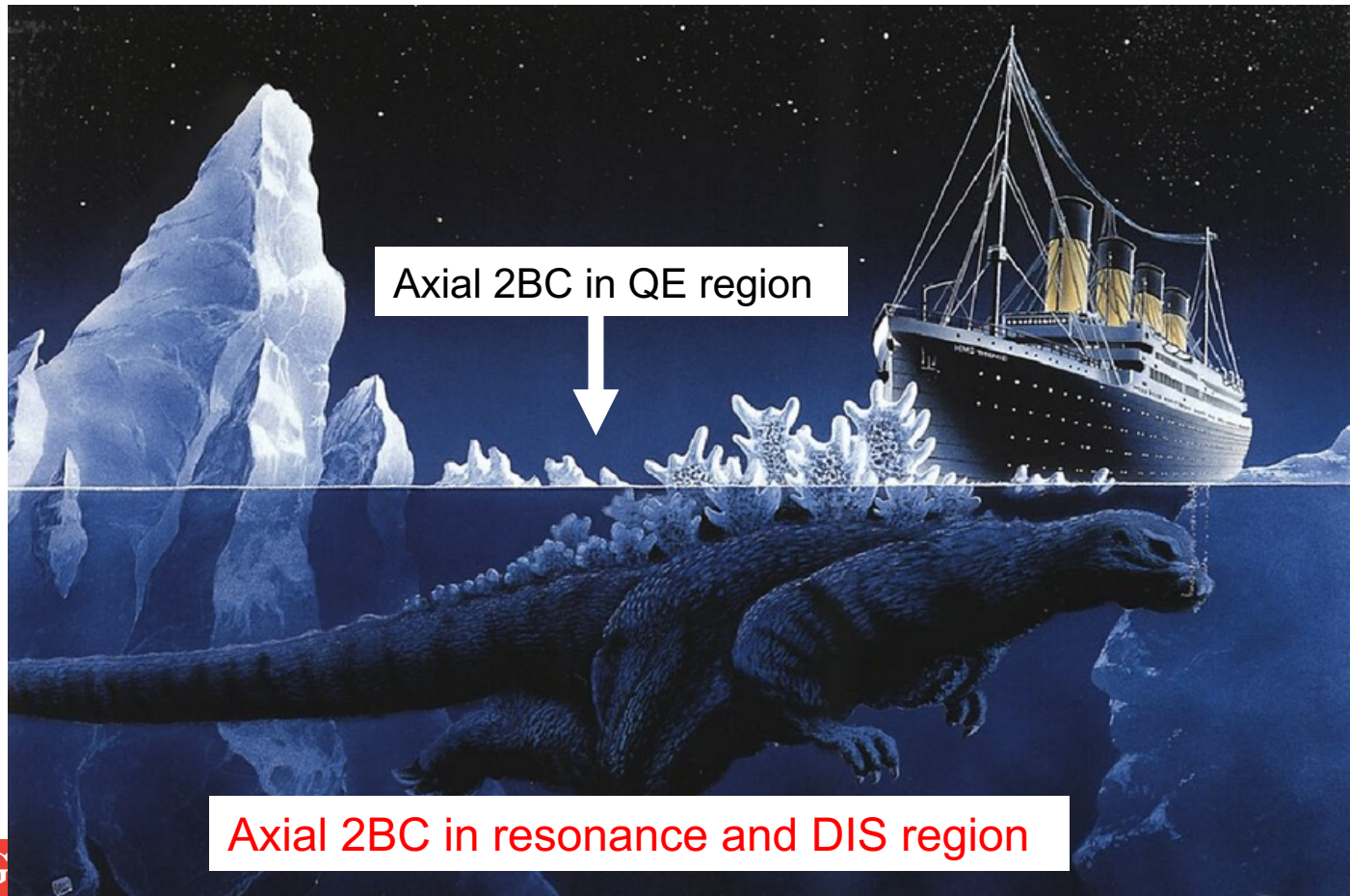
3. Beyond QE peak

Axial 2-body current in QE region may be a tip of the iceberg...

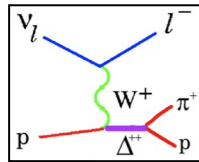


3. Beyond QE peak

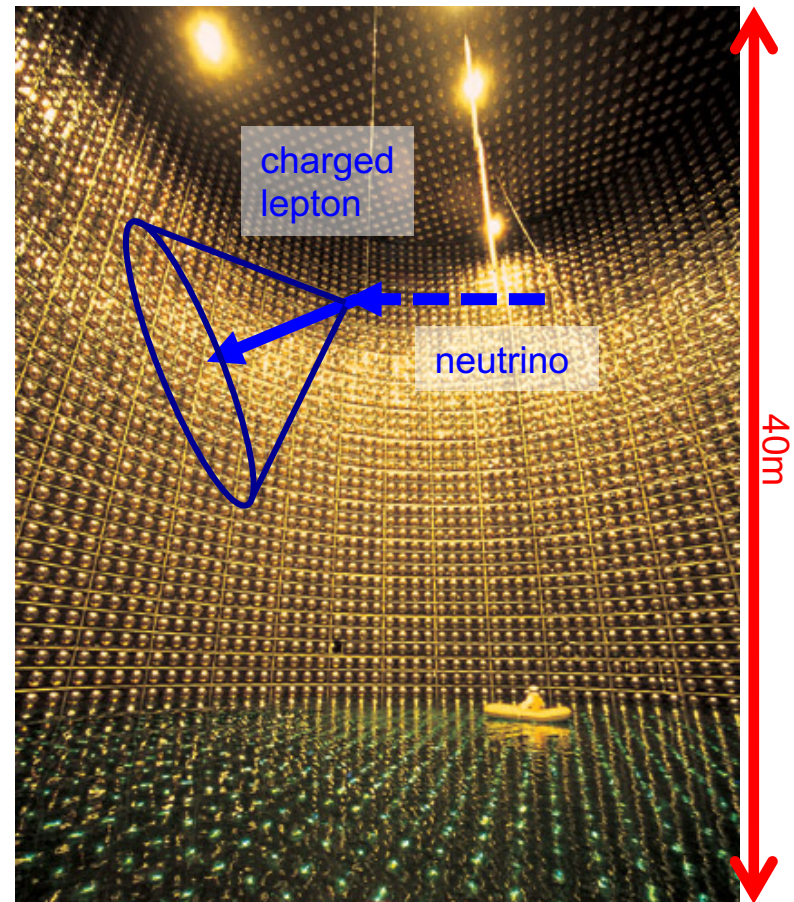
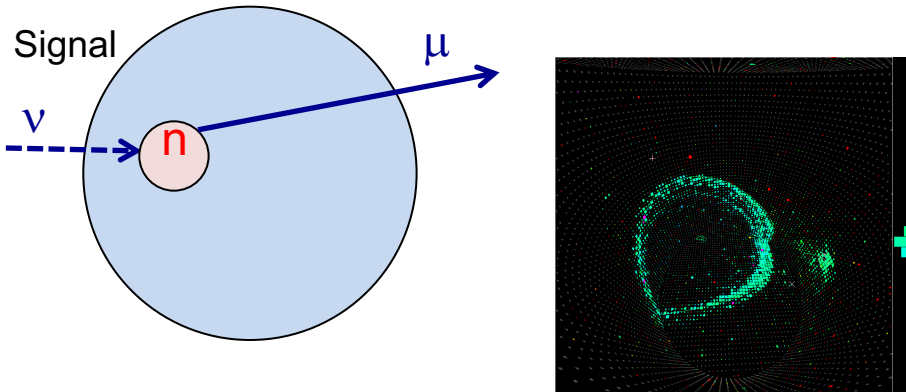
Axial 2-body current in QE region may be a tip of the iceberg..., or maybe a tip of gozilla!



3. non-QE background (resonance pion production)

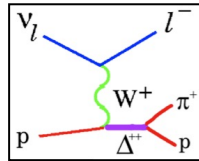


non-QE background \rightarrow shift spectrum



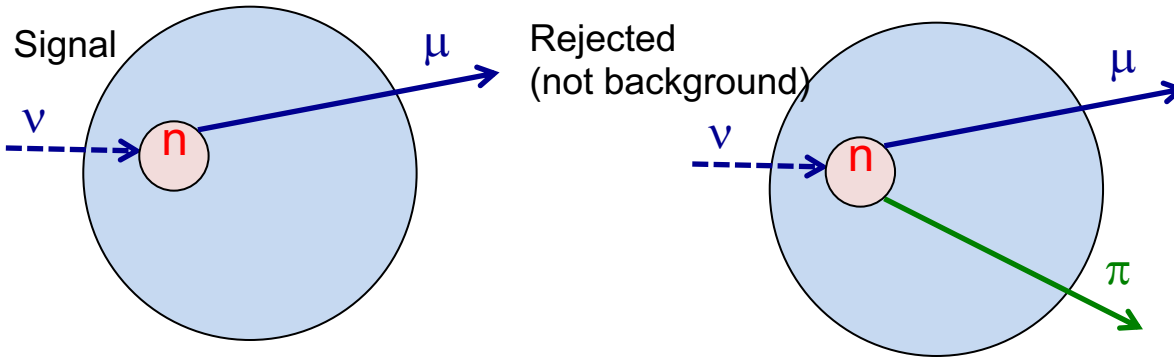
Typical neutrino oscillation detector

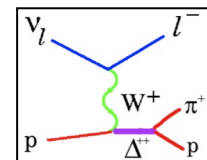
- Big and dense, to maximize interaction rate
- Coarsely instrumented, to minimize cost (not great detector to measure hadrons)



3. non-QE background (resonance pion production)

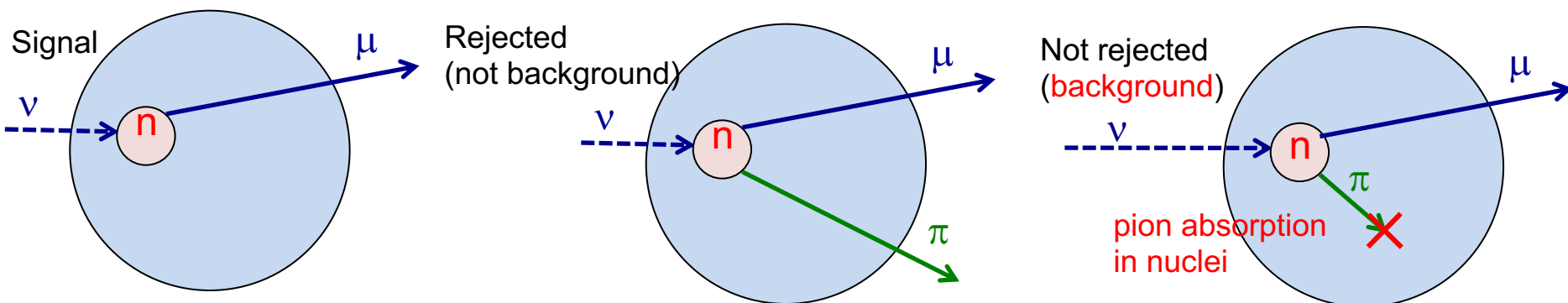
non-QE background \rightarrow shift spectrum

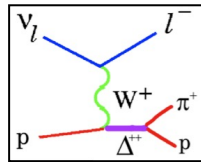




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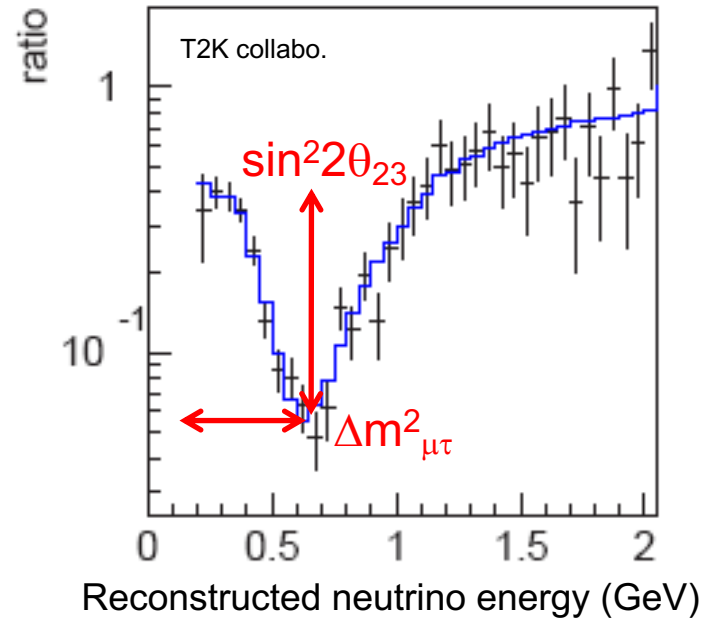
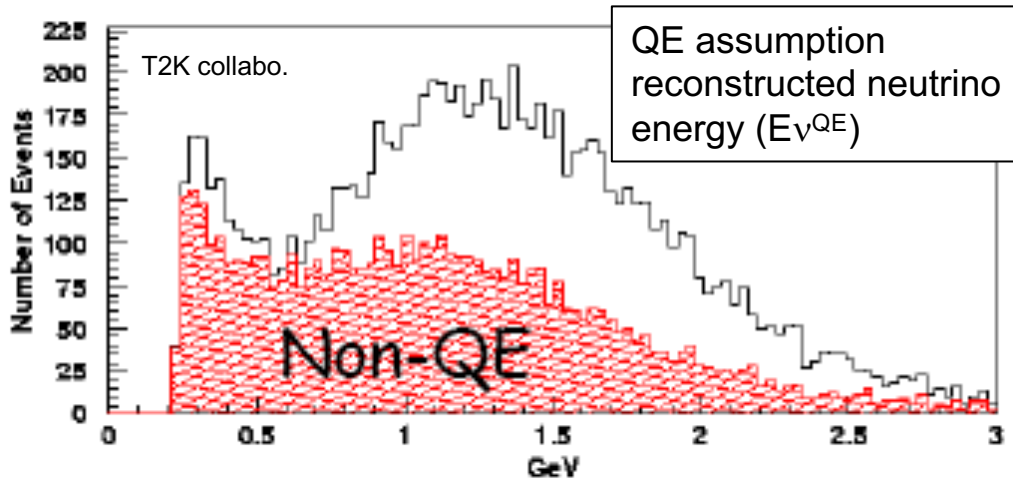
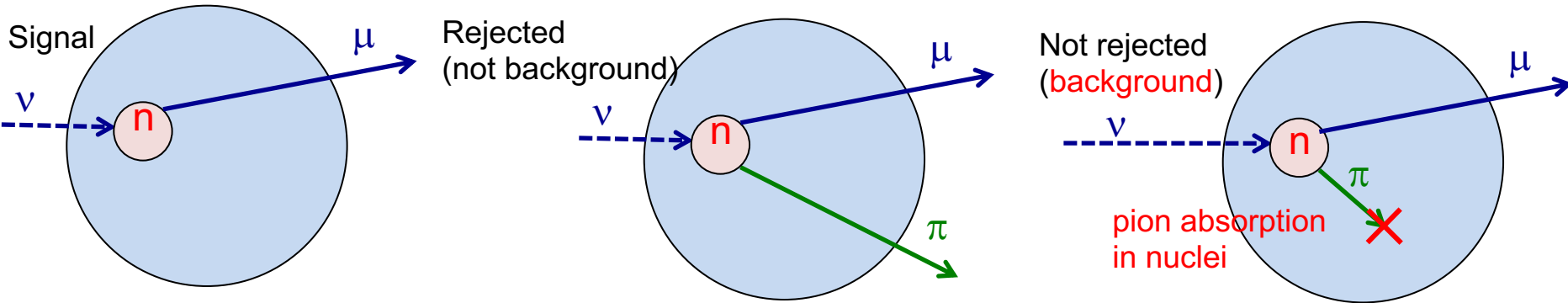
non-QE background \rightarrow shift spectrum



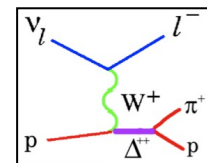


3. non-QE background (resonance pion production)

non-QE background → shift spectrum



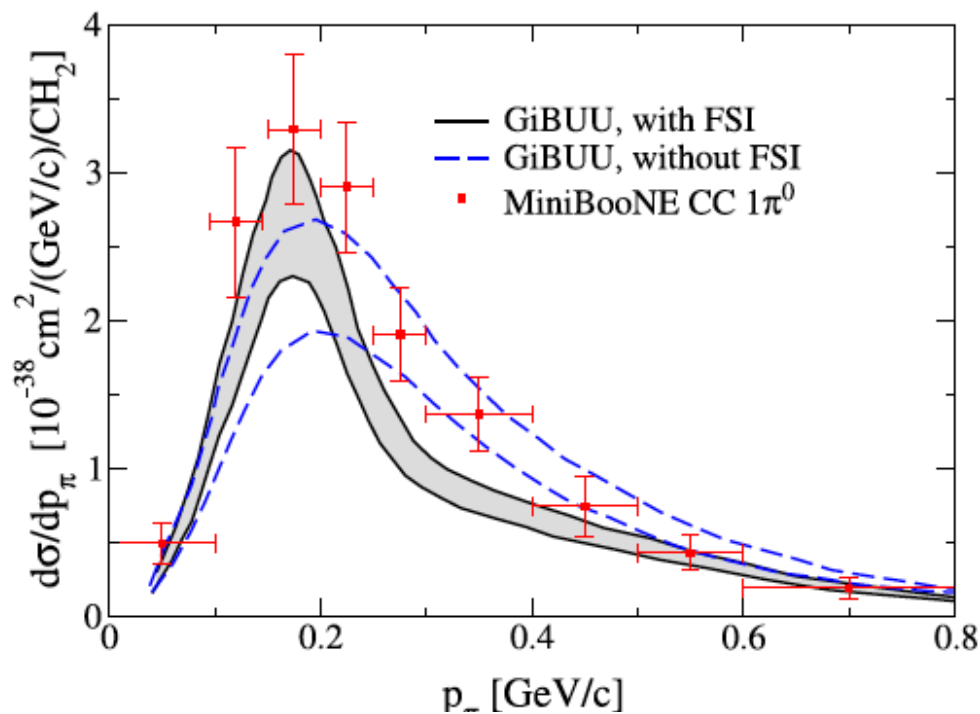
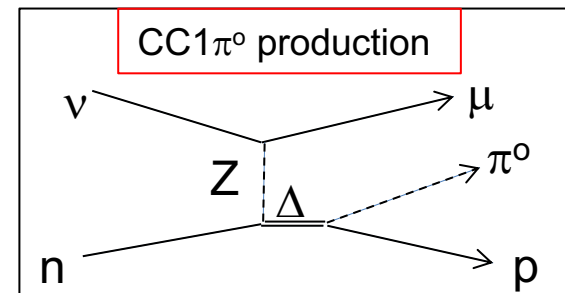
Solution: you need a good prediction of neutrino-induced pion production (hard)



3. Pion puzzle

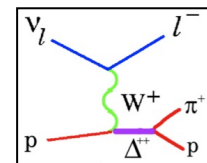
Final state interaction

- Cascade model as a standard of the community
- Advanced models are not available for event-by-event simulation



ex) Giessen BUU transport model

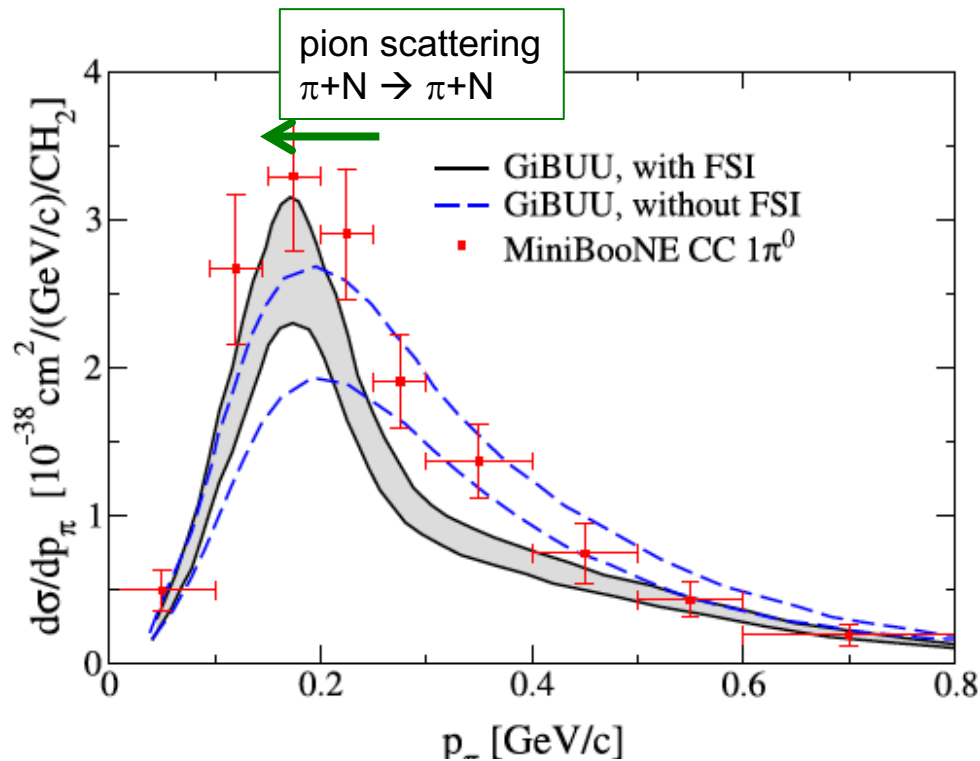
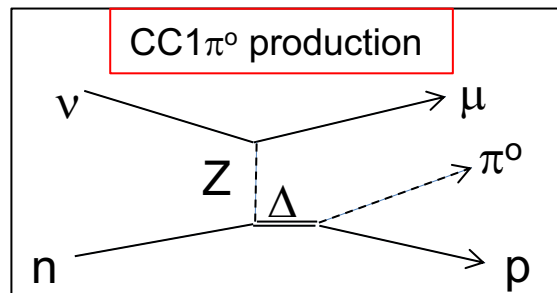
- Developed for heavy ion collision, and now used to calculate final state interactions of pions in nuclear media



3. Pion puzzle

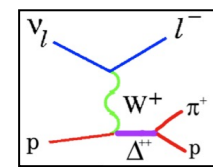
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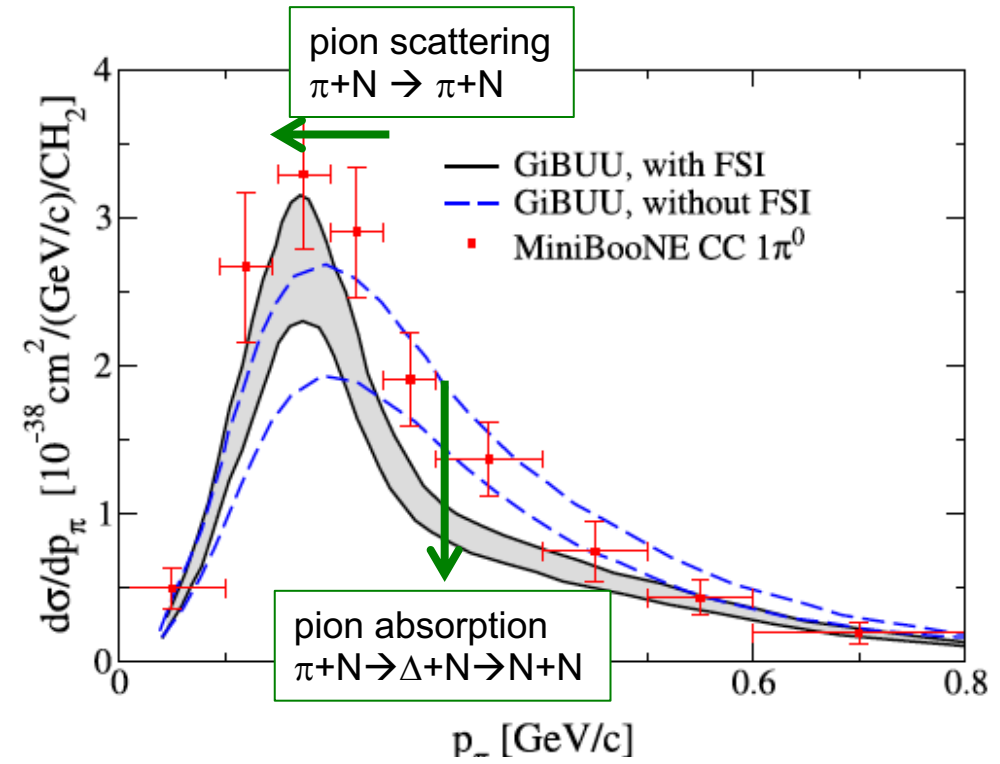
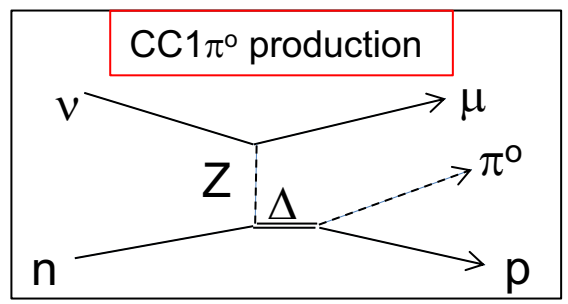
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3. Pion puzzle

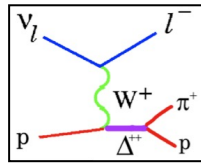
Final state interaction

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ex) Giessen BUU transport model

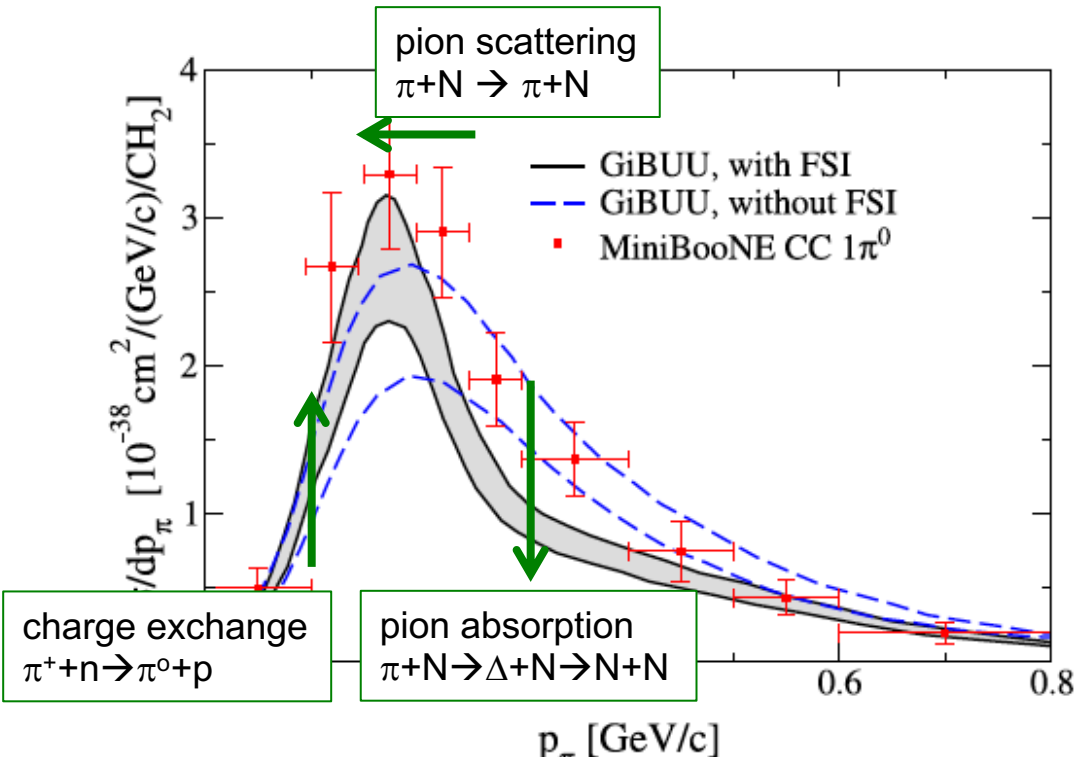
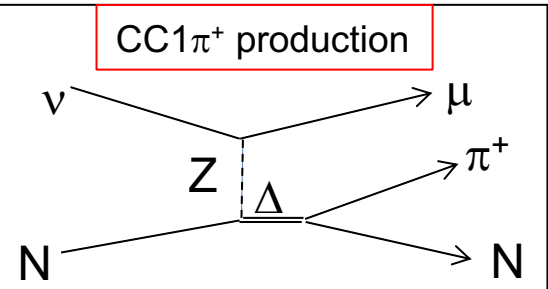
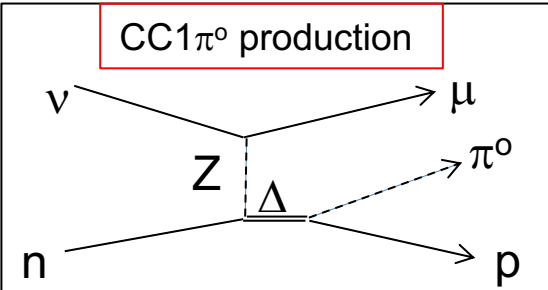
- Developed for heavy ion collision, and now used to calculate final state interactions of pions in nuclear media



3. Pion puzzle

Final state interaction

- Cascade model as a standard of the community
- Advanced models are not available for event-by-event simulation

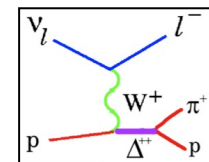


ex) Giessen BUU transport model

- Developed for heavy ion collision, and now used to calculate final state interactions of pions in nuclear media

You need to predict both
 1. all pion production channels
 2. all final state interaction

3. Pion puzzle

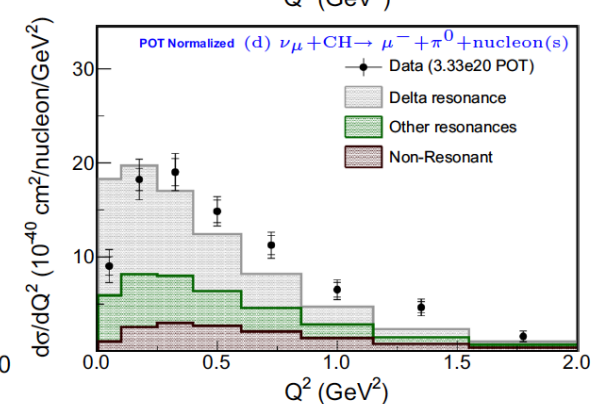
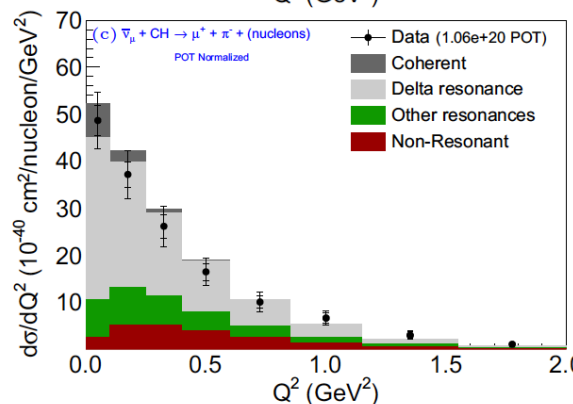
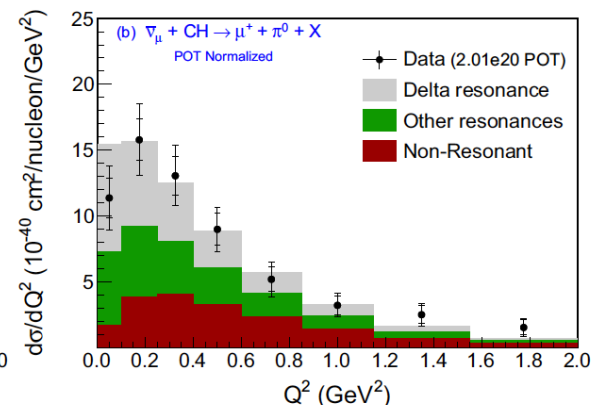
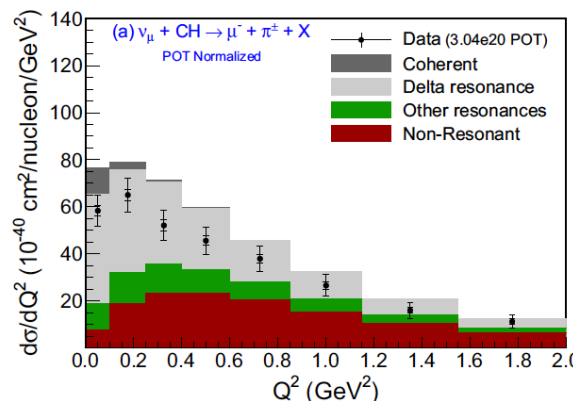


Data tension – internal: MINERvA pion data

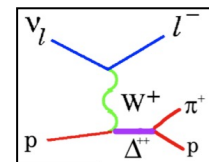
- It is extremely difficult to tune pion and/or FSI parameters to fit all pion data
- $\nu_\mu CC\pi^\pm$, low Q^2 suppression, over-predicted
- $\nu_\mu CC\pi^0$, strong low Q^2 suppression
- $\bar{\nu}_\mu CC\pi^-$, no low Q^2 suppression
- $\bar{\nu}_\mu CC\pi^0$, low Q^2 suppression, under-predicted

The study relies of available knobs in the generator

It looks the simulation doesn't have good knobs to tune



- You need to predict both
1. all pion production channels
 2. all final state interaction

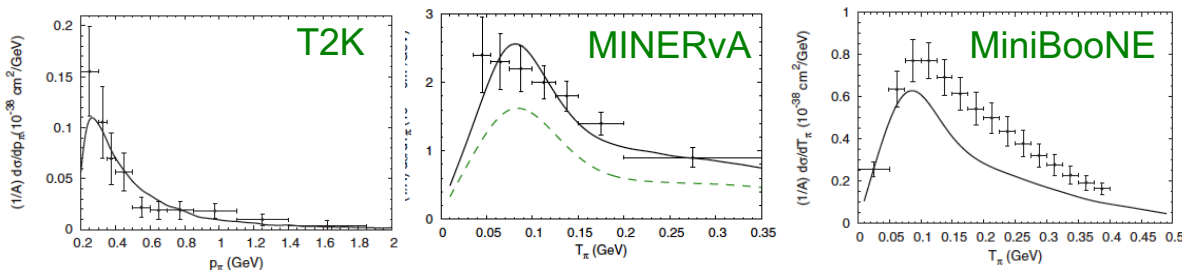


3. Data tension

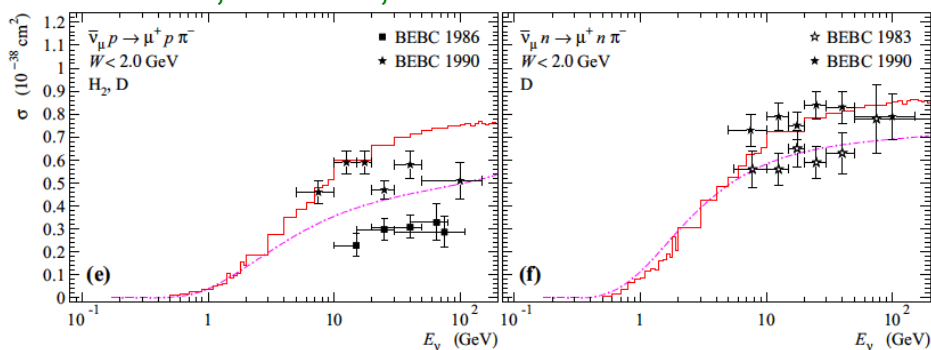
Data tension – external

- Tension between different experiments
- Tension between different targets
- Tension between different analyses

MiniBooNE vs. T2K vs. MINERvA (GiBUU)

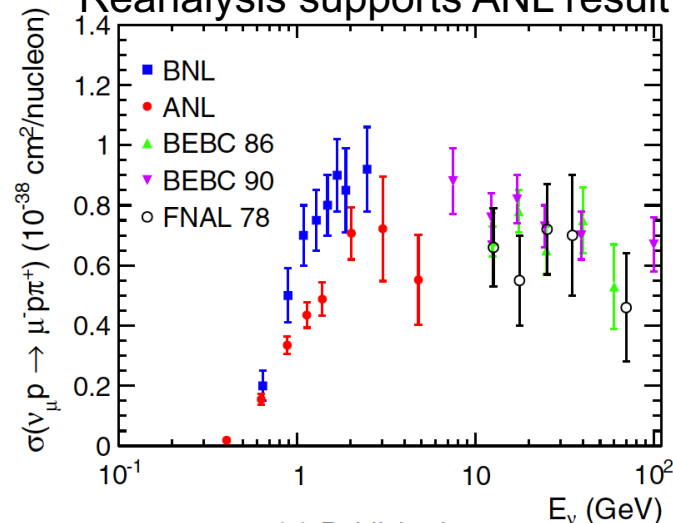


BEBC, H vs. D, 1983 vs. 1986 vs. 1990

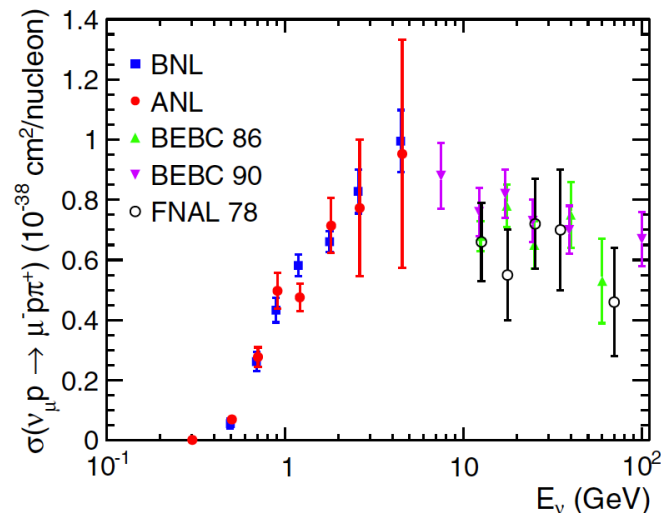


ANL vs. BNL data

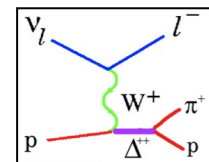
Reanalysis supports ANL result



(a) Published



(b) This analysis



3. MiniBooNE low-energy excess

Short-baseline anomalies

- Collection of data suggesting existence of sterile neutrino in 1eV scale
- MiniBooNE has the single highest significance in all anomalies
- Resonance related backgrounds (NC π^0 , NC γ) are dominant at low-E ν_e candidate
- Important for future ν_e appearance oscillation experiments, HyperK and DUNE)

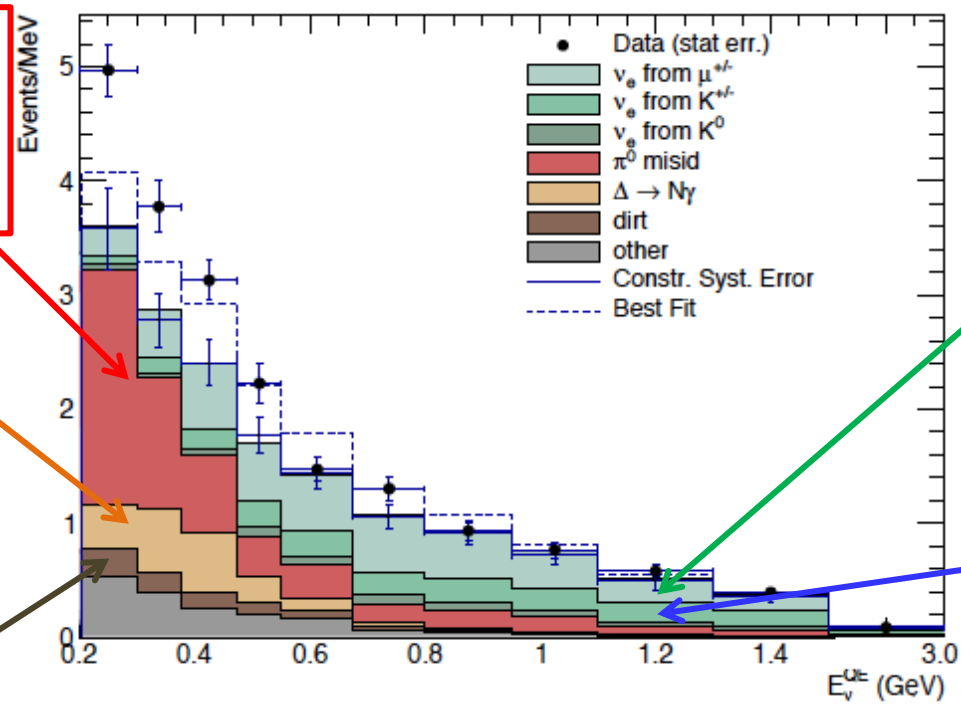
Asymmetric π^0 decay is constrained from measured NC π^0 rate ($\pi^0 \rightarrow \gamma$)

Δ resonance rate is constrained from measured NC π^0 rate

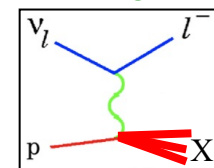
dirt rate is measured from dirt data sample

ν_e from μ decay is constrained from ν_μ CCQE measurement

ν_e from K decay is constrained from SciBooNE high energy ν_μ event measurement



1. Neutrino interaction physics - introduction
2. Charged-Current Quasi-Elastic (CCQE) interaction
3. Neutrino baryonic resonance interaction
- 4. Neutrino shallow- and deep-inelastic scatterings**
5. Conclusions



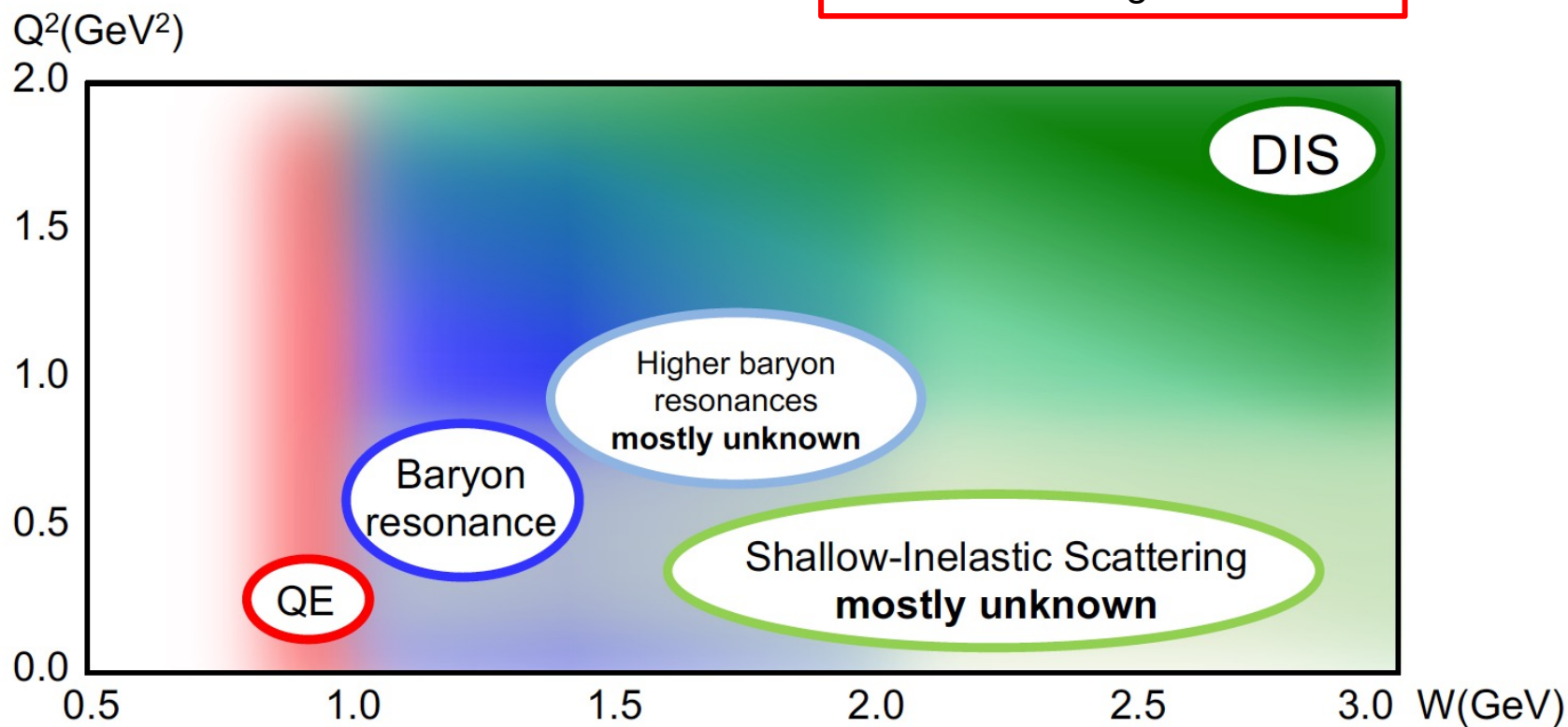
4. Shallow Inelastic Scattering (SIS)

Cross section

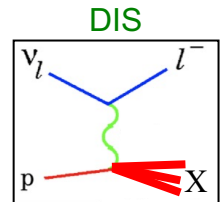
- Higher resonances and hadron dynamics
- Quark-Hadron duality (low Q^2 , low W DIS)
- Nuclear dependent DIS

Neutrino experiments around 1-10 GeV are not quite DIS yet

- Shallow \rightarrow low Q^2
- Inelastic \rightarrow large W

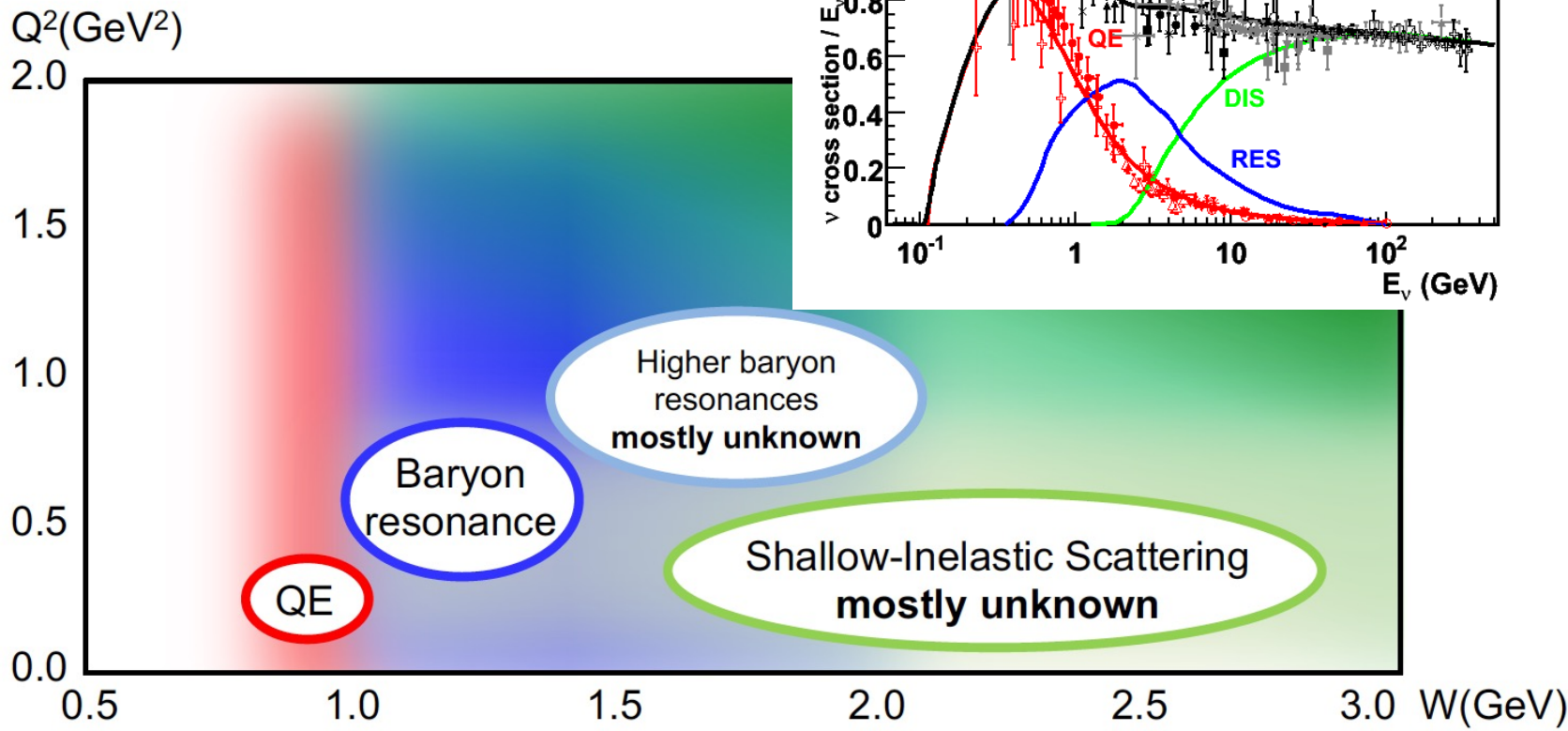


4. Shallow Inelastic Scattering (SIS)



Cross section

- Higher resonances and hadron dynamics
- Quark-Hadron duality (low Q^2 , low W DIS)
- Nuclear dependent DIS



4. Higher baryonic resonances

Cross section

- Higher resonances and hadron dynamics
- Quark-Hadron duality (low Q^2 , low W DIS)
- Nuclear dependent DIS

DCC model

- Channels are coupled (πN , $\pi\pi N$, etc), total amplitude is conserved
- Most of axial form factors are unknown

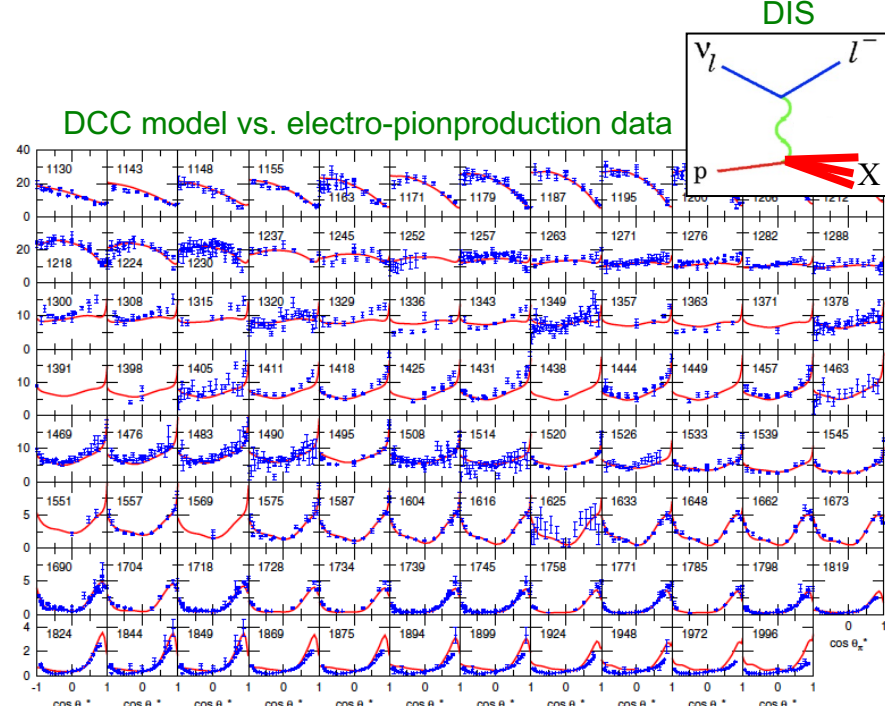
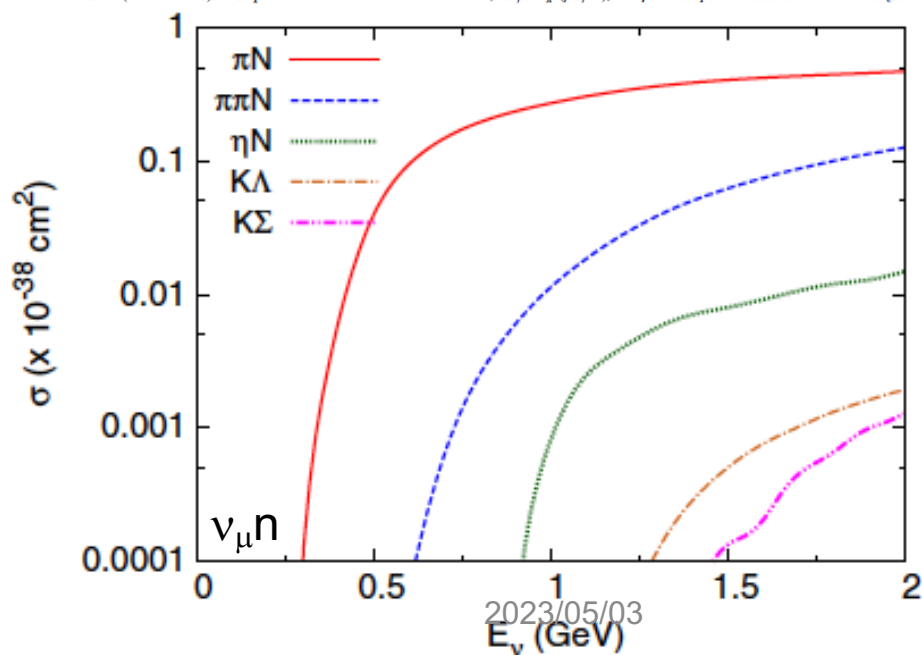
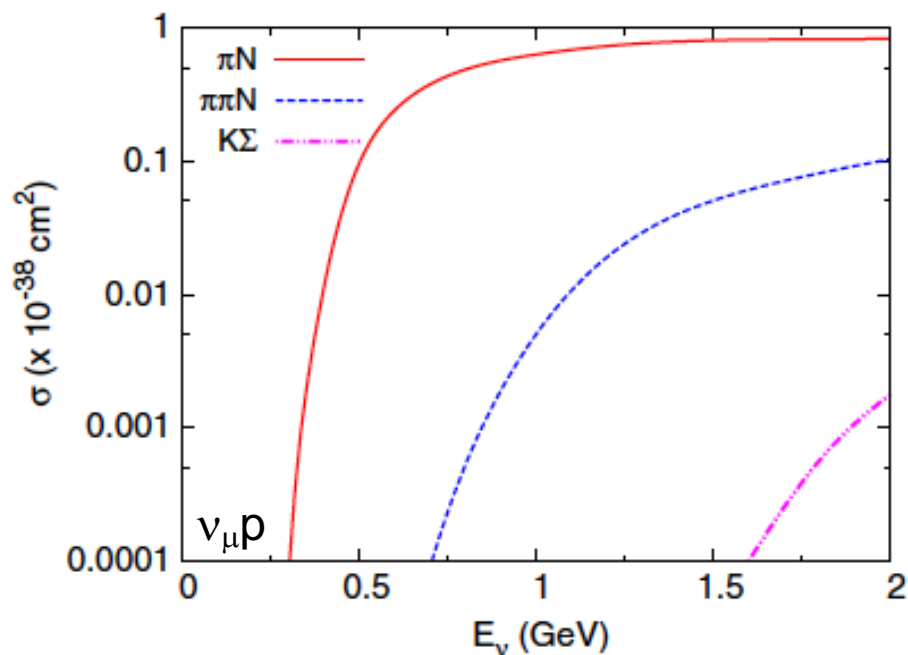
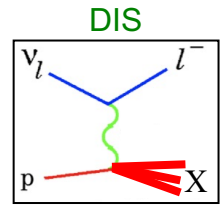


FIG. 8 (color online). Unpolarized differential cross sections, $d\sigma/d\Omega_*^*$ ($\mu\text{b}/\text{sr}$), for $\gamma n \rightarrow \pi^- p$. The data are from Refs. [55–78].

4. Quark-Hadron duality

Nachtmann variable $\xi = \frac{2x}{\left(1 + \sqrt{1 + \frac{4x^2 M^2}{Q^2}}\right)}$



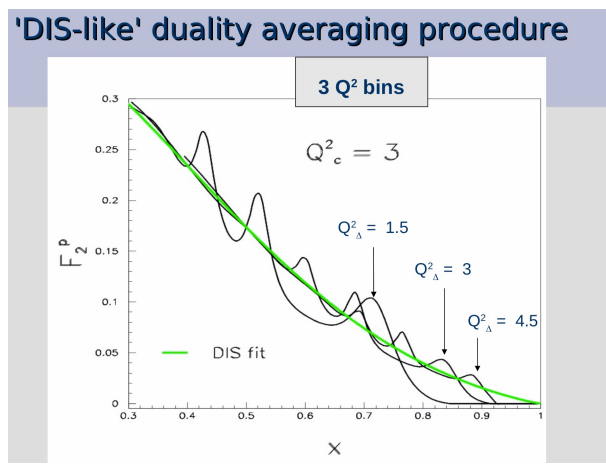
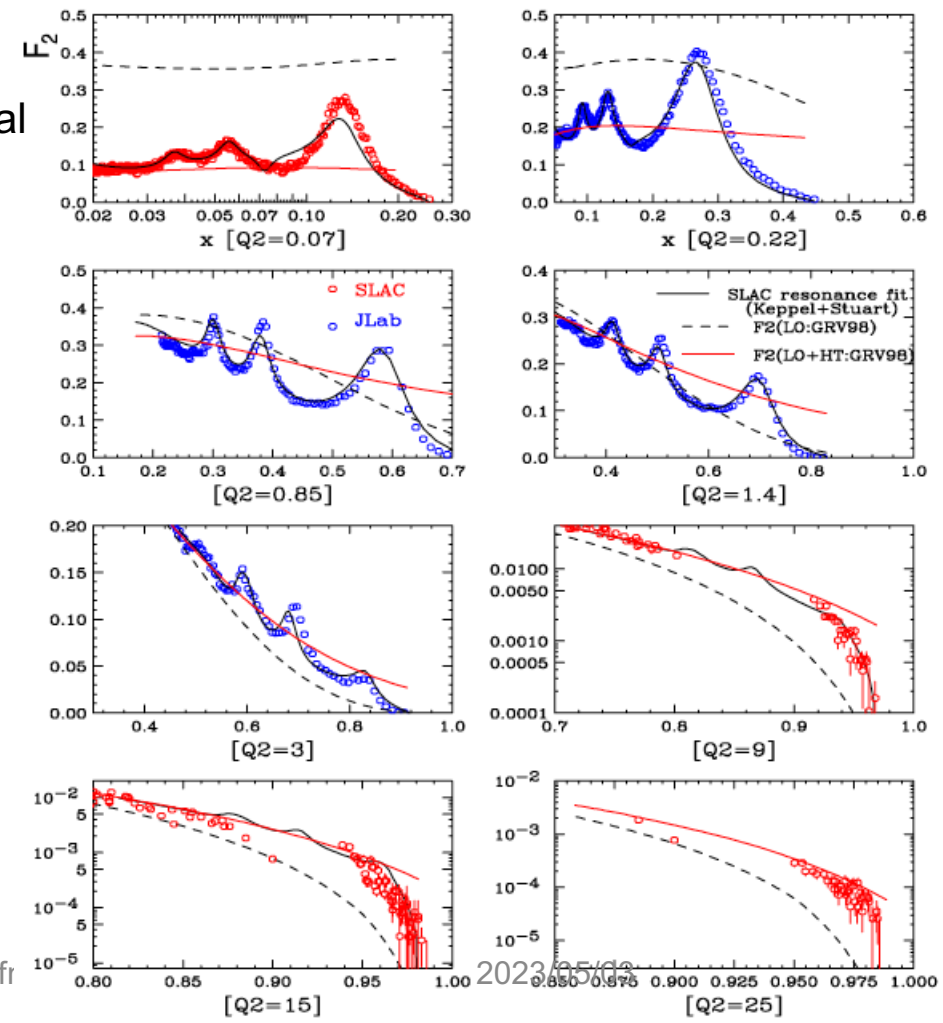
Cross section

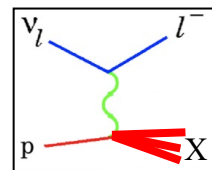
- Higher resonances and hadron dynamics
- Quark-Hadron duality (low Q^2 , low W DIS)
- Nuclear dependent DIS

Bodek-Yang correction is a phenomenological model to reproduce duality-like behavior, accepted by all neutrino simulation

DIS \neq Bjorken limit
 DIS = Q^2 average of all resonances

Proton F2 function GRV98-BY correction vs. data





4. Nuclear dependent DIS

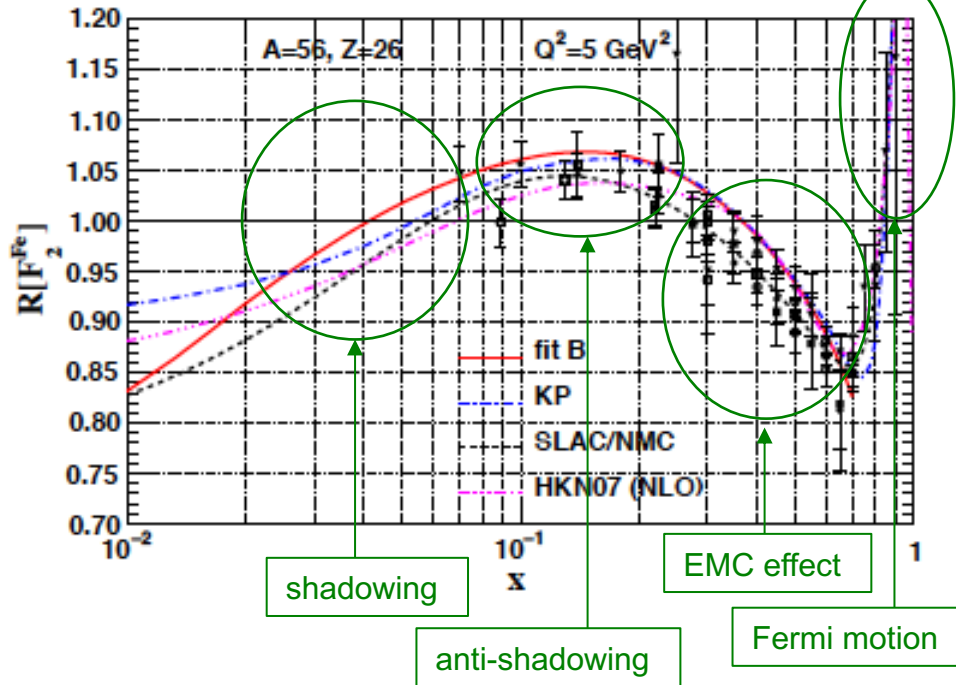
Cross section

- Higher resonances and hadron dynamics
- Quark-Hadron duality (low Q^2 , low W DIS)
- Nuclear dependent DIS

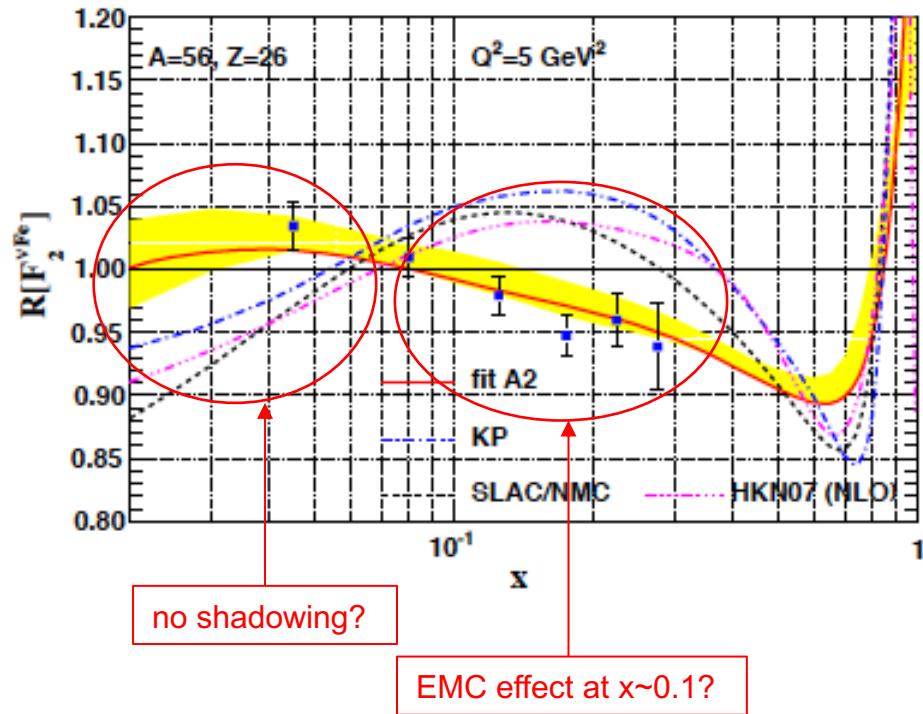
Nuclear PDF

- Shadowing, EMC effect, Fermi motion
- Likely due to nucleon dynamics in nucleus
- Various models describe charged lepton data
- Neutrino data look very different

e^+ -Fe nuclear correction factor



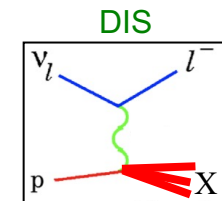
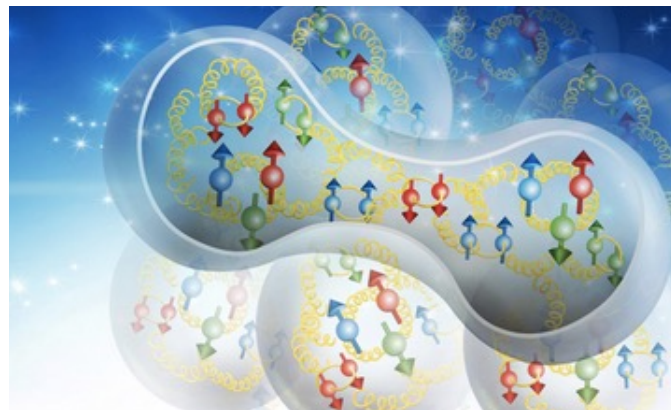
ν -Fe nuclear correction factor



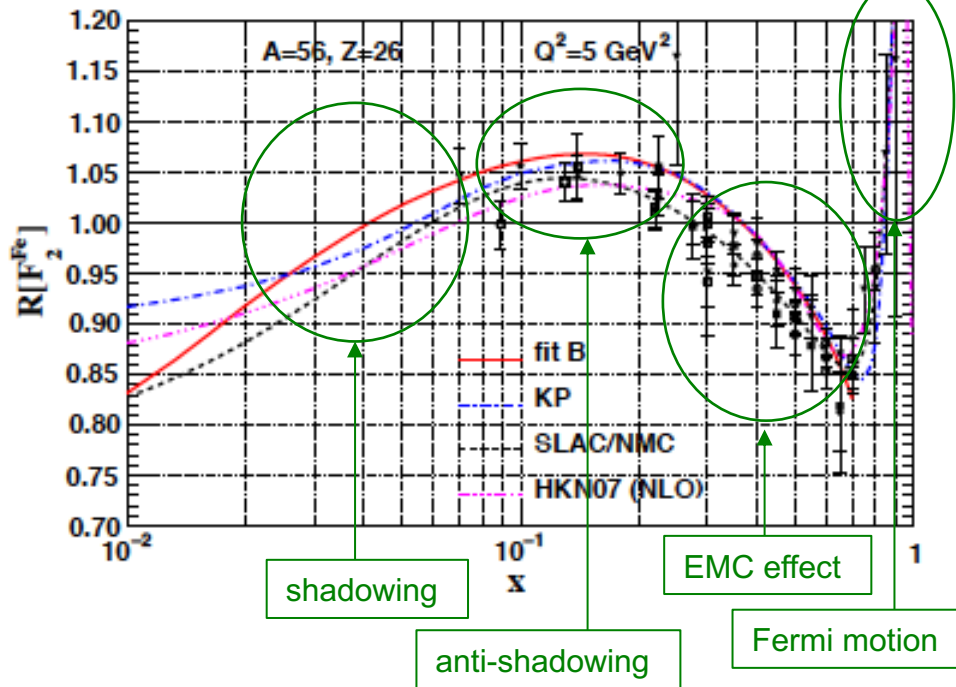
4. EMC effect

Nuclear dependent DIS

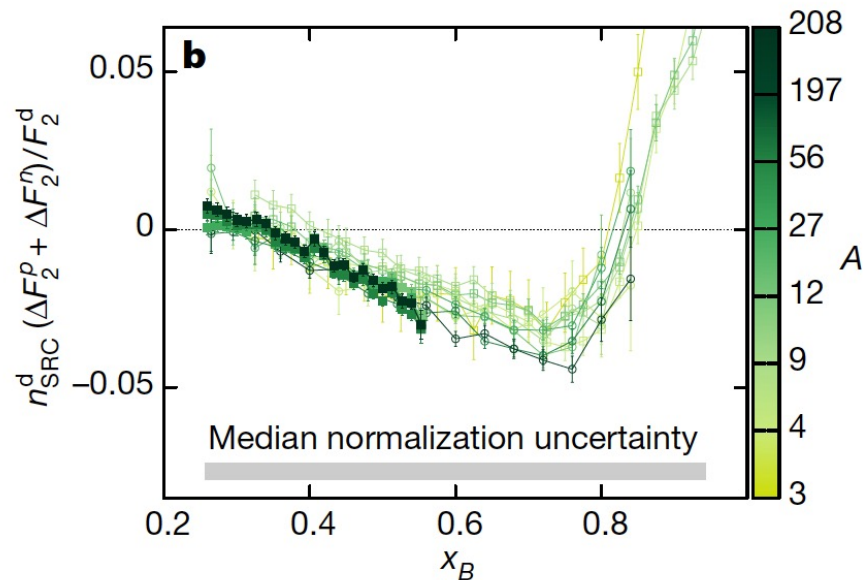
- First observed by the EMC experiment
- Structure function depends on nuclei
- Quarks feel presence of other quarks in other nucleons



e^\pm -Fe nuclear correction factor



EMC effect can be modeled from the amount of correlated pairs in nuclei (CLAS in JLab).



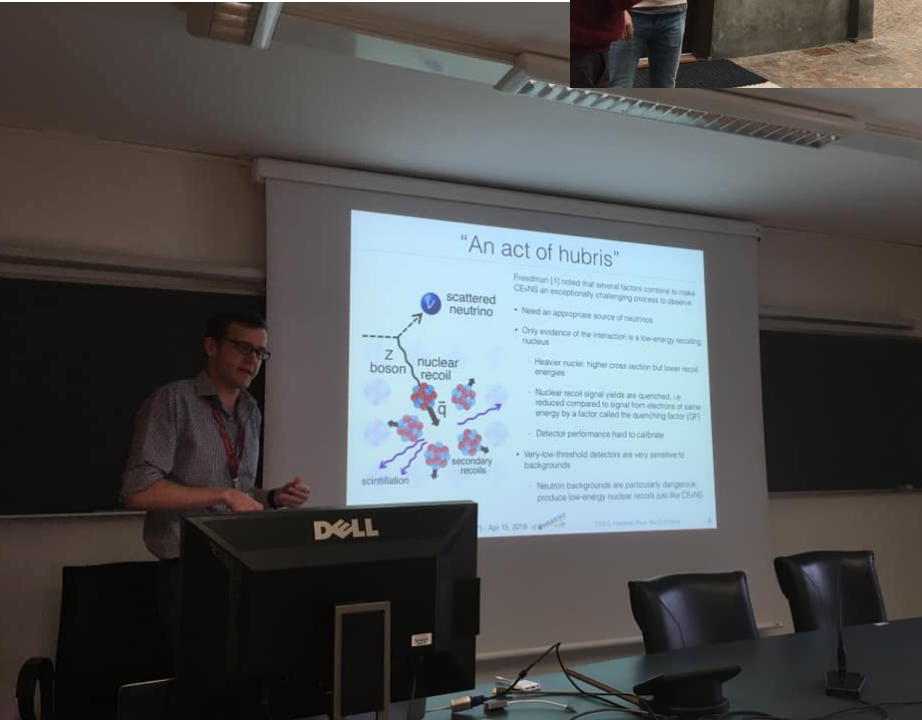
4. Atomic nuclei as laboratories for BSM physics

ECT* workshop, 15 Apr. 15-19 2019, Trento, Italy

<https://www.ectstar.eu/workshops/atomic-nuclei-as-laboratories-for-bsm-physics/>

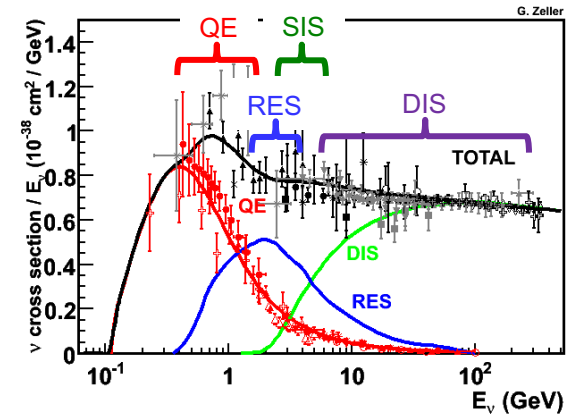
Topics include;

- Neutrino interactions
- EMC effect
- $0\nu\beta\beta$
- Direct dark matter
- etc



4. NuInt conference series

<https://nustec.fnal.gov/nuint-conference-series/>



The main conference in the neutrino interaction physics community

- Every ~18 months
- The next one will be Spring 2024 in São Paulo (Brazil)

NuINT 2022

The 13th International Workshop on Neutrino-Nucleus Interactions
in the Few GeV Regions

October 24 to 29, 2022 (OFFLINE)

Hoam Faculty House
Seoul National University
Seoul, Korea



4. NuSTEC

NuSTEC Workshop on Electron Scattering 2022 (Tel Aviv)
<https://indico.fnal.gov/event/50863/>

Neutrino Scattering Theory-Experiment Collaboration
<https://nustec.fnal.gov/>

NuSTEC CTGWG seminar and CEWG meetings
<https://indico.fnal.gov/category/990/>



Cross-experiment WG (CEWG)

Cross-theory-generator WG (CTGWG)

Workshops and Schools



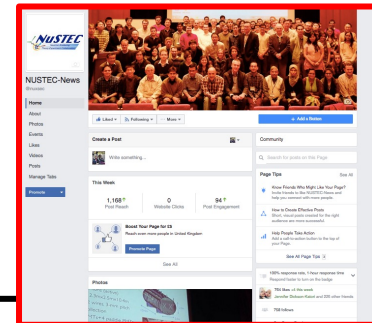
NuSTEC

Publications

Long-term community planning

Outreach

NuSTEC-News Facebook page
<https://www.facebook.com/nuxsec>



NuSTEC, PPNP100(2018)1
<https://www.sciencedirect.com/science/article/pii/S0146641018300061>



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Conclusion

1 to 10 GeV neutrino interaction measurements are crucial to successful next-generation neutrino oscillation experiments (DUNE, Hyper-K)

Nucleon correlation physics drastically change neutrino cross sections, both size and shape.

Recent new models and theories show nucleon correlation physics is important in many sub-fields of particle physics.

Neutrino interaction physics beyond QE region is confusing.

Future neutrino interaction measurements should focus on high-statistics neutrino hadron production measurements. This is the key to understand neutrino interaction models and nuclear effect.

Thank you for your attention!

Backup

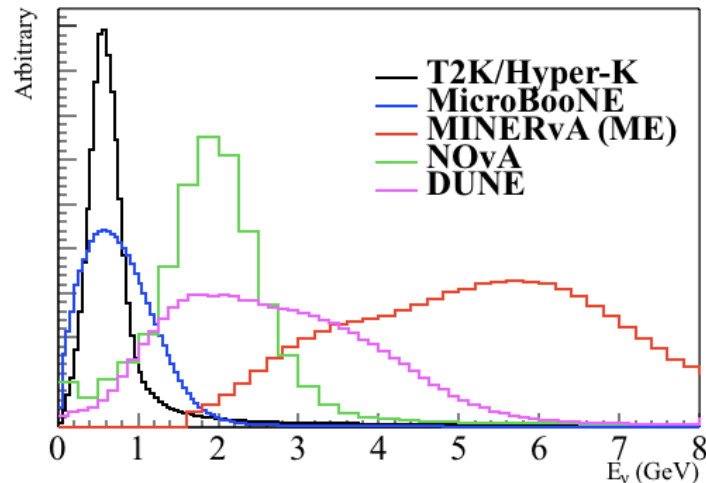
1. Next generation neutrino oscillation experiments

Neutrino oscillation experiments

- Past: K2K, MiniBooNE, MINOS, DeepCore
- Present to Future: T2K, NOvA, PINGU, ORCA, Hyper-Kamiokande, DUNE...

We don't know the energy of incoming neutrinos...

- We need to simulate all physics from $E_\nu=0$ to $E_\nu \sim \text{few GeV}$
- We need to simulate all physics from $\omega, |\vec{q}|=0$ to $\omega, |\vec{q}| \sim \text{few GeV}$



Two rules of neutrino interaction physics

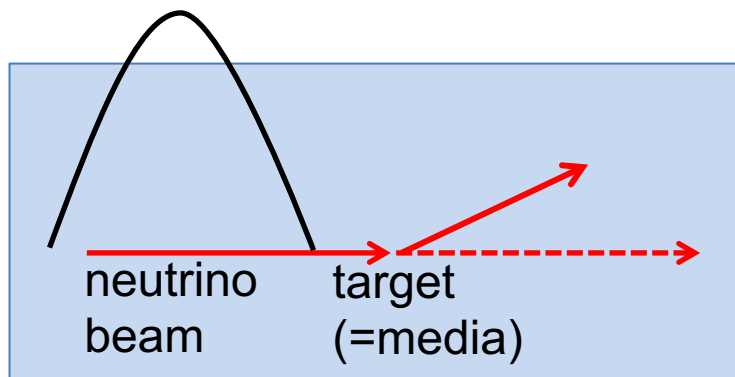
1. Neutrinos cannot choose kinematic

2. Neutrino kinematics are not fully determined

1. Typical neutrino detectors

Neutrino scattering

- Wideband beam
- observables are inclusive



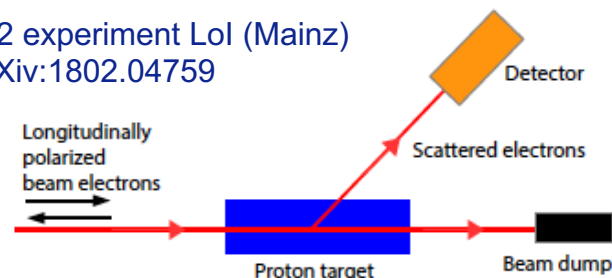
Incomplete kinematics

- Large mass, coarse instrumentation
- No one measures neutrino energy directly
- **Reconstructing kinematics (E_ν , Q^2 , W , x , y , ...)** in 1-10 GeV depends on interaction models

Electron scattering

- well defined energy, well known flux
- reconstruct energy-momentum transfer
- kinematics is completely fixed

P2 experiment Lol (Mainz)
arXiv:1802.04759



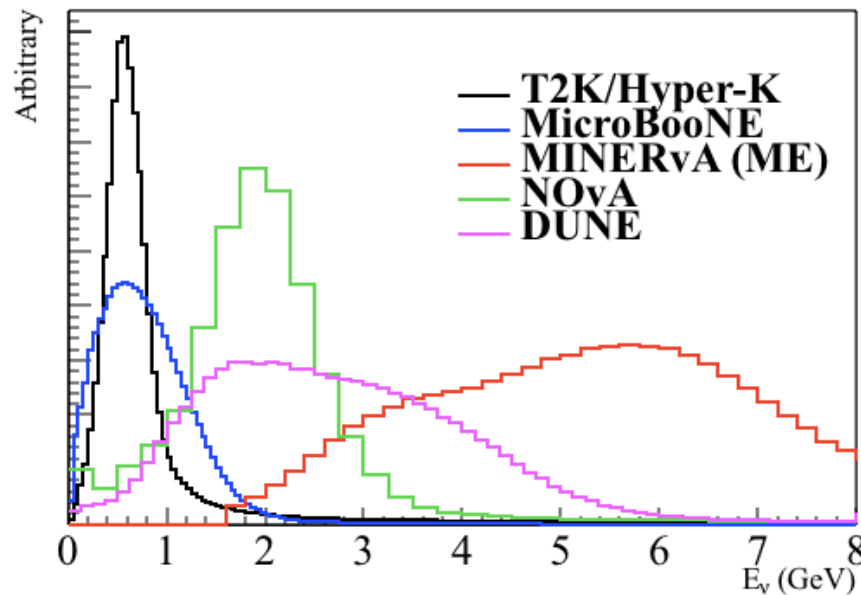
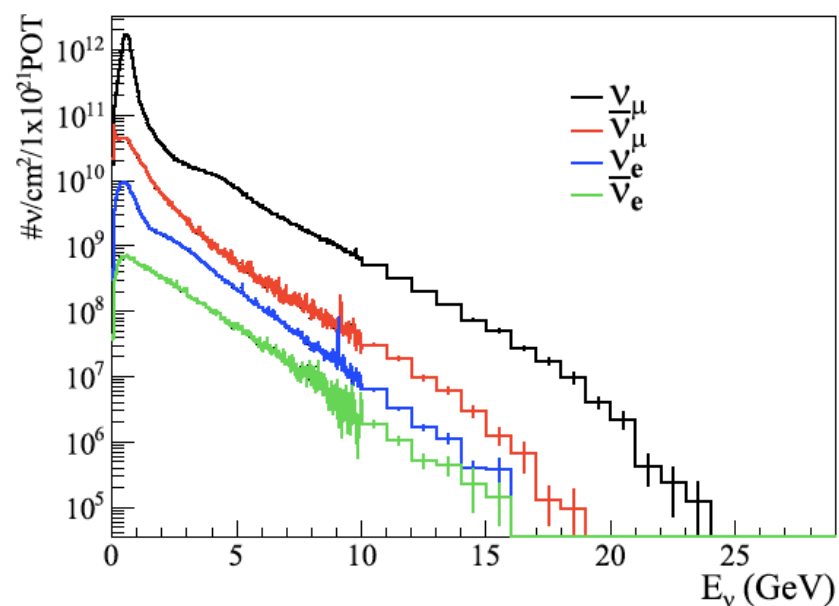
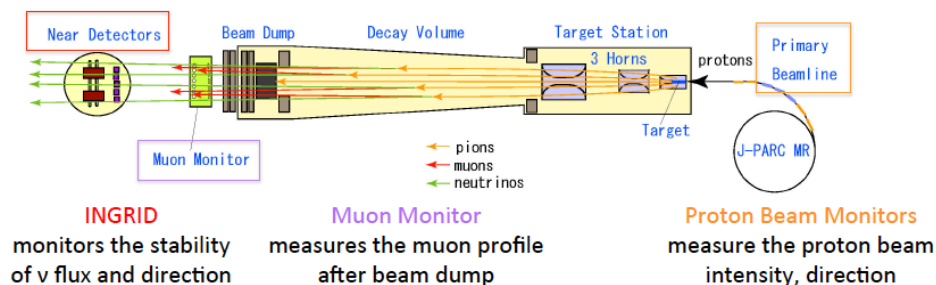
Two rules of neutrino interaction physics

1. Neutrinos cannot choose kinematic
2. **Neutrino kinematics are not fully determined**

1. Typical neutrino beams for oscillation experiments

e.g.) J-PARC neutrino beam (T2K)

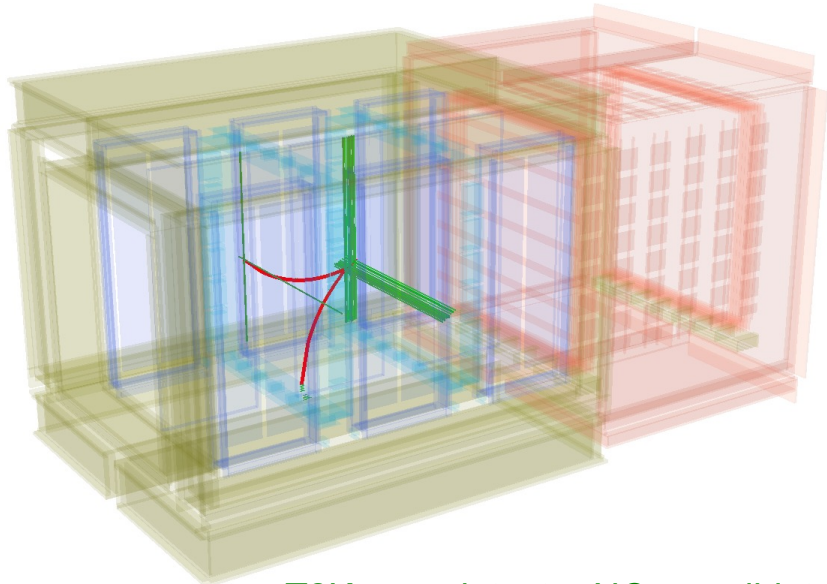
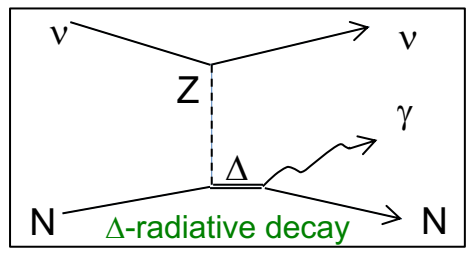
- pion decay-in-flight (high flux)
- off-axis beam (narrow band)
- but has components up to ~ 10 GeV
- typical beam 1-10 GeV
- $\sim 4\%$ normalization error (best case)



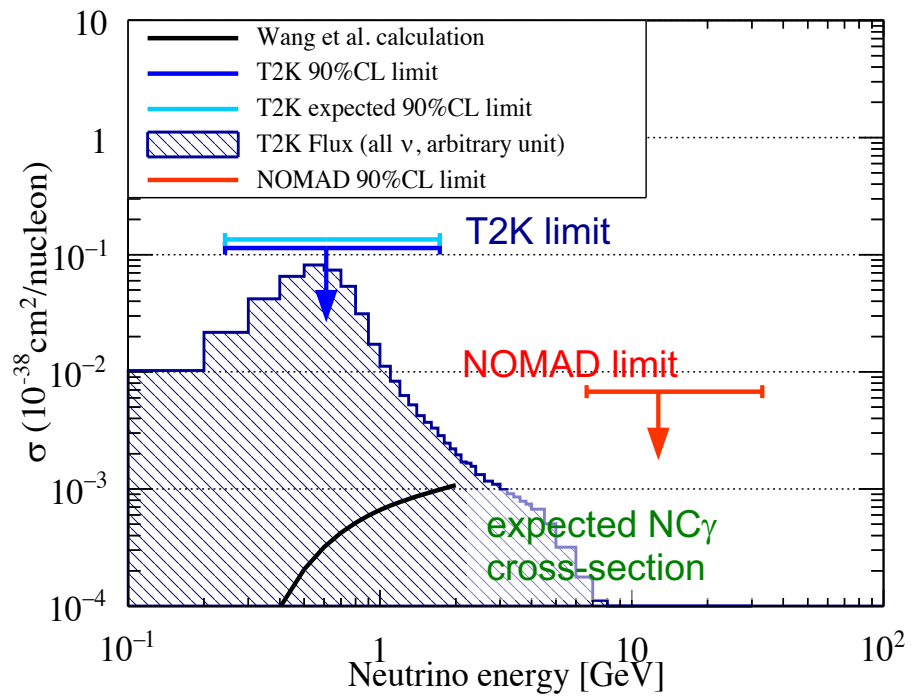
$$P_{\mu \rightarrow e}(L/E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (\text{eV}^2) \frac{L(\text{km})}{E(\text{GeV})} \right)$$

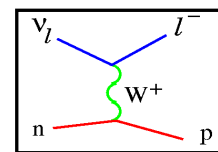
4. T2K Neutrino NC single photon production (NC_γ)

Neutrino induced NC single photon production (NC_γ) process is not experimentally identified. NC_γ is misID background for every electron-neutrino appearance oscillation experiment. T2K and NOMAD set limits on this process, but $\sim x3$ higher cross-section can explain all MiniBooNE excess.



T2K near detector NC_γ candidate

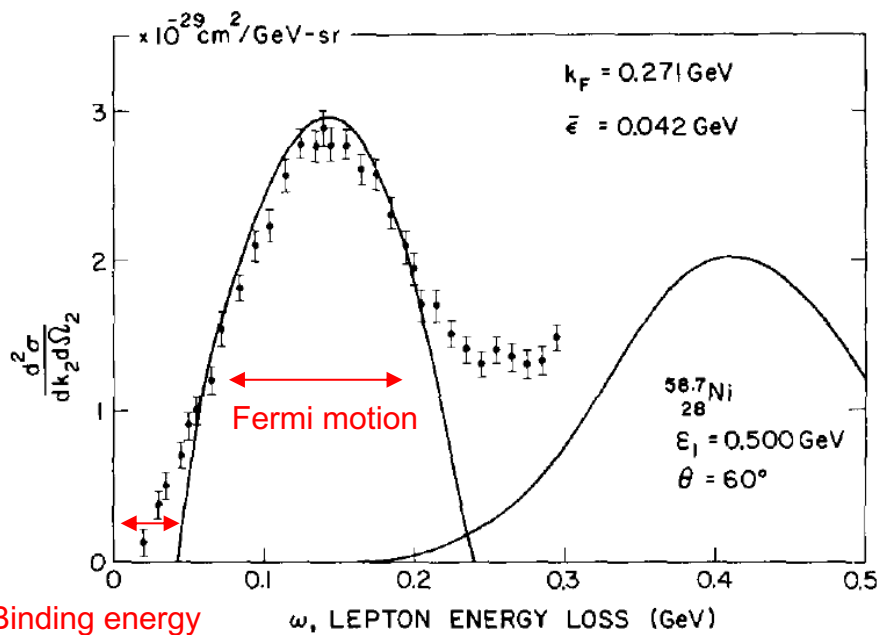




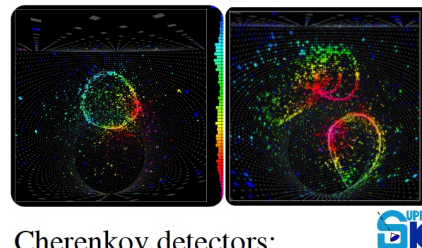
2. Fermi motion

Fermi motion

- Measured energy is smeared from the true energy if you assume nucleon at rest
- High resolution detector can measure all outgoing hadrons
 - initial nucleon momentum can be reconstructed (no Fermi motion smearing)

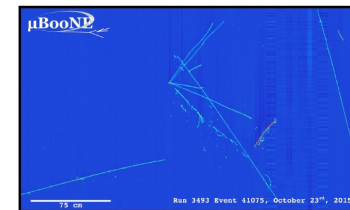


Binding energy



Cherenkov detectors:
 Assuming QE interaction
 Using lepton only

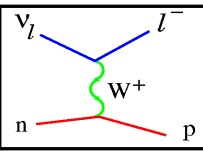
$$E_{QE} = \frac{2M\epsilon + 2ME_l - m_l^2}{2(M - E_l + |k_l| \cos \theta_l)}$$



Tracking detectors:
 Calorimetric sum
 Using All detected particles

$$E_{cal} = E_l + E_p^{kin} + \epsilon$$

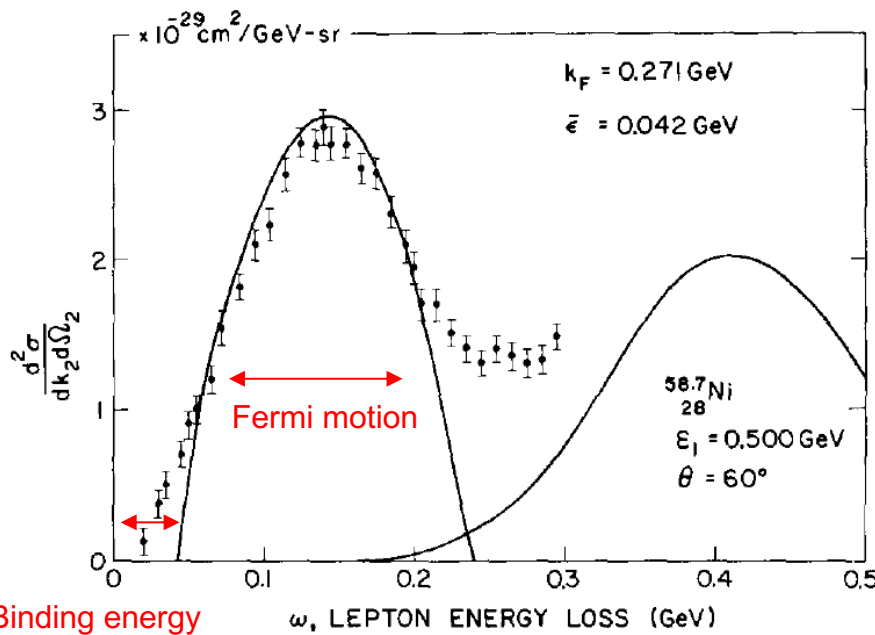
[1p0π]



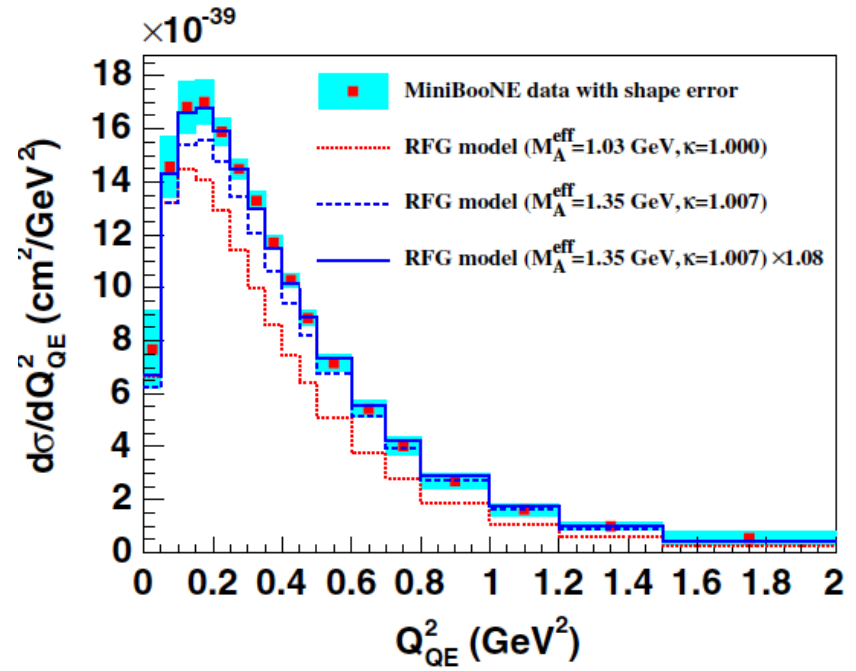
2. Pauli blocking

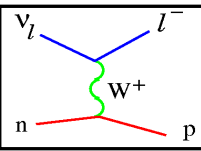
Pauli blocking

- Low momentum transfer reaction is forbidden.
- data show more suppression than what Pauli blocking can → RPA(?)
- In the global Fermi model, Pauli blocking looks unphysical



Binding energy

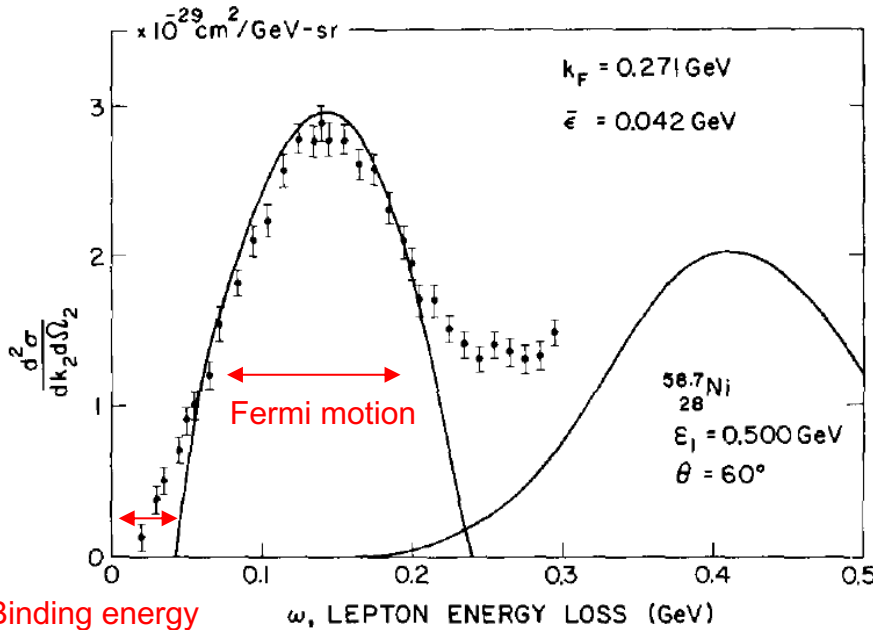




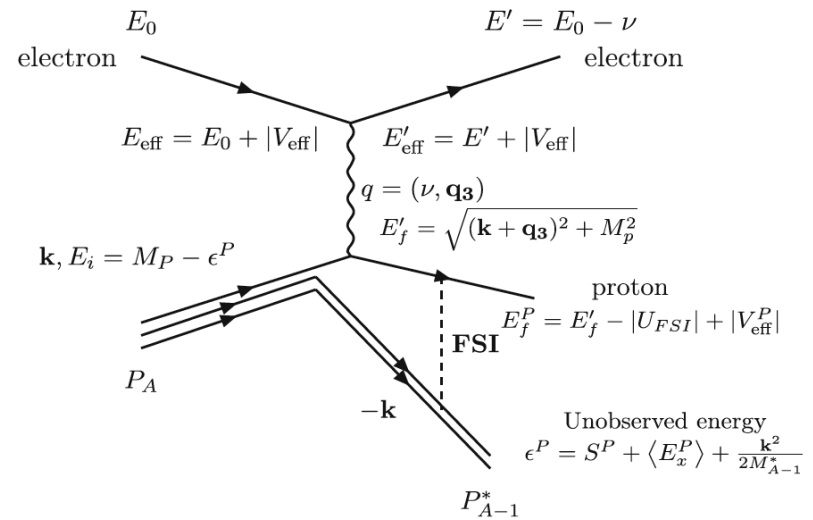
2. Nuclear Shell structure and binding energy

Binding energy ~ unobserved energy

- Energy to cost to release 1 nucleon, not constant
- Separation energy + excitation energy + recoil energy
 - Separation energy: energy to release 1 nucleon from the shell (~15 MeV, depends)
 - Excitation energy: energy used to excite leftover target nucleus (~1 MeV)
 - Recoil energy: kinetic energy of recoil target nucleus (~2-3 MeV)



Electron scattering on proton



2. Nucleon correlations in neutrino physics

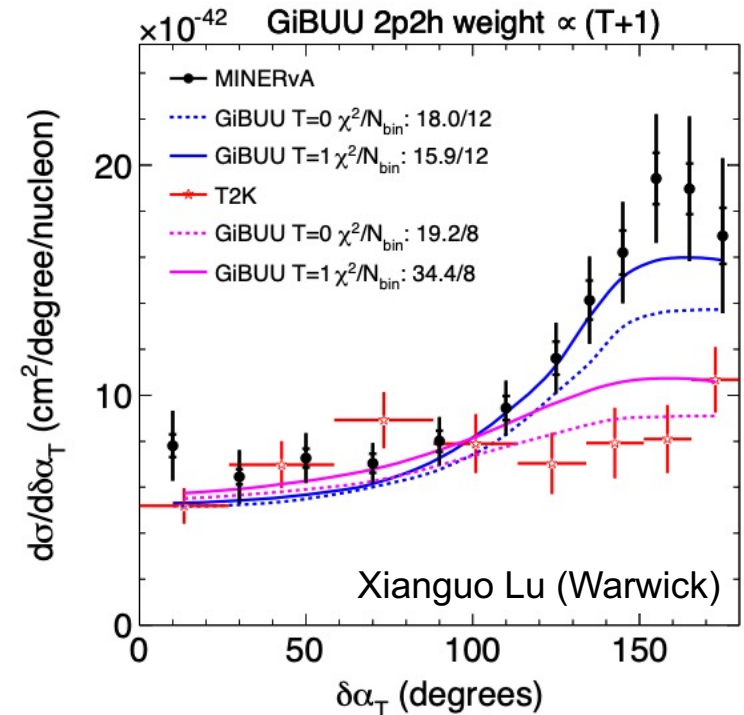
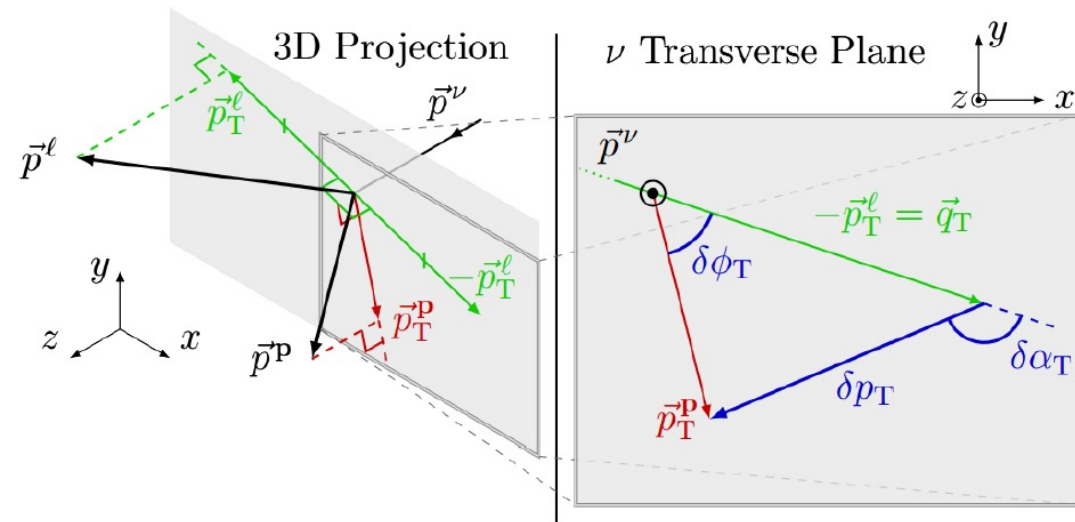
We want to constrain nuclear model from neutrino data

- Final state hadron measurement is the key

1 muon + 1 proton sample

- 5 dof (mu E and $\cos\theta$, proton E and $\cos\theta$, mu-p opening angle).
- Low statistics, and these are converted to 3 kinematic variables.

Data prefer advanced nuclear models, but it's not easy to identify which 2p2h model is right



Importance of axial 2BC is understood qualitatively, we need more quantitative understanding

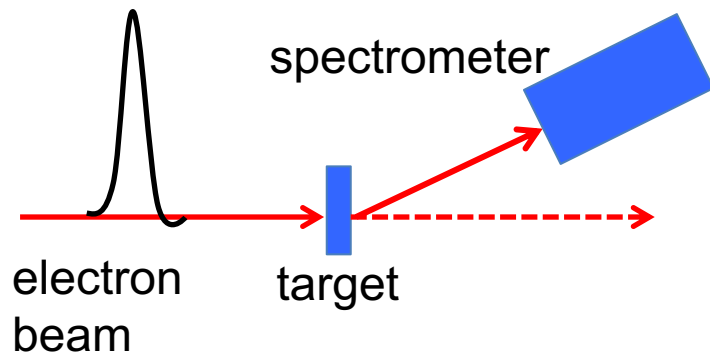
2. New paradigm of lepton scattering experiments

Flux-averaged differential cross-section

- Incomplete kinematics, reconstruction of E_ν , Q^2 , q^3 , W , x , y , ... depends on models

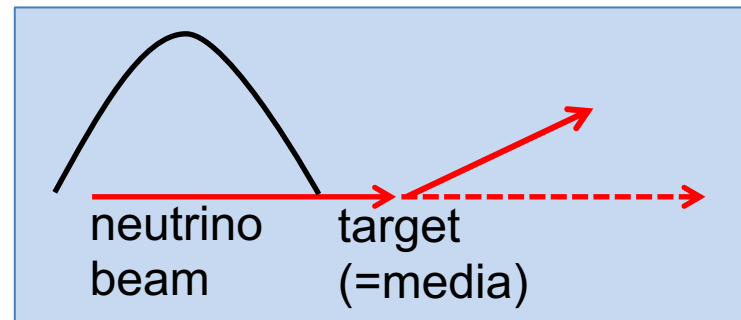
Electron scattering

- well defined energy, well known flux
- reconstruct energy-momentum transfer
- measure each process



Neutrino scattering

- Wideband beam (unknown E_ν)
- cannot fix kinematics
- inclusive measurement (CCQE, RES...)



2. New paradigm of lepton scattering experiments

Flux-averaged differential cross-section

- Incomplete kinematics, reconstruction of E_ν , Q^2 , q^3 , W , x , y , ... depends on models

Electron scattering

- well defined energy, well known flux
- reconstruct energy-momentum transfer
- measure each process

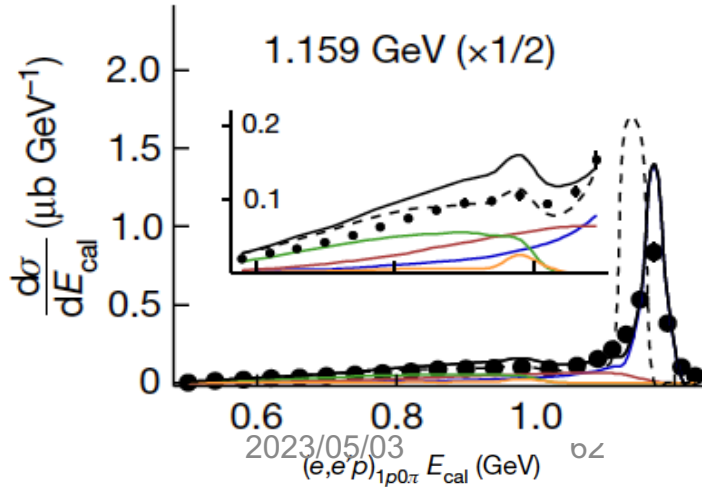
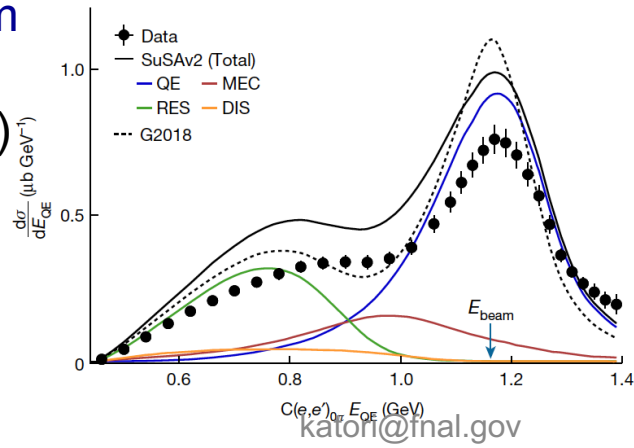
Neutrino experiment don't reconstruct E_ν (and Q^2) with great precision

Reconstructed beam electron energy spectrum by

- QE kinematic (HyperK)
- Calorimetric (DUNE)

$$E_\nu^{QE} = \frac{ME_\nu - 0.5m_\mu^2}{M - E_\mu + p_\mu \cos\theta}$$

$$E_\nu^{Cal} = E_\mu + \sum^{all} E_{had}^i$$



2. New paradigm of lepton scattering experiments

Flux-averaged differential cross-section

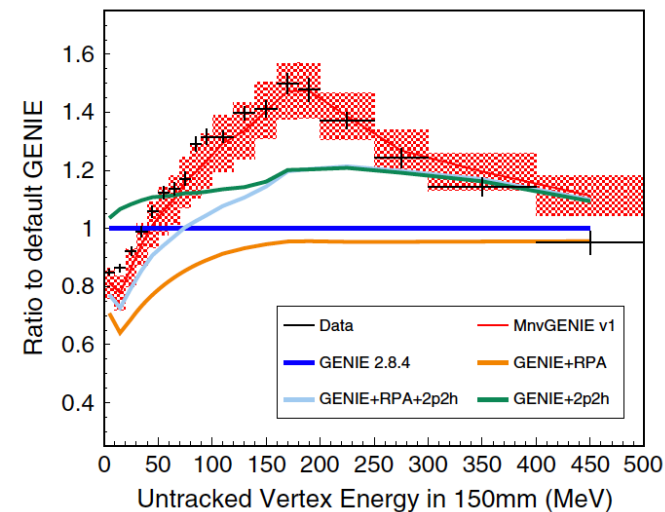
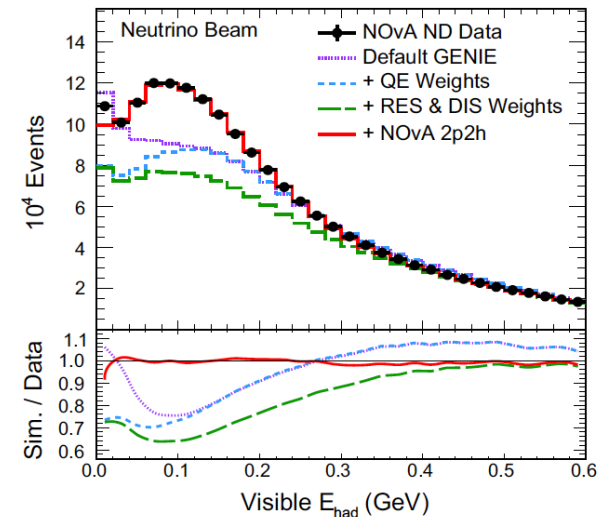
- Incomplete kinematics, reconstruction of E_ν , Q^2 , q_3 , W , x , y , ... depends on models
- New kinematic variables from hadrons

Visible hadronic energy deposit: E_{had} , E_{avail}

- Sum of all hadron energy deposit
- Strongly correlated to energy transfer (q_0 or ω or ν)
- Sensitive to 2p2h

Vertex activity

- Some of all hadronic activities around the vertex
- Low energy nucleons (=2 nucleon emission)



2. New paradigm of lepton scattering experiments

These studies suggest no nuclear models fit neutrino data without tuning

Flux-averaged differential cross-section

- Incomplete kinematics, reconstruction of E_ν , Q^2 , q_3 , W , x , y , ... depends on models
- New kinematic variables from hadrons

Visible hadronic energy deposit: E_{had} , E_{avail}

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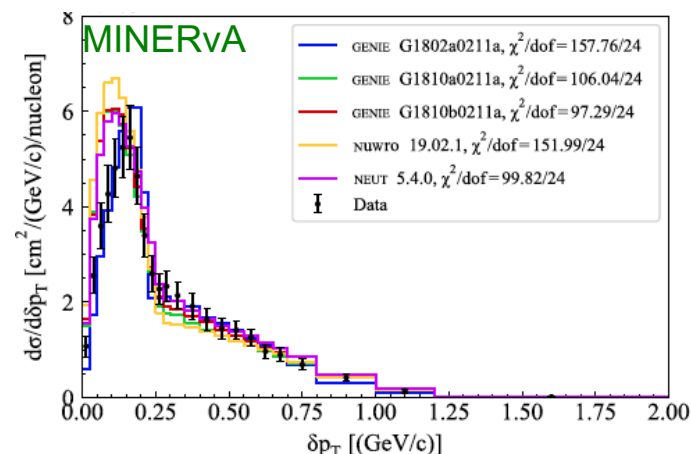
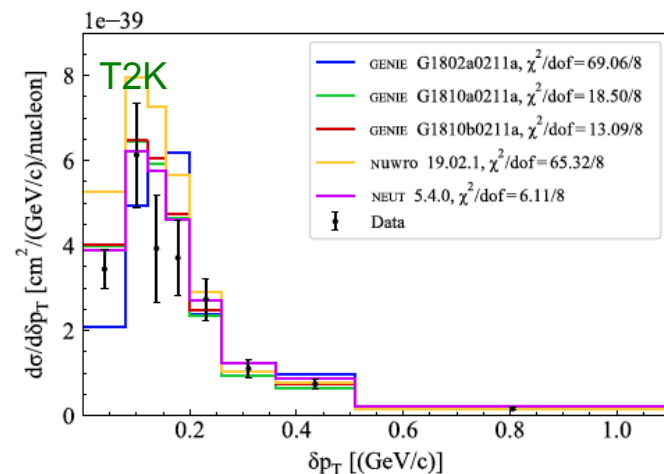
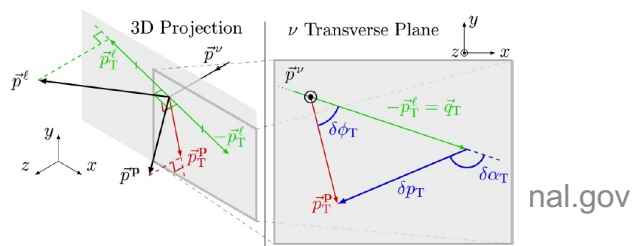
Vertex activity

- Some of all hadronic activities around the vertex
- Low energy nucleons (=2 nucleon emission)

Transverse kinematic Imbalance (TKI) variables

$\delta P_T \sim$ nucleon momentum distribution

$\delta \alpha_T \sim$ FSI



2. Generator implementation is our bottleneck

Flux-averaged differential cross-section

- Incomplete kinematics, reconstruction of E_ν , Q^2 , q^3 , W , x , y ,... depends on models
- New kinematic variables from hadrons

Hadron variables

- Visible hadronic energy deposit: E_{had} , E_{avail}
- Vertex activity
- Transverse kinematic Imbalance (TKI) variables

Hadrons are affected by FSIs

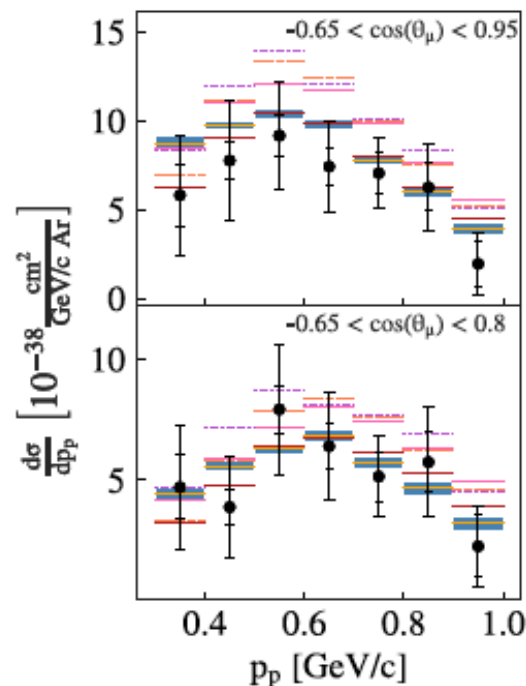
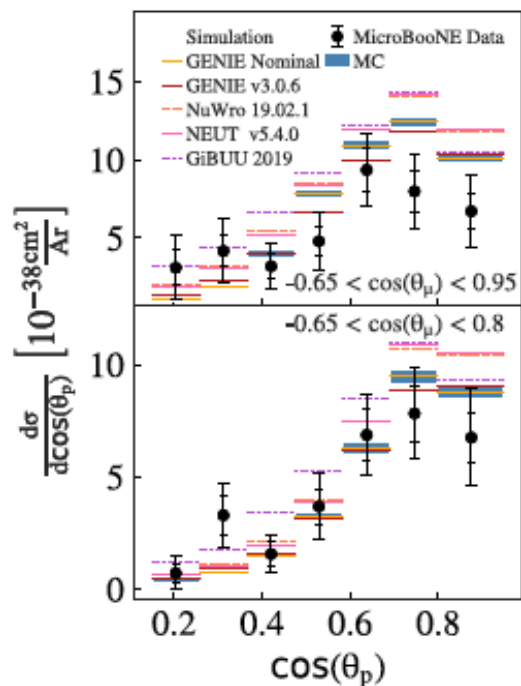
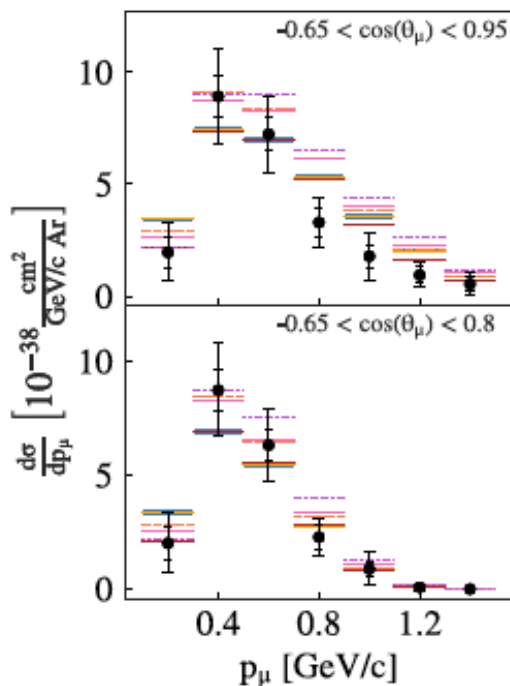
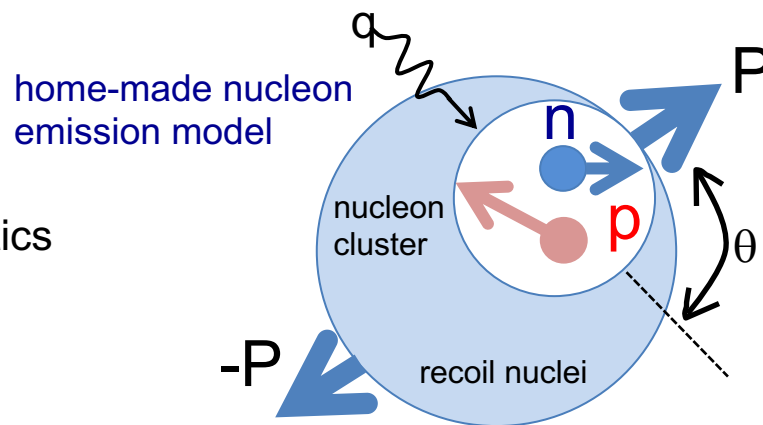
- Without implementing in generators, theoretical nuclear models cannot be compared with data
- Generator implementation is continuously a problem of our community

2. Nucleon correlations in neutrino physics

There is a strong belief in experimental community that hadron final states tell everything about 2p2h...

We need prediction of hadronic final states from theorists

- double differential cross-section = lepton kinematics
- final hadron multiplicity/kinematics = home-made



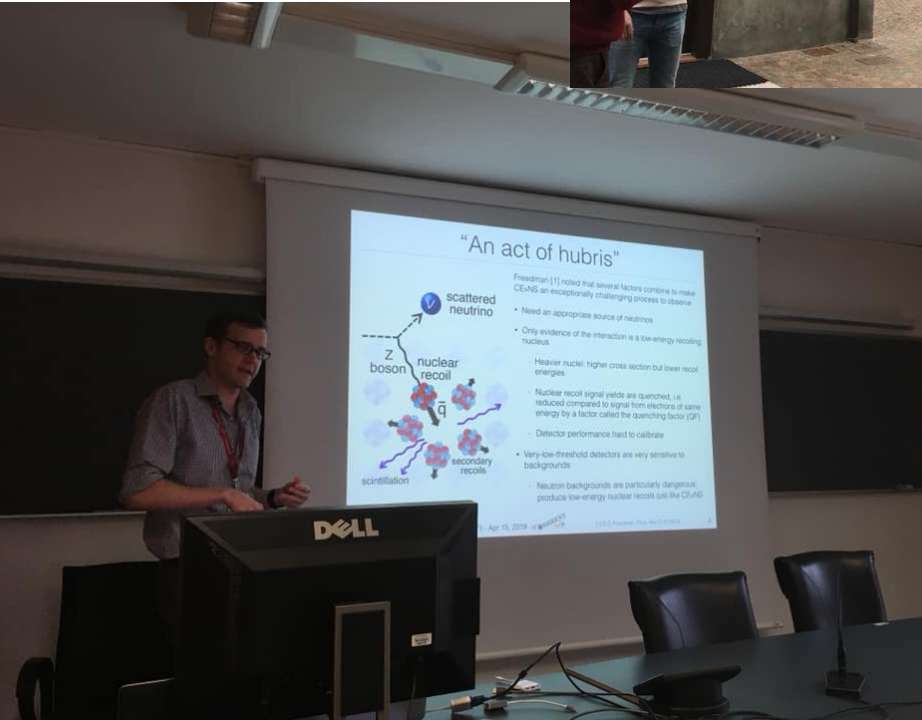
4. Atomic nuclei as laboratories for BSM physics

ECT* workshop, 15 Apr. 15-19 2019, Trento, Italy

<https://www.ectstar.eu/workshops/atomic-nuclei-as-laboratories-for-bsm-physics/>

Topics include;

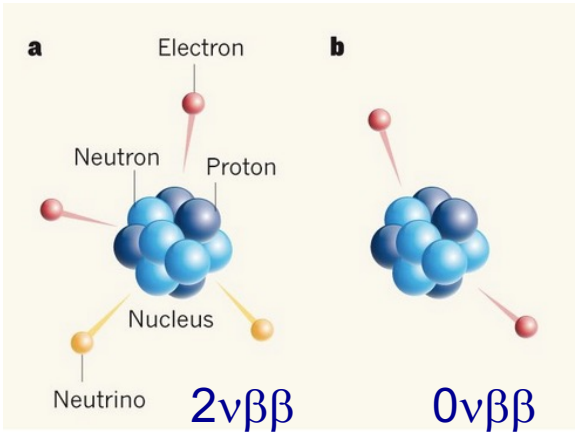
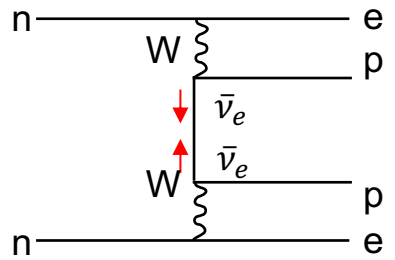
- Neutrino interactions
- $0\nu\beta\beta$
- Direct dark matter
- EMC effect
- etc



4. Neutrino-less double beta decay ($0\nu\beta\beta$)

Majorana particle

- double beta decay ($2\nu\beta\beta$) is the second order nuclear process, possible only for few elements (^{82}Se , ^{76}Ge , ^{100}Mo , ^{130}Te , ^{136}Xe , etc)
- $0\nu\beta\beta$ is the lepton number violation process (BSM process)
- Expected half-life, $\tau(0\nu\beta\beta) > 10^{27}$ yrs ($\gg 10^{10}$ yrs \sim lifetime of universe)

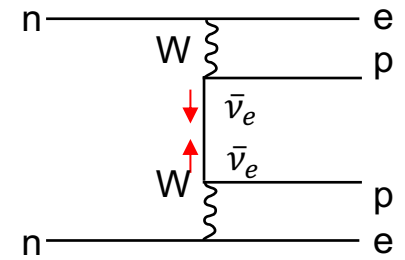


Measured half-life of $0\nu\beta\beta$ process is related to effective Majorana mass ($m_{\beta\beta}^2$)

- Phase space
- Nuclear matrix element
- effective g_A

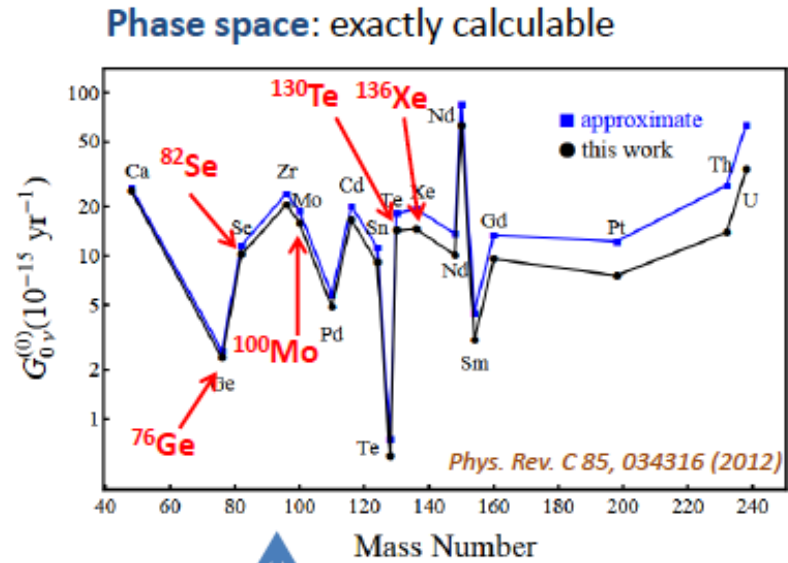
$$1/\tau = G(Q,Z) g_A^4 |M_{nucl}|^2 m_{\beta\beta}^2$$

4. Neutrino-less double beta decay ($0\nu\beta\beta$)

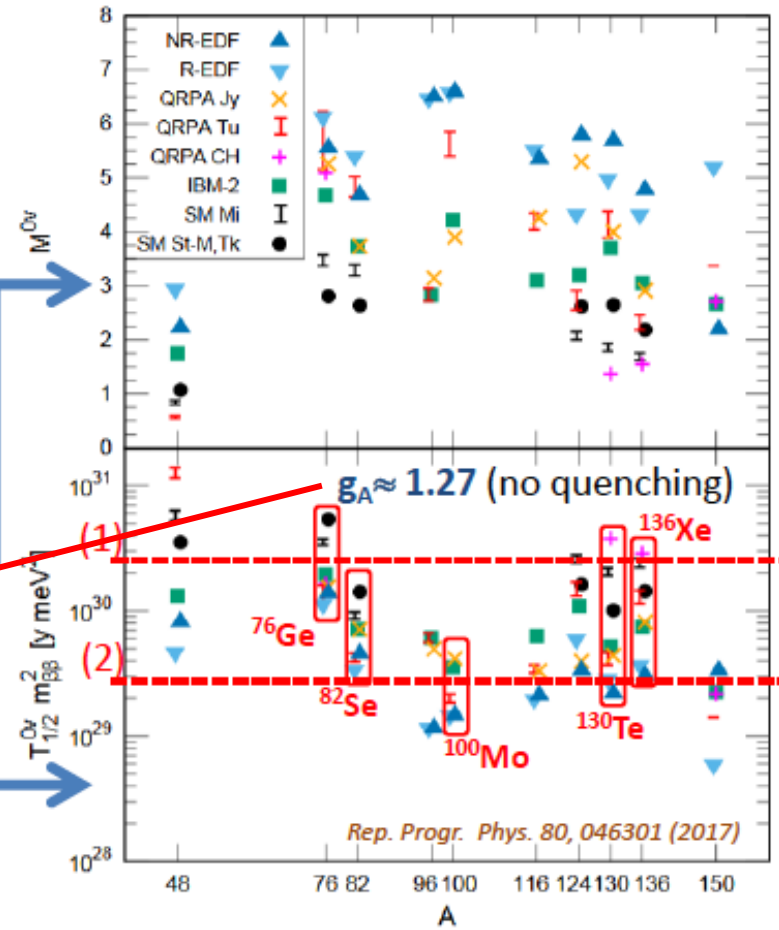


Majorana particle

- Measured half-life of $0\nu\beta\beta$ process is related to effective Majorana mass ($m_{\beta\beta}^2$)



Nuclear matrix elements: several models



$$1/\tau = G(Q,Z) g_A^4 |M_{nucl}|^2 m_{\beta\beta}^2$$

Nuclear physics gives large systematics

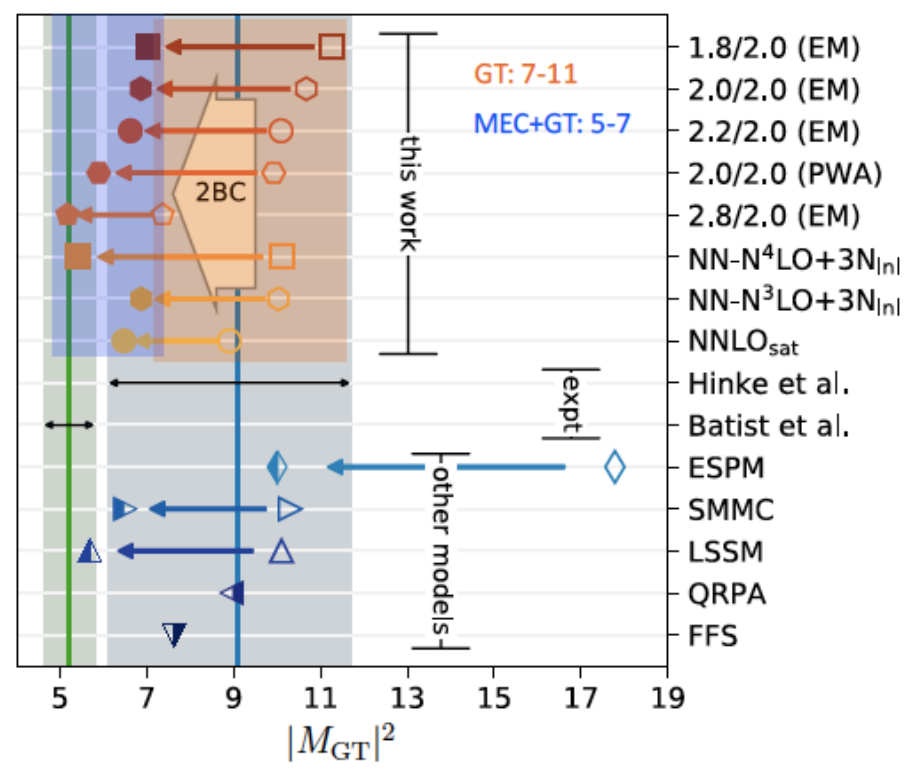
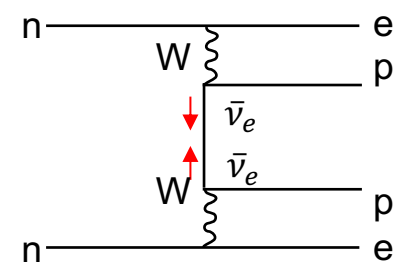
- Nuclear matrix element calculation
- Nuclear quenching of g_A



4. Nucleon correlation and $0\nu\beta\beta$

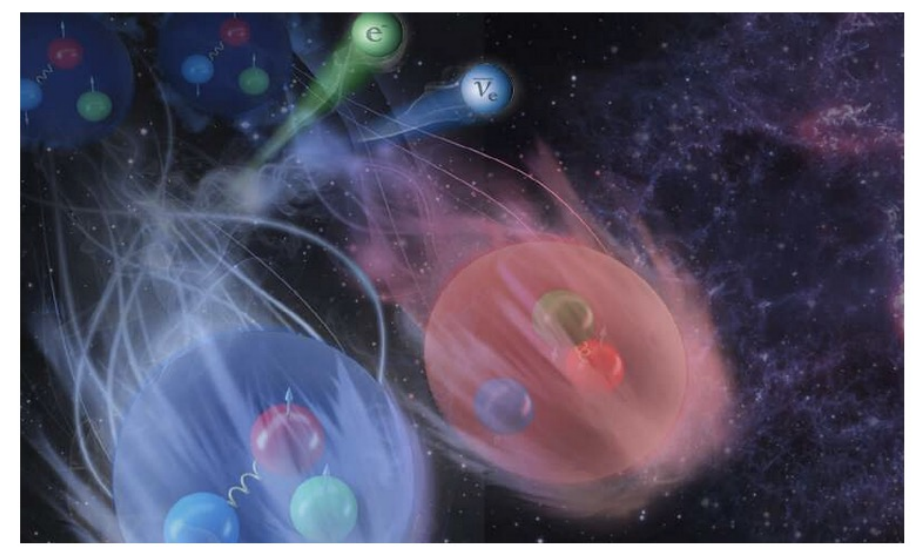
Beta decay quenching

- Axial coupling looks smaller in nuclei
- Ab initio calculation shows matrix element is suppressed due to nucleon 2-body current (2BC)
- Another uncertainty of $0\nu\beta\beta$



Physicists solve a beta-decay puzzle with advanced nuclear models

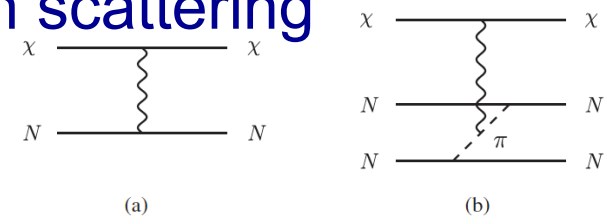
by Oak Ridge National Laboratory



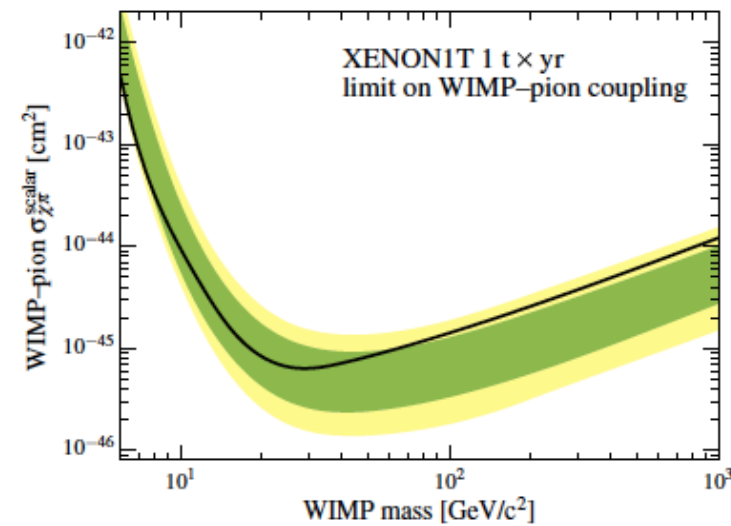
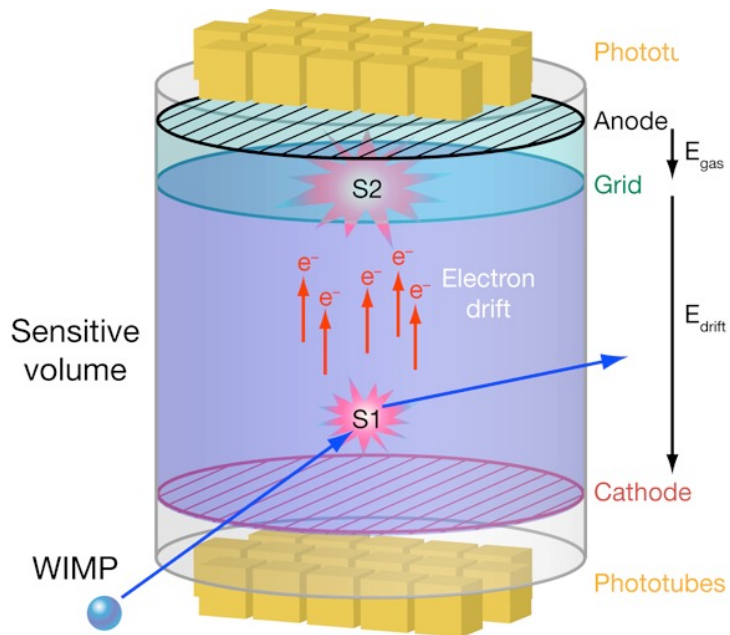
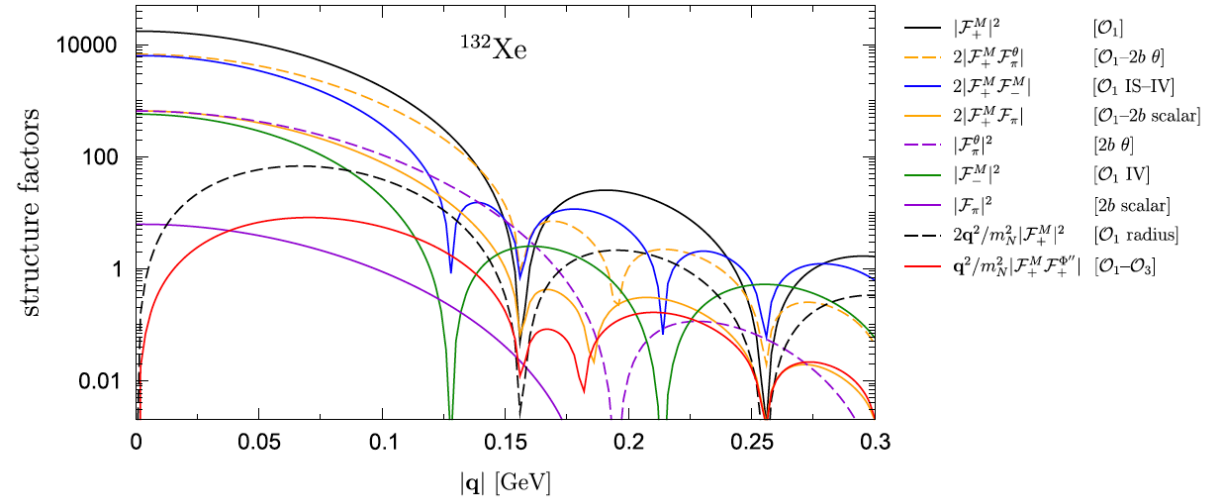
4. Nucleon correlation and Direct WIMP-pion scattering

Direct WIMP-pion interaction

- WIMP interaction depends on nuclear structure function
- Chiral effective field theory including 2-body current
- Assuming leading contributions are zero, WIMP-pion scattering can be studied.



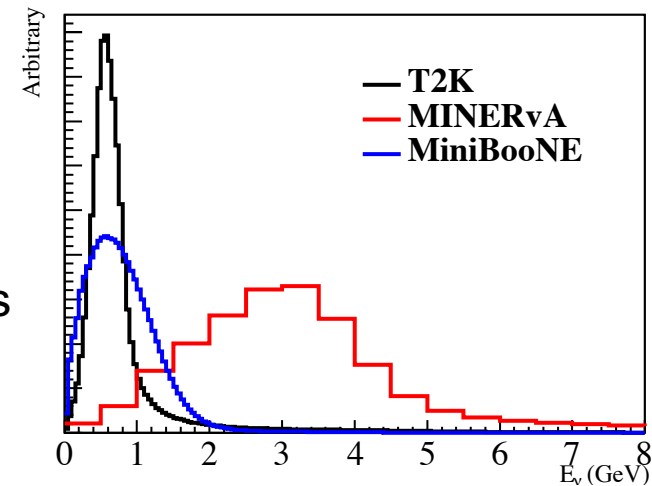
$$\frac{d\sigma_{\chi N}^{SI}}{dq^2} = \frac{1}{4\pi v^2} \left| \left(c_+^M - \frac{\mathbf{q}^2}{m_N^2} \dot{c}_+^M \right) \mathcal{F}_+^M(\mathbf{q}^2) + c_\pi \mathcal{F}_\pi(\mathbf{q}^2) + c_\pi^\theta \mathcal{F}_\pi^\theta(\mathbf{q}^2) + \left(c_-^M - \frac{\mathbf{q}^2}{m_N^2} \dot{c}_-^M \right) \mathcal{F}_-^M(\mathbf{q}^2) + \frac{\mathbf{q}^2}{2m_N^2} [c_+^{\Phi''} \mathcal{F}_+^{\Phi''}(\mathbf{q}^2) + c_-^{\Phi''} \mathcal{F}_-^{\Phi''}(\mathbf{q}^2)] \right|^2,$$



5. MiniBooNE

Mineral oil (CH_2) Cherenkov detector

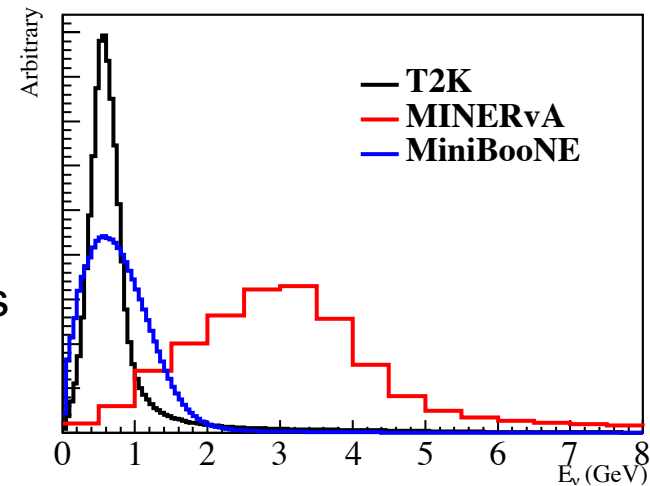
- 4π coverage, $\langle E \rangle \sim 800$ MeV beam up to 2 GeV
- Measure Cherenkov radiations from charged particles
- Some calorimetric (scintillation)



5. MiniBooNE

Mineral oil (CH₂) Cherenkov detector

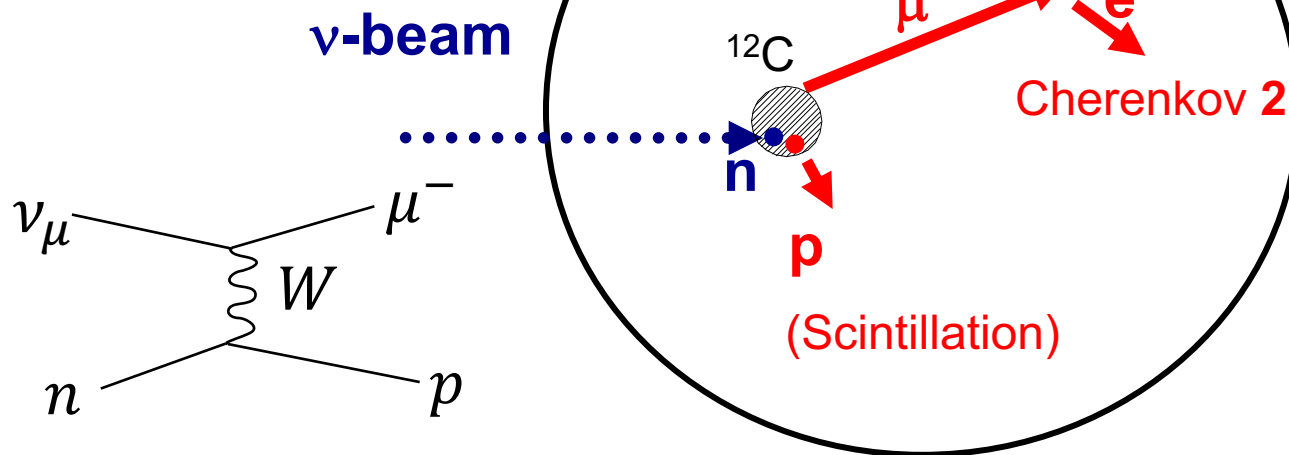
- 4π coverage, <E>~800 MeV beam up to 2 GeV
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- Some calorimetric (scintillation)



MiniBooNE CCQE measurement

MiniBooNE detector

(spherical Cherenkov detector)

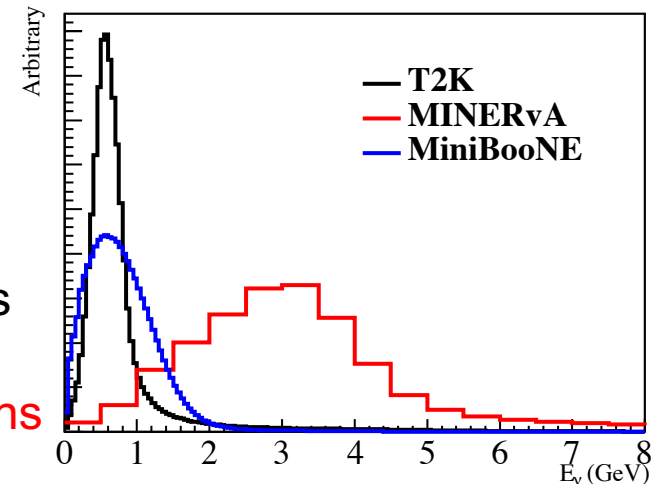


muon like Cherenkov light and subsequent decayed electron (Michel electron) like Cherenkov light are the signal of CCQE event

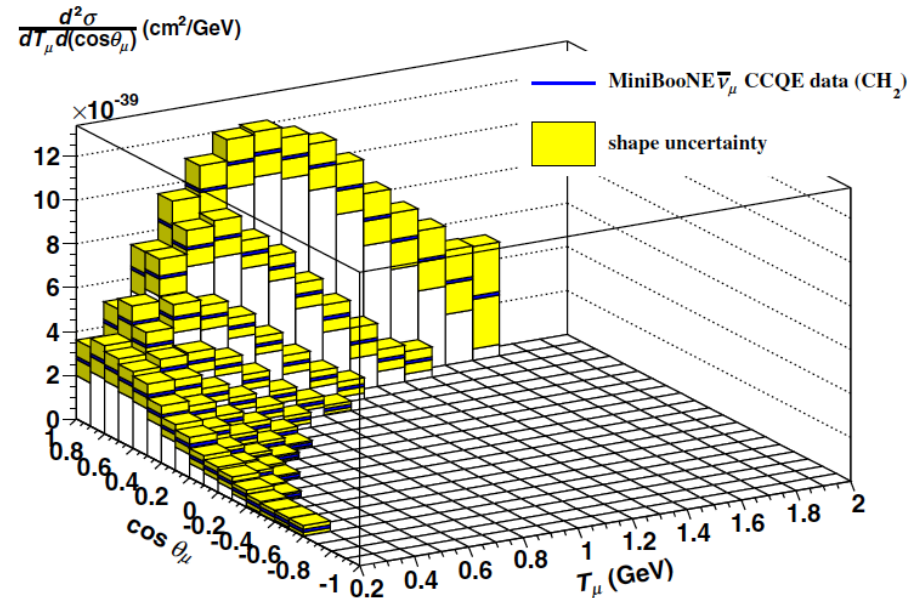
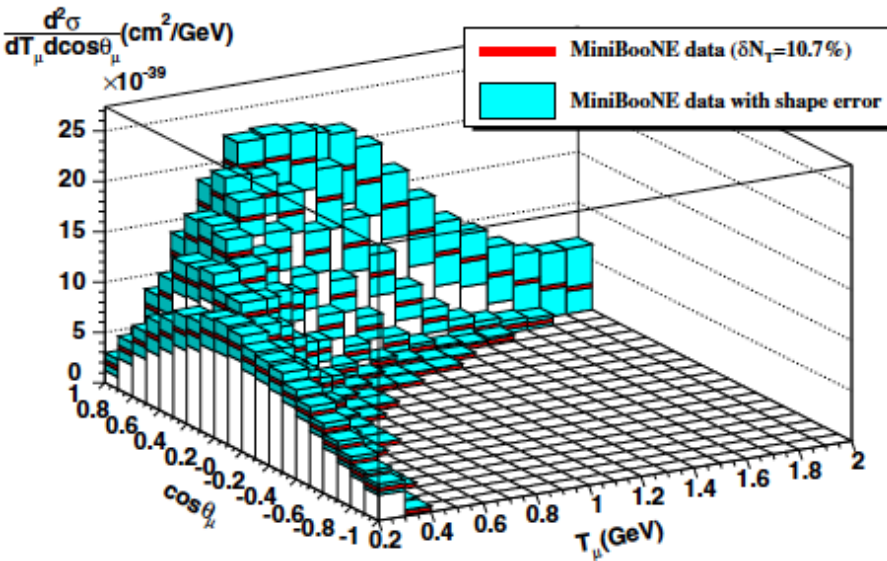
5. MiniBooNE

Mineral oil (CH₂) Cherenkov detector

- 4π coverage, <E>~800 MeV beam up to 2 GeV
- Measure Cherenkov radiations from charged particles
- Some calorimetric (scintillation)
- Measured first **flux-integrated differential cross sections**
- Solved **CCQE puzzle**



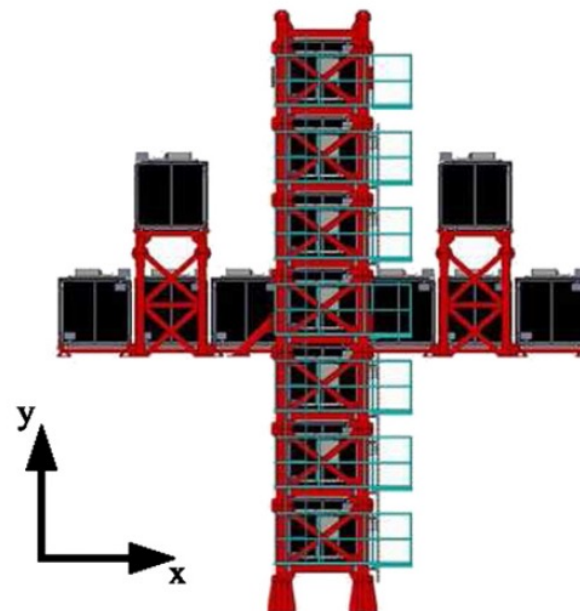
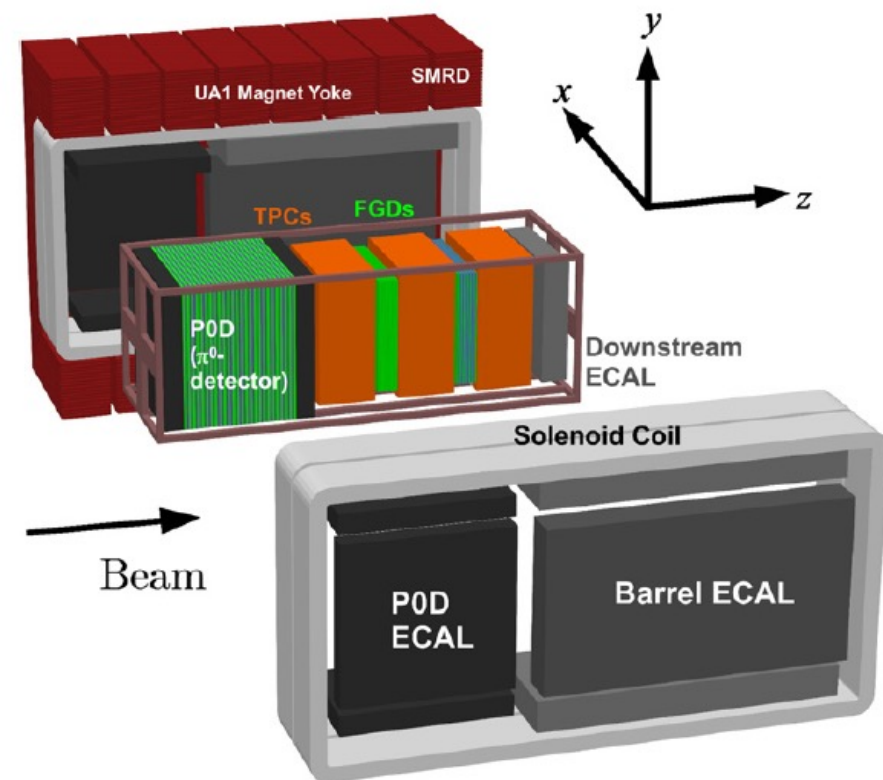
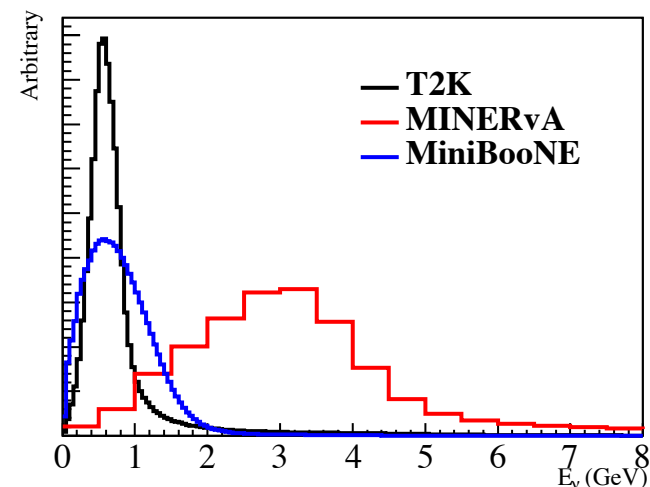
neutrino and anti-neutrino CCQE-like double differential cross sections



5. T2K near detector complex

INGRID, FGD, P0D, ECal, TPC, SMRD

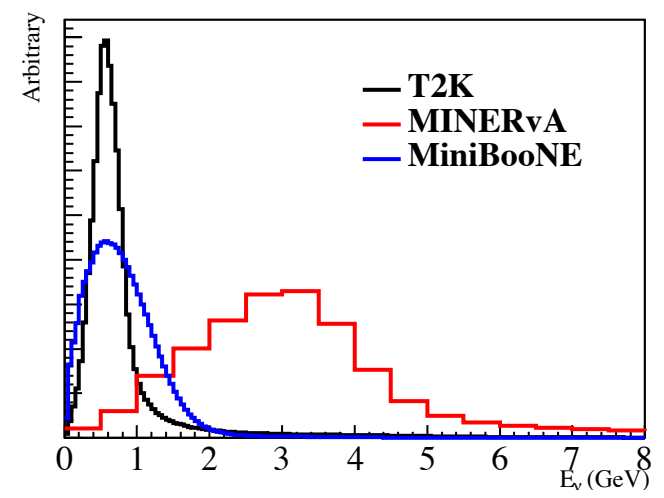
- Plastic scintillation trackers (except gas TPC)
- 0.2T magnet for momentum measurement
- $\langle E \rangle \sim 600$ MeV off-axis beam
- variety of targets (CH, H₂O, Pb, Ar)
- Limited coverage (combination of sub-detectors)



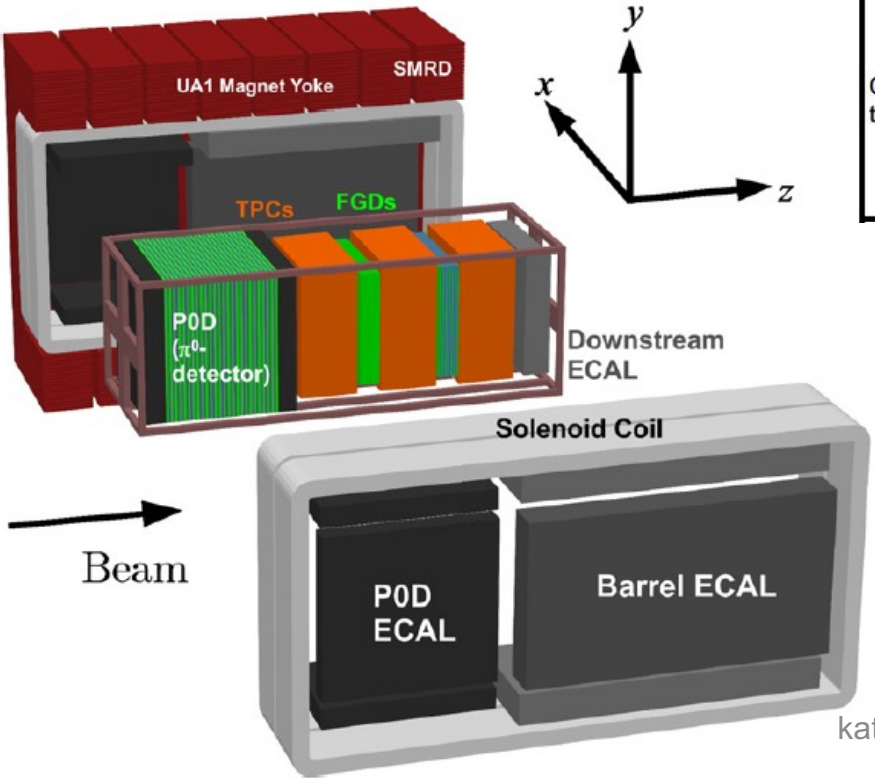
5. T2K near detector complex

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- Plastic scintillation trackers (except gas TPC)
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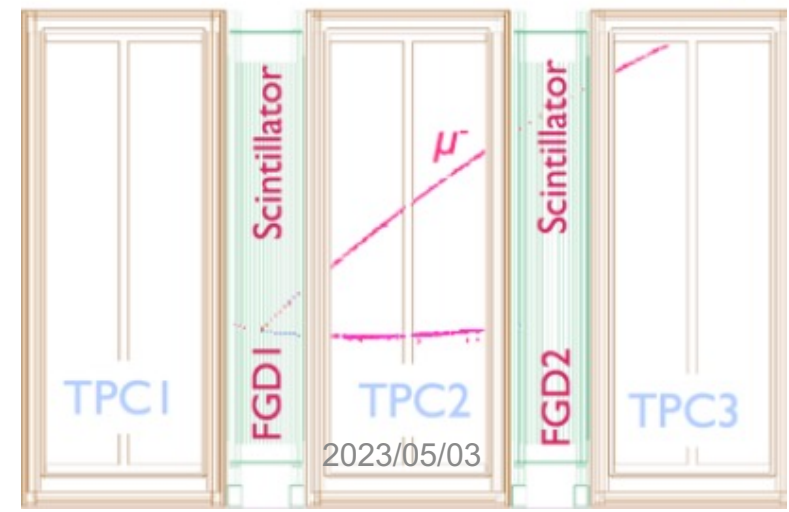


neutrino CC0 π double differential cross sections



	(1)	(2)	(3)	(4a)	(4b)
CCQE topology					
	ECAL+SMRD	ECAL+SMRD	ECAL+SMRD	ECAL+SMRD	ECAL+SMRD
	FGD	FGD	FGD	FGD	FGD
	TPC	TPC	TPC	TPC	TPC
	ECAL+SMRD	ECAL+SMRD	ECAL+SMRD	ECAL+SMRD	ECAL+SMRD

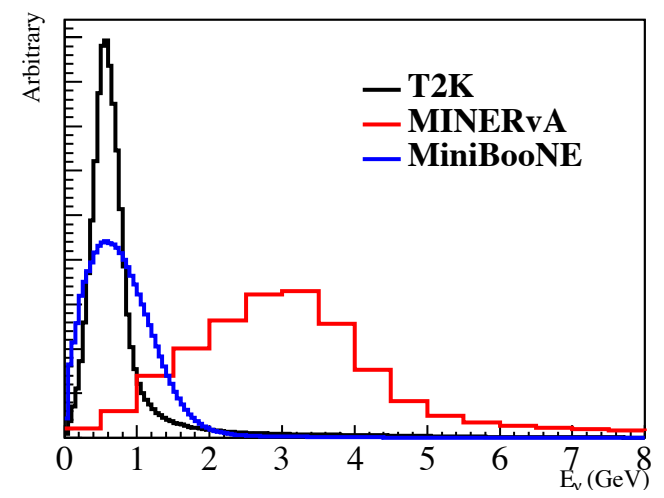
Run #: 4200 Evt #: 24083 Time: Sun 2010-03-21 22:33:25 JST



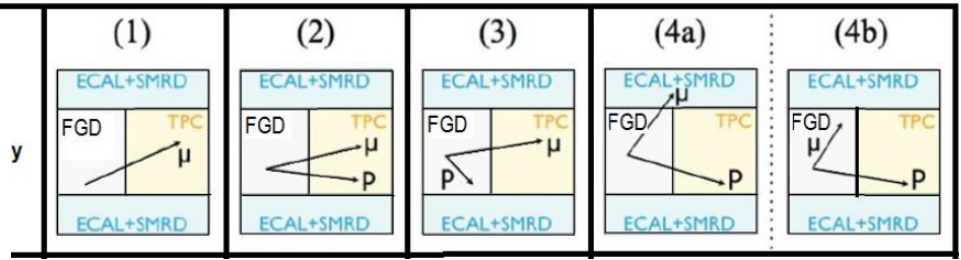
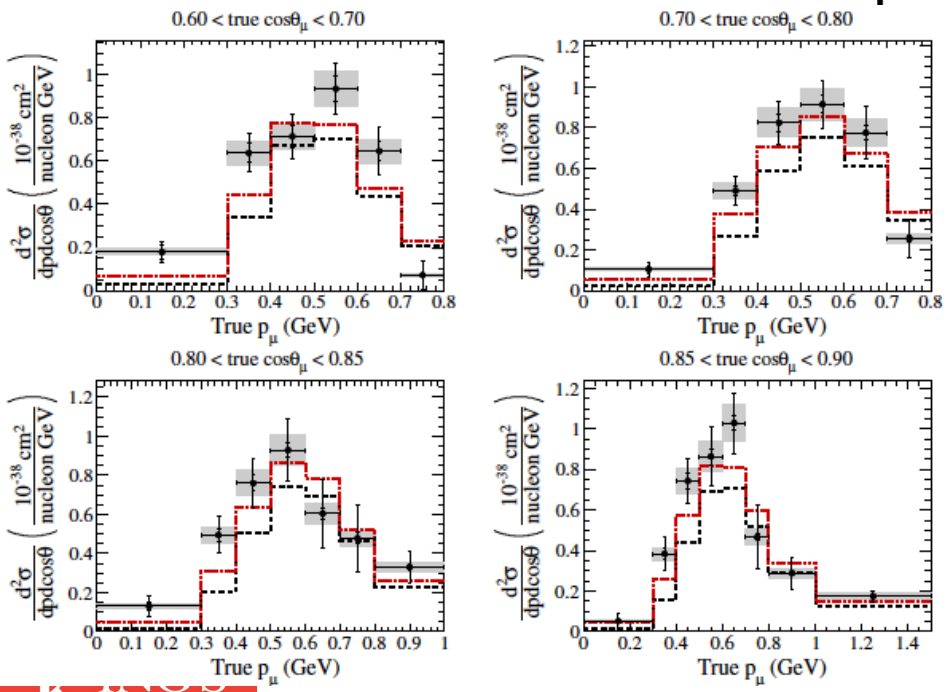
5. T2K near detector complex

INGRID, FGD, P0D, ECal, TPC, SMRD

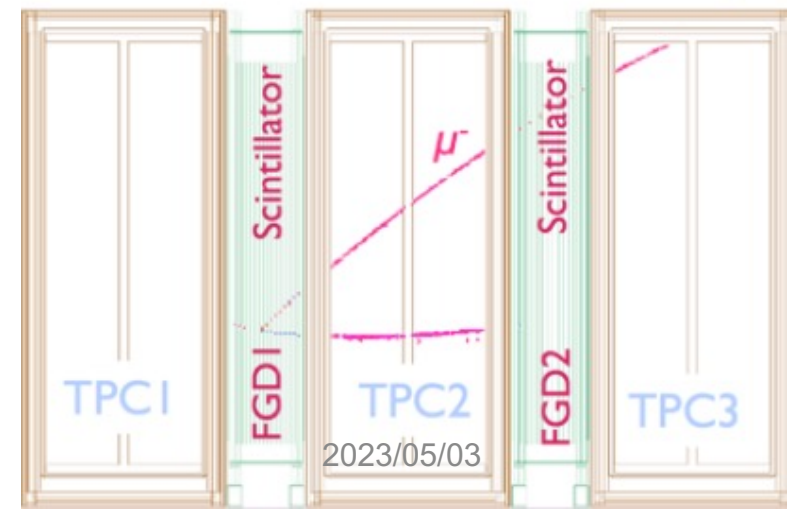
- Plastic scintillation trackers (except gas TPC)
- 0.2T magnet for momentum measurement
- $\langle E \rangle \sim 600$ MeV off-axis beam
- variety of targets (CH, H₂O, Pb, Ar)
- Limited coverage (combination of sub-detectors)



neutrino CC0 π double differential cross sections



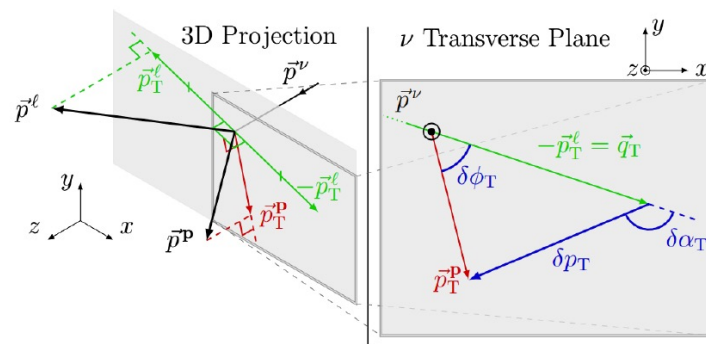
Run #: 4200 Evt #: 24083 Time: Sun 2010-03-21 22:33:25 JST



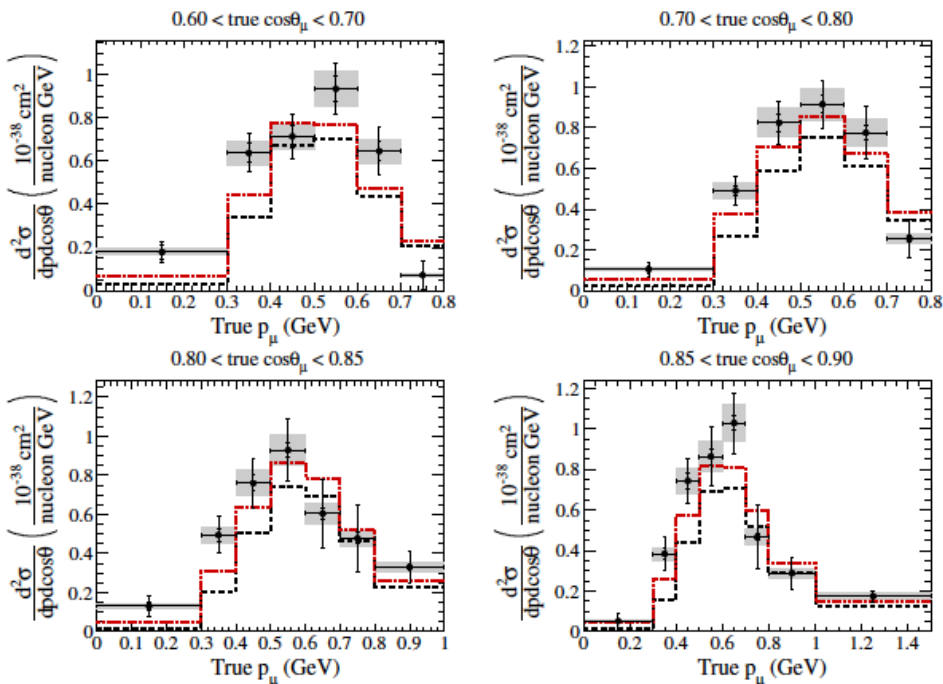
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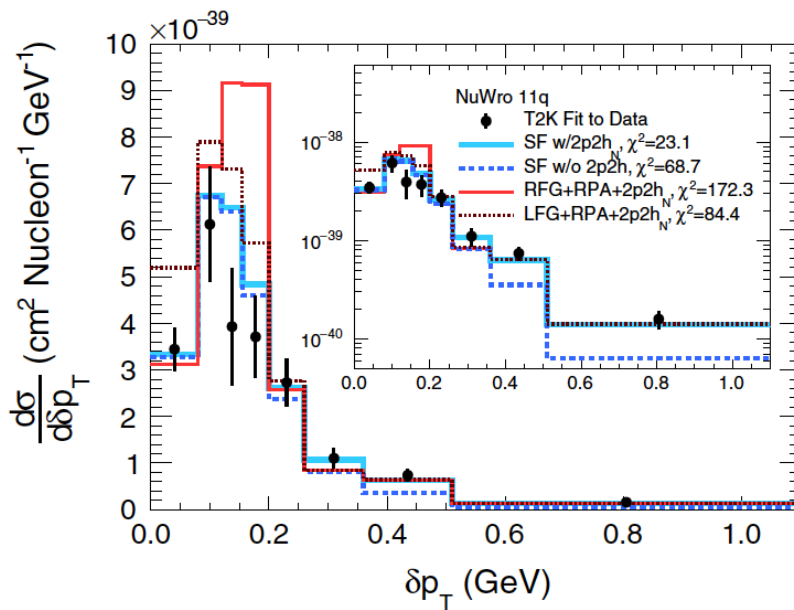
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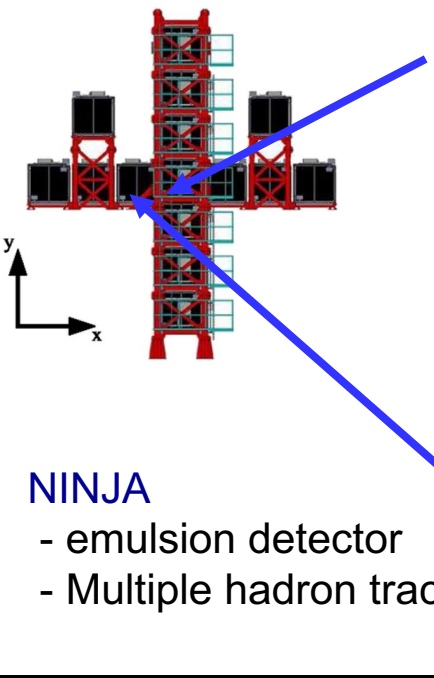
neutrino CC0 π double differential cross sections



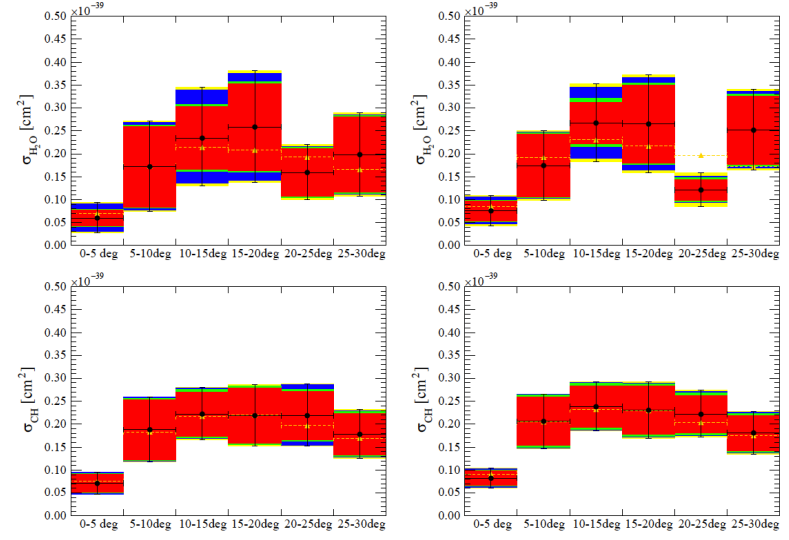
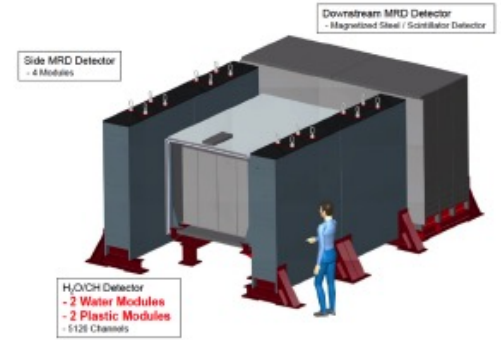
neutrino CC0 π 1p differential cross sections



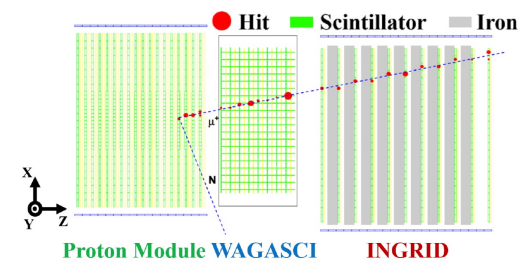
5. T2K near detector complex



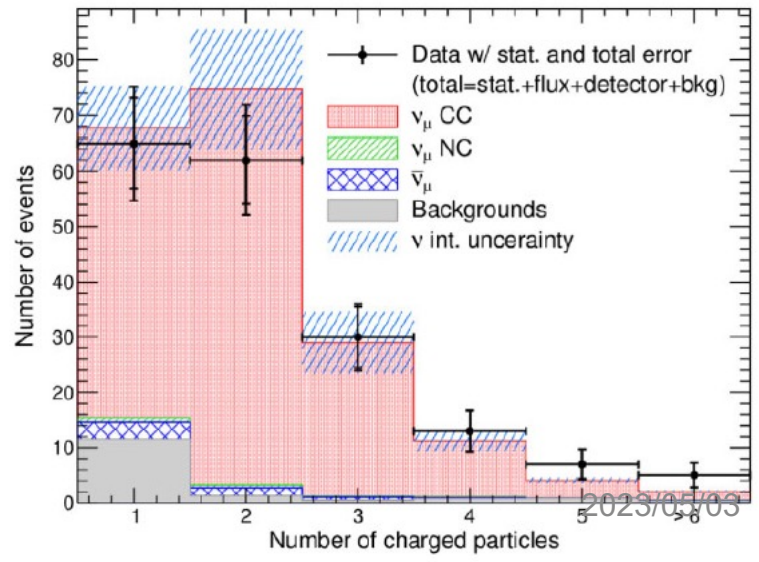
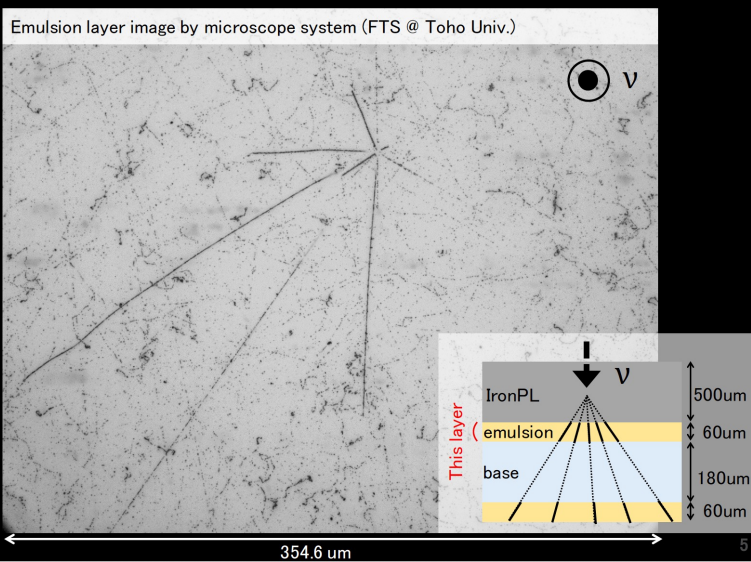
WAGASCI
- water target



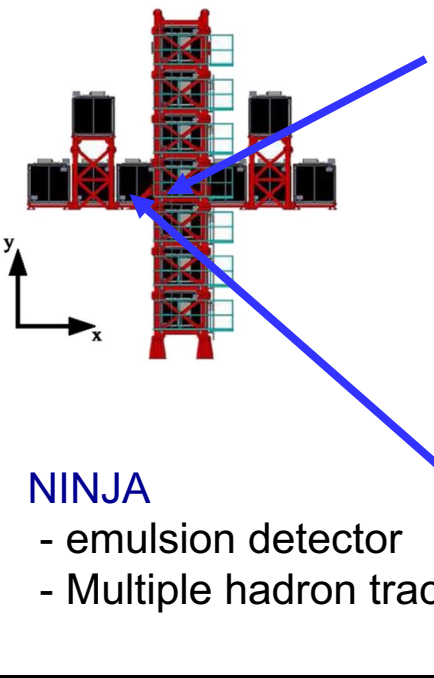
NINJA
- emulsion detector
- Multiple hadron tracks



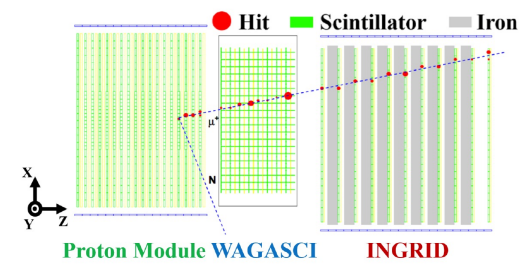
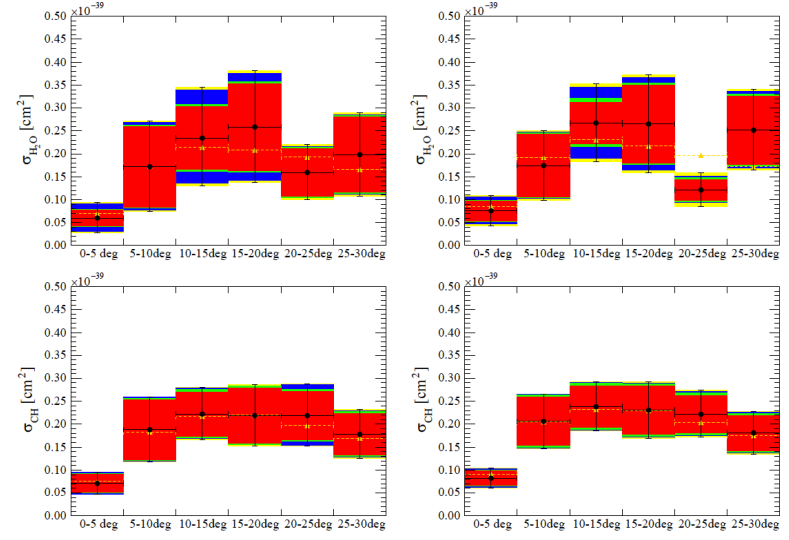
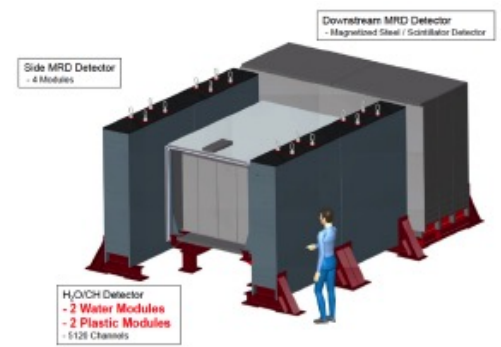
An example of ν - iron interaction (2016 NINJA iron target run)



5. T2K near detector complex

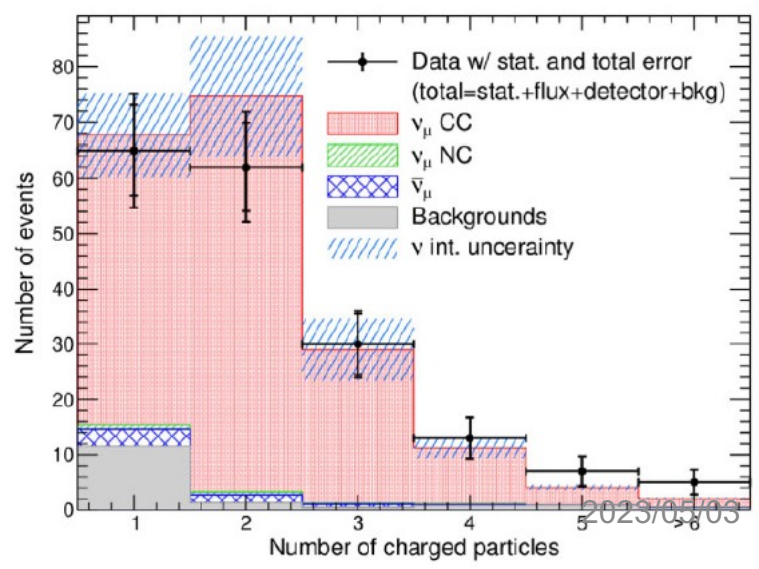
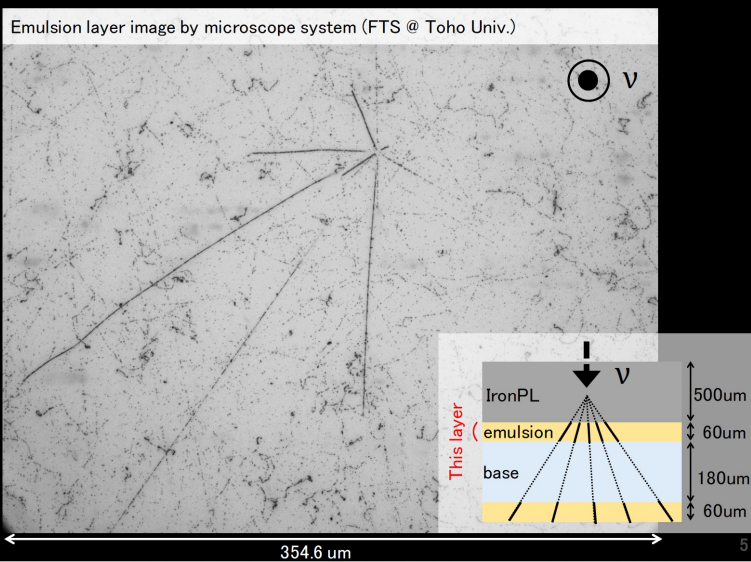


WAGASCI
- water target



NINJA
- emulsion detector
- Multiple hadron tracks

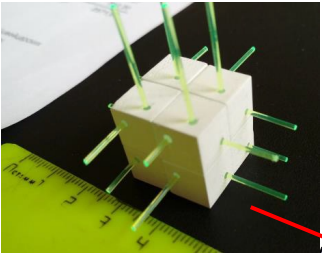
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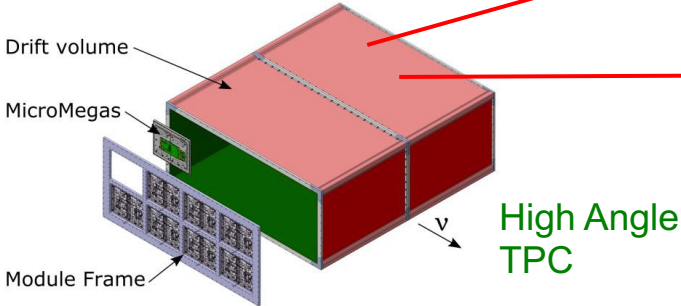
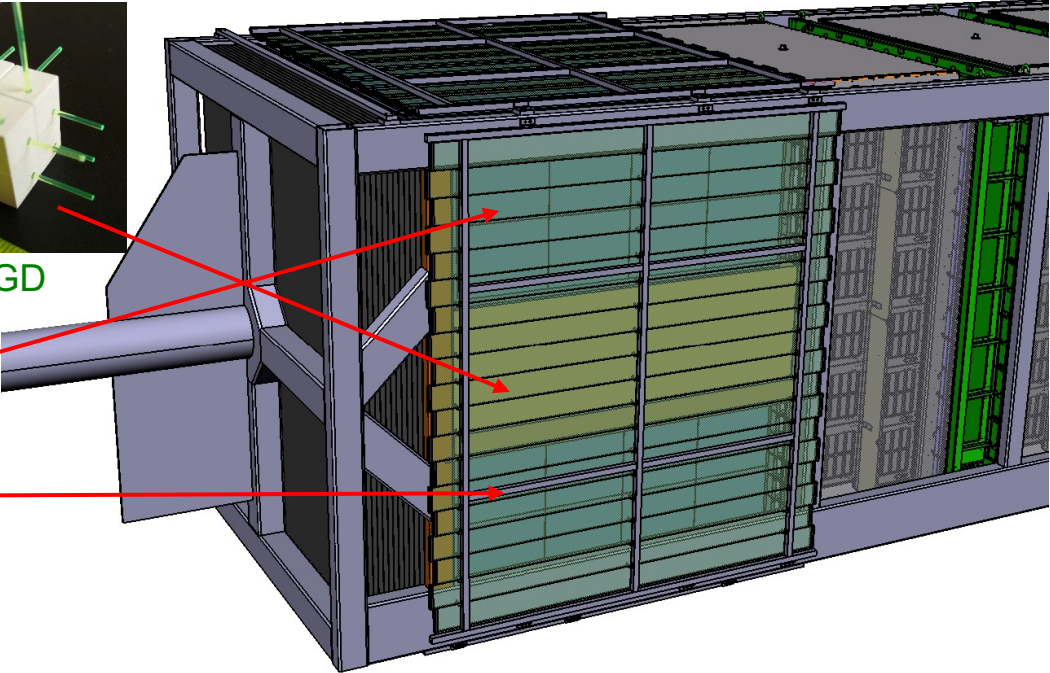
5. ND280 Upgrade

ND280 Upgrade

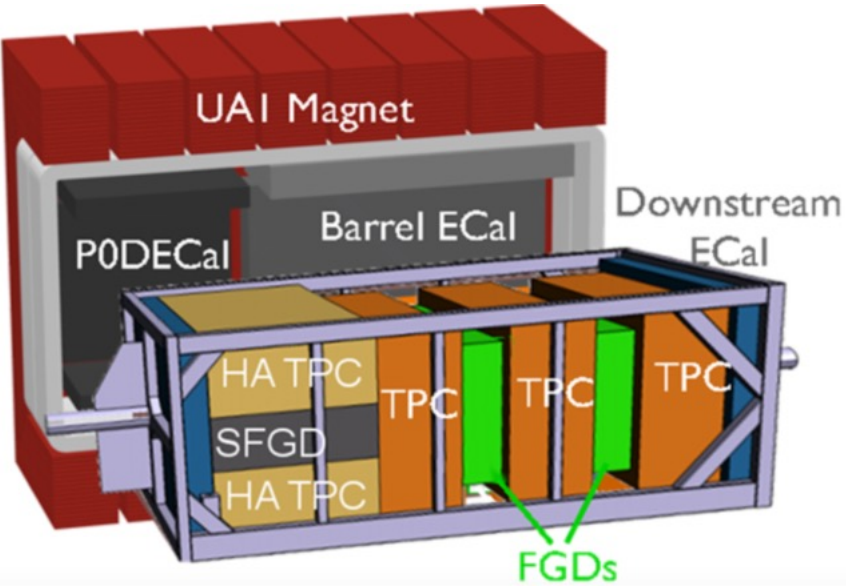
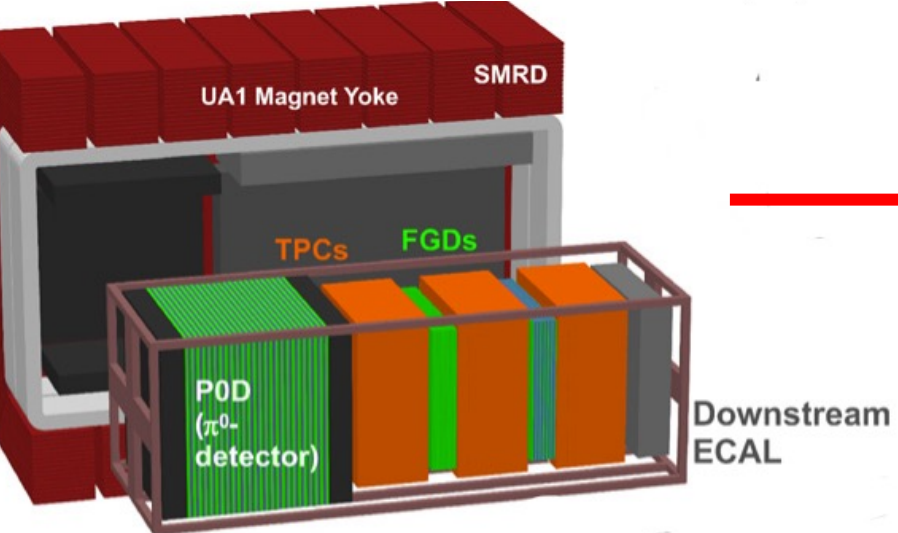
- Out: P0D detector
- In: High Angle TPC (HATPC)
- In: SuperFGD



SuperFGD



High Angle TPC



5. ND280 Upgrade

ND280 Upgrade

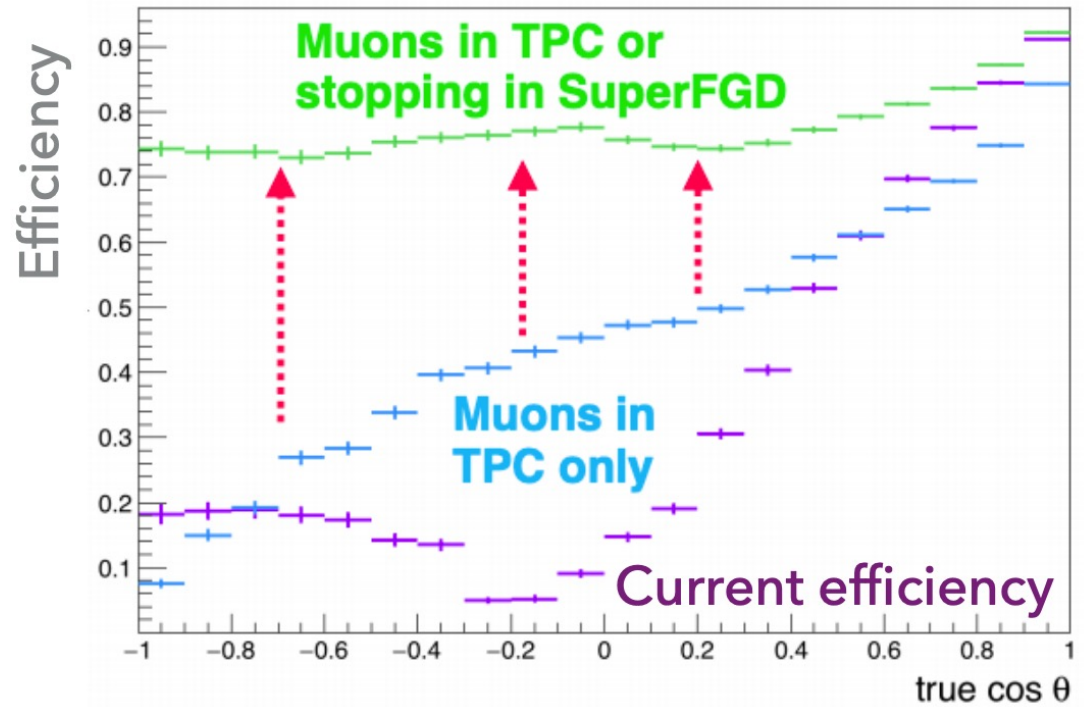
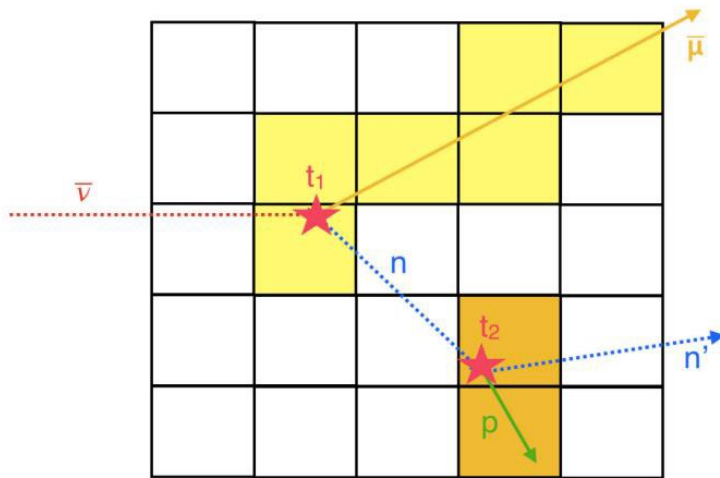
- Out: POD detector
- In: High Angle TPC (HATPC)
- In: SuperFGD

4 π coverage

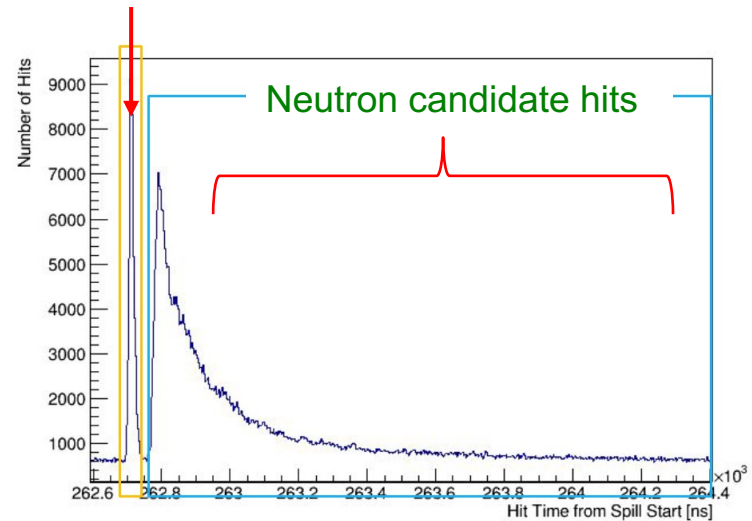
- It matches with Hyper-K phase space

Neutron tagging

- SuperFGD beam test at LANL
- ToF to measure energy



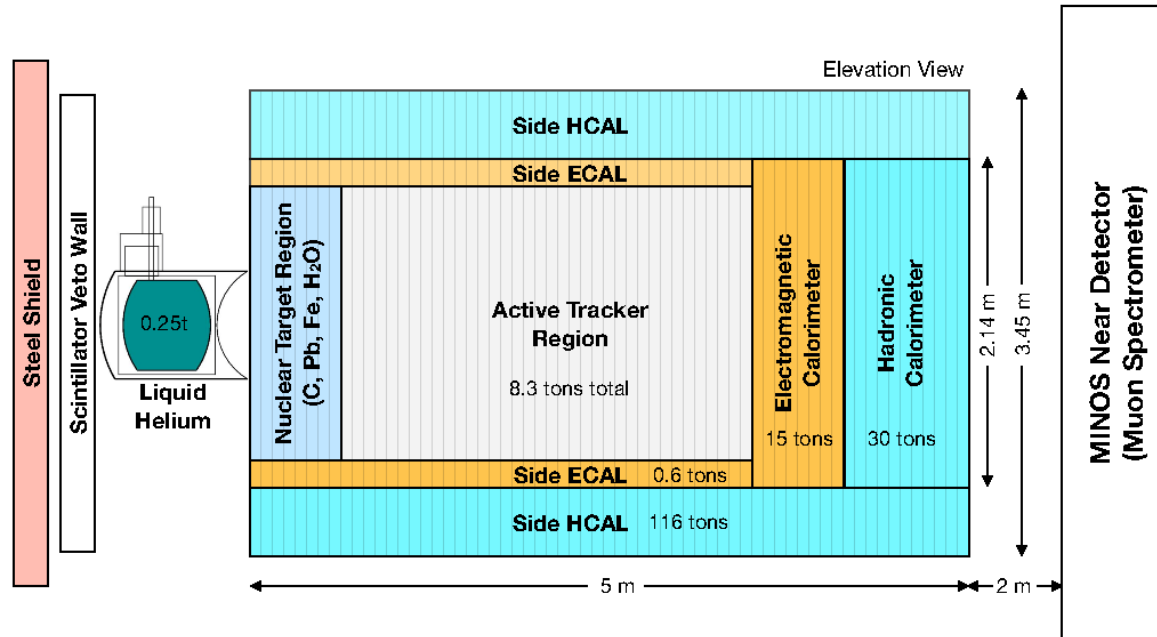
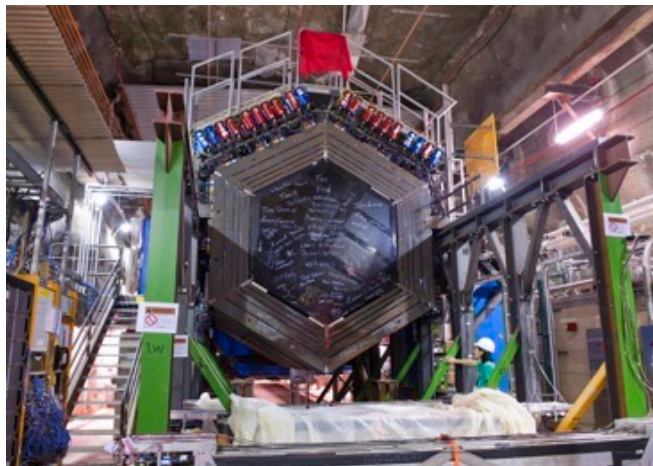
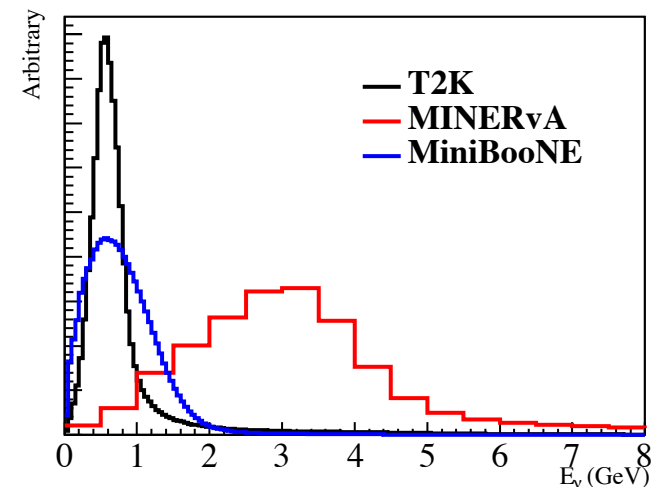
t_0 (γ -ray)



5. MINERvA

Scintillation tracker

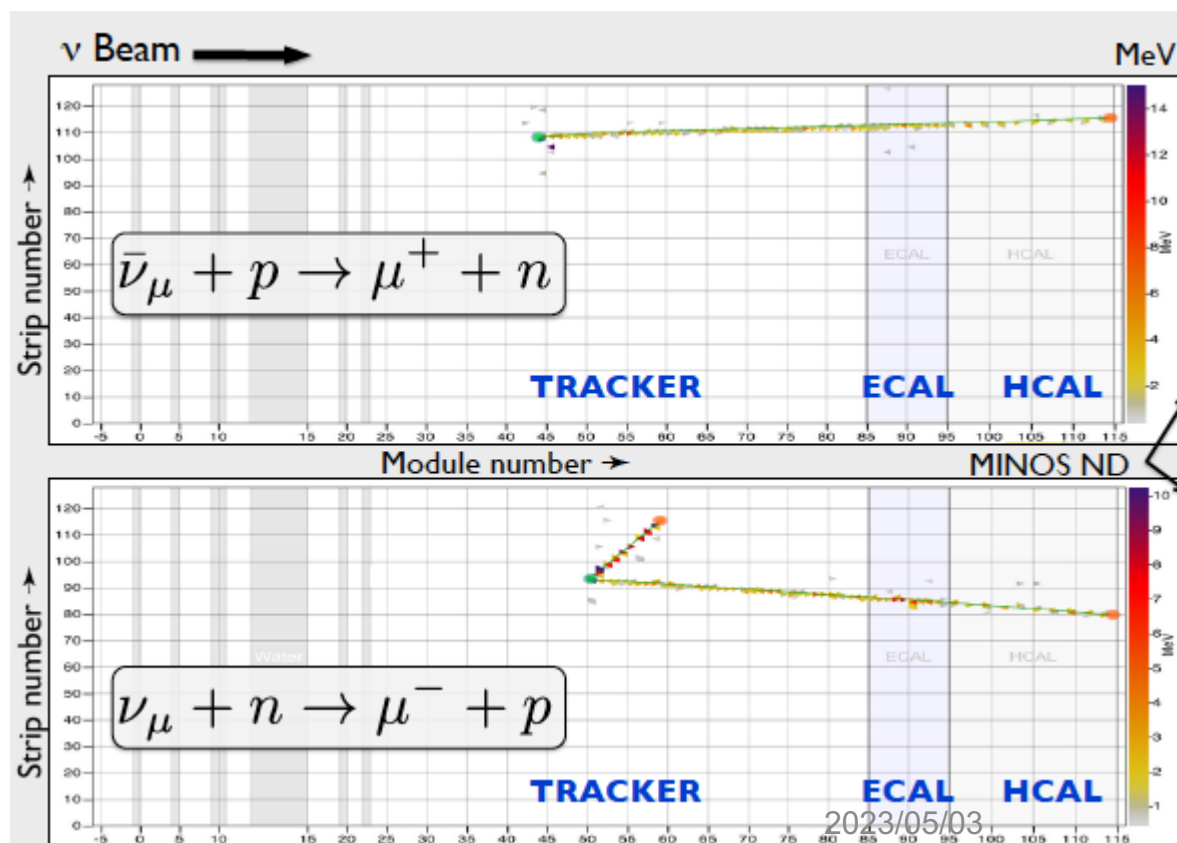
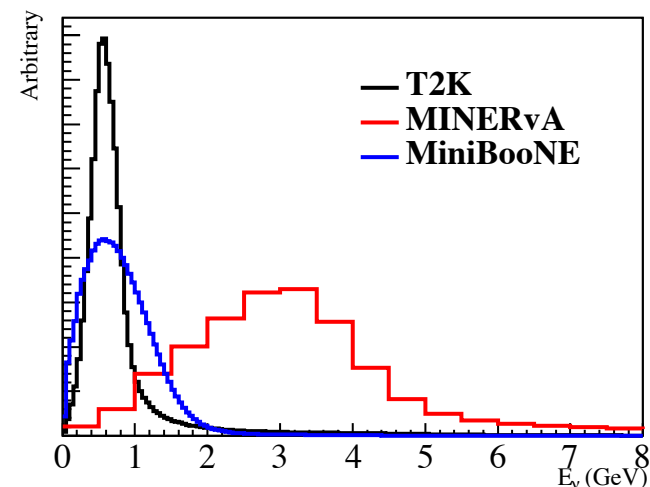
- $\langle E \rangle \sim 3.5$ GeV on-axis beam
- variety of targets (CH, Pb, Fe)
- Small acceptance due to MINOS ND
- charge separation by MINOS ND
- internal flux constraint (DIS, ν -e)



5. MINERvA

Scintillation tracker

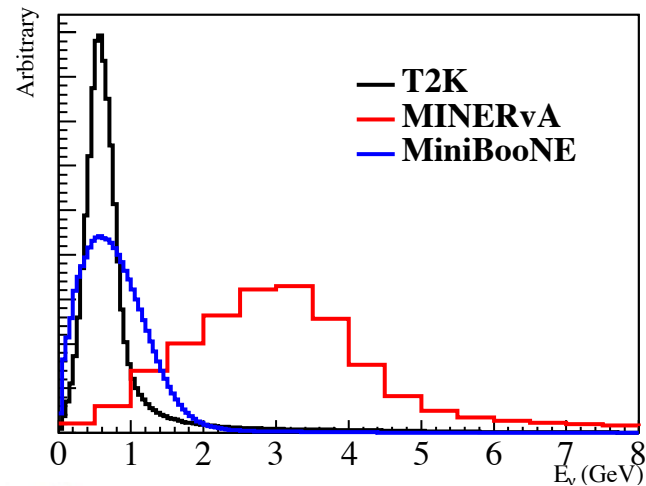
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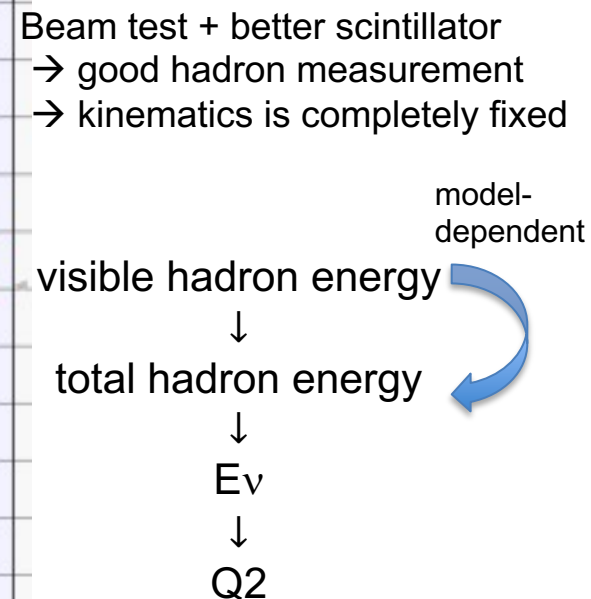
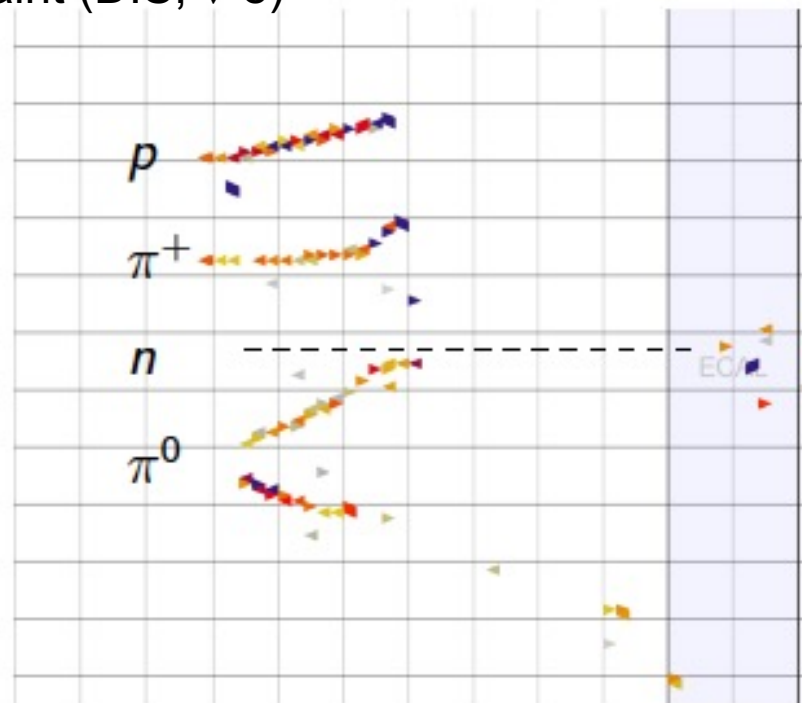


Kinetic energy

Kinetic energy

0

Total energy



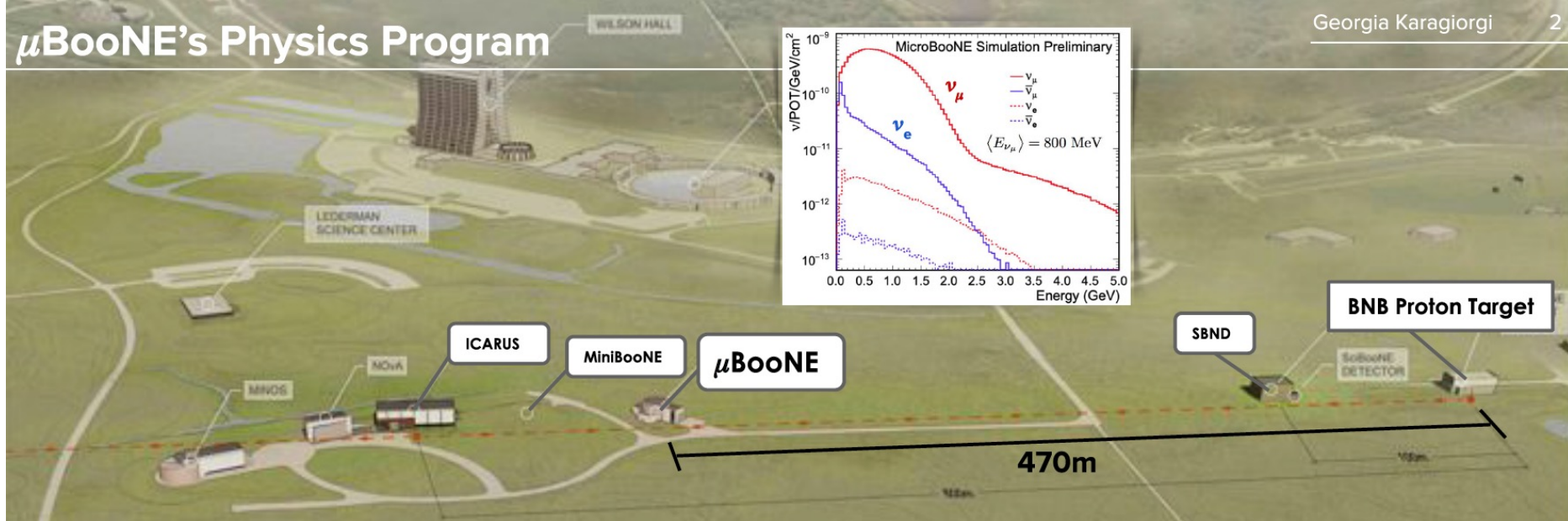
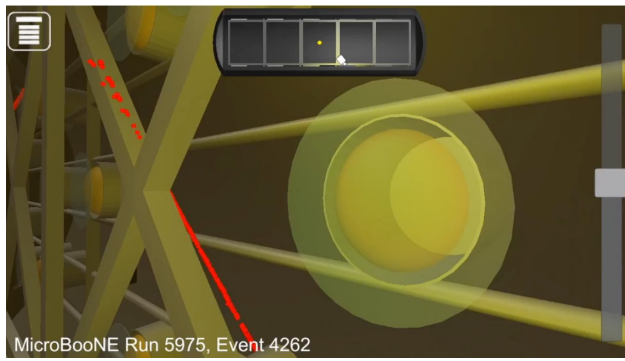
On average, we see *available* hadronic energy $E_{avail} \neq q_0$:

$$E_{avail} = \sum (\text{Proton and } \pi^\pm \text{ KE}) + (\text{Total } E \text{ of other particles except neutrons})$$

5. MicroBooNE

86ton LArTPC

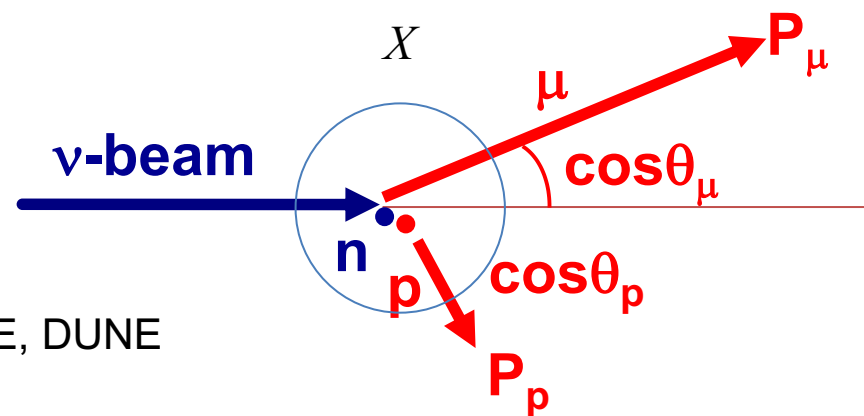
- $\langle E \rangle \sim 800$ MeV BNB on-axis beam
- Single phase LArTPC, 3-wire-plane reading
- 3mm pitch
- photon detection system
- ArgoNeuT, LArIAT, SBND, ICARUS, protoDUNE, DUNE



5. MicroBooNE

86ton LArTPC

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MicroBooNE CC $\mu+p$ differential cross section

Outgoing proton kinematics are measured to reconstruct Fermi motion

Multiple Coulomb scattering to estimate escaping muon energy

Large cosmic ray background, but mostly understood

Low statistics for hadron measurements

