

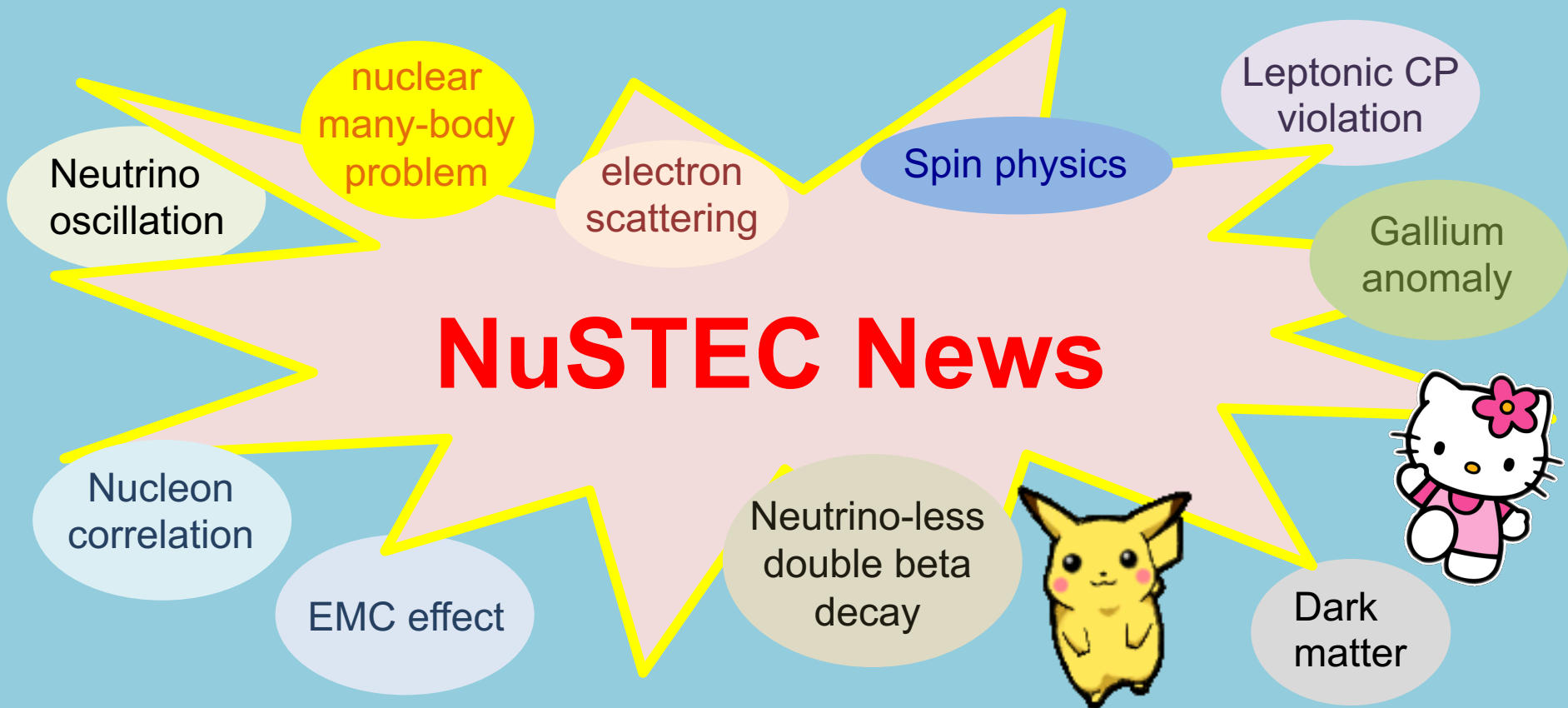
Nuclear Physics for **Beyond the Standard Model** **Neutrino Physics**

outline

1. Neutrino interaction physics - introduction
2. Charged-Current Quasi-Elastic (CCQE) interaction
3. Neutrino baryonic resonance interaction
4. Neutrino shallow- and deep-inelastic scatterings
5. Conclusion

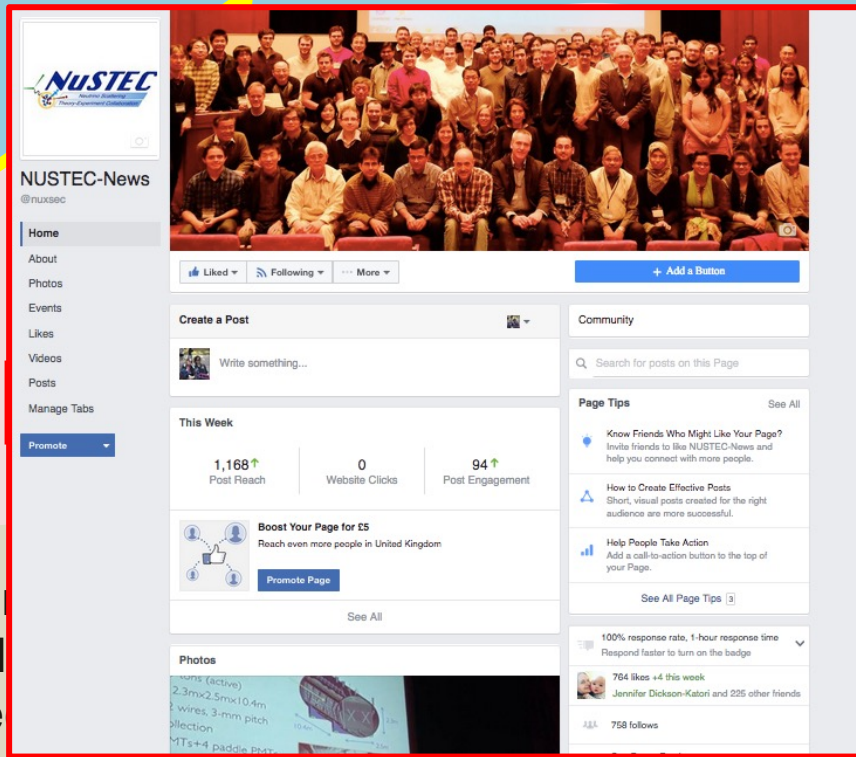
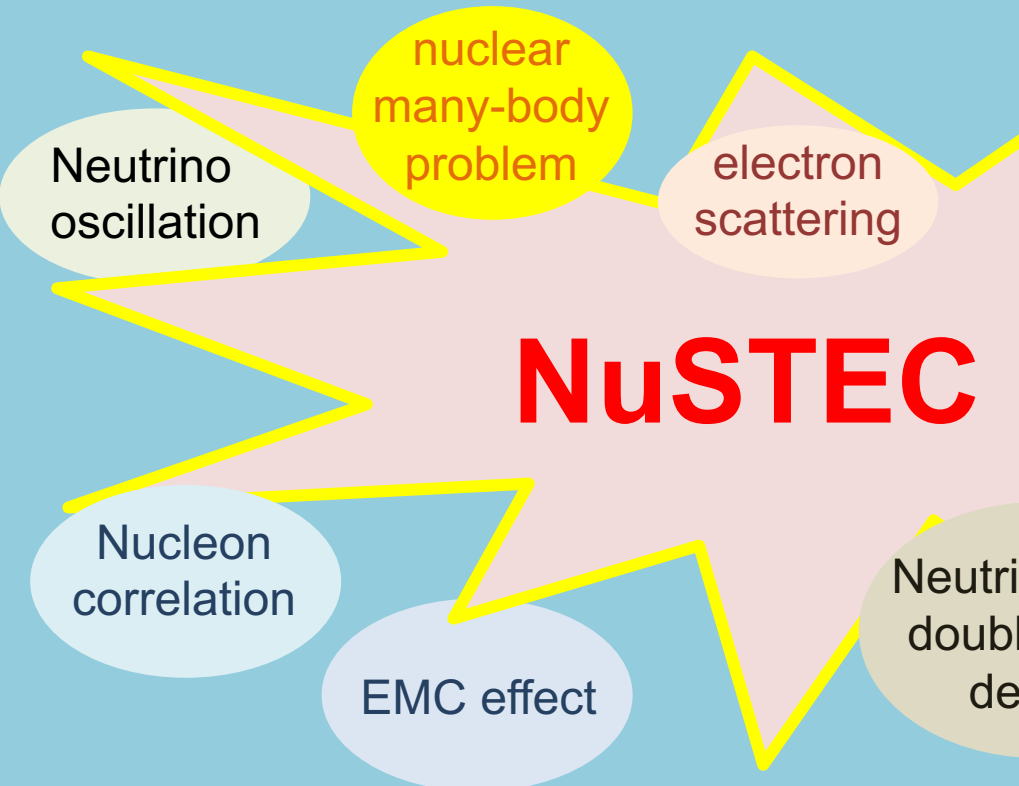
Teppei Katori  @teppeikatori
King's College London
ICN seminar, UNAM, Nov. 17, 2021

Fun Timely Intellectual Adorable!



Subscribe “NuSTEC News”
E-mail to listserv@fnal.gov, Leave the subject line blank, Type "subscribe nustec-news firstname lastname"
(or just send e-mail to me, katori@FNAL.GOV)
like “@nuxsec” on Facebook page, use hashtag #nuxsec

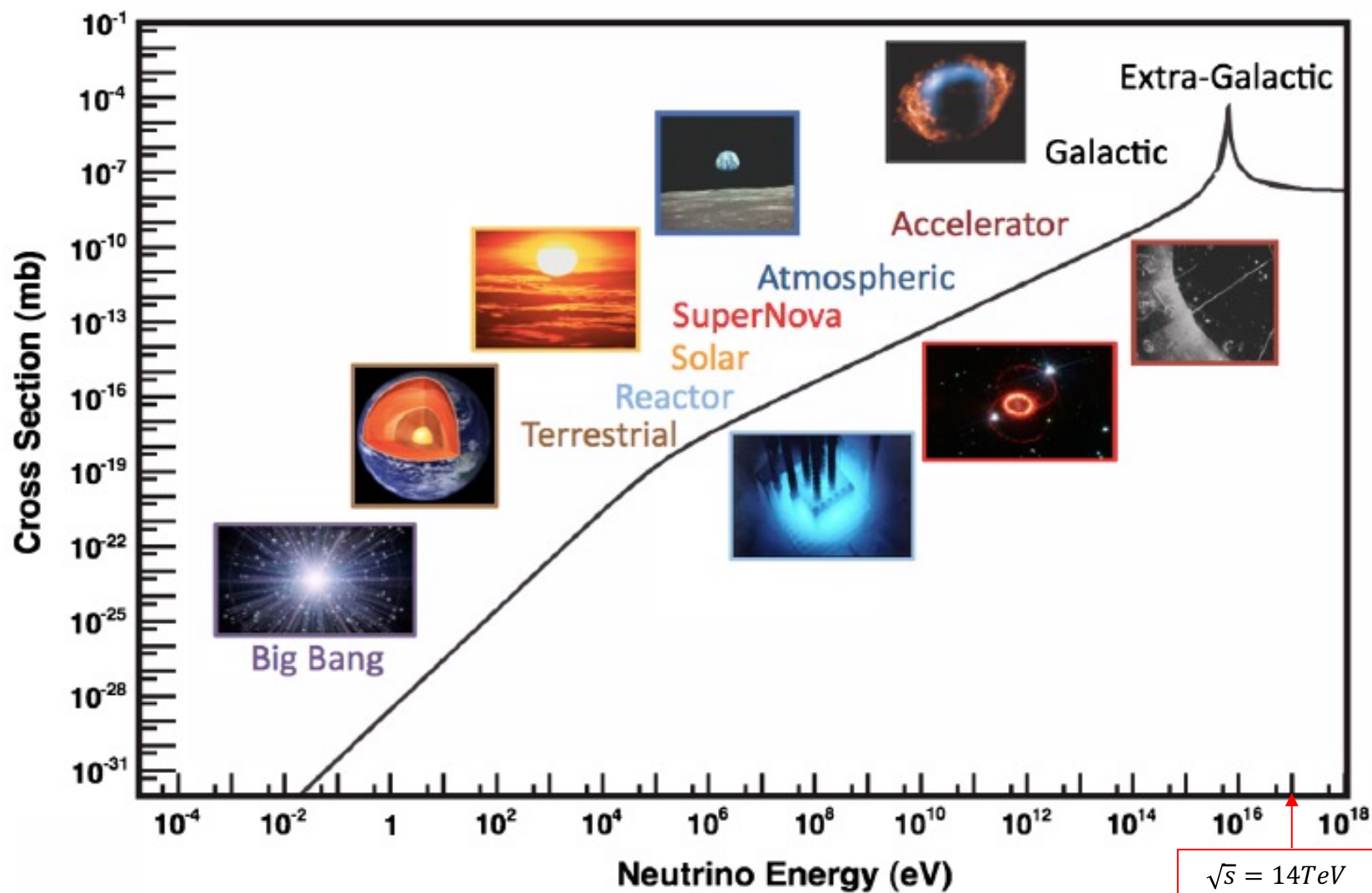
Fun Timely Intellectual Adorable!



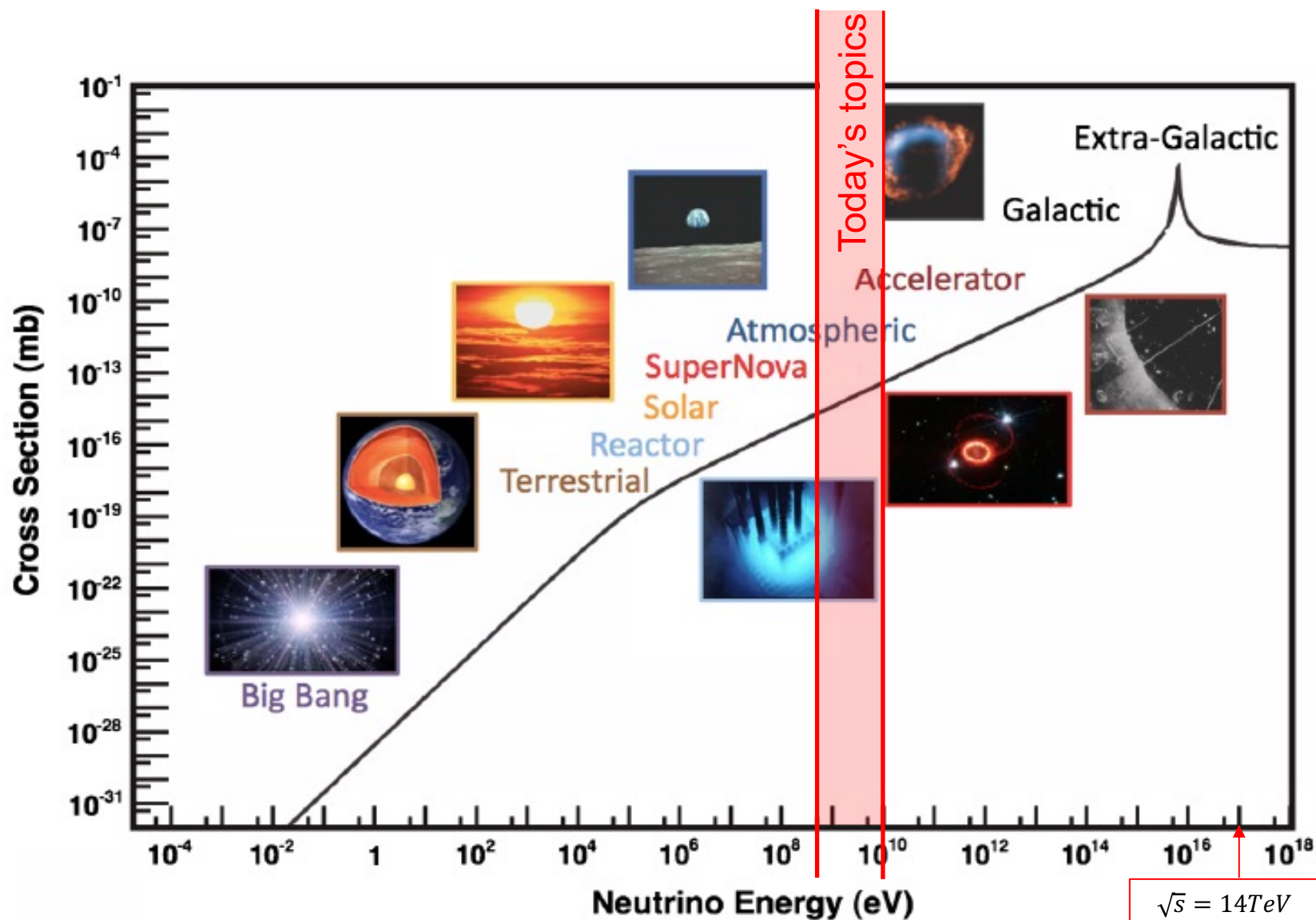
Subscribe "NuSTEC News"
E-mail to listserv@fnal.gov, Leave the subject line blank, Type "subscribe nustec-news firstname lastname"
(or just send e-mail to me, katori@fnal.gov)
like "[@nuxsec](https://www.facebook.com/nuxsec)" on Facebook page, use hashtag [#nuxsec](https://www.facebook.com/nuxsec)

- 1. Neutrino interaction physics - introduction**
2. Charged-Current Quasi-Elastic (CCQE) interaction
3. Neutrino baryonic resonance interaction
4. Neutrino shallow- and deep-inelastic scatterings
5. Conclusions

1. From eV to EeV: Neutrino cross sections across energy scales



1. From eV to EeV: Neutrino cross sections across energy scales

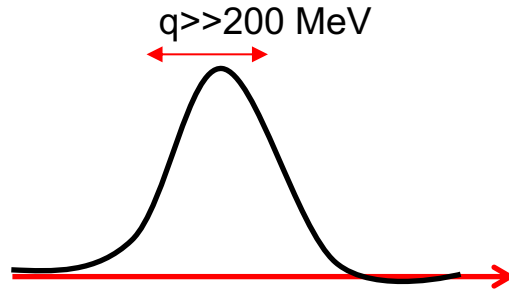


1. Neutrino interaction physics around 1-10 GeV

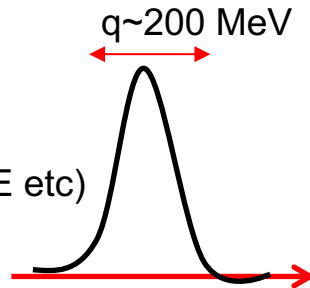
Size of wave packet \sim momentum transfer (\sim energy)

$$\hbar c = 197 \text{ MeV} \cdot \text{fm} \rightarrow 200 \text{ MeV} \sim 1 \text{ fm (size of nucleon)}$$

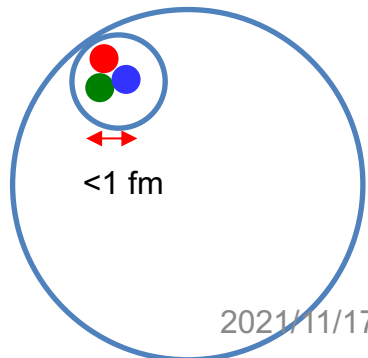
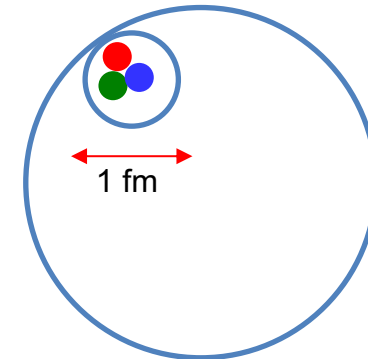
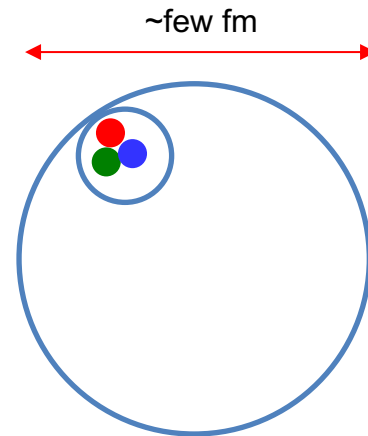
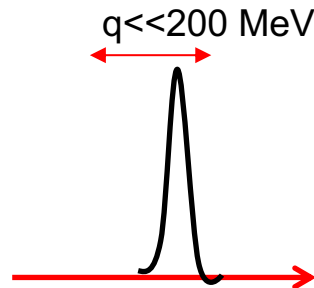
$\ll 1$ GeV neutrino beam
(solar neutrinos, etc)



~ 1 GeV neutrino beam
(T2K, NOvA, HyperK, DUNE etc)



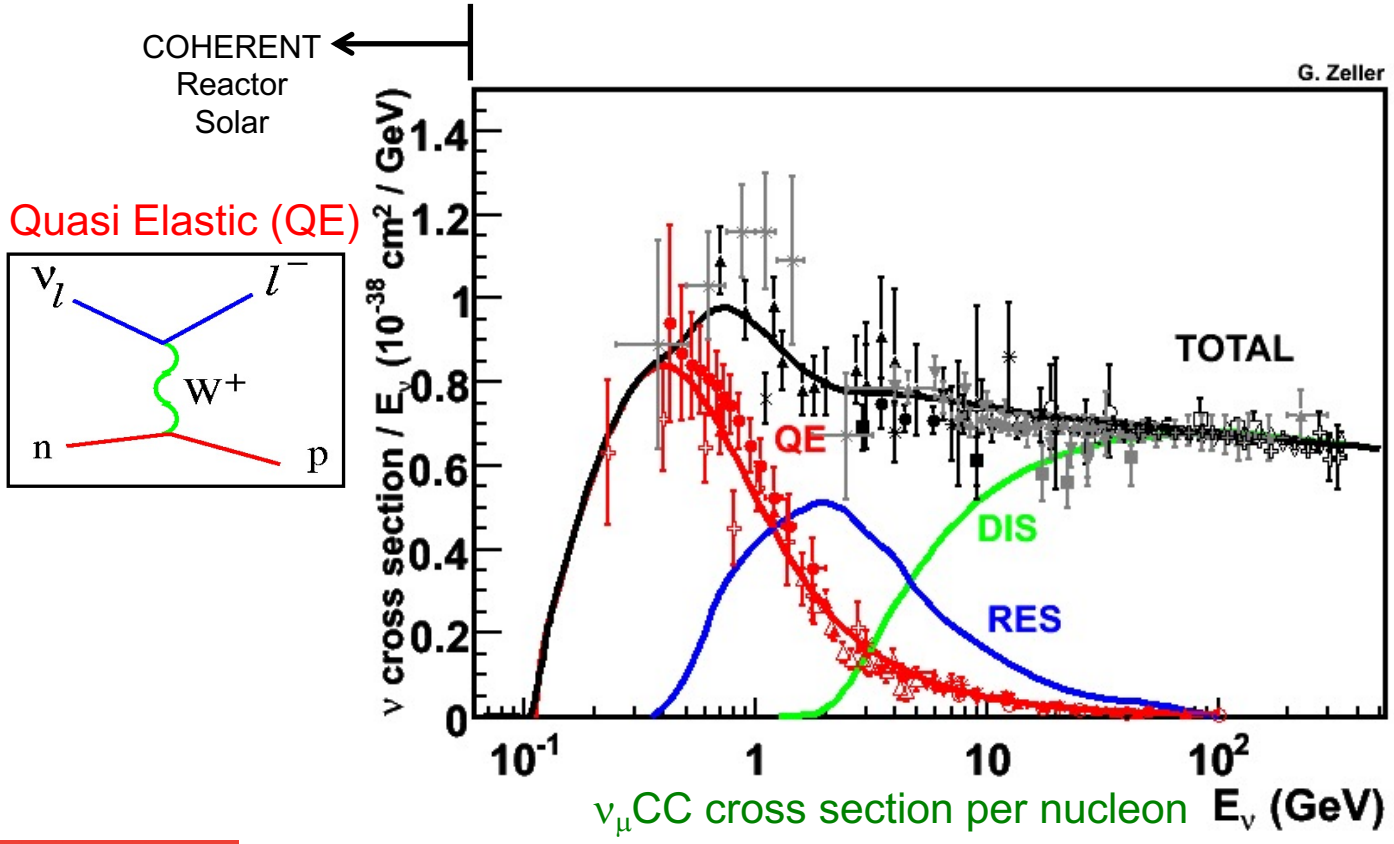
$\gg 1$ GeV neutrino beam
(LHC, astrophysical)



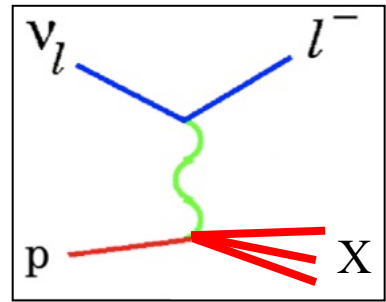
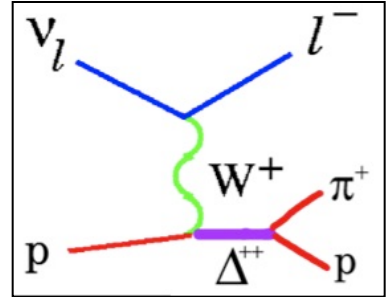
1. Neutrino interaction physics around 1-10 GeV

Neutrino interaction physics around 1-10 GeV

- degree of freedom change from nucleus → nucleon → parton
- There is no such thing (they all interfere)



baryonic
RESONance



Deep Inelastic
Scattering (DIS)

1. Next goal of high energy physics

Establish Neutrino Standard Model (ν SM)

- SM + 3 active massive neutrinos

Unknown parameters of ν SM

1. Dirac CP phase
 2. θ_{23} ($\theta_{23}=40^\circ$ and 50° are same for $\sin 2\theta_{23}$, but not for $\sin\theta_{23}$)
 3. normal mass ordering $m_1 < m_2 < m_3$ or inverted mass ordering $m_3 < m_1 < m_2$
 4. Dirac or Majorana
 5. Majorana phase
 6. Absolute neutrino mass
- } not relevant to neutrino oscillation experiment(?)

We need higher precision neutrino experiments around 1-10 GeV.

Low energy beam (~1 GeV)

- shorter baseline (lower flux reduction)
- lower neutrino production
- lower interaction rate
- kinematic energy reconstruction

High energy beam (~few GeV)

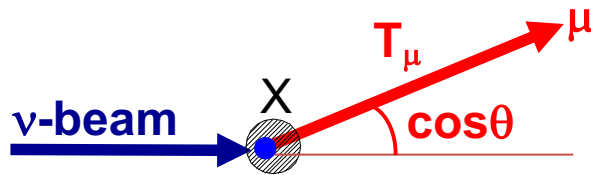
- longer baseline (higher flux reduction)
- higher neutrino production
- higher interaction rate
- calorimetric energy reconstruction

$$P_{\mu \rightarrow e}(L/E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L(km)}{E(GeV)} \right)$$

1. Next goal of high energy physics

Kinematics energy reconstruction

- problem: it assume 2-body neutrino interaction with single nucleon



$$E_\nu^{QE} = \frac{ME_\nu - 0.5m_\mu^2}{M - E_\mu + p_\mu \cos\theta}$$

Low energy beam (~1 GeV)

- shorter baseline (lower flux reduction)
- lower neutrino production
- lower interaction rate
- kinematic energy reconstruction

Calorimetric energy reconstruction

- problem: you have to measure energy deposit from all outgoing particles

$$E_\nu^{Cal} = E_\mu + \sum_{i=1}^{all} E_{had}^i$$

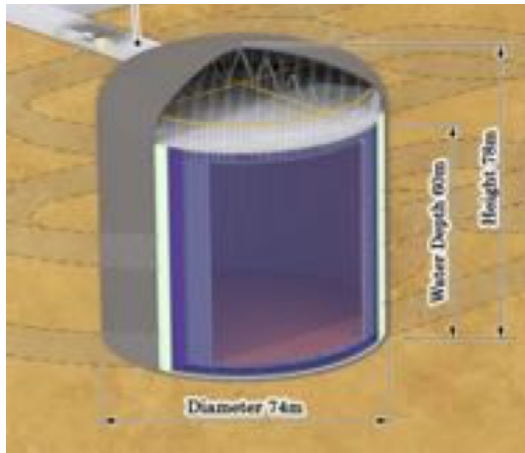
High energy beam (~few GeV)

- longer baseline (higher flux reduction)
- higher neutrino production
- higher interaction rate
- calorimetric energy reconstruction

1. Next goal of high energy physics

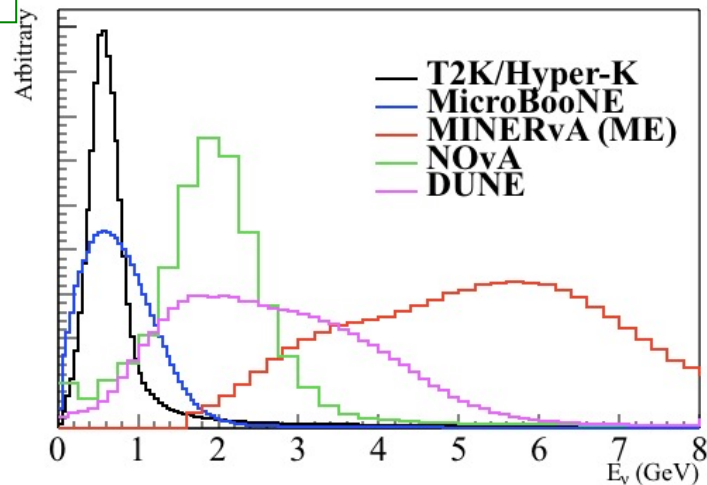
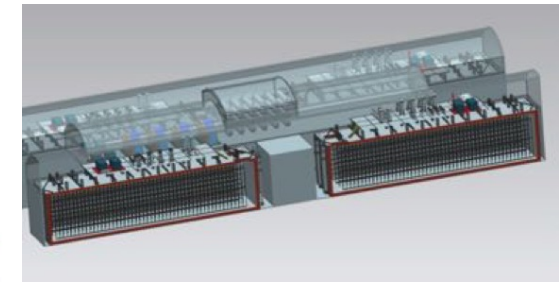
Hyper-Kamiokande (Japan)

- Water target
- Narrow band 0.6 GeV
- Low spatial resolution
- High time resolution



DUNE (USA)

- Argon target
- wide band 1-4 GeV
- High spatial resolution
- Low time resolution



Low energy beam (~1 GeV)

- shorter baseline (lower flux reduction)
- lower neutrino production
- lower interaction rate
- kinematic energy reconstruction

High energy beam (~few GeV)

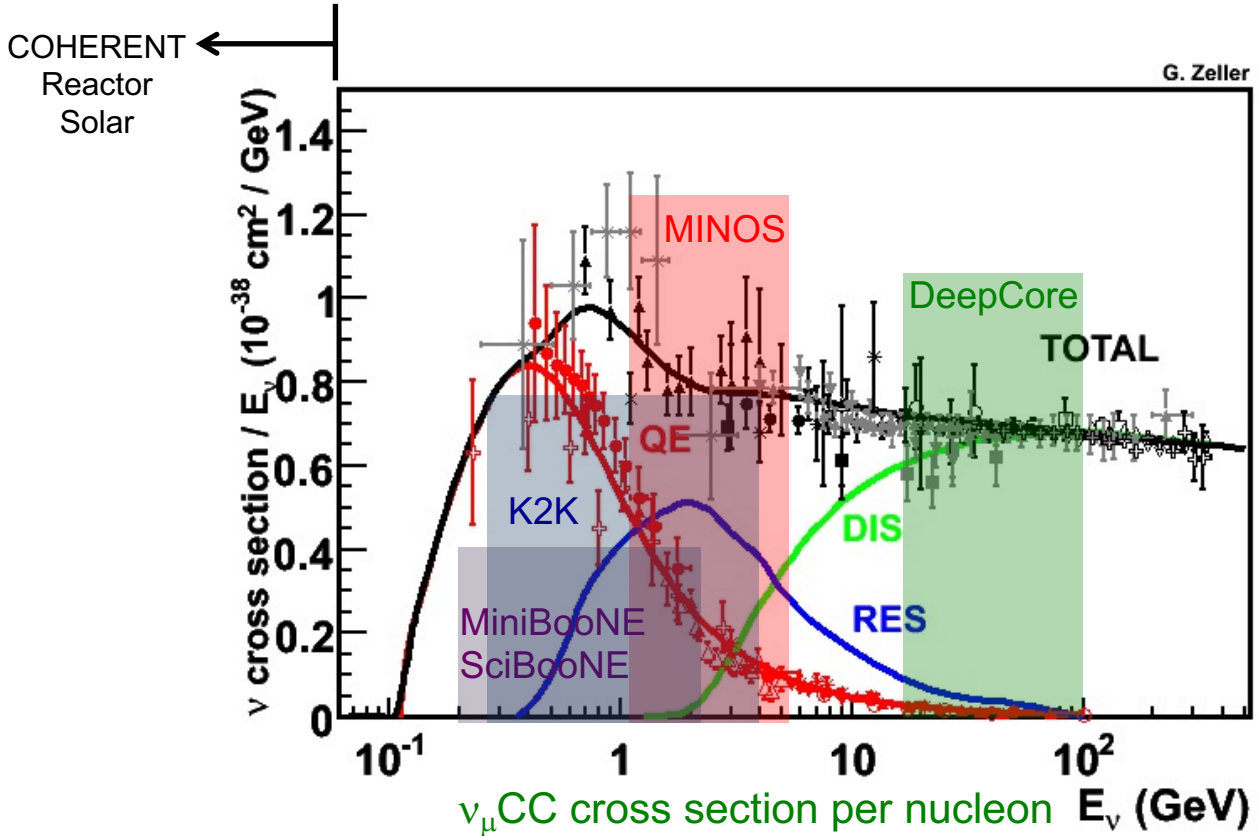
- longer baseline (higher flux reduction)
- higher neutrino production
- higher interaction rate
- calorimetric energy reconstruction

$$P_{\mu \rightarrow e}(L/E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L(km)}{E(GeV)} \right)$$

1. Next generation neutrino oscillation experiments

Neutrino oscillation experiments

- Past: K2K, MiniBooNE, MINOS, DeepCore
- Present to Future: T2K, NOvA, IceCube-Upgrade, ORCA, Hyper-Kamiokande, DUNE...

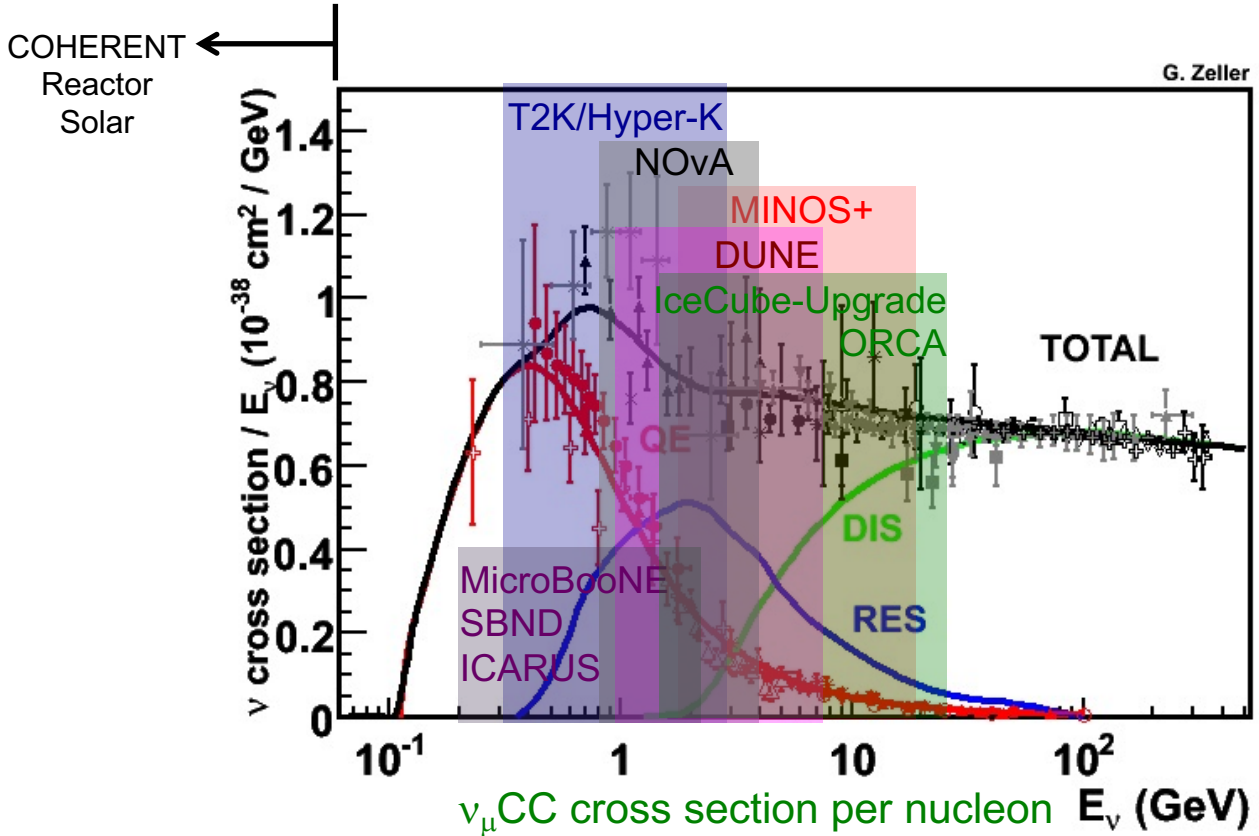


$$P_{\mu \rightarrow e}(L/E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L(km)}{E(GeV)} \right)$$

1. Next generation neutrino oscillation experiments

Neutrino oscillation experiments

- Past: K2K, MiniBooNE, MINOS, DeepCore
- Present to Future: T2K, NOvA, IceCube-Upgrade, ORCA, Hyper-Kamiokande, DUNE...



$$P_{\mu \rightarrow e}(L/E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L(km)}{E(GeV)} \right)$$

1. Neutrino cross-section formula

Cross-section

- product of Leptonic and Hadronic tensor

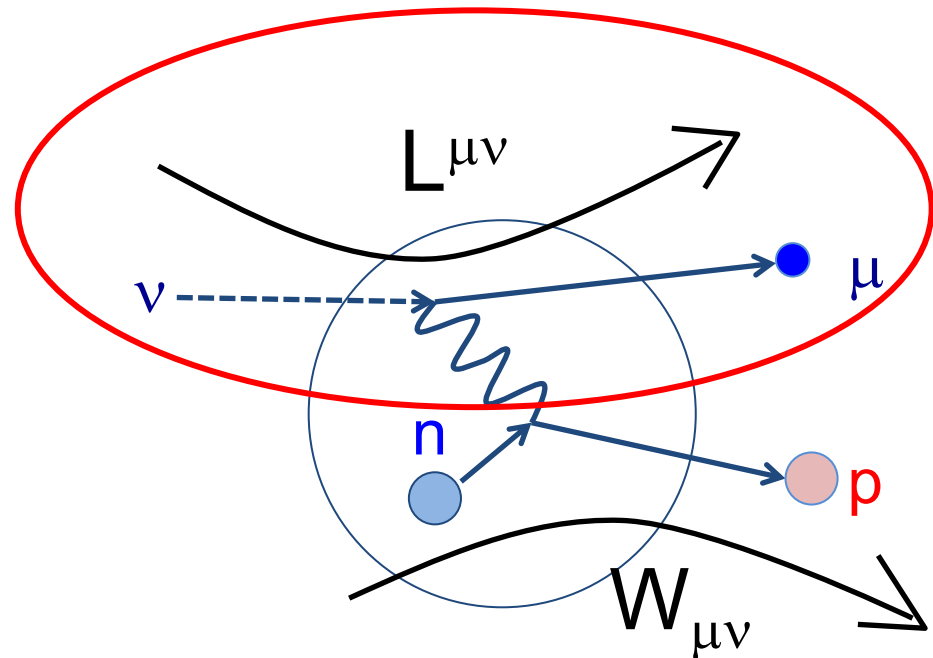
$$d\sigma \sim L^{\mu\nu}W_{\mu\nu}$$

Leptonic tensor

→ the Standard Model (easy)

Hadronic tensor

→ nuclear physics (hard)



1. Neutrino cross-section formula

Cross-section

- product of Leptonic and Hadronic tensor

$$d\sigma \sim L^{\mu\nu}W_{\mu\nu}$$

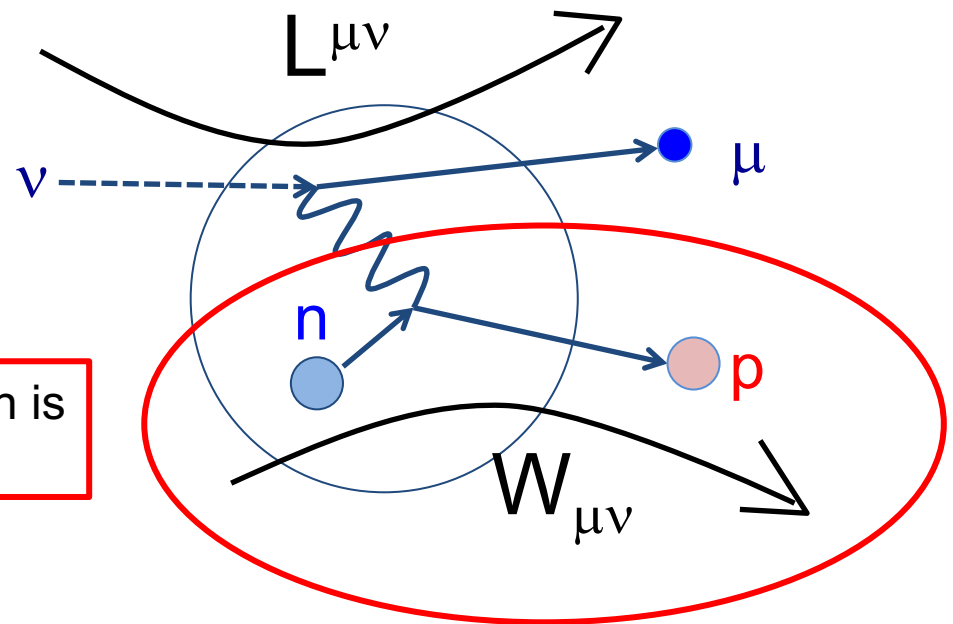
Leptonic tensor

→ the Standard Model (easy)

Hadronic tensor

→ nuclear physics (hard)

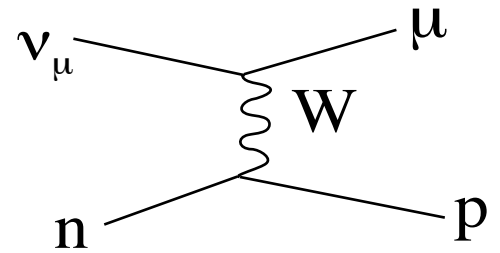
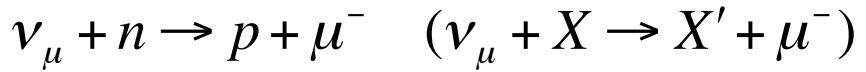
All complication of neutrino cross-section is how to model the hadronic tensor part



1. Neutrino interaction physics - introduction
- 2. Charged-Current Quasi-Elastic (CCQE) interaction**
3. Neutrino baryonic resonance interaction
4. Neutrino shallow- and deep-inelastic scatterings
5. Conclusions

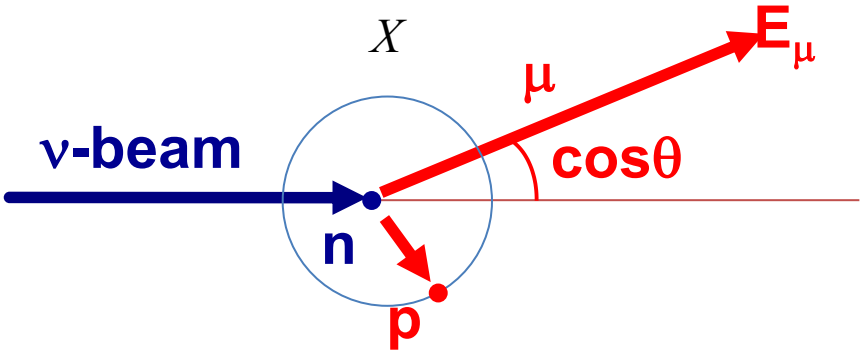
2. Charged Current Quasi-Elastic scattering (CCQE)

The simplest and the most abundant interaction around ~1 GeV.



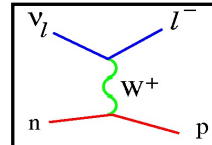
Neutrino energy is reconstructed from the observed lepton kinematics
 “QE assumption”

1. assuming neutron at rest
2. assuming interaction is CCQE



$$E_\nu^{QE} = \frac{ME_\nu - 0.5m_\mu^2}{M - E_\mu + p_\mu \cos\theta}$$

CCQE is the single most important channel of neutrino oscillation physics
 T2K, NOvA, microBoonE, Hyper-Kamiokande...etc



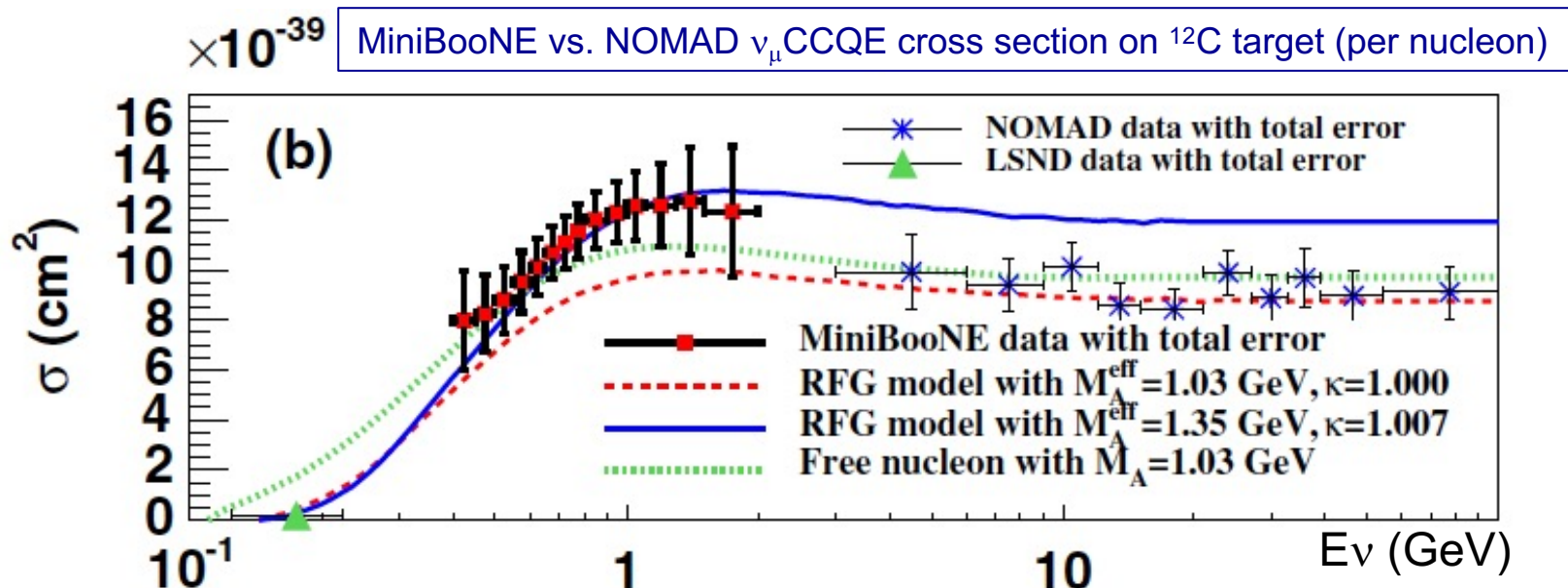
2. CCQE puzzle

Simplest channel, but both shape and normalization disagree

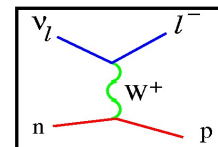
1. low Q^2 suppression \rightarrow Low forward efficiency? (detector?)
2. high Q^2 enhancement \rightarrow Axial mass > 1.0 GeV? (physics?)
3. large normalization \rightarrow Beam simulation is wrong? (flux?)

CCQE interaction on nuclear targets are precisely measured by electron scattering

- Lepton universality = precise prediction for neutrino CCQE cross-section...?



2. Solution of CCQE puzzle

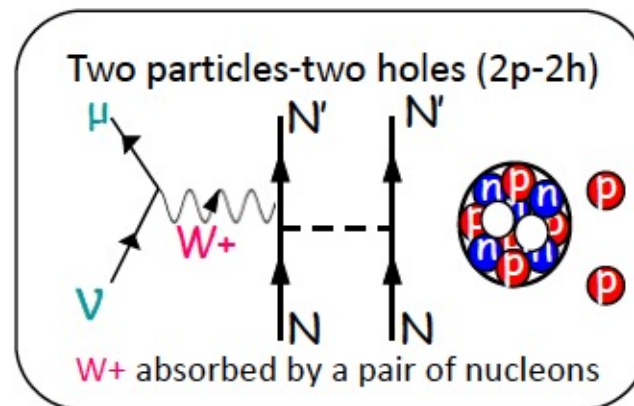
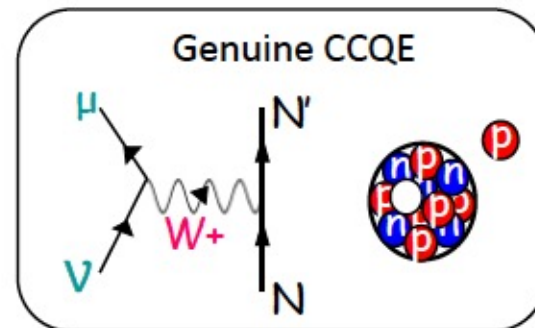
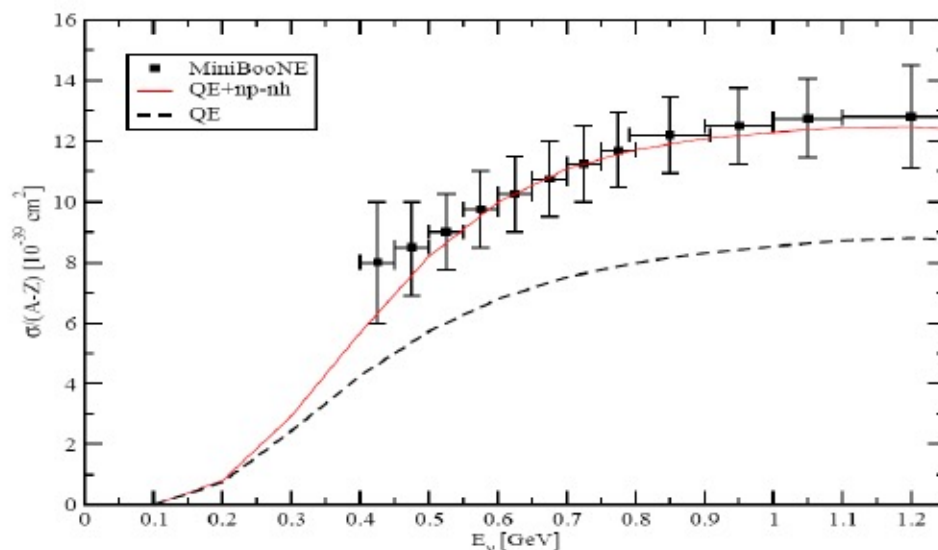


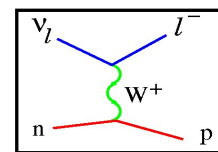
Presence of 2-body current

- Martini et al showed 2p-2h effect can add up 30-40% more cross section!

An explanation of this puzzle

Inclusion of the multinucleon emission channel (np-nh)





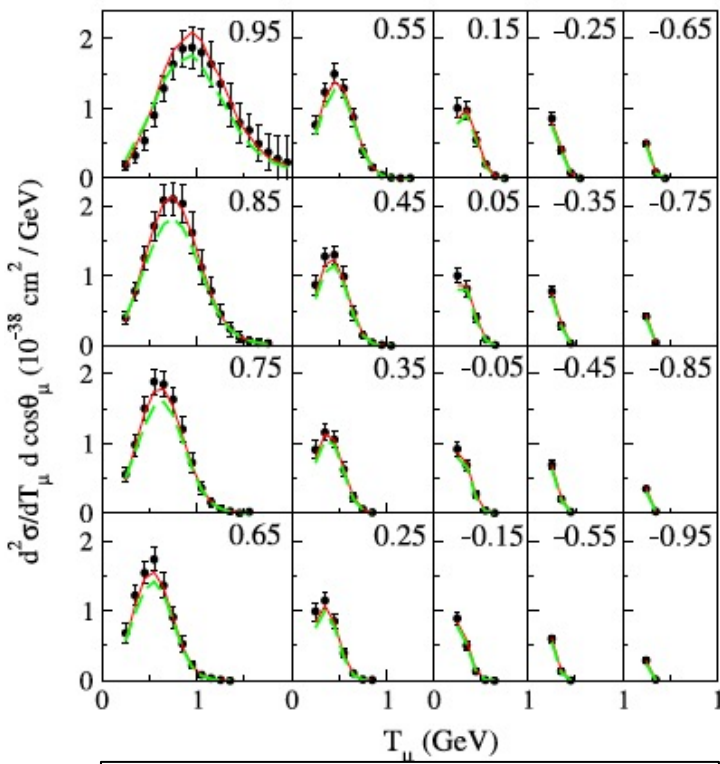
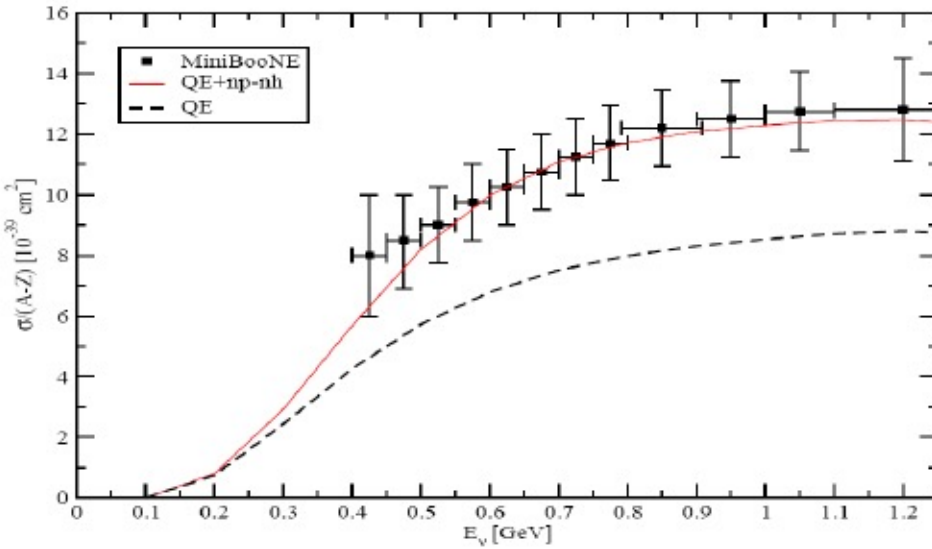
2. Solution of CCQE puzzle

Presence of 2-body current

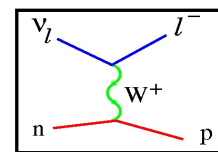
- Martini et al showed 2p-2h effect can add up 30-40% more cross section!
- consistent result is obtained by Nieves et al

An explanation of this puzzle

Inclusion of the multinucleon emission channel (np-nh)



Valencia model vs. MiniBooNE CCQE double differential cross-section data

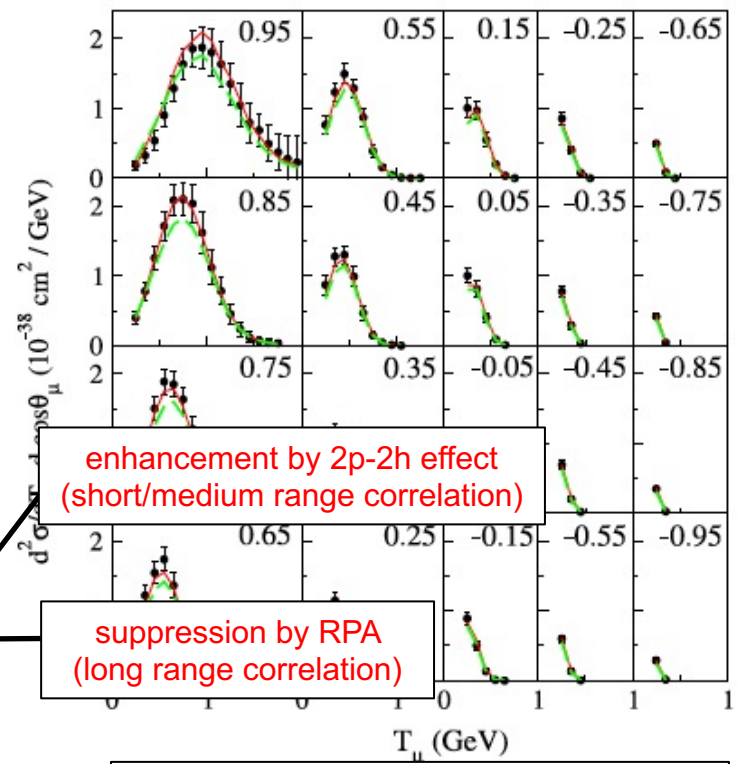
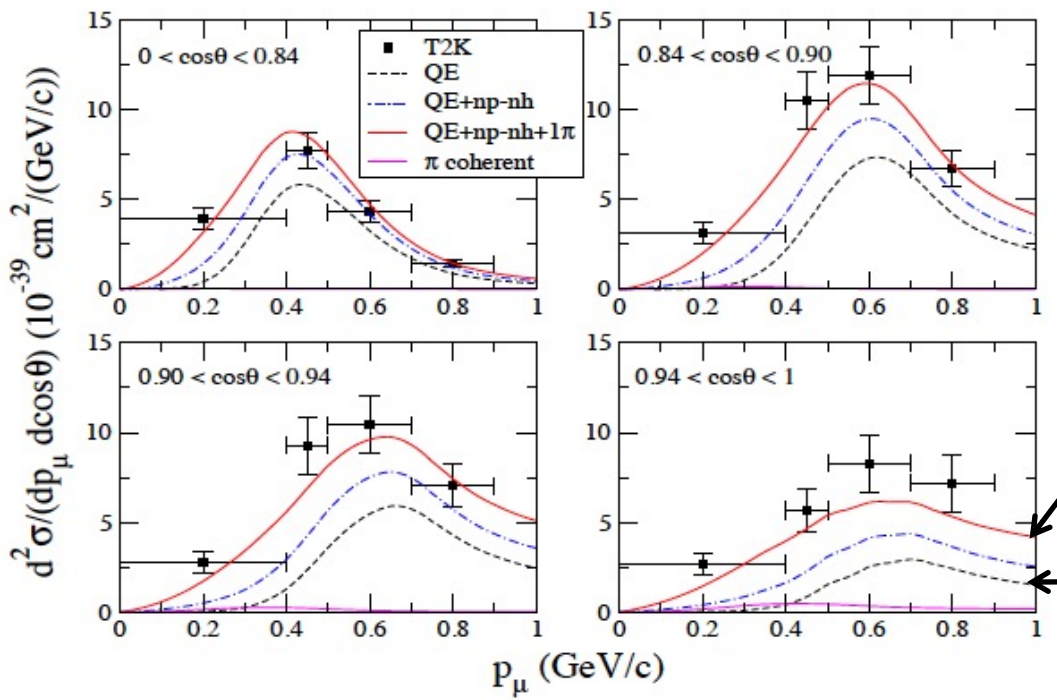


2. Solution of CCQE puzzle

Presence of 2-body current

- Martini et al showed 2p-2h effect can add up 30-40% more cross section!
- consistent result is obtained by Nieves et al
- The model can explain T2K data simultaneously

Martini model vs. T2K CC double differential cross-section data



Valencia model vs. MiniBooNE CCQE double differential cross-section data

2. Solution of CCQE puzzle

Many phenomenological models agree **qualitatively** with MiniBooNE CCQE-like double differential data.

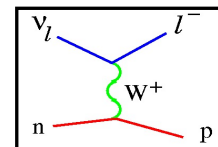
Martini – RPA+2p2h

Nieves – Valencia 2p2h model

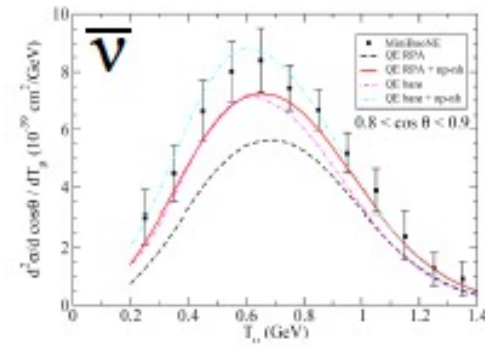
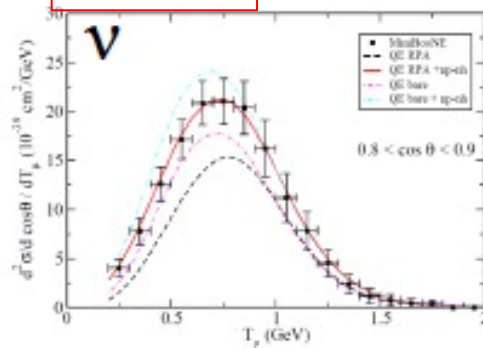
SuSA – Superscaling+MEC

Giusti – Relativistic Green's function

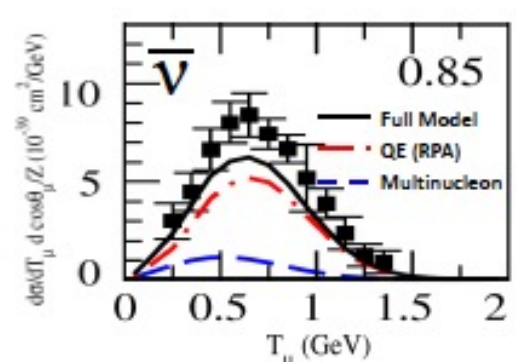
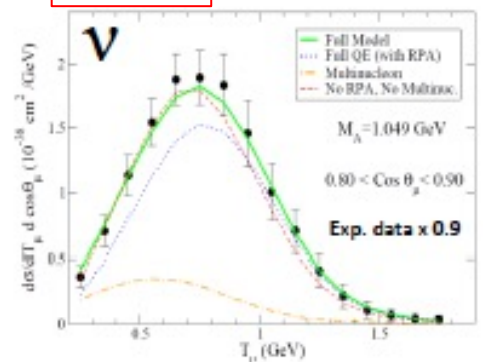
Butkevich – RDWIA+MEC



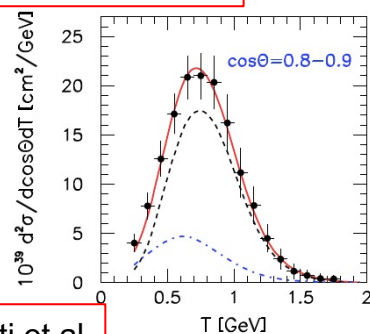
Martini et al



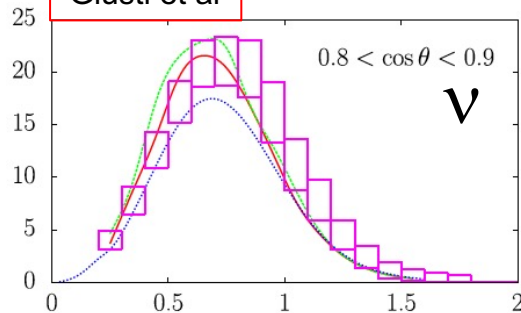
Valencia



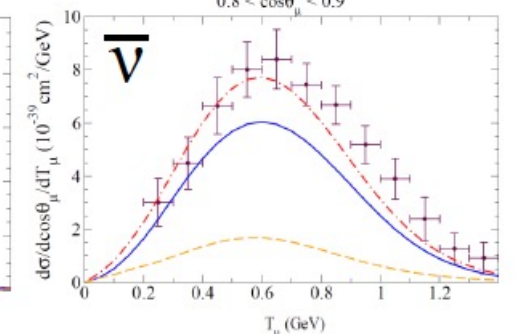
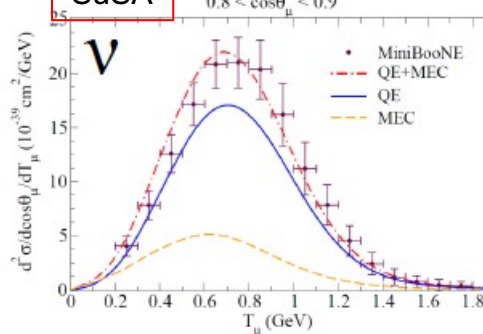
Butkevich et al

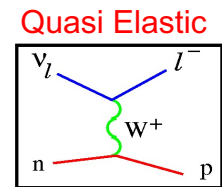


Giusti et al



SuSA





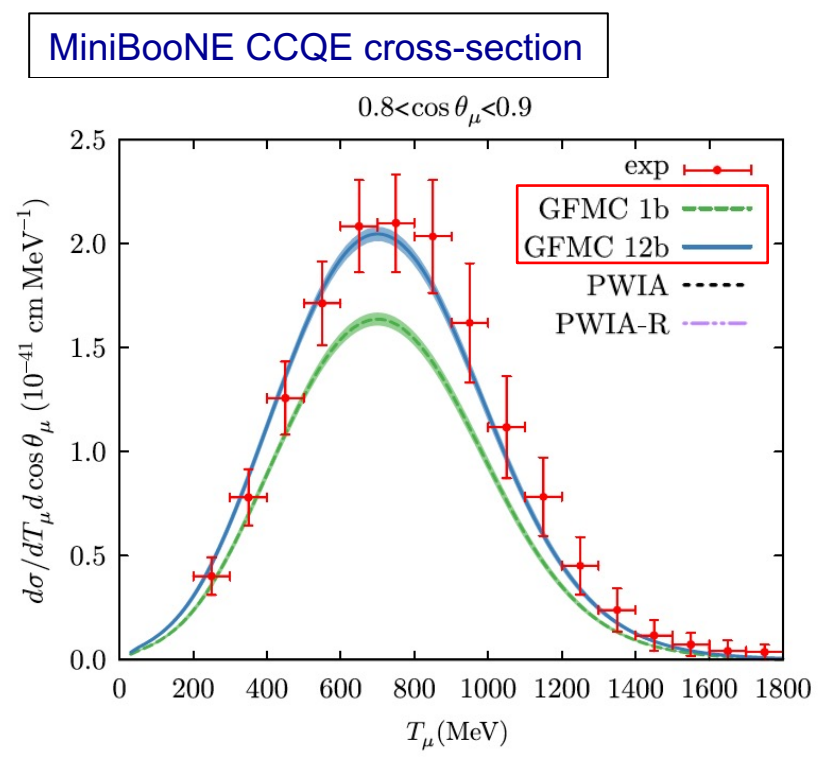
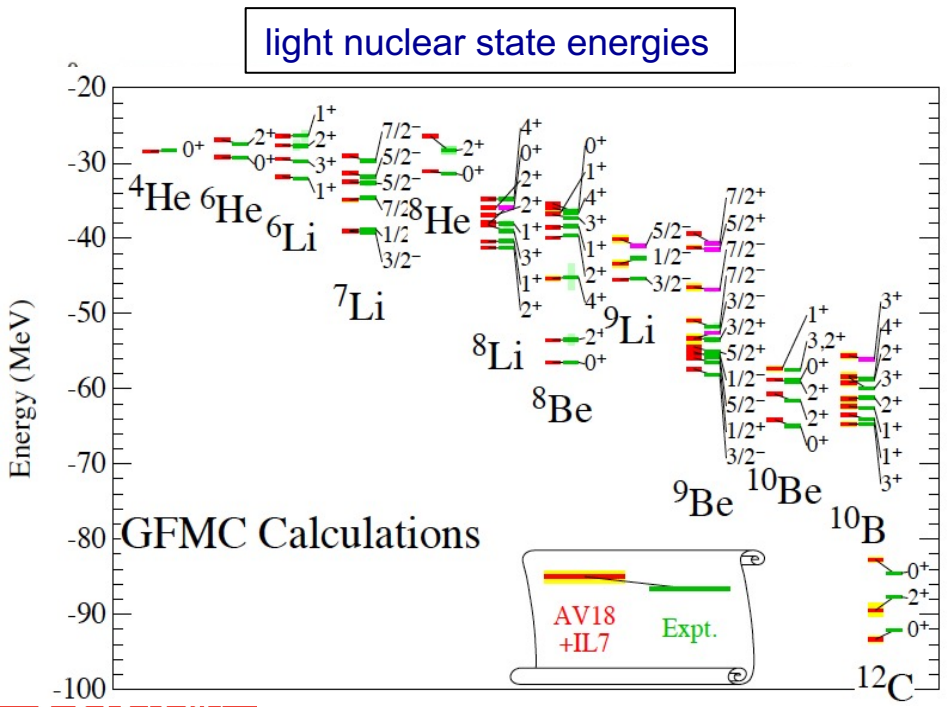
2. Nucleon correlations in neutrino physics

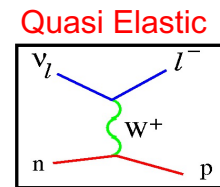
Ab-initio calculation

- Quantum Monte Carlo (QMC)
- Predicts energy levels of all light nuclei
- Consistent result with phenomenological models
- Ground state includes correct nucleon correlations

$$|\Psi_V\rangle = \mathcal{S} \prod_{i < j}^A \left[1 + \boxed{U_{ij}} + \sum_{k \neq i, j}^A \boxed{\tilde{U}_{ijk}^{TNI}} \right] |\Psi_J\rangle$$

2N potential (Av18) 3N potential (IL7)





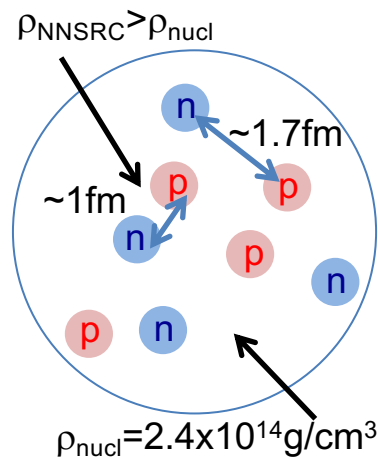
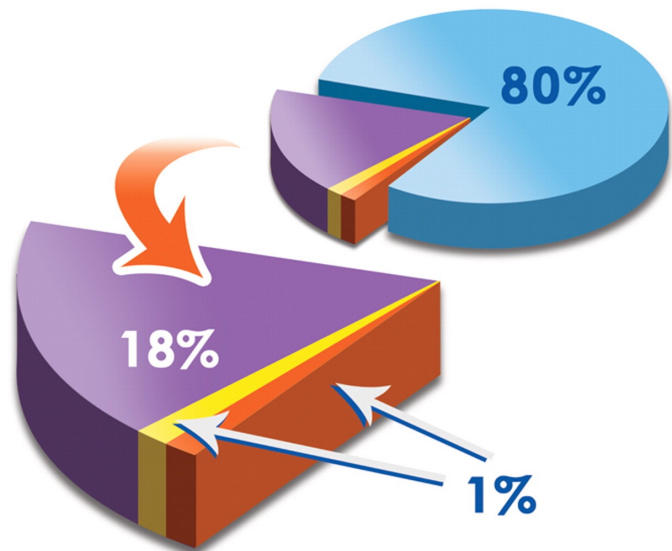
2. Nucleon correlations in neutrino physics

Ab-initio calculation

- Quantum Monte Carlo (QMC)
- Predicts energy levels of all light nuclei
- Consistent result with phenomenological models
- Ground state includes correct nucleon correlations

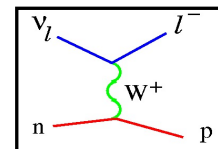
$$|\Psi_V\rangle = \mathcal{S} \prod_{i < j}^A \left[1 + \boxed{U_{ij}} + \sum_{k \neq i, j}^A \boxed{\tilde{U}_{ijk}^{TNI}} \right] |\Psi_J\rangle$$

2N potential (Av18) 3N potential (IL7)



Physics of nucleon correlation

- neutrino interaction
- EMC effect
- $0\nu\beta\beta$
- Direct WIMP detection
- etc

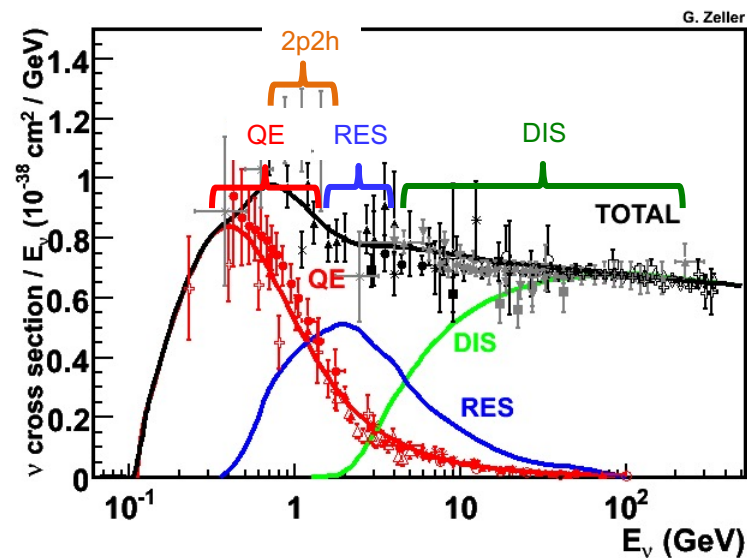
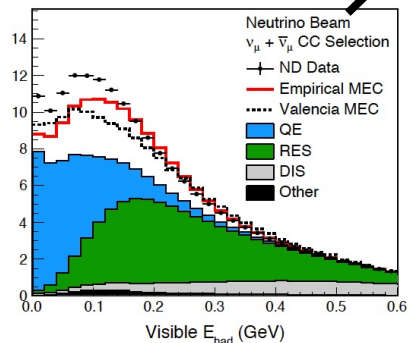
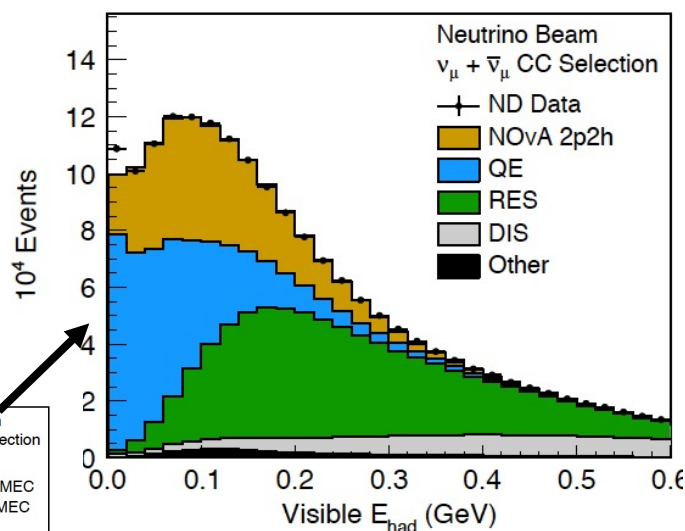


2. Nucleon correlations in neutrino physics

2-particle 2-hole (2p2h) effect

- Essential to describe data
- The biggest topic in nuxsec community (T2K, NOvA, MINERvA, MicroBooNE, etc)
- 2p2h models in generators don't describe data well?
- High resolution detector (LArTPC, emulsion, etc) can find what is going on?

NOvA near detector data-MC comparison after fit



2. Nucleon correlations in neutrino physics

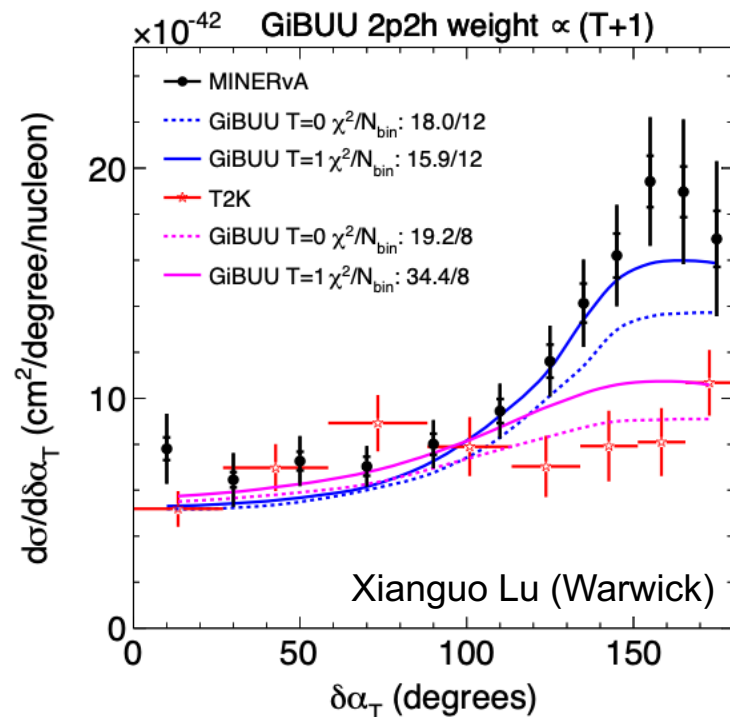
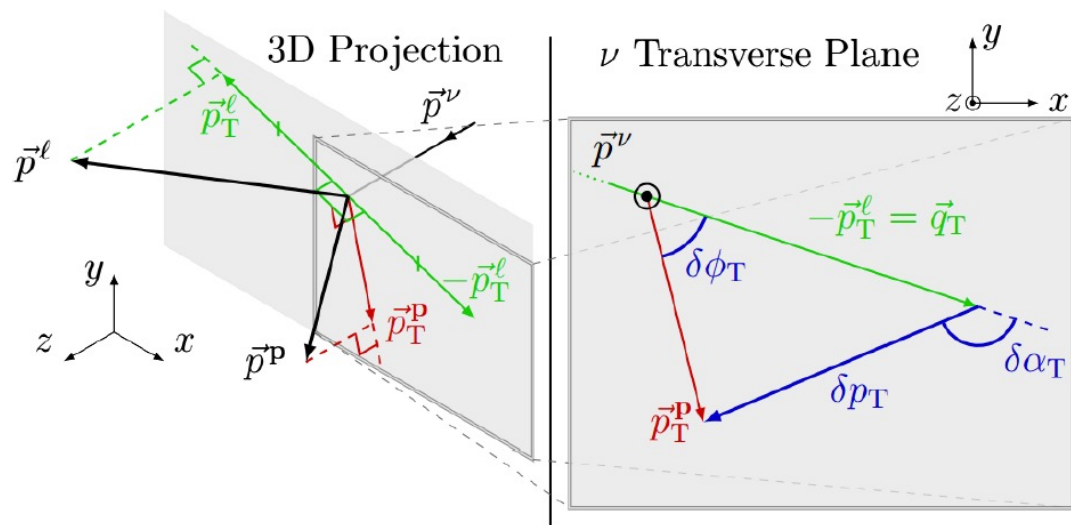
We want to constrain nuclear model from neutrino data

- Final state hadron measurement is the key

1 muon + 1 proton sample

- 5 dof (mu E and $\cos\theta$, proton E and $\cos\theta$, mu-p opening angle).
- Low statistics, and these are converted to 3 kinematic variables.

Data prefer advanced nuclear models, but it's not easy to identify 2p2h model



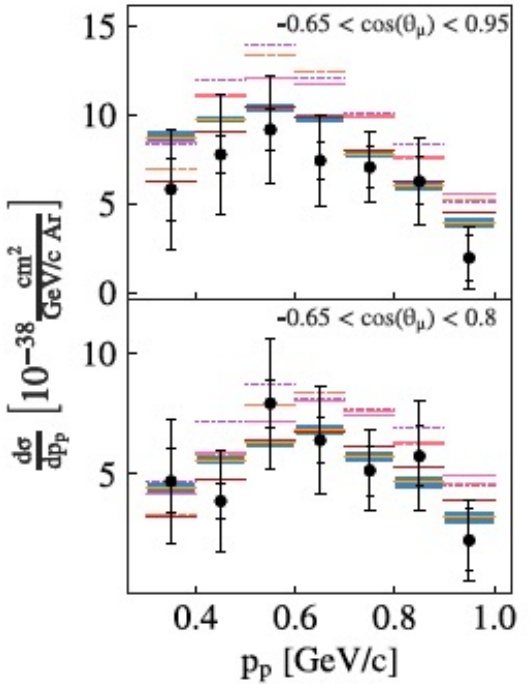
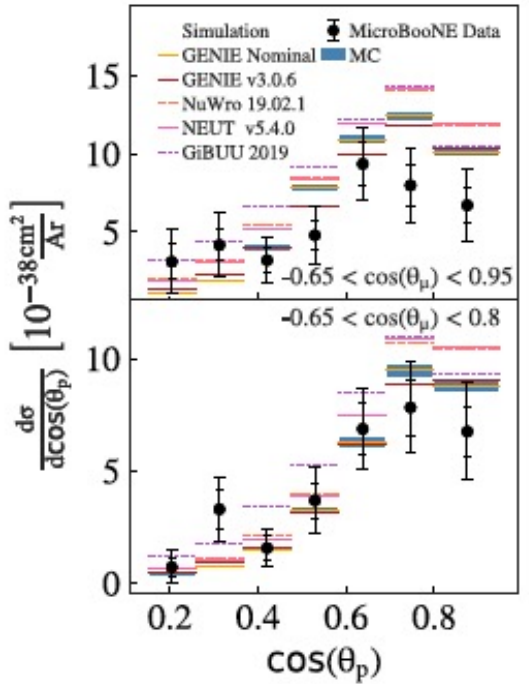
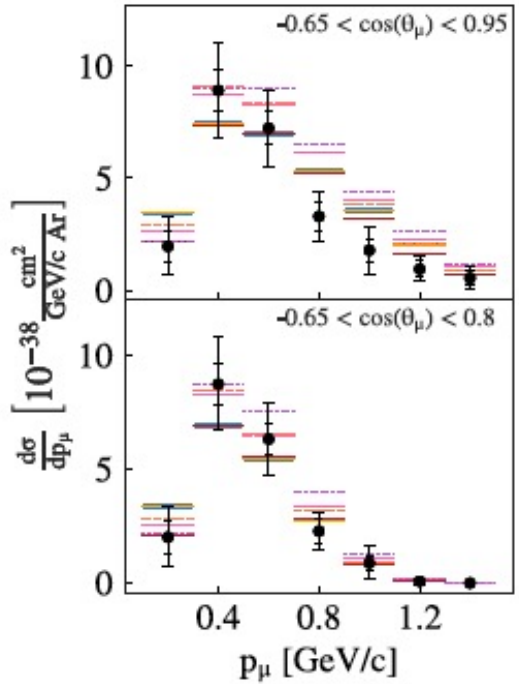
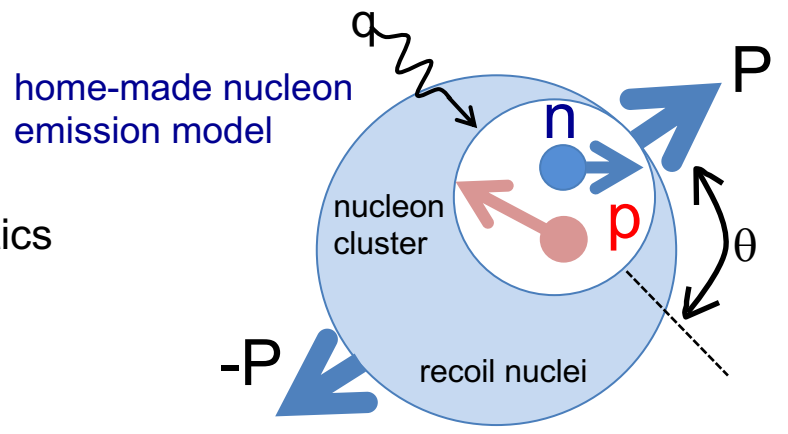
Importance of axial 2BC is understood qualitatively, we need more quantitative understanding!
 (how much 2-body current? Any characteristic shapes in kinematic variables?)

2. Nucleon correlations in neutrino physics

There is a strong belief in experimental community that hadron final states tell everything about 2p2h...

We need prediction of hadronic final states from theorists

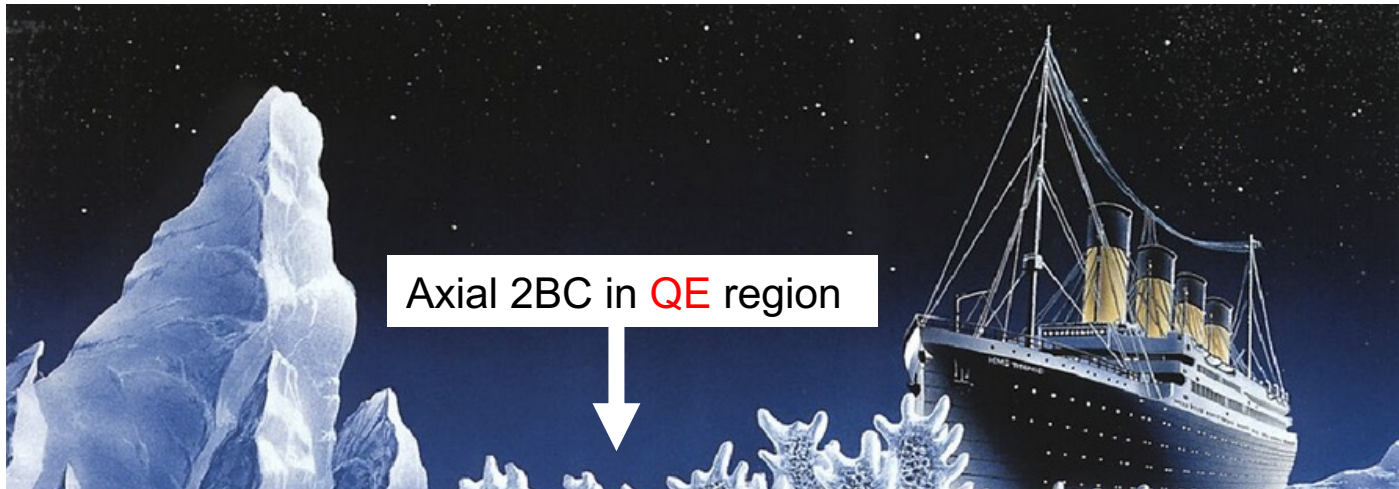
- double differential cross-section = lepton kinematics
- final hadron multiplicity/kinematics = home-made



1. Neutrino interaction physics - introduction
2. Charged-Current Quasi-Elastic (CCQE) interaction
- 3. Neutrino baryonic resonance interaction**
4. Neutrino shallow- and deep-inelastic scatterings
5. Conclusions

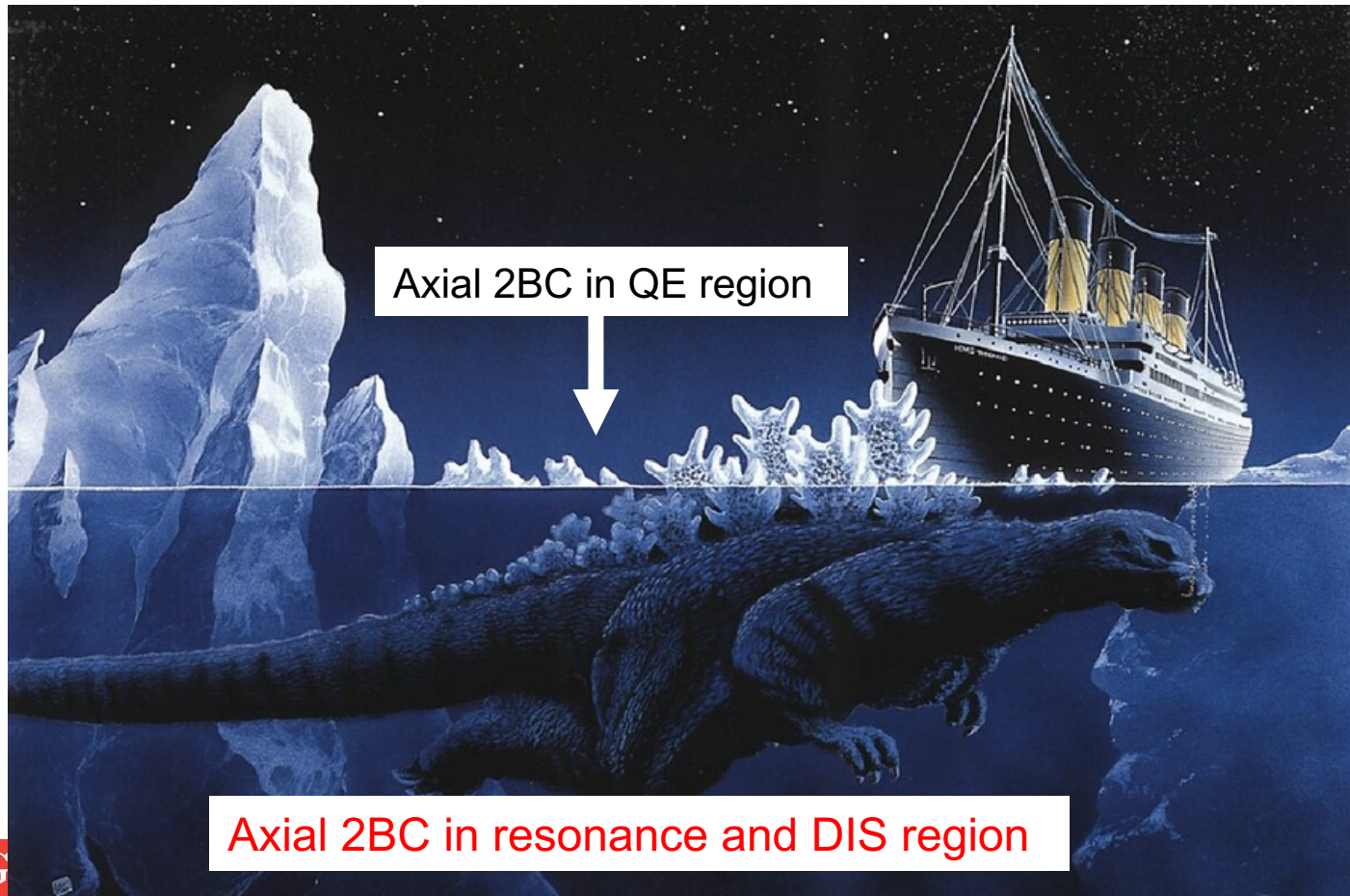
3. Beyond QE peak

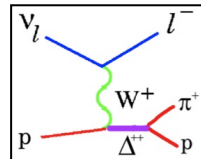
Axial 2-body current in QE region may be a tip of the iceberg...



3. Beyond QE peak

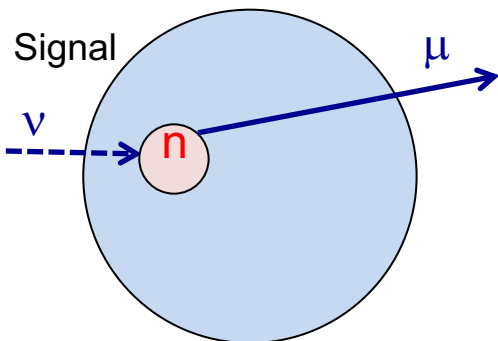
Axial 2-body current in QE region may be a tip of the iceberg..., or maybe a tip of gozilla!





3. non-QE background (resonance pion production)

non-QE background \rightarrow shift spectrum



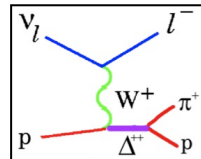
Neutrino energy is reconstructed from the observed lepton kinematics

“QE assumption”

1. assuming neutron at rest
2. assuming interaction is CCQE

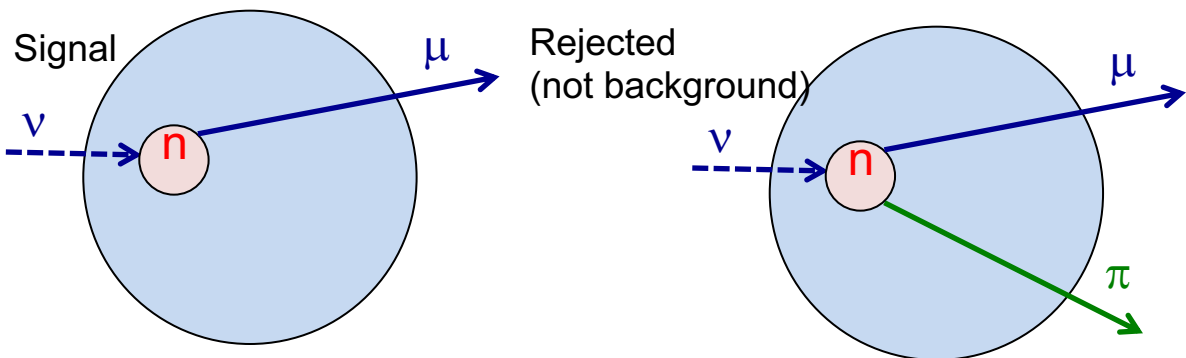
Typical neutrino oscillation detector

- Big and dense, to maximize interaction rate
- Coarsely instrumented, to minimize cost (not great detector to measure hadrons)



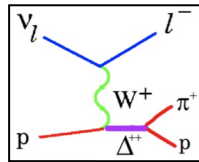
3. non-QE background (resonance pion production)

non-QE background \rightarrow shift spectrum



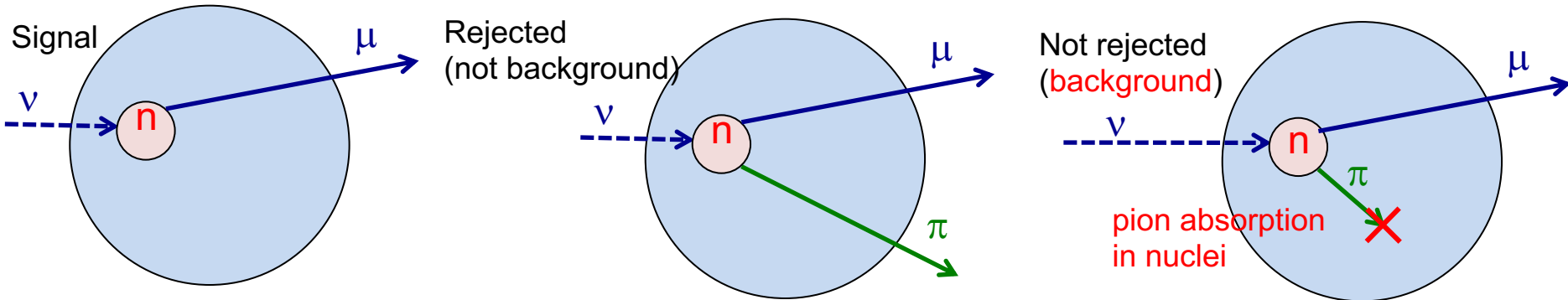
Typical neutrino oscillation detector

- Big and dense, to maximize interaction rate
- Coarsely instrumented, to minimize cost (not great detector to measure hadrons)



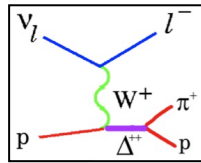
3. non-QE background (resonance pion production)

non-QE background \rightarrow shift spectrum



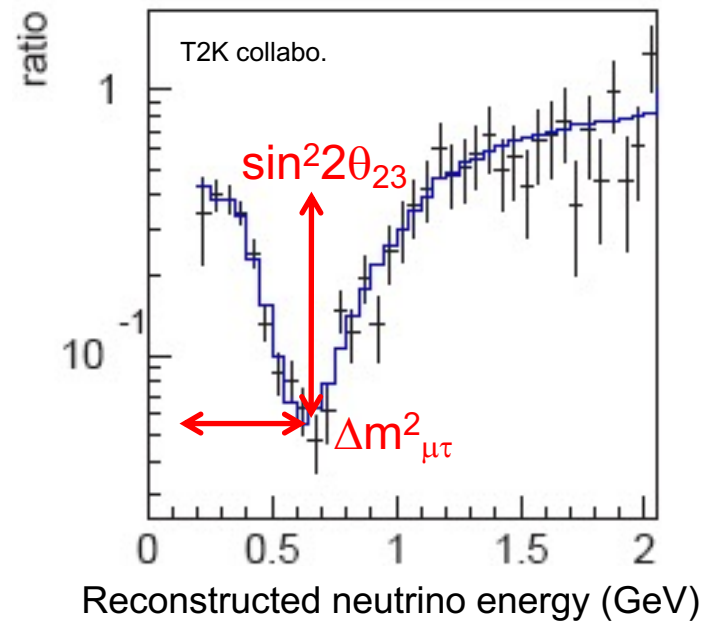
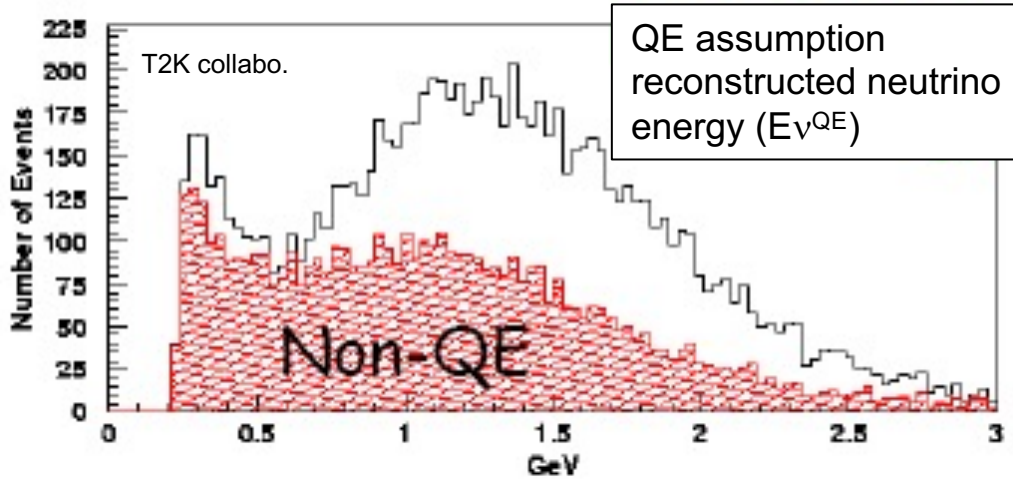
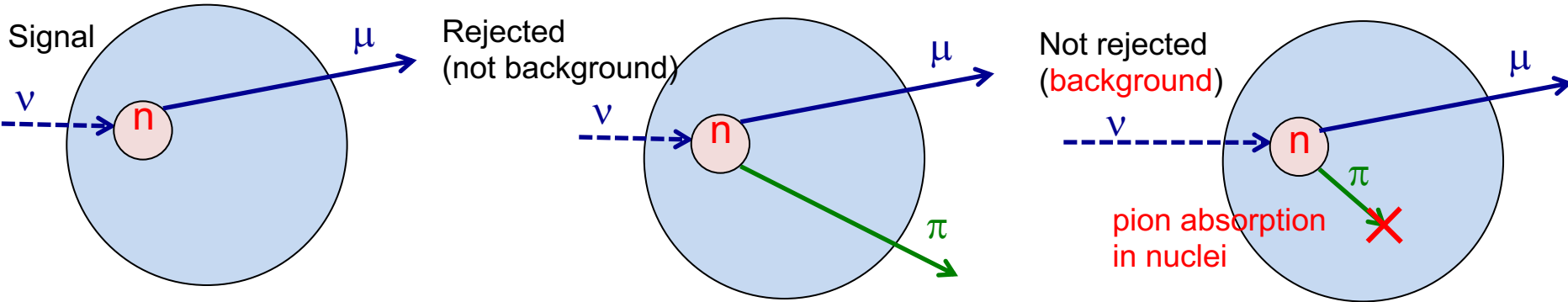
Typical neutrino oscillation detector

- Big and dense, to maximize interaction rate
- Coarsely instrumented, to minimize cost (not great detector to measure hadrons)

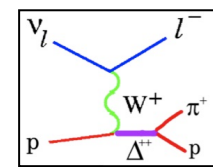


3. non-QE background (resonance pion production)

non-QE background \rightarrow shift spectrum



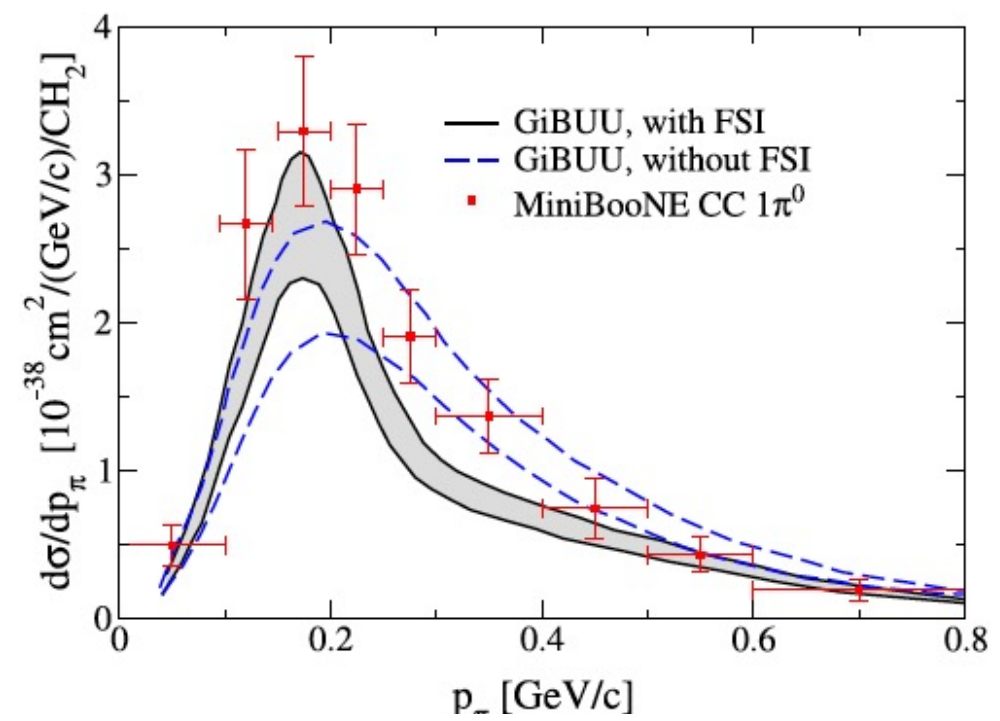
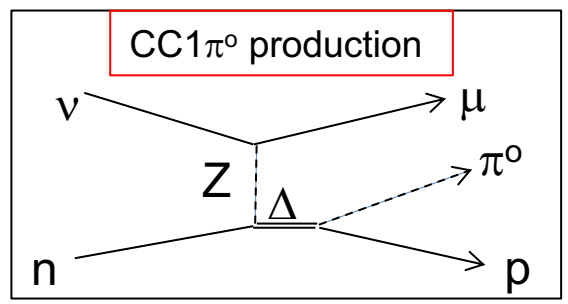
Solution: you need a good prediction of out-going hadron final states (hard)



3. Pion puzzle

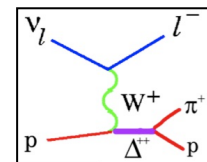
Final state interaction

- Cascade model as a standard of the community
- Advanced models are not available for event-by-event simulation



ex) Giessen BUU transport model

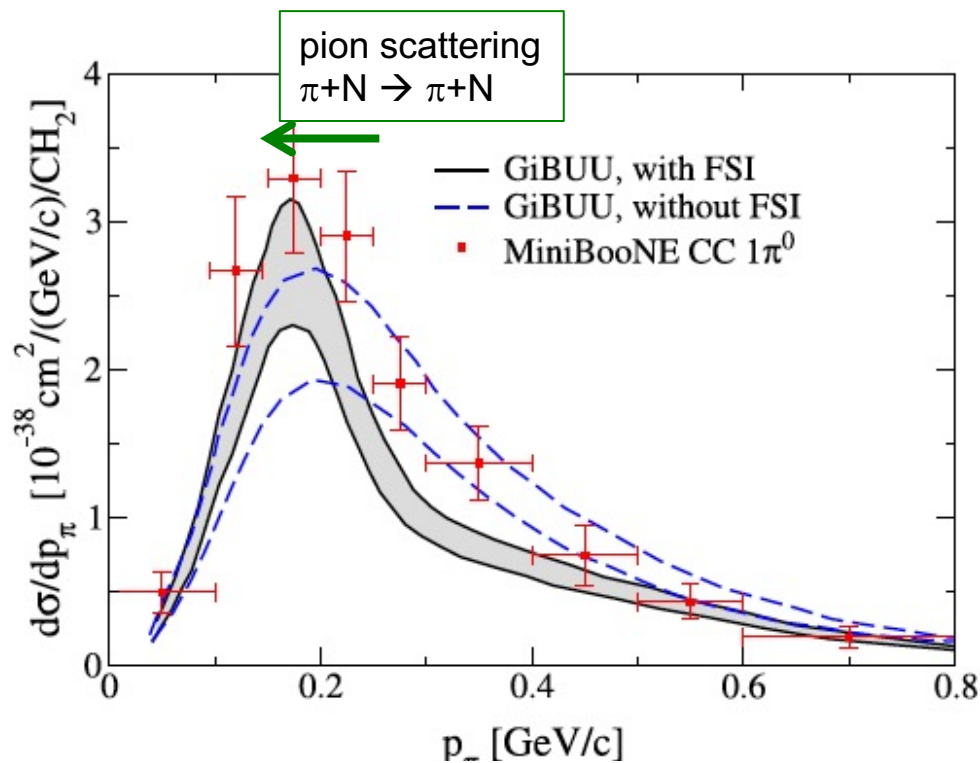
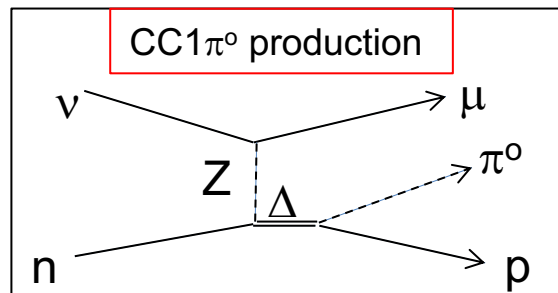
- Developed for heavy ion collision, and now used to calculate final state interactions of pions in nuclear media



3. Pion puzzle

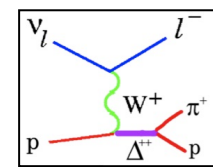
Final state interaction

- Cascade model as a standard of the community
- Advanced models are not available for event-by-event simulation



ex) Giessen BUU transport model

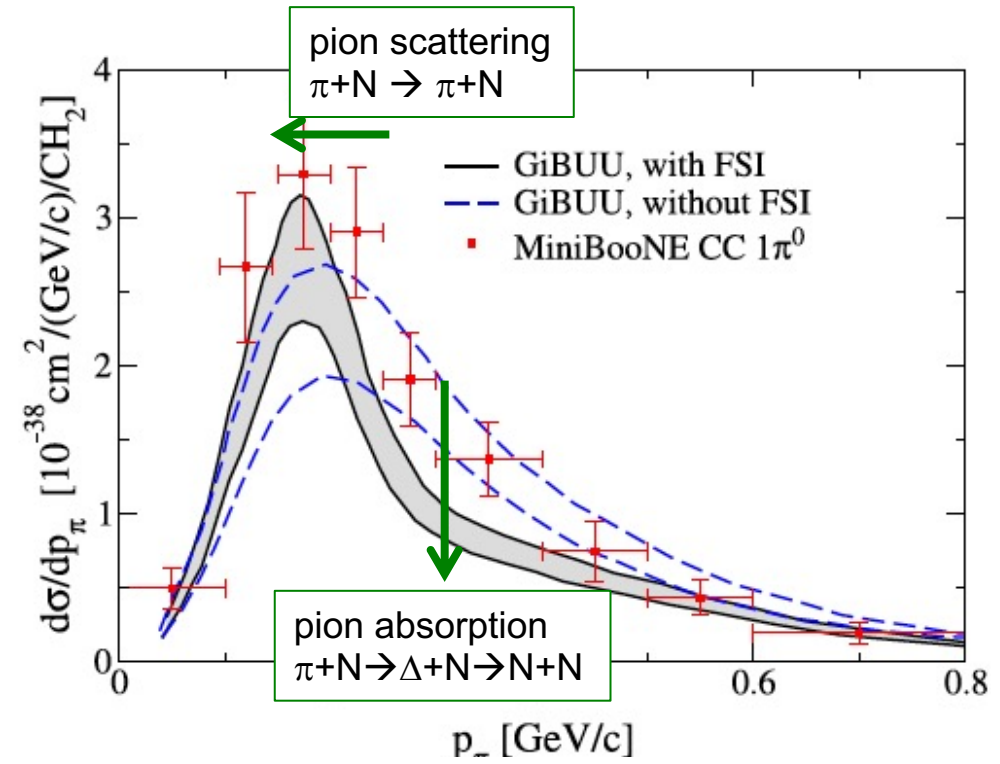
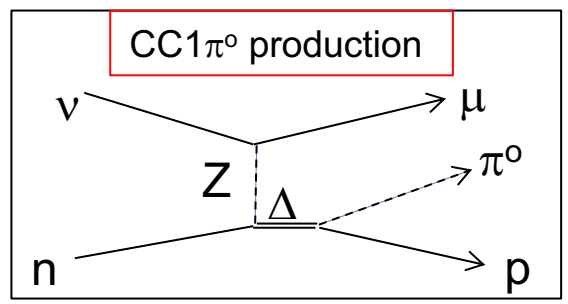
- Developed for heavy ion collision, and now used to calculate final state interactions of pions in nuclear media



3. Pion puzzle

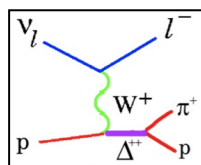
Final state interaction

- Cascade model as a standard of the community
- Advanced models are not available for event-by-event simulation



ex) Giessen BUU transport model

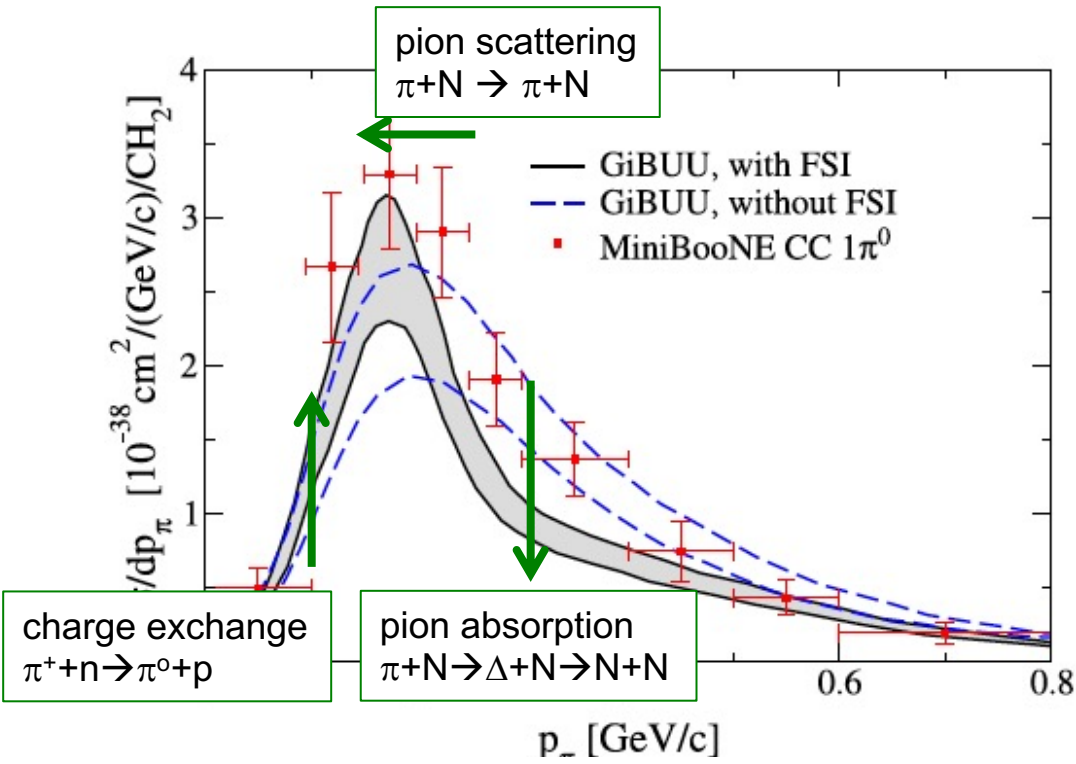
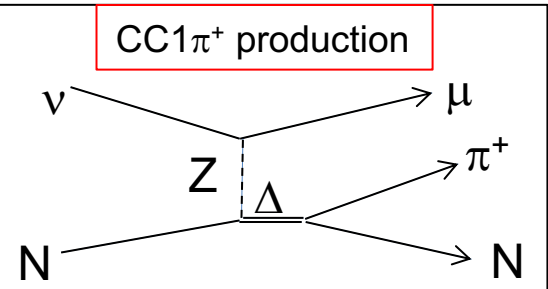
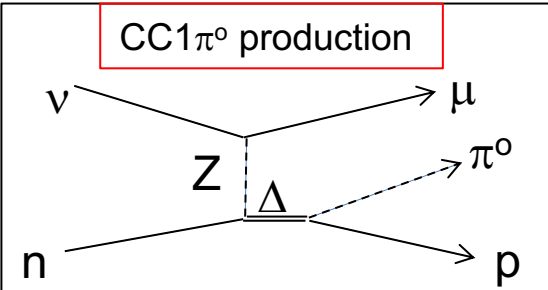
- Developed for heavy ion collision, and now used to calculate final state interactions of pions in nuclear media



3. Pion puzzle

Final state interaction

- Cascade model as a standard of the community
- Advanced models are not available for event-by-event simulation



ex) Giessen BUU transport model

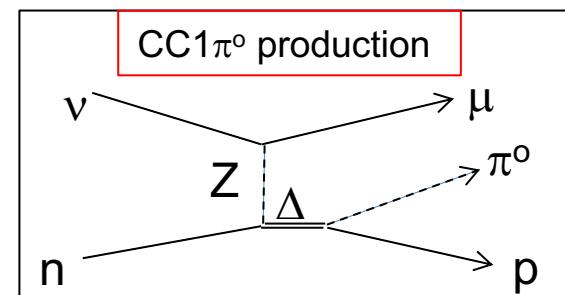
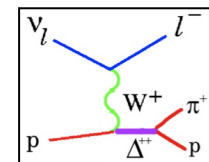
- Developed for heavy ion collision, and now used to calculate final state interactions of pions in nuclear media

You need to predict both
 1. pion production model
 2. final state interaction

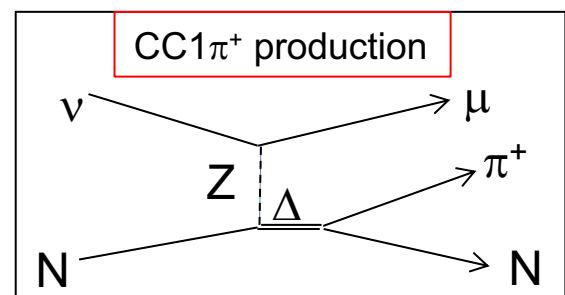
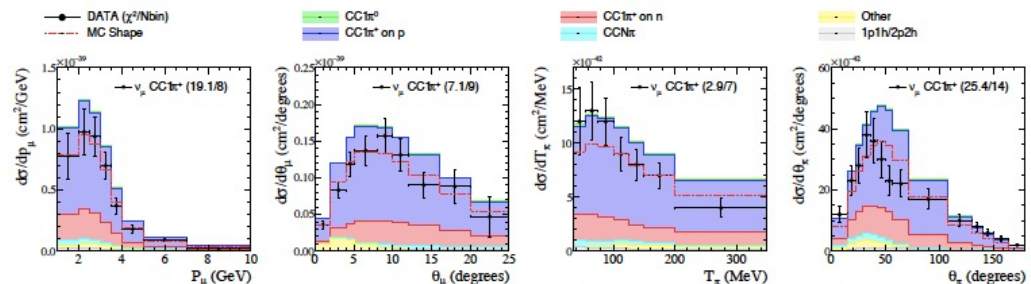
3. Pion puzzle

MINERvA try to fit 4 different data set to tune MC

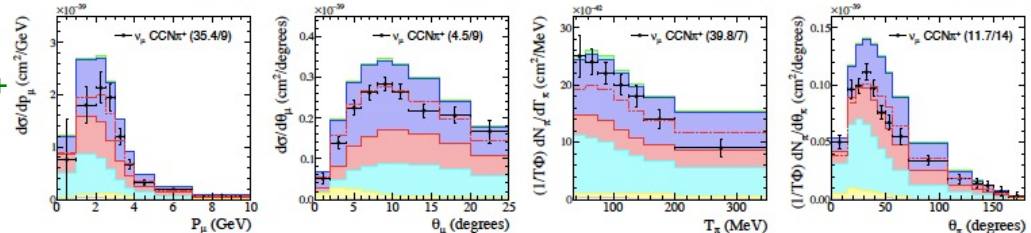
- Strong tensions between data set
- Both cross section and FSI models need to be improved



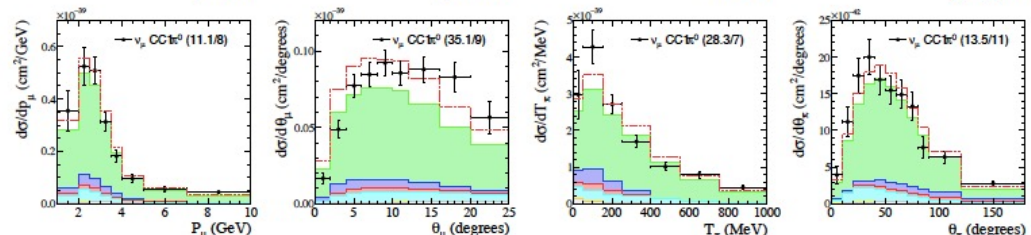
$\nu_{\mu} \text{CC1}\pi^+$



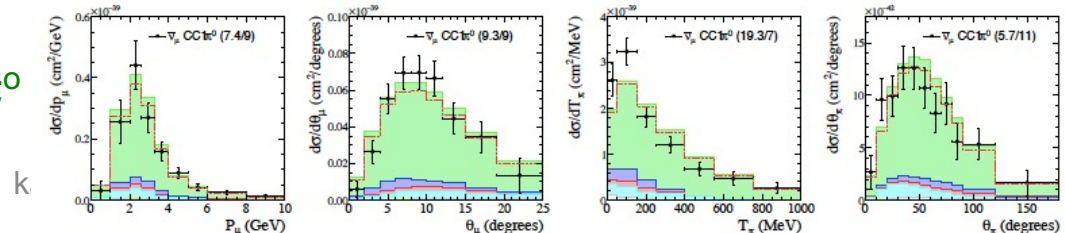
$\nu_{\mu} \text{CCN}\pi^+$



$\nu \text{CC1}\pi^0$

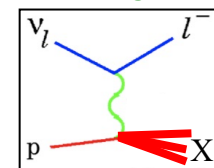


$\bar{\nu} \text{CC1}\pi^0$



You need to predict both
1. pion production model
2. final state interaction

1. Neutrino interaction physics - introduction
2. Charged-Current Quasi-Elastic (CCQE) interaction
3. Neutrino baryonic resonance interaction
- 4. Neutrino shallow- and deep-inelastic scatterings**
5. Conclusions



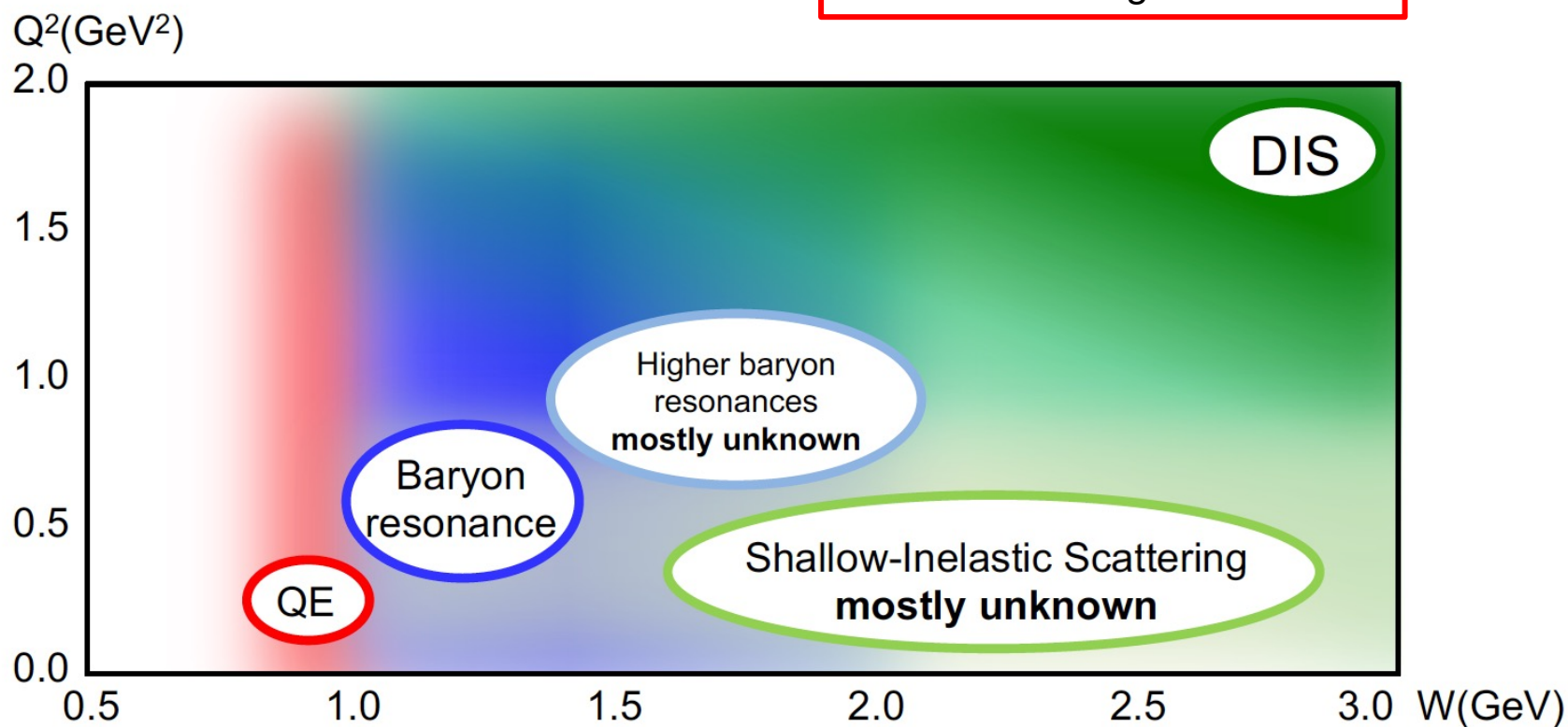
4. Shallow Inelastic Scattering (SIS)

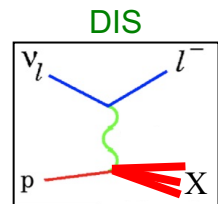
Cross section

- Higher resonances and hadron dynamics
- Quark-Hadron duality (low Q^2 , low W DIS)
- Nuclear dependent DIS

Neutrino experiments around 1-10 GeV are not quite DIS yet

- Shallow \rightarrow low Q^2
- Inelastic \rightarrow large W





4. Shallow Inelastic Scattering (SIS)

Cross section

- Higher resonances and hadron dynamics
- Quark-Hadron duality (low Q^2 , low W DIS)
- Nuclear dependent DIS

$Q^2(\text{GeV}^2)$

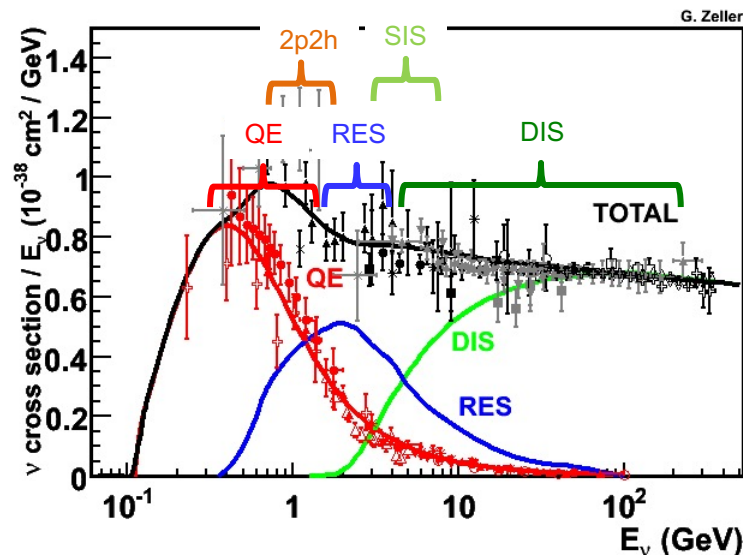
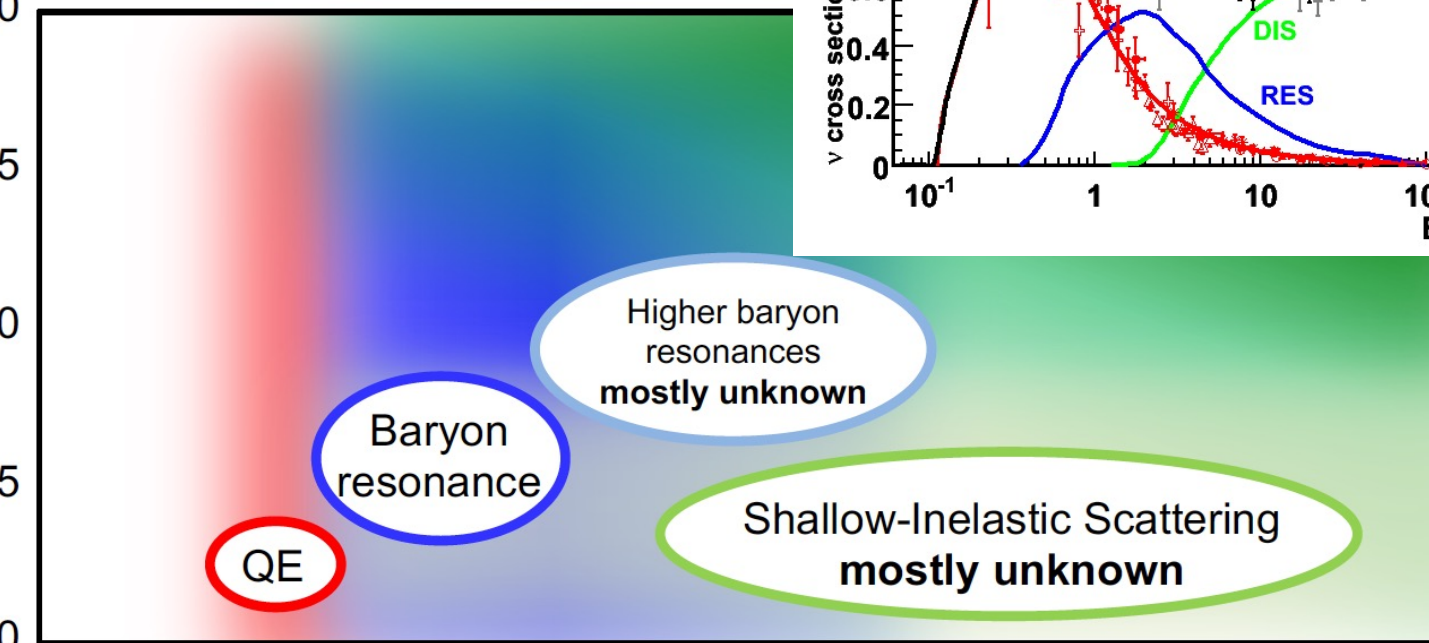
2.0

1.5

1.0

0.5

0.0



G. Zeller

4. Higher baryonic resonances

Cross section

- Higher resonances and hadron dynamics
- Quark-Hadron duality (low Q^2 , low W DIS)
- Nuclear dependent DIS

DCC model

- Channels are coupled (πN , $\pi\pi N$, etc), total amplitude is conserved
- Most of axial form factors are unknown

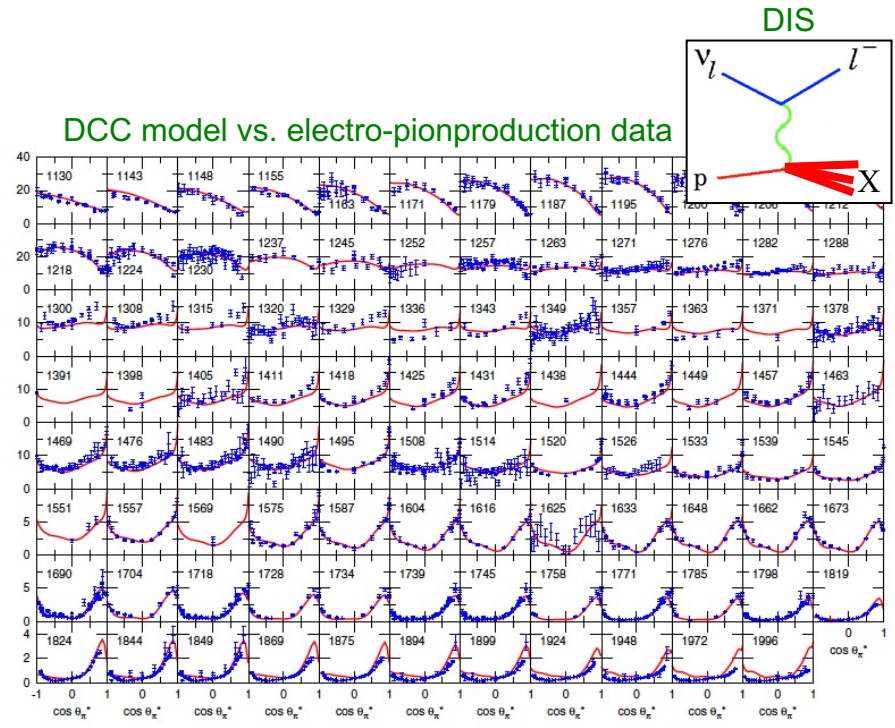
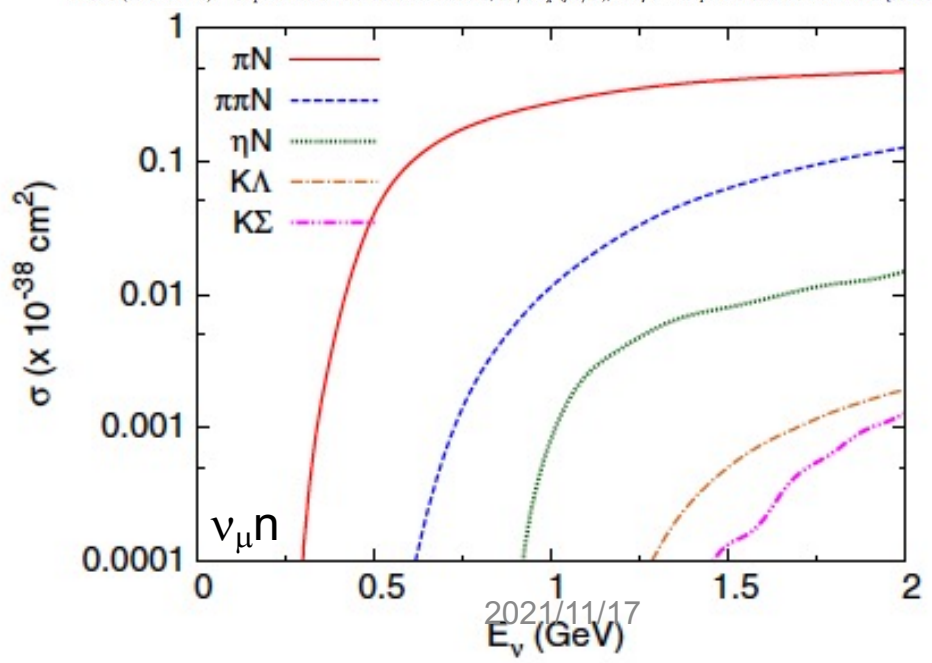
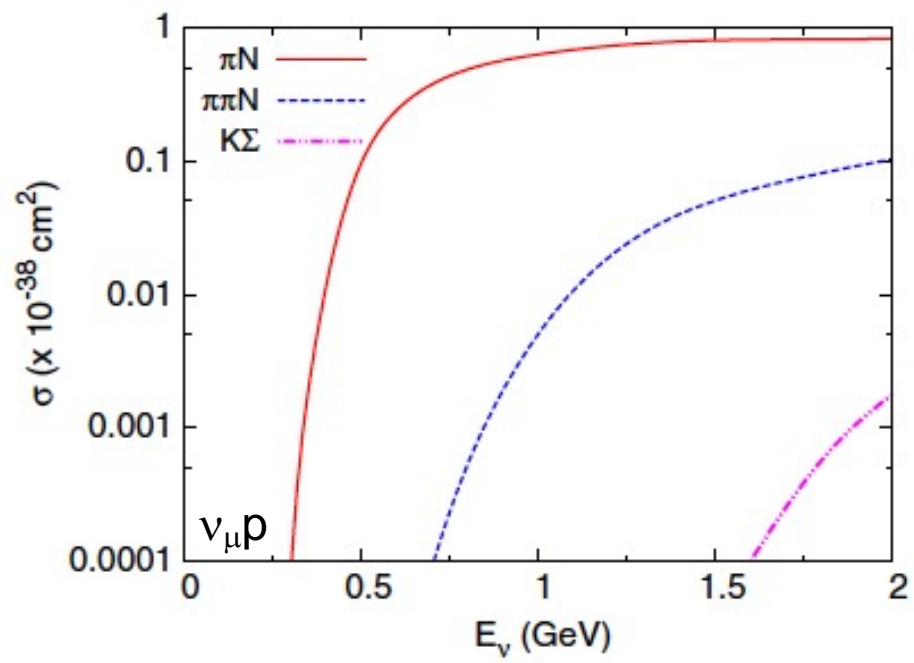
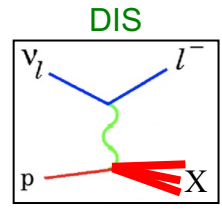


FIG. 8 (color online). Unpolarized differential cross sections, $d\sigma/d\Omega_*^2$ ($\mu\text{b/sr}$), for $\gamma n \rightarrow \pi^- p$. The data are from Refs. [55–78].

4. Quark-Hadron duality

Nachtmann variable $\xi = \frac{2x}{\left(1 + \sqrt{1 + \frac{4x^2 M^2}{Q^2}}\right)}$



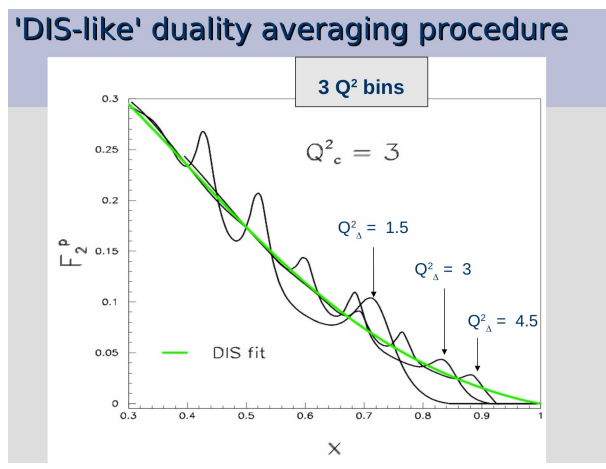
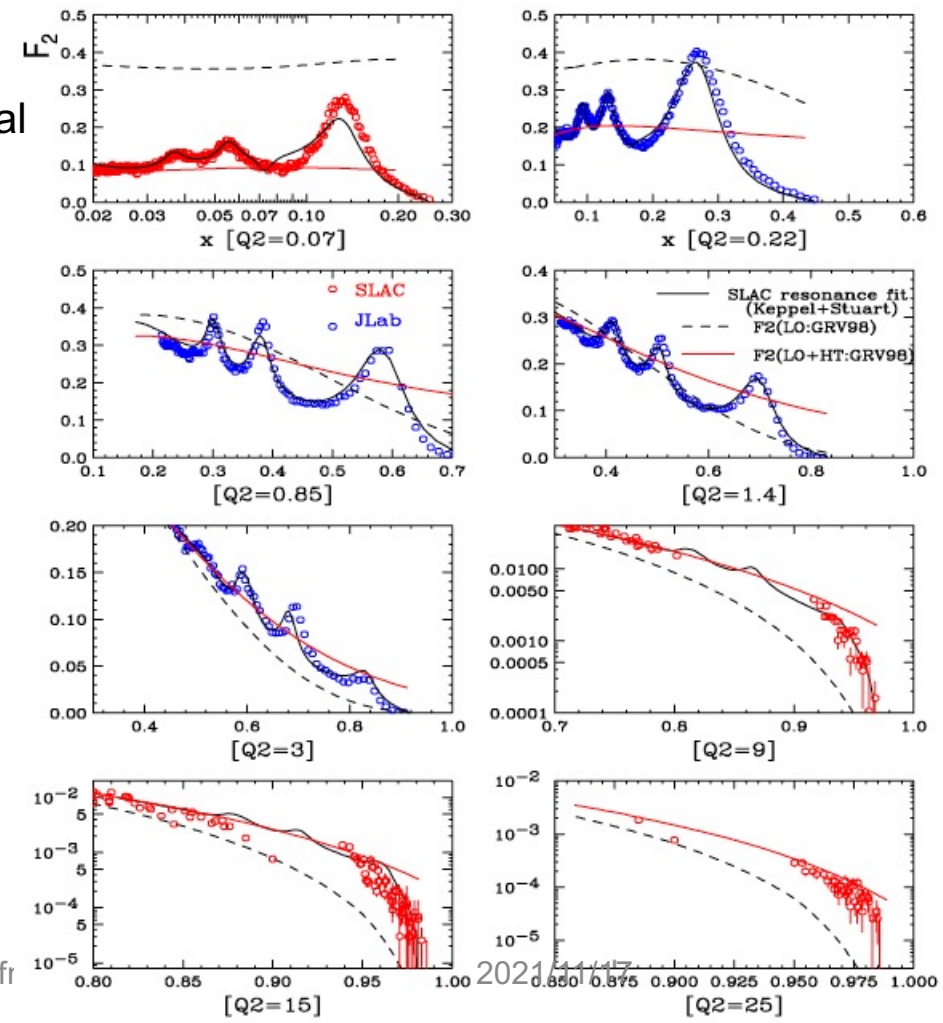
Cross section

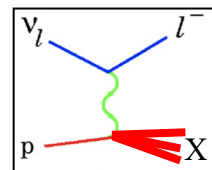
- Higher resonances and hadron dynamics
- Quark-Hadron duality (low Q^2 , low W DIS)
- Nuclear dependent DIS

Bodek-Yang correction is a phenomenological model to reproduce duality-like behavior, accepted by all neutrino simulation

DIS \neq Bjorken limit
 DIS = Q^2 average of all resonances

Proton F2 function GRV98-BY correction vs. data





4. Nuclear dependent DIS

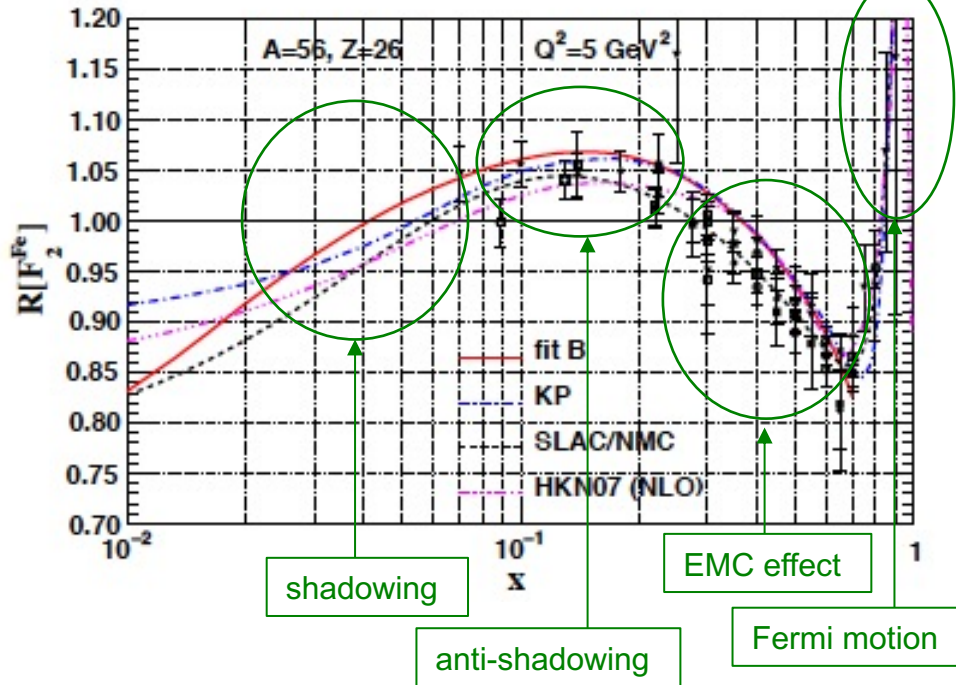
Cross section

- Higher resonances and hadron dynamics
- Quark-Hadron duality (low Q^2 , low W DIS)
- Nuclear dependent DIS

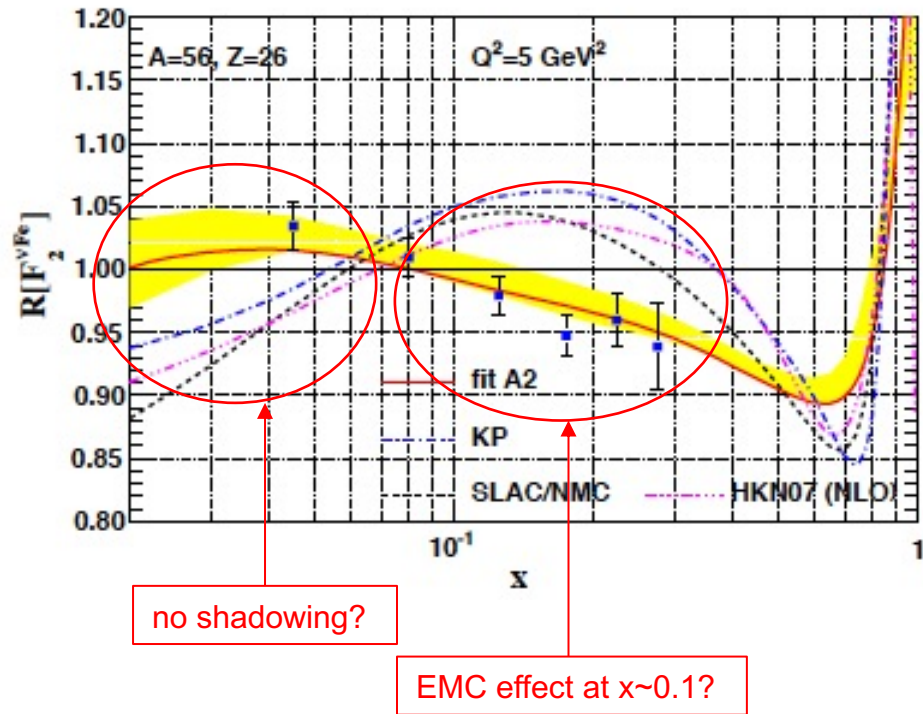
Nuclear PDF

- Shadowing, EMC effect, Fermi motion
- Likely due to nucleon dynamics in nucleus
- Various models describe charged lepton data
- Neutrino data look very different

e^+ -Fe nuclear correction factor



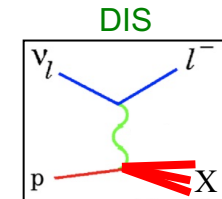
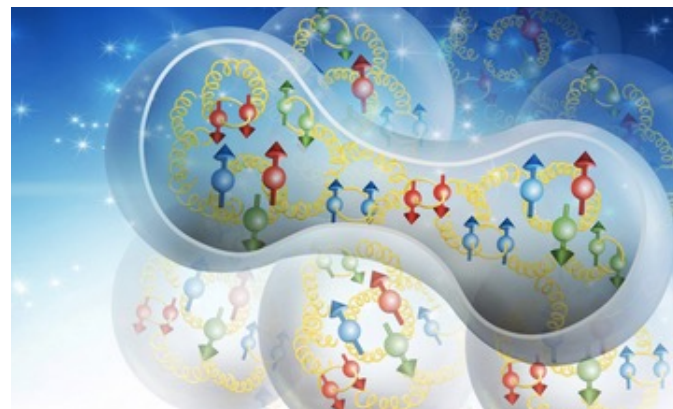
ν -Fe nuclear correction factor



4. EMC effect

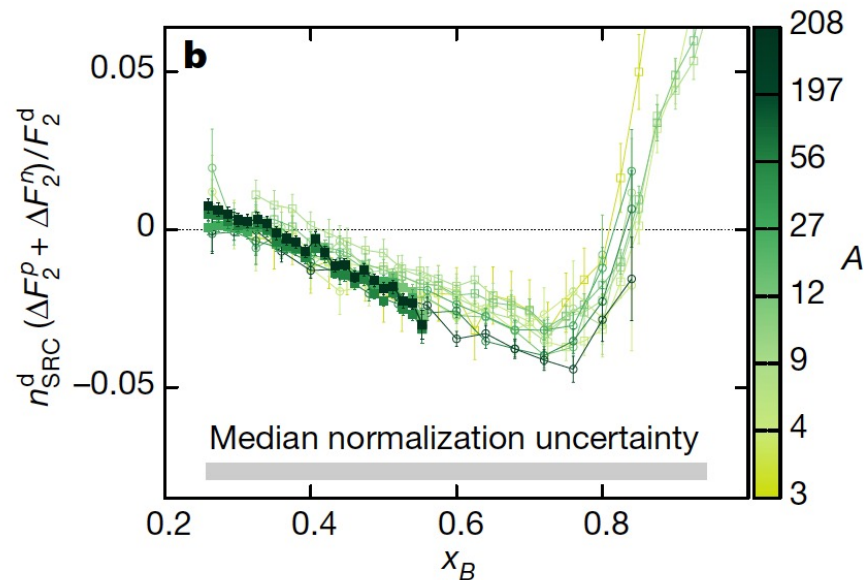
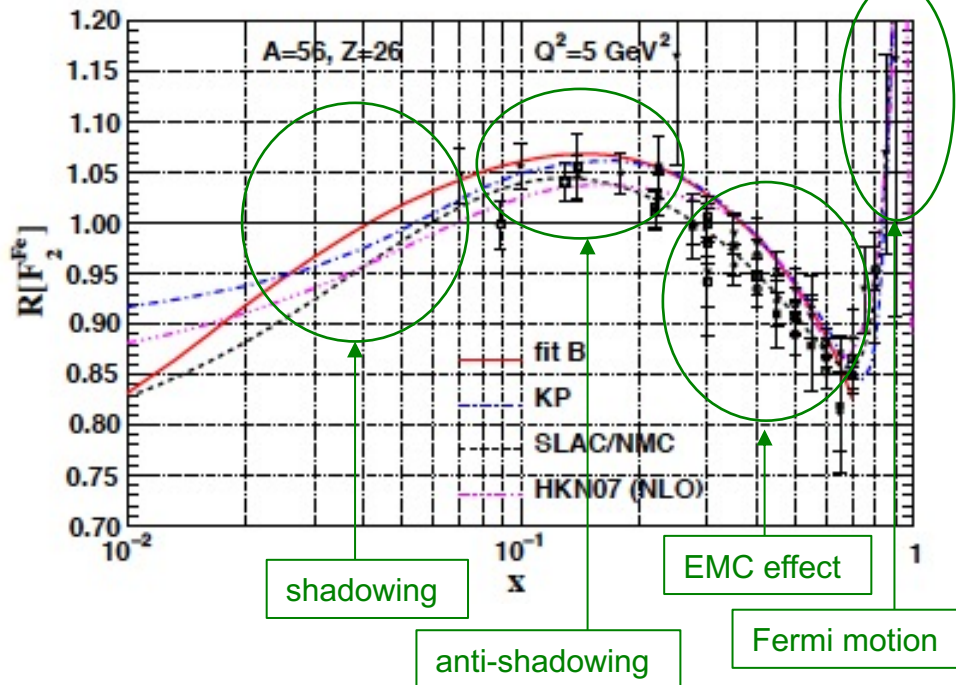
Nuclear dependent DIS

- First observed by the EMC experiment
- Structure function depends on nuclei
- Quarks feel presence of other quarks in other nucleons



EMC effect can be modeled from the amount of correlated pairs in nuclei (CLAS in JLab).

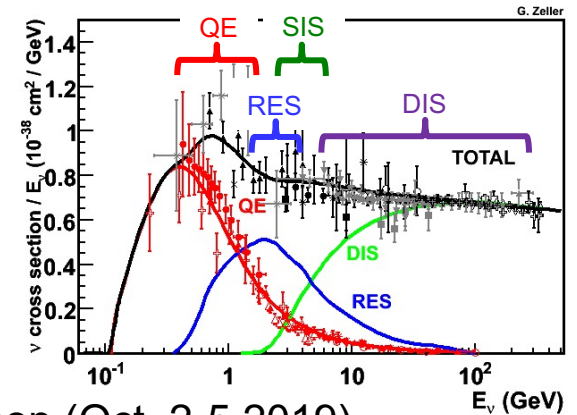
e^\pm -Fe nuclear correction factor



4. NuSTEC workshops

NuSTEC Workshop on
**Neutrino-Nucleus Pion Production
 in the Resonance Region**
 2019 October 2-5
 The University of Pittsburgh, USA
 nustec.fnal.gov/pion19

NuSTEC pion workshop (Oct. 2-5 2019)
 the University of Pittsburgh, USA
<https://nustec.fnal.gov/pion19/>



2018 October 11-13
 Gran Sasso Science Institute, Italy
 ν**S&DIS** workshop
 Neutrino Shallow- and Deep-
 inelastic Scattering workshop
 nustec.fnal.gov/nuSDIS18

NuSTEC SIS workshop (Oct. 11-13 2018)
 L'Aquila, Italy
<https://nustec.fnal.gov/nuSDIS18/>
 Summary paper (<https://arxiv.org/abs/1907.13252>)

Neutrino hadron production channels

- Background for oscillation measurement at HyperK
- Signal and background for oscillation measurement in DUNE
- Background for proton decay
- Signal and background for BSM physics search

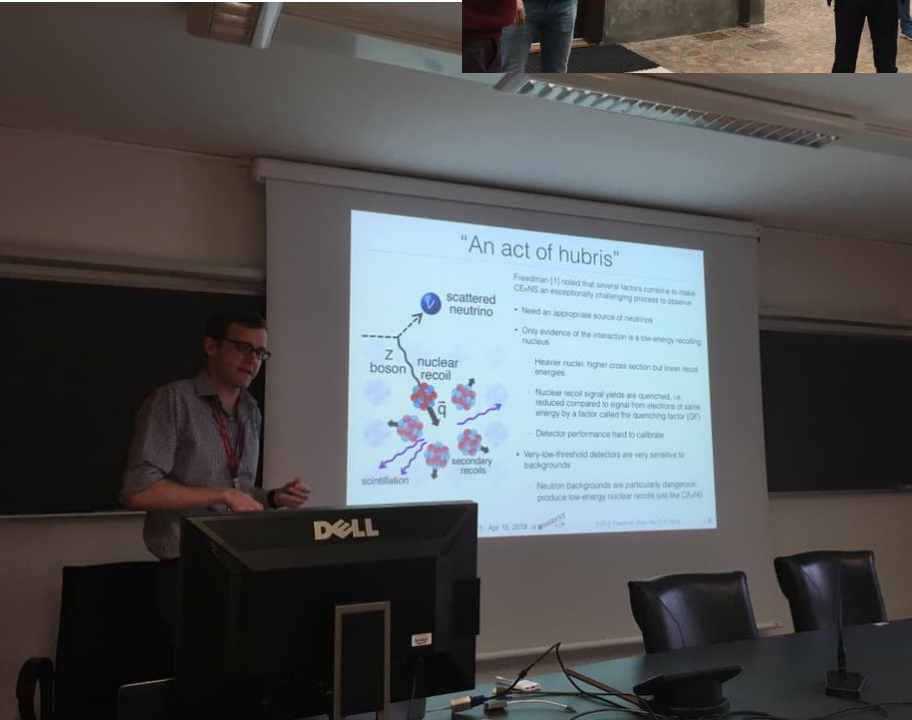
4. Atomic nuclei as laboratories for BSM physics

ECT* workshop, 15 Apr. 15-19 2019, Trento, Italy

<http://www.ectstar.eu/node/4436>

Topics include;

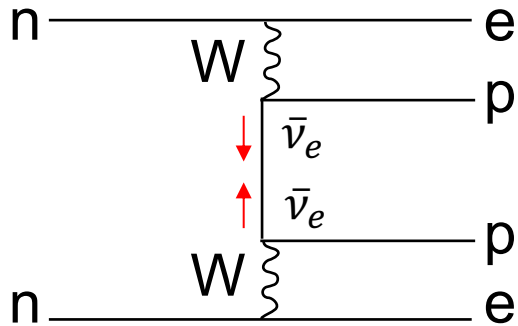
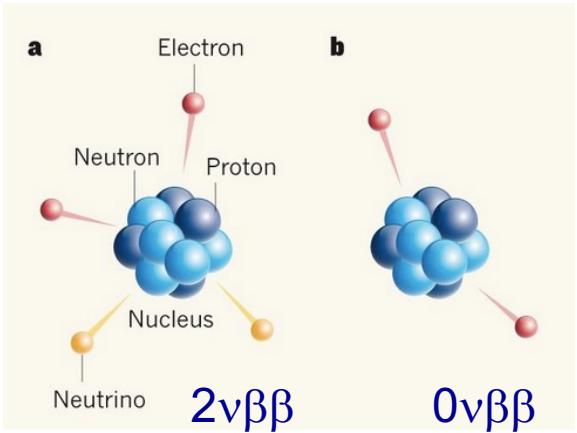
- 2p2h
- EMC effect
- $0\nu\beta\beta$
- dark matter
- etc



4. Neutrino-less double beta decay ($0\nu\beta\beta$)

Majorana particle

- double beta decay ($2\nu\beta\beta$) is the second order nuclear process, possible only for few elements (^{82}Se , ^{76}Ge , ^{100}Mo , ^{130}Te , ^{136}Xe , etc)
- $0\nu\beta\beta$ is the lepton number violation process (BSM process)
- Expected half-life, $\tau(0\nu\beta\beta) > 10^{27}$ yrs ($\gg 10^{10}$ yrs \sim life of universe)

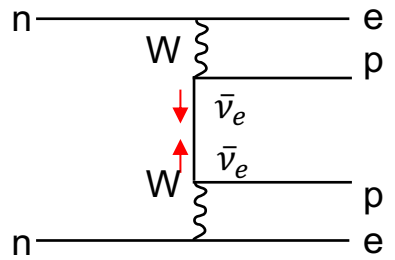


$$1/\tau = G(Q,Z) g_A^4 |M_{nucl}|^2 m_{\beta\beta}^2$$

Measured half-life of $0\nu\beta\beta$ process is related to effective Majorana mass ($m_{\beta\beta}^2$)

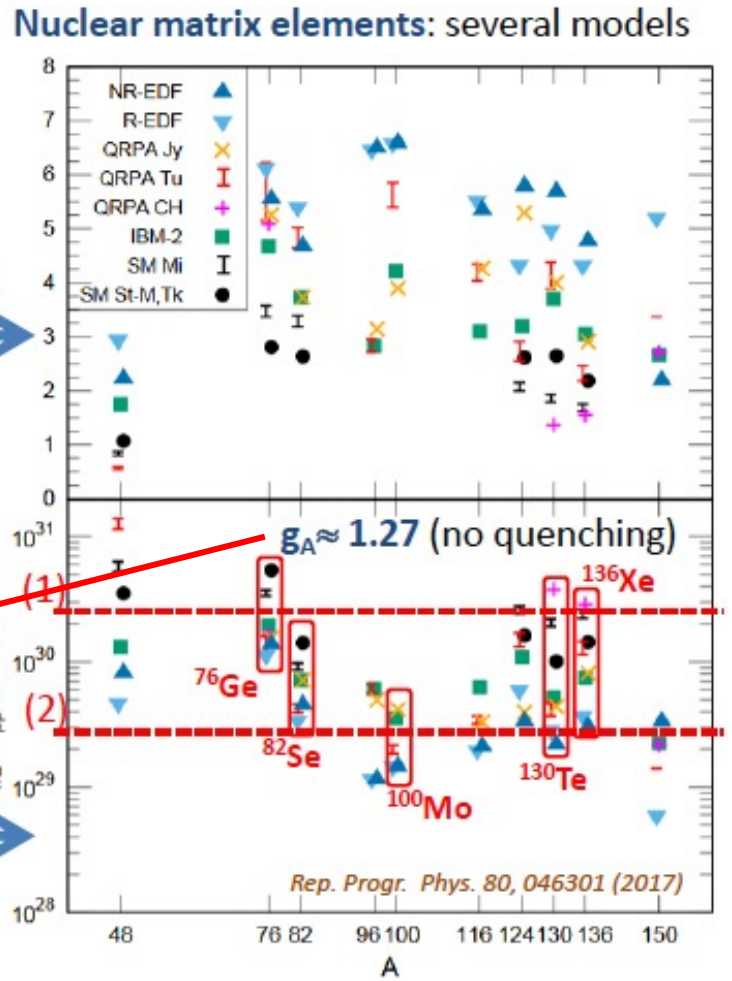
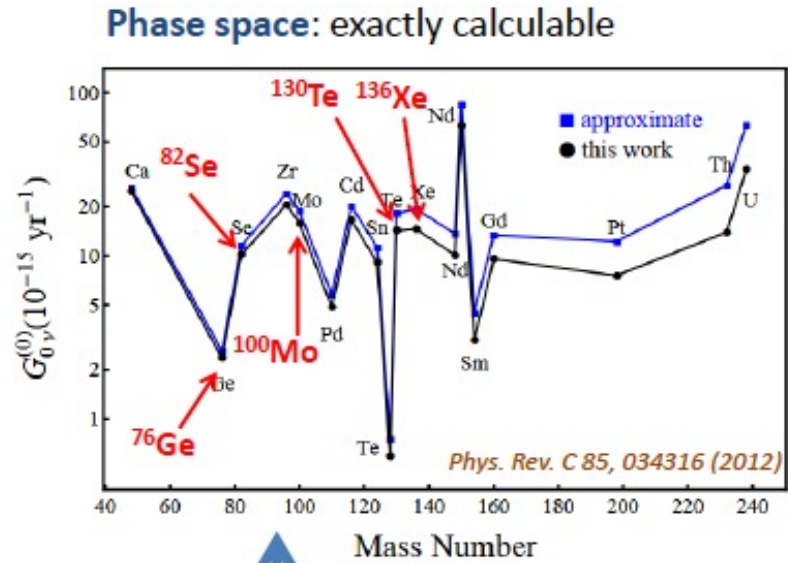
- Phase space
- Nuclear matrix element
- effective g_A

4. Neutrino-less double beta decay ($0\nu\beta\beta$)



Majorana particle

- Measured half-life of $0\nu\beta\beta$ process is related to effective Majorana mass ($m_{\beta\beta}^2$)



$$1/\tau = G(Q,Z) g_A^4 |M_{nucl}|^2 m_{\beta\beta}^2$$

Nuclear physics gives large systematics

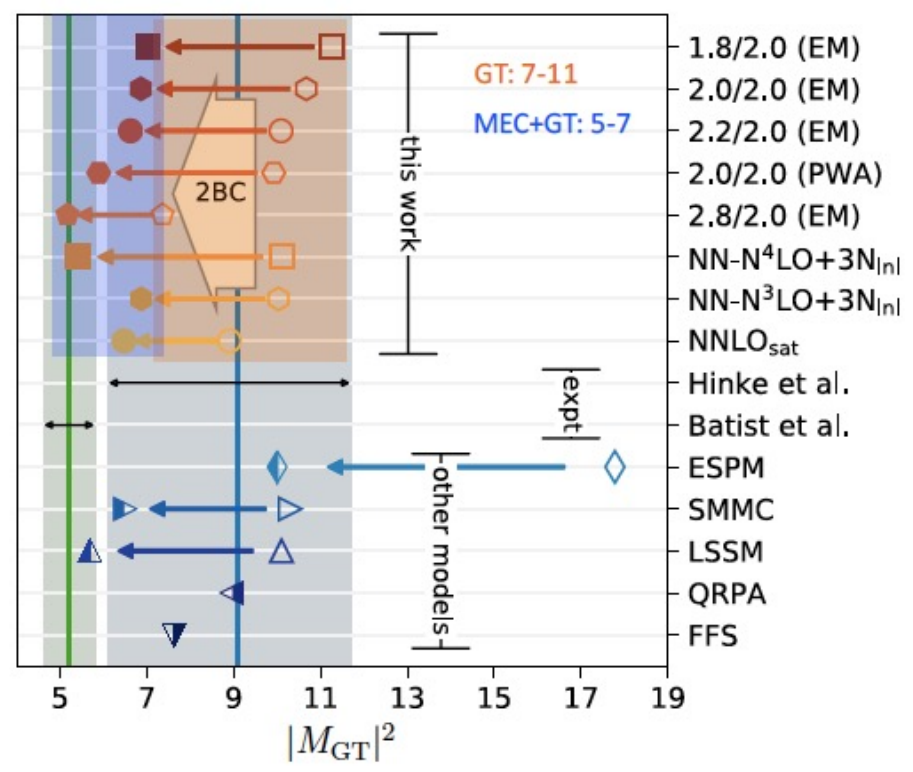
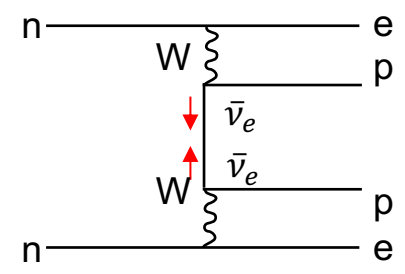
- Nuclear matrix element calculation
- Nuclear quenching of g_A



4. Neutrino-less double beta decay ($0\nu\beta\beta$)

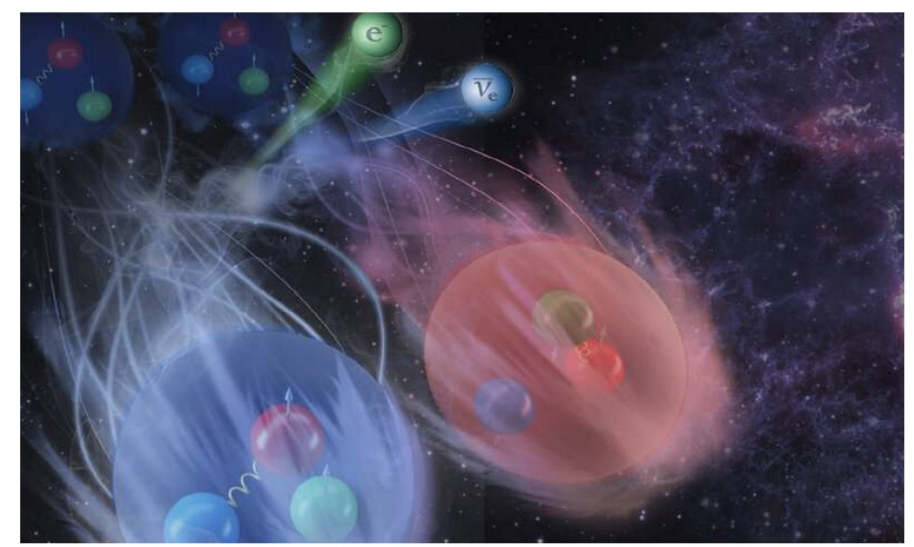
Beta decay quenching

- Axial coupling looks smaller in nuclei
- Ab initio calculation shows matrix element is suppressed due to nucleon 2-body current (2BC)
- Another uncertainty of $0\nu\beta\beta$



Physicists solve a beta-decay puzzle with advanced nuclear models

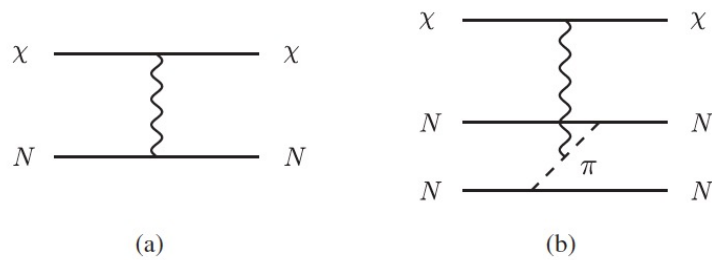
by Oak Ridge National Laboratory



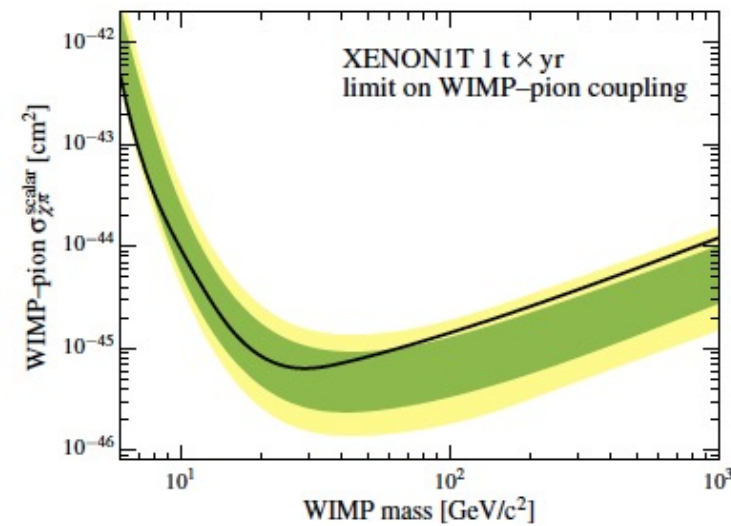
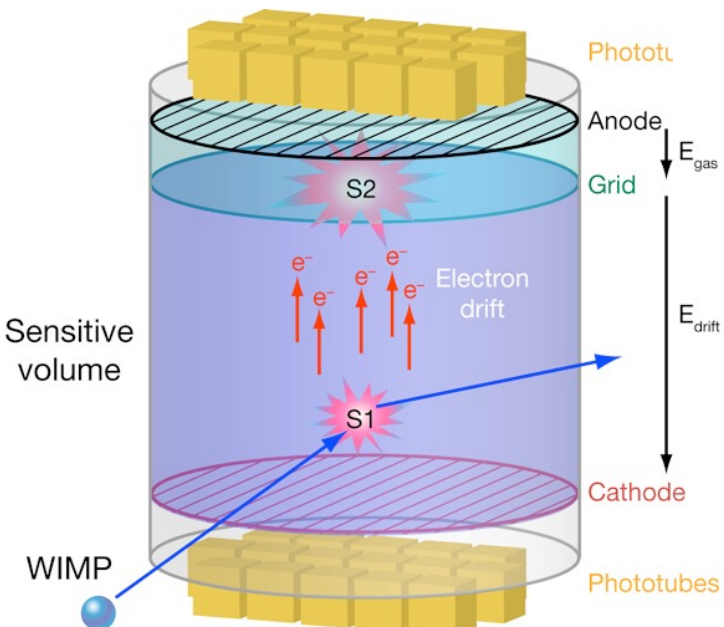
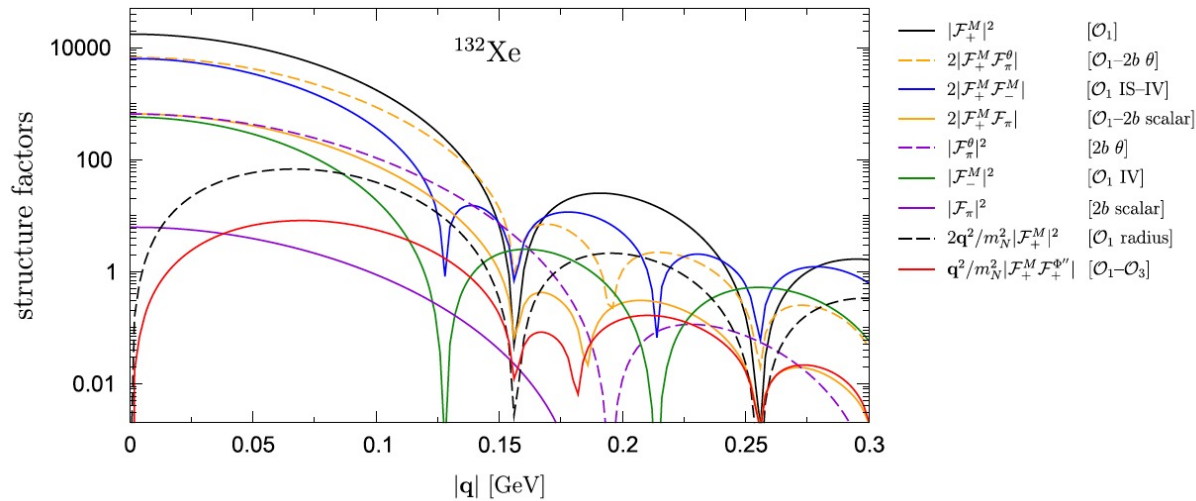
4. Direct WIMP-pion scattering

Nuclear structure function

- WIMP interaction depends on nuclear structure function
- Chiral effective field theory including 2-body current
- Assuming leading contributions are zero, WIMP-pion scattering can be studied.



$$\frac{d\sigma_{\chi N}^{SI}}{dq^2} = \frac{1}{4\pi v^2} \left| \left(c_+^M - \frac{\mathbf{q}^2}{m_N^2} \dot{c}_+^M \right) \mathcal{F}_+^M(\mathbf{q}^2) + c_\pi \mathcal{F}_\pi(\mathbf{q}^2) + c_\pi^\theta \mathcal{F}_\pi^\theta(\mathbf{q}^2) + \left(c_-^M - \frac{\mathbf{q}^2}{m_N^2} \dot{c}_-^M \right) \mathcal{F}_-^M(\mathbf{q}^2) + \frac{\mathbf{q}^2}{2m_N^2} [c_+^{\Phi''} \mathcal{F}_+^{\Phi''}(\mathbf{q}^2) + c_-^{\Phi''} \mathcal{F}_-^{\Phi''}(\mathbf{q}^2)] \right|^2,$$



Subscribe "NuSTEC News"

E-mail to listserv@fnal.gov, Leave the subject line blank, Type "subscribe nustec-news firstname lastname"

(or just send e-mail to me, katori@FNAL.GOV)

like "@nuxsec" on Facebook page, use hashtag #nuxsec

Conclusion

1 to 10 GeV neutrino interaction measurements are crucial to successful next-generation neutrino oscillation experiments (DUNE, Hyper-K)

Nucleon correlation physics drastically change neutrino cross sections, both size and shape.

Recent new models and theories show nucleon correlation physics is important in many sub-fields of particle physics.

Neutrino interaction physics beyond QE region is confusing.

Future neutrino interaction measurements should focus on high-statistics neutrino hadron production measurements. This is the key to understand neutrino interaction models and nuclear effect.

Thank you for your attention!

Backup

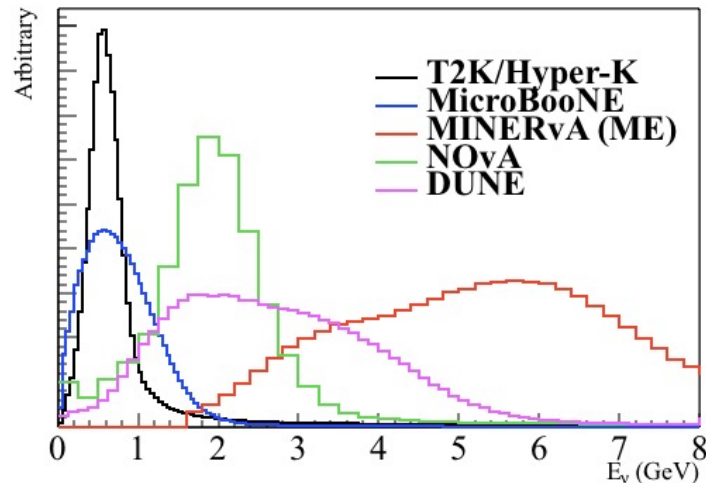
1. Next generation neutrino oscillation experiments

Neutrino oscillation experiments

- Past: K2K, MiniBooNE, MINOS, DeepCore
- Present to Future: T2K, NOvA, PINGU, ORCA, Hyper-Kamiokande, DUNE...

We don't know the energy of incoming neutrinos...

- We need to simulate all physics from $E_\nu=0$ to $E_\nu \sim \text{few GeV}$
- We need to simulate all physics from $\omega, |\vec{q}|=0$ to $\omega, |\vec{q}| \sim \text{few GeV}$



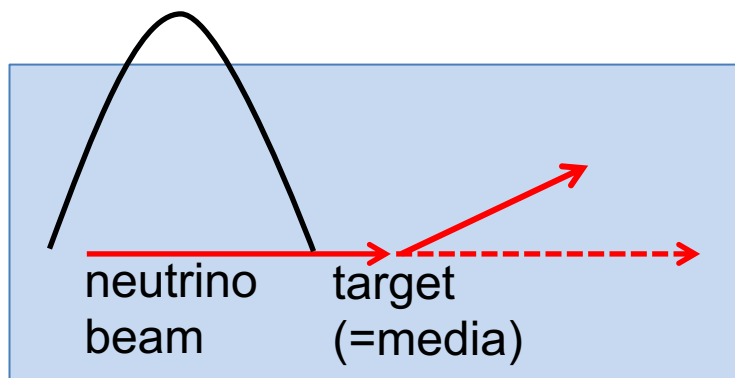
Two rules of neutrino interaction physics

1. Neutrinos cannot choose kinematic
2. Neutrino kinematics are not fully determined

1. Typical neutrino detectors

Neutrino scattering

- Wideband beam
- observables are inclusive



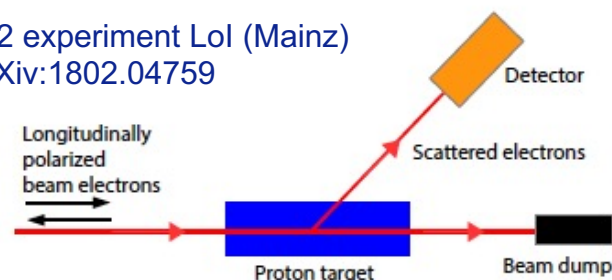
Incomplete kinematics

- Large mass, coarse instrumentation
- No one measures neutrino energy directly
- **Reconstructing kinematics (E_ν , Q^2 , W , x , y , ...)** in 1-10 GeV depends on interaction models

Electron scattering

- well defined energy, well known flux
- reconstruct energy-momentum transfer
- kinematics is completely fixed

P2 experiment Lol (Mainz)
arXiv:1802.04759



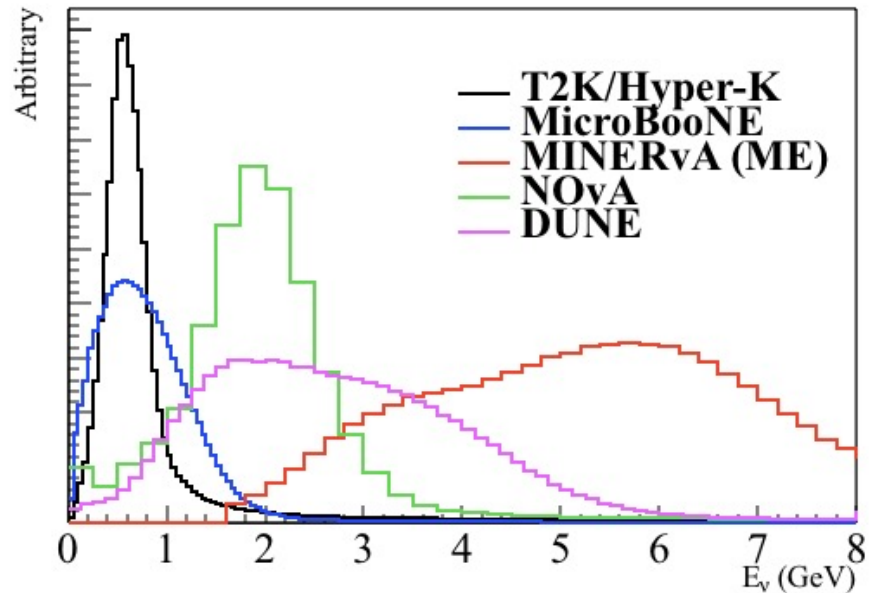
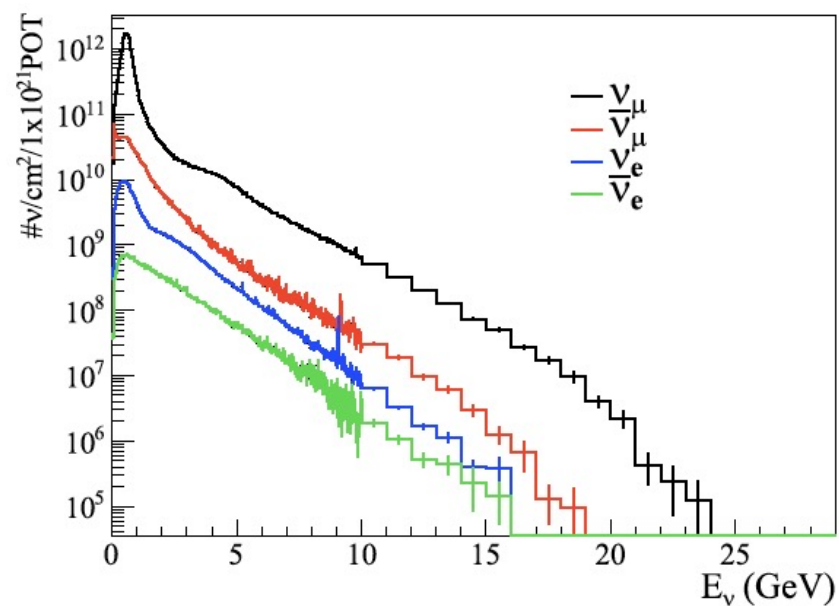
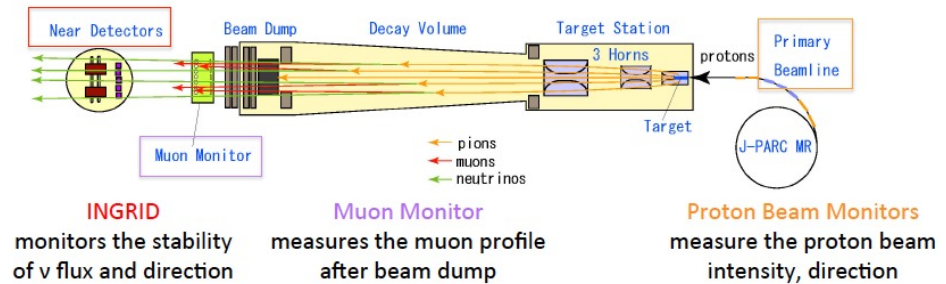
Two rules of neutrino interaction physics

1. Neutrinos cannot choose kinematic
2. **Neutrino kinematics are not fully determined**

1. Typical neutrino beams for oscillation experiments

e.g.) J-PARC neutrino beam (T2K)

- pion decay-in-flight (high flux)
- off-axis beam (narrow band)
- but has components up to ~ 10 GeV
- typical beam 1-10 GeV
- $\sim 4\%$ normalization error (best case)



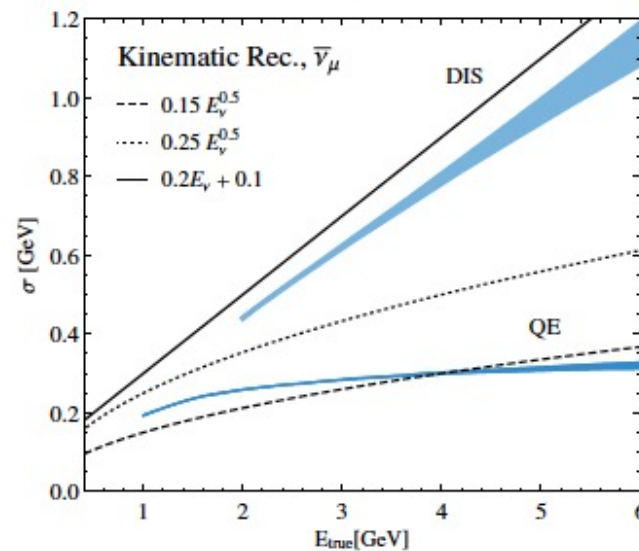
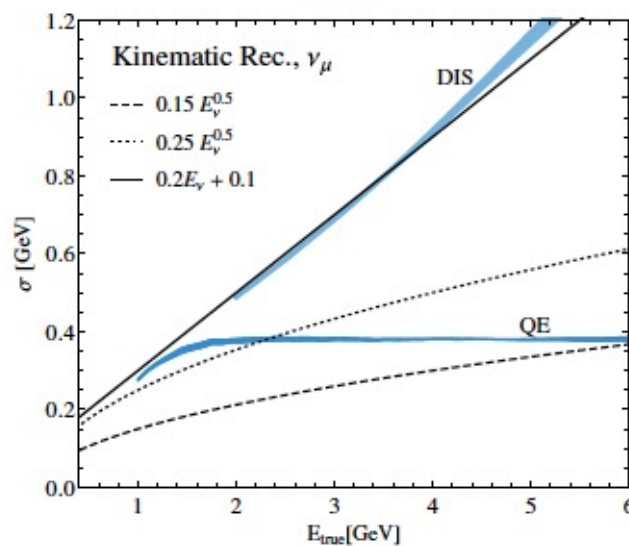
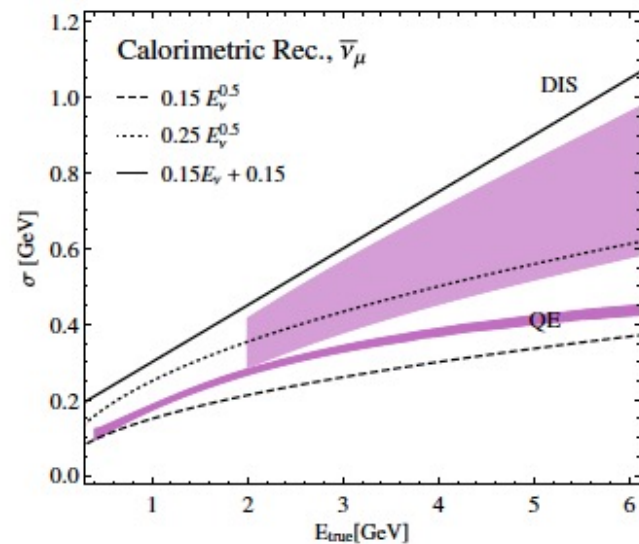
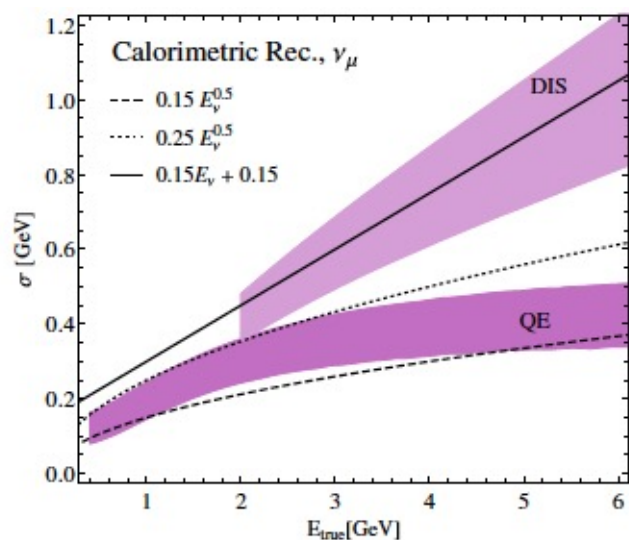
$$P_{\mu \rightarrow e}(L/E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (\text{eV}^2) \frac{L(\text{km})}{E(\text{GeV})} \right)$$

2. Kinematic E reconstruction vs calorimetric E reconstruction

Calorimetric energy reconstruction suffers invisible hadrons (=neutrons)

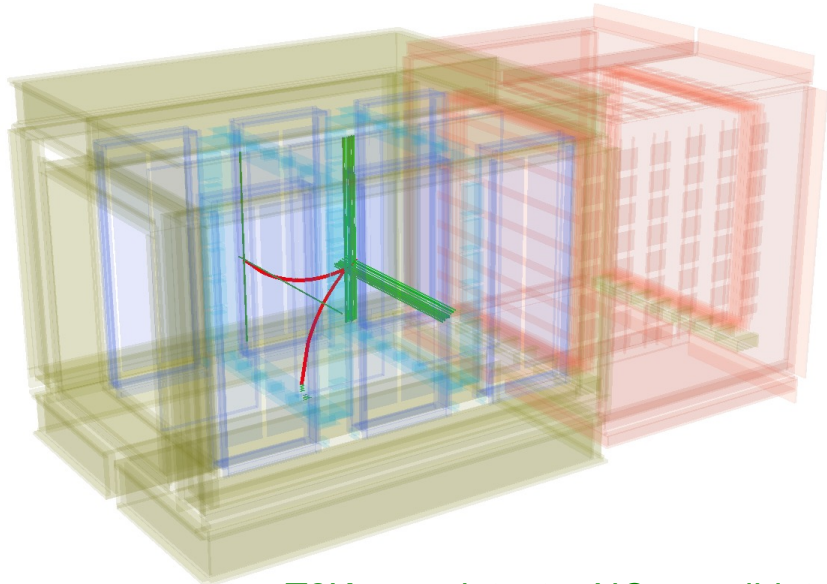
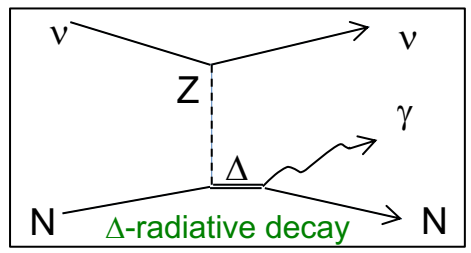
It largely depends on **neutrino interaction and hadron simulation**

- multiplicity
 - kinematics
 - nuclear effect
 - re-scattering
 - charge exchange
 - baryonic resonance
 - nucleon correlation
- etc

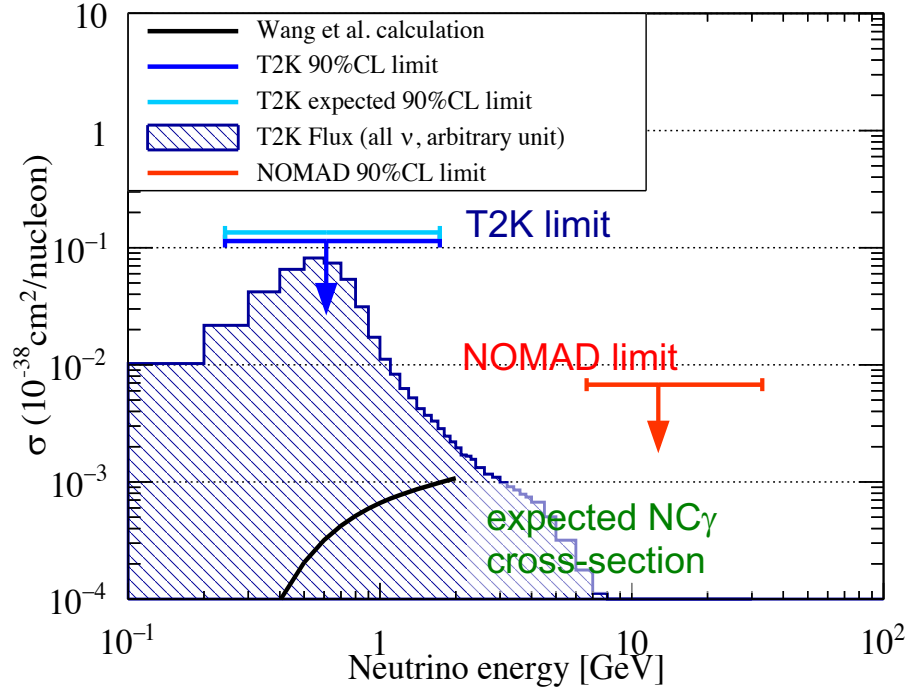


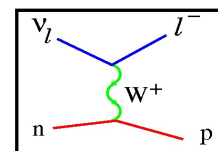
4. T2K Neutrino NC single photon production (NC_γ)

Neutrino induced NC single photon production (NC_γ) process is not experimentally identified. NC_γ is misID background for every electron-neutrino appearance oscillation experiment. T2K and NOMAD set limits on this process, but $\sim x3$ higher cross-section can explain all MiniBooNE excess.



T2K near detector NC_γ candidate

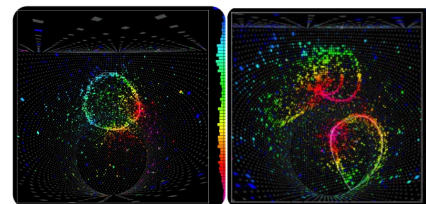
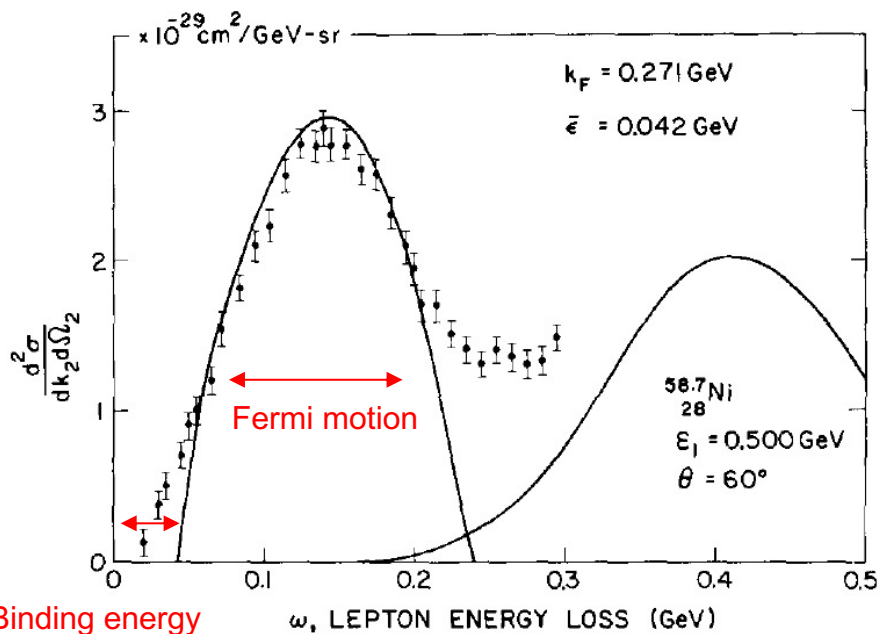




2. Fermi motion

Fermi motion

- Measured energy is smeared from the true energy if you assume nucleon at rest
- High resolution detector can measure all outgoing hadrons
 - initial nucleon momentum can be reconstructed (no Fermi motion smearing)

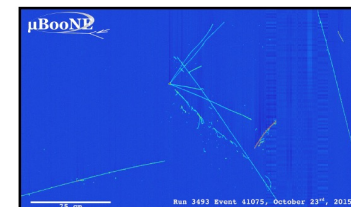


Cherenkov detectors:

Assuming QE interaction

Using lepton only

$$E_{QE} = \frac{2M\epsilon + 2ME_l - m_l^2}{2(M - E_l + |k_l| \cos \theta_l)}$$



Tracking detectors:

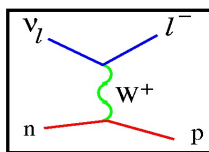
Calorimetric sum

Using All detected particles

$$E_{cal} = E_l + E_p^{kin} + \epsilon$$

[1p0π]

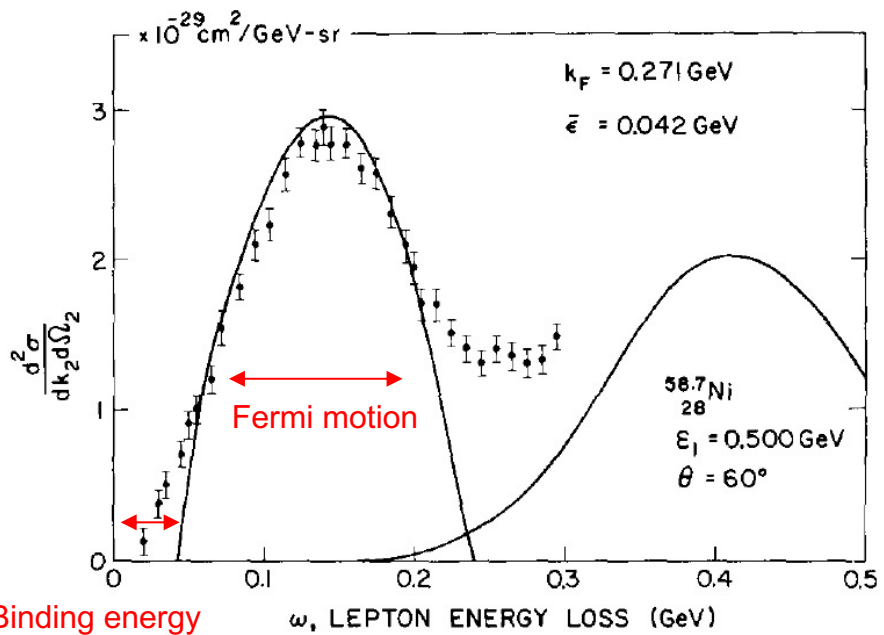
Binding energy



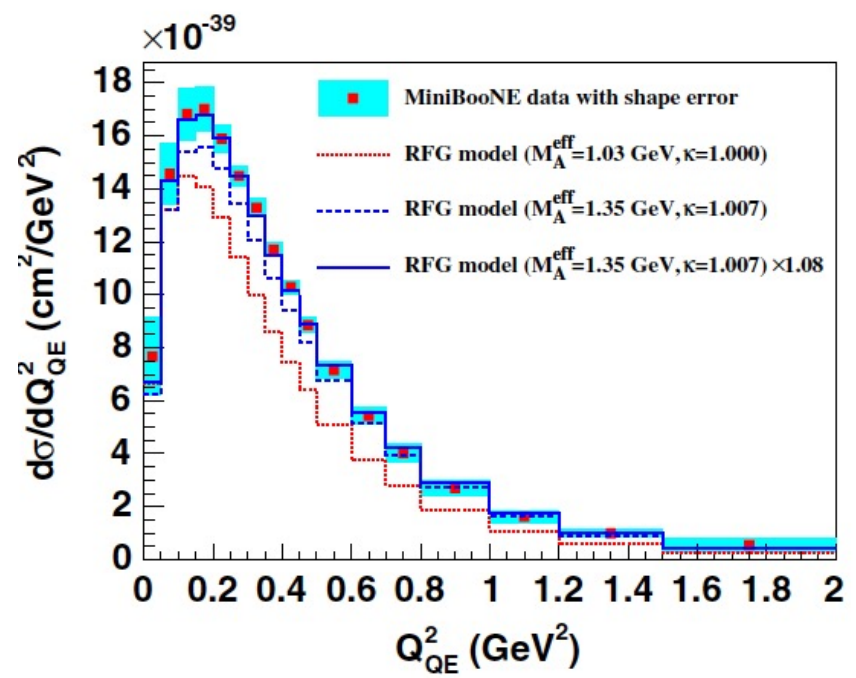
2. Pauli blocking

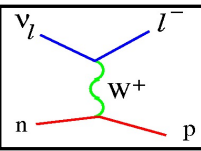
Pauli blocking

- Low momentum transfer reaction is forbidden.
- data show more suppression than what Pauli blocking can → RPA(?)
- In the global Fermi model, Pauli blocking looks unphysical



Binding energy

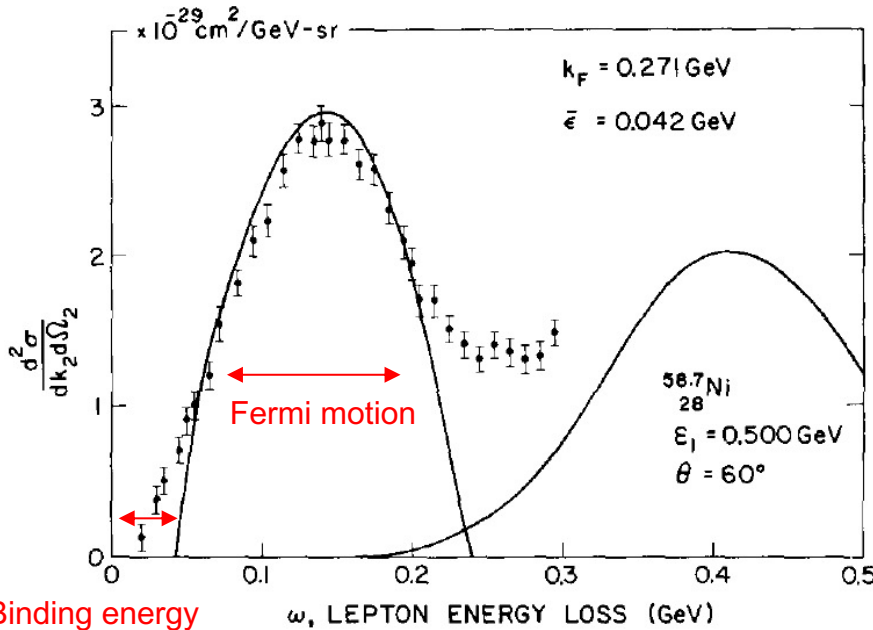




2. Nuclear Shell structure and binding energy

Binding energy ~ unobserved energy

- Energy to cost to release 1 nucleon, not constant
- Separation energy + excitation energy + recoil energy
 - Separation energy: energy to release 1 nucleon from the shell (~15 MeV, depends)
 - Excitation energy: energy used to excite leftover target nucleus (~1 MeV)
 - Recoil energy: kinetic energy of recoil target nucleus (~2-3 MeV)



Binding energy

Electron scattering on proton

