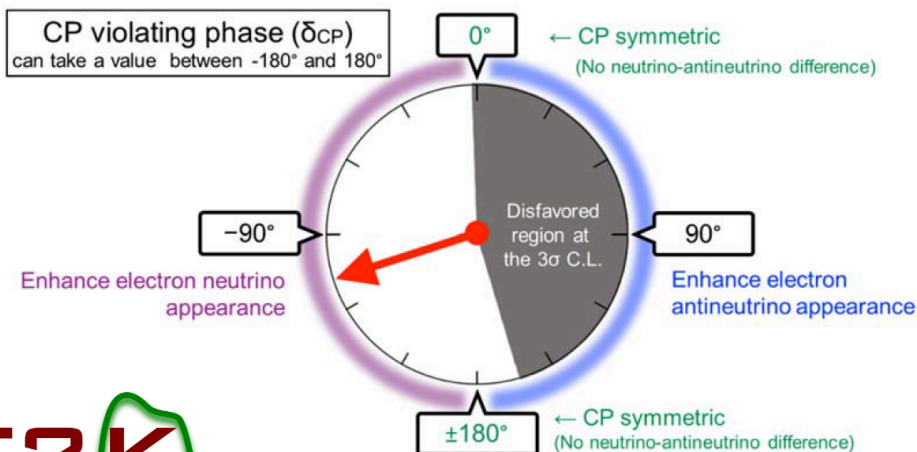


Constraint on the matter–antimatter symmetry-violating phase in neutrino oscillations

outline

1. Neutrino oscillation experiments
2. Neutrino beam
3. Neutrino detector
4. Neutrino interaction physics
5. Oscillation result
6. SK-Gd, T2K-Upgrade, Hyper-Kamiokande
7. Conclusion



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2020/04/22

1



1. Neutrino oscillation experiments

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7. Conclusion

Neutrino Physics is center stage!

BREAKTHROUGH PRIZE



2016 Fundamental Physics Breakthrough Prize

- Koichiro Nishikawa (K2K and T2K)
- Atsuto Suzuki (KamLAND)
- Kam-Biu Luk (Daya Bay)
- Yifang Wang (Daya Bay)
- Art McDonald (SNO)
- Yoichiro Suzuki (Super-Kamiokande)
- Takaaki Kajita (Super-Kamiokande)



The Nobel Prize in Physics 2015

Takaaki Kajita, Arthur B. McDonald

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The Nobel Prize in Physics 2015



Photo © Takaaki Kajita

Takaaki Kajita

Prize share: 1/2



Photo: K. McFarlane,
Queen's University
/SNOLAB

Arthur B. McDonald

Prize share: 1/2

The Nobel Prize in Physics 2015 was awarded jointly to Takaaki Kajita and Arthur B. McDonald "*for the discovery of neutrino oscillations, which shows that neutrinos have mass*"

1. Next goal of high energy physics

Establish Neutrino Standard Model (ν SM)

- SM + 3 active massive neutrinos

Unknown parameters of ν SM

1. Dirac CP phase
 2. θ_{23} ($\theta_{23}=40^\circ$ and 50° are same for $\sin 2\theta_{23}$, but not for $\sin \theta_{23}$)
 3. normal mass ordering $m_1 < m_2 < m_3$ or inverted mass ordering $m_3 < m_1 < m_2$
 4. Dirac or Majorana
 5. Majorana phases
 6. Absolute neutrino mass
- } not relevant to neutrino oscillation experiments

We need higher precision neutrino oscillation experiments around 1-10 GeV.

1. Neutrino oscillations

Flavour eigenstate = superposition of Hamiltonian eigenstate (mass eigenstate in vacuum)

$$|\nu_\alpha\rangle = \sum_i U_{\alpha i} |\nu_i\rangle, \quad U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix}$$

U is PMNS (Pontecovo-Maki-Nakagawa-Sakata) Matrix

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\frac{\alpha_{21}}{2}} & 0 \\ 0 & 0 & e^{i\frac{\alpha_{31}}{2}} \end{pmatrix}$$

Here $c_{23} = \cos\theta_{23}$ and $s_{23} = \sin\theta_{23}$, δ_{CP} = Dirac phase, α_i = Majorana phase.

Oscillation probability of neutrino flavor ν_α to ν_β with energy E , baseline L is

$$P_{\alpha \rightarrow \beta}(L, E) = 1 - 4 \sum_{i>j} \text{Re}(U_{\alpha i}^* U_{\beta i}^* U_{\alpha j} U_{\beta j}) \sin^2 \left(\frac{\Delta m_{ij}^2 L}{4E} \right) + 2 \sum_{i>j} \text{Im}(U_{\alpha i}^* U_{\beta i}^* U_{\alpha j} U_{\beta j}) \sin \left(\frac{\Delta m_{ij}^2 L}{2E} \right)$$

After adjusting units, the 2-flavor oscillation can be written in a simple formula.

$$P_{\mu \rightarrow e}(L / E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L(km)}{E(GeV)} \right)$$

1. Neutrino oscillations

Keep the first order of CP violation for muon neutrino to electron neutrino oscillation

Jarlskog invariant

$$J_{CP,I} = \frac{1}{8} \cos\theta_{13} \sin(2\theta_{12}) \sin(2\theta_{23}) \sin(2\theta_{13}) \sin\delta_{CP}$$

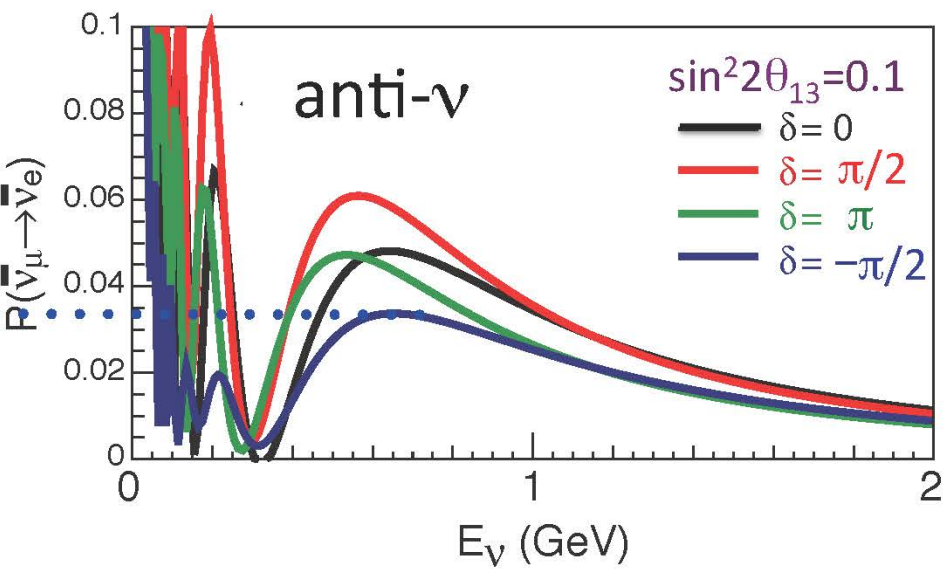
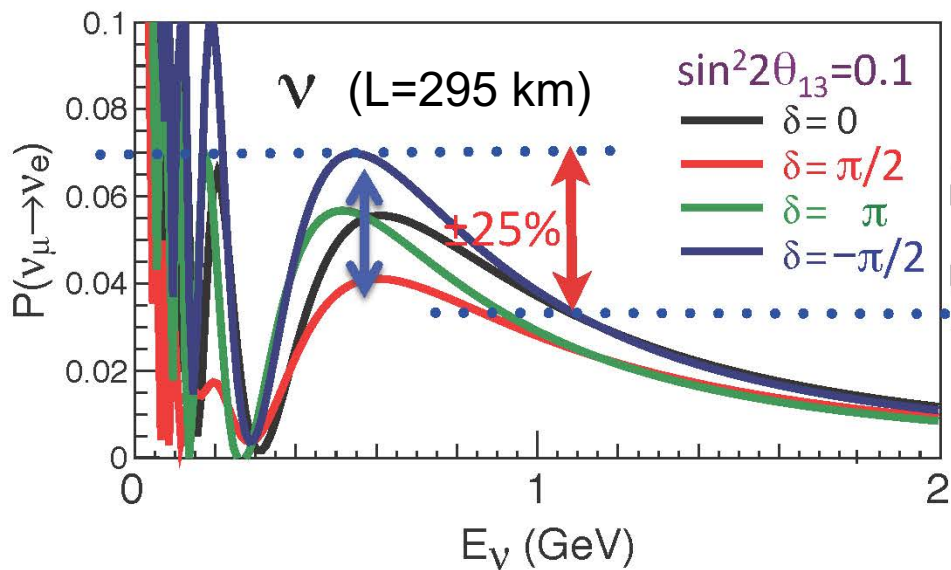
$$P(\nu_{\mu} \rightarrow \nu_e) \approx \sin^2(2\theta_{13}) \sin^2\theta_{23} \sin^2\left(\frac{1.27\Delta m_{32}^2 L}{E}\right)$$

$$\mp \frac{1.27\Delta m_{21}^2 L}{E} 8J_{CP} \sin^2\left(\frac{1.27\Delta m_{32}^2 L}{E}\right)$$

- Neutrino
+ Antineutrino

(1) If there is no CP violation,
 $P(\nu_{\mu} \rightarrow \nu_e)$ and $P(\bar{\nu}_{\mu} \rightarrow \bar{\nu}_e)$ are
the same

(2) Expected oscillation probability
to measure δ_{CP} is small



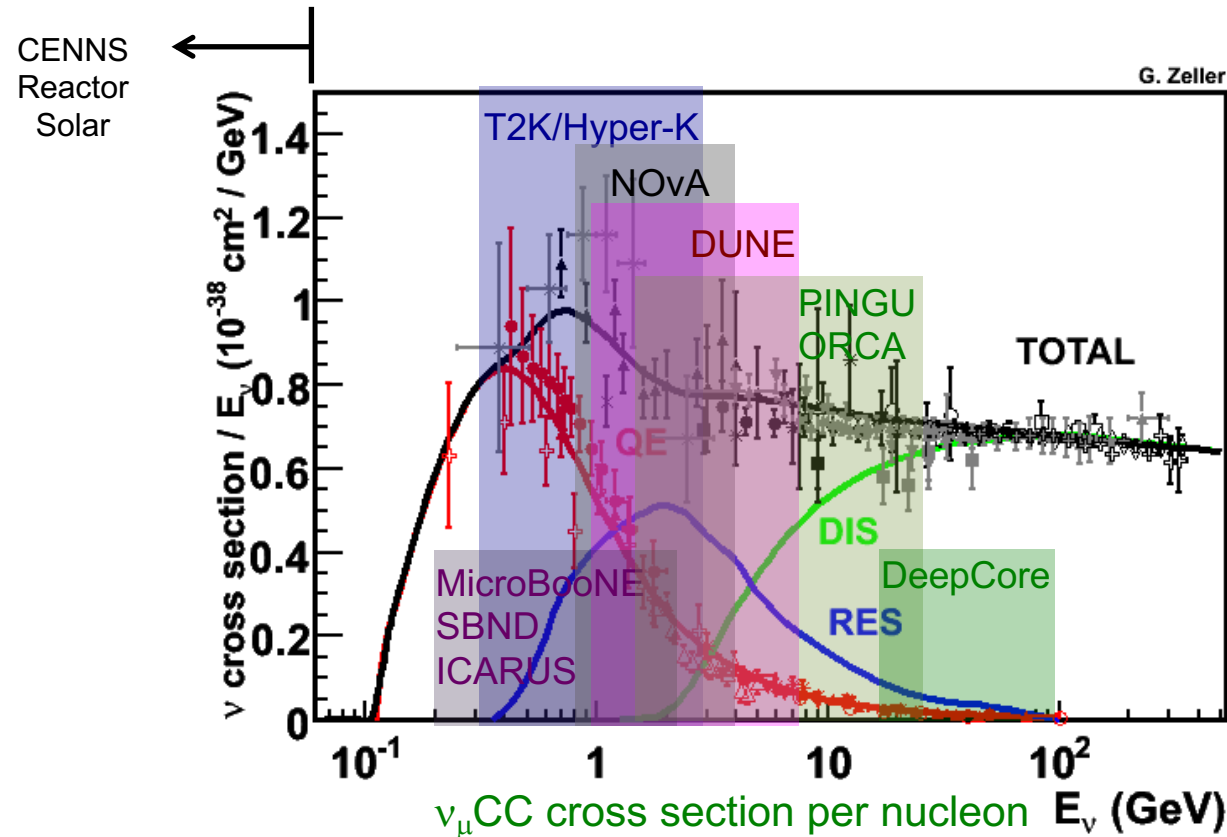
$$P_{\mu \rightarrow e}(L/E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L(km)}{E(GeV)} \right)$$

1. Next generation neutrino oscillation experiments

Neutrino oscillation experiments

- Present: T2K, NOvA, MicroBooNE, DeepCore...
- Future: Hyper-Kamiokande, DUNE, PINGU, ORCA...

- Most of experiments are 1-10 GeV
 - Cross-section data are poor



For $\Delta m^2_{32} \sim 2.5 \times 10^{-3} \text{ eV}^2$,
the first oscillation max is
 $L/E \sim 500 \text{ km/GeV}$

ex) T2K
 $L \sim 300 \text{ km}$
 $E \sim 0.6 \text{ GeV}$

ex) NOvA
 $L \sim 800 \text{ km}$
 $E \sim 2 \text{ GeV}$

$$P_{\mu \rightarrow e}(L / E) = \sin^2 2 \theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L (km)}{E (GeV)} \right)$$

1. Neutrino oscillation analysis

Data

Create neutrino beam

Measure muon neutrinos at the near detector

Measure electron neutrinos (and muon neutrinos) at the far detector

Simulation

Simulate neutrino flux at the near detector and the far detector

Simulate neutrino interaction at the near detector and the far detector

Simulate detector response at the near detector and far detector

Apply oscillation formula to the far event distribution

Oscillation analysis

Compare data and simulation to find oscillation parameters
(use near detector information to constrain systematic errors)

$$P_{\mu \rightarrow e}(L / E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L(km)}{E(GeV)} \right)$$

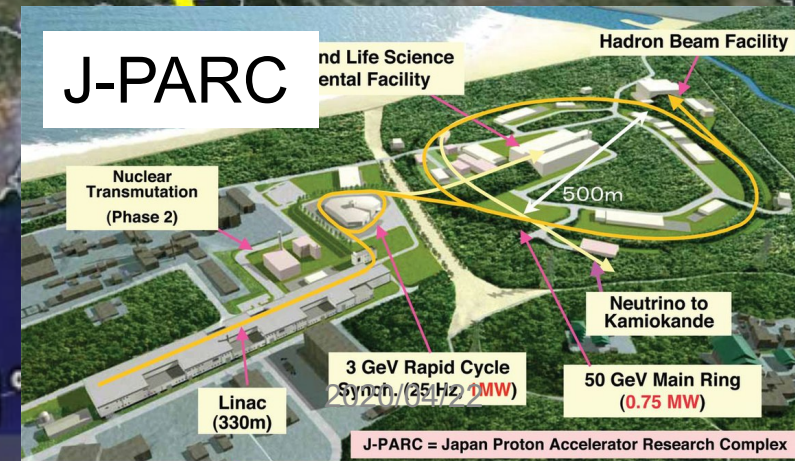
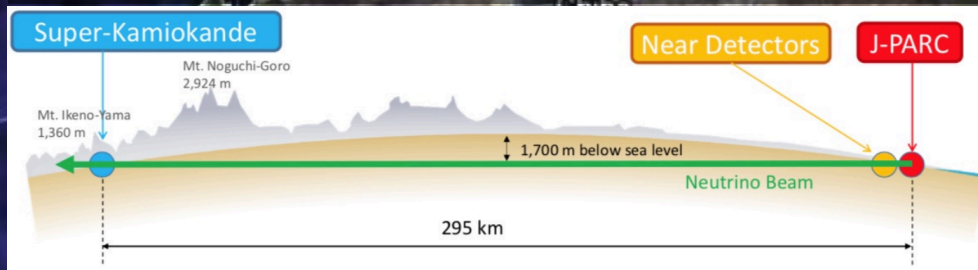
Super-Kamiokande detector
(far detector)

T2K

295km

Neutrino beam

T2K (Tokai to Kamioka)
experiment



1. Neutrino oscillation experiments

2. Neutrino beam

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2. J-PARC

LINAC
- 400 MeV



RCS (Rapid Cycling Synchrotron)
- 3 GeV



To
Kamioka

Neutrino

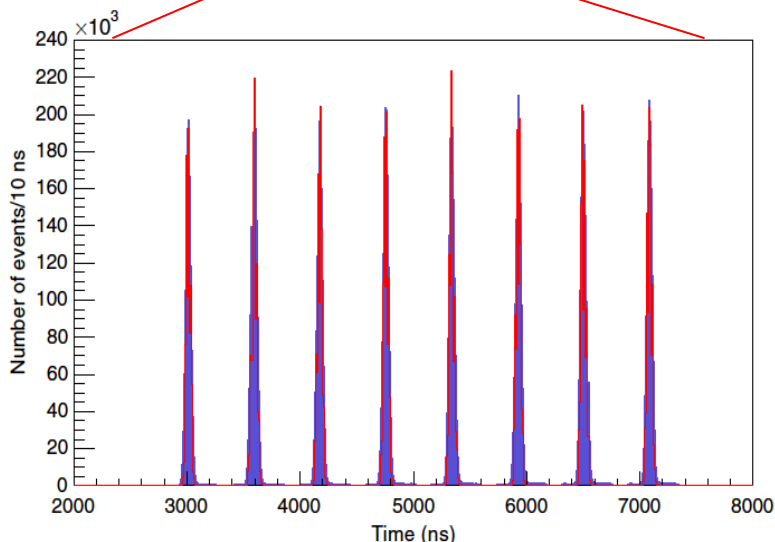
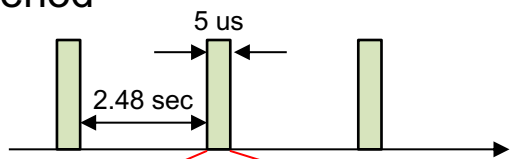
Main Ring
- 30 GeV



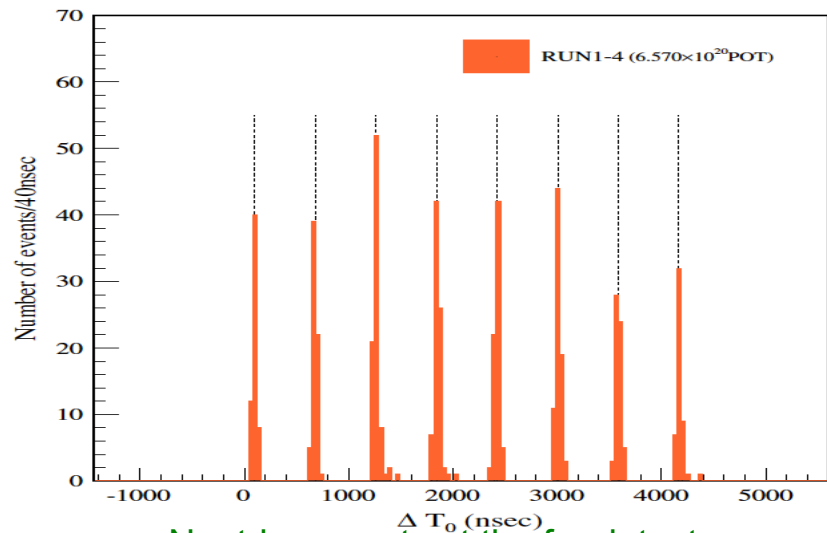
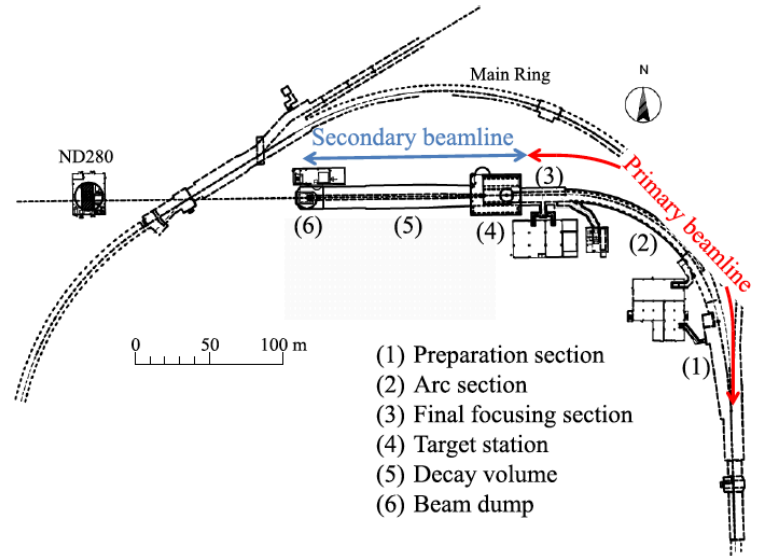
2. Neutrino beamline

Primary beamline

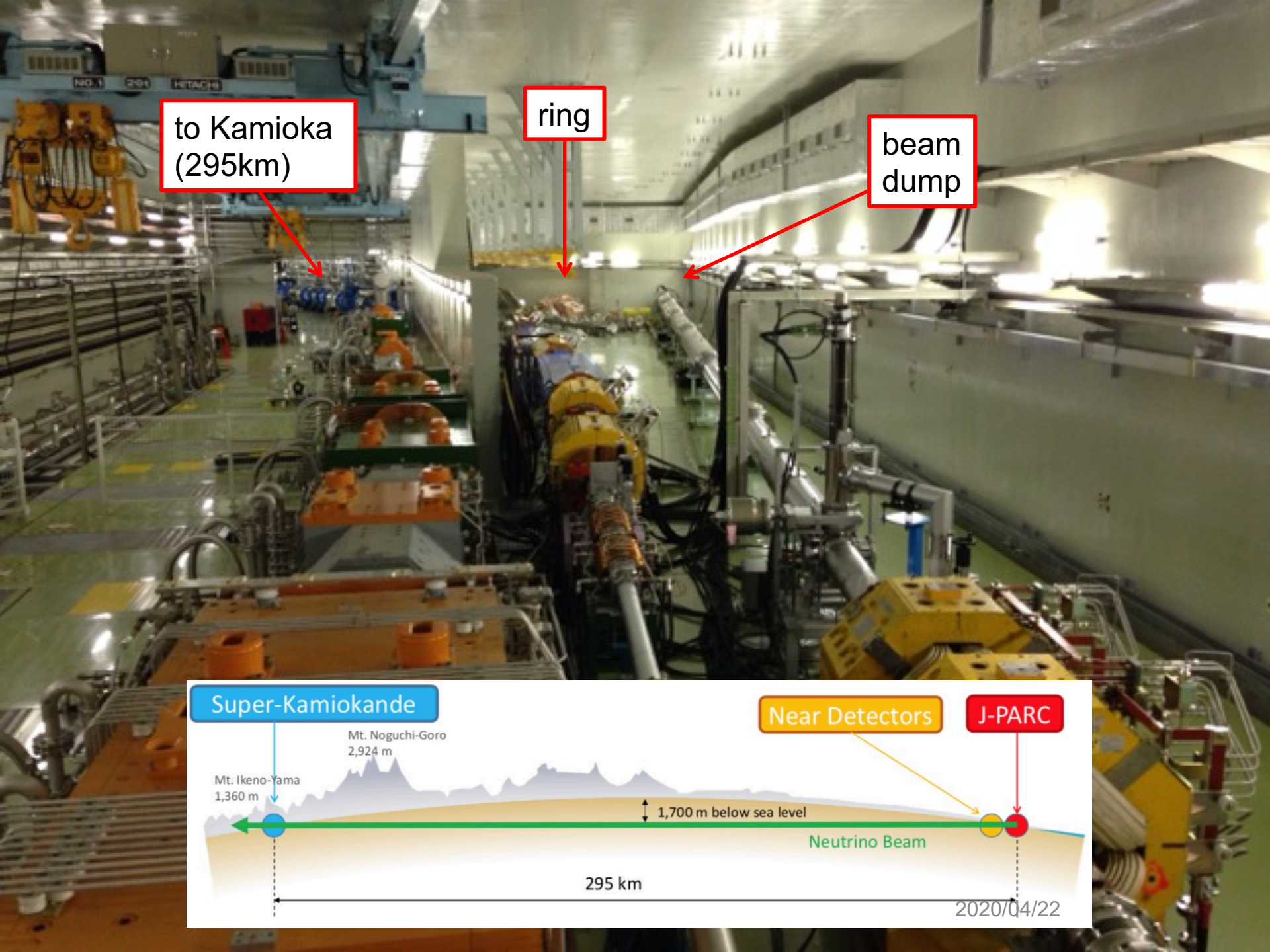
- 30 GeV protons are extracted from MR by superconducting magnets
- 1 pulse contains 8 bunches in ~5us, about ~2.5E14 ppp (protons per pulse) with 2.48 sec period



Neutrino events at the near detector



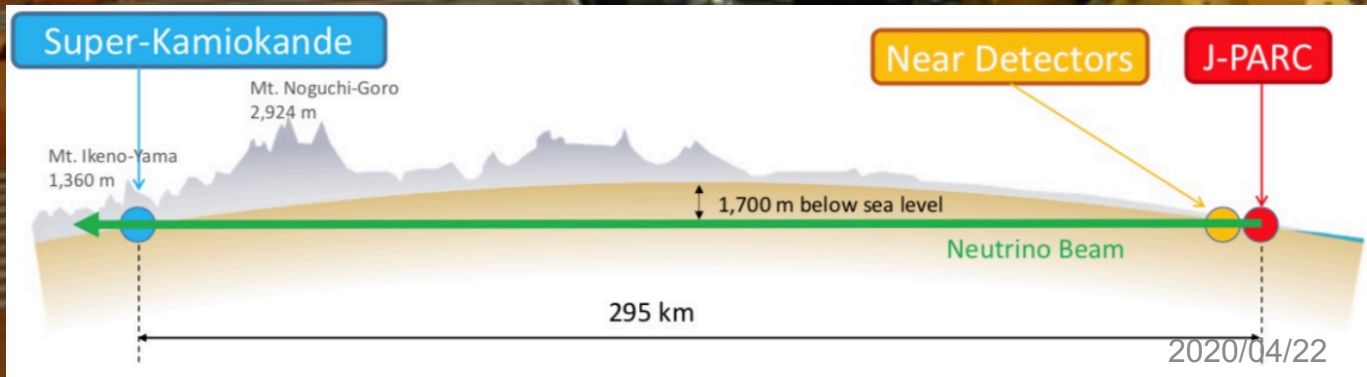
Neutrino events at the far detector



to Kamioka
(295km)

ring

beam
dump

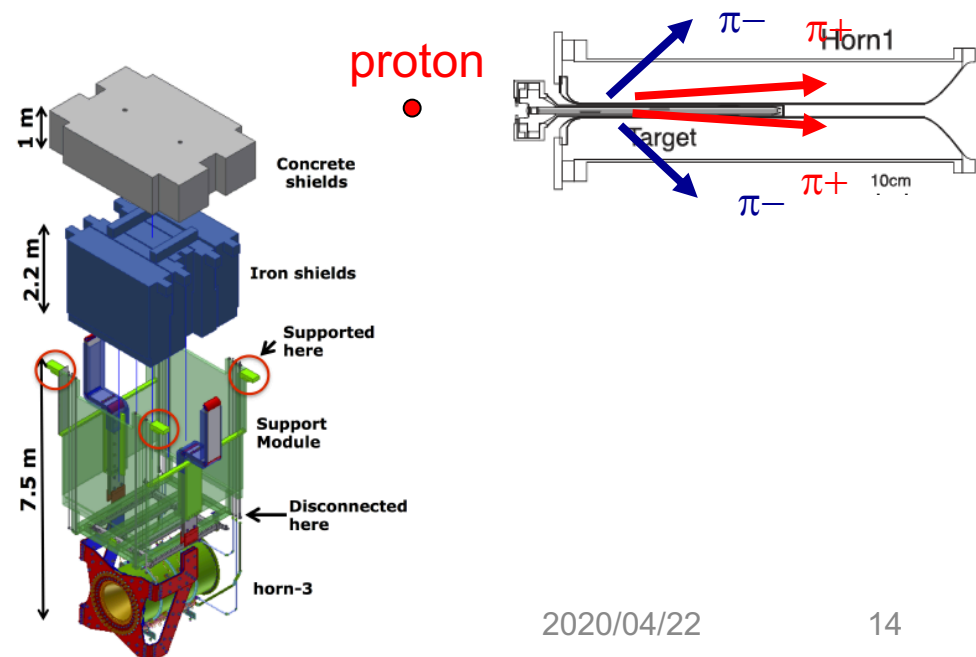
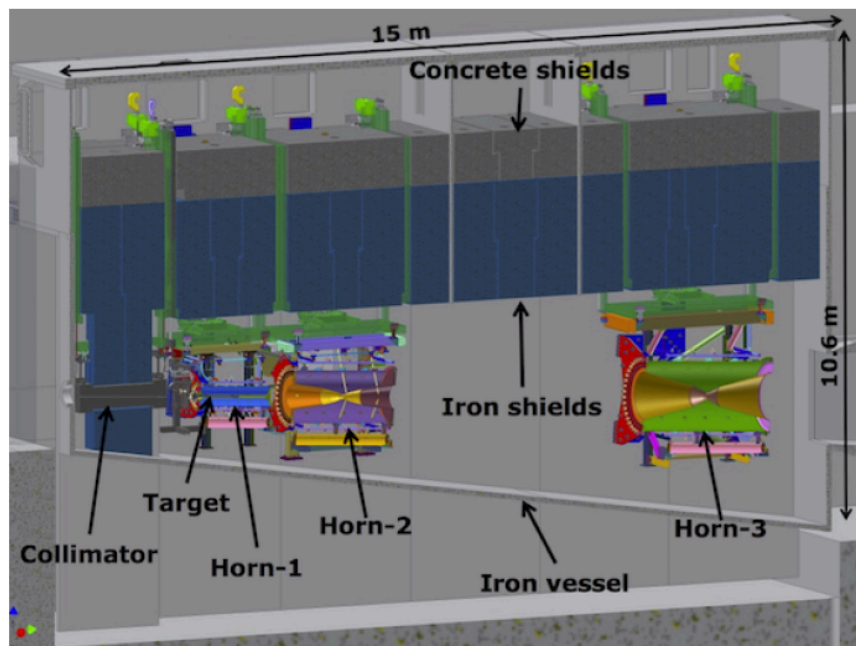
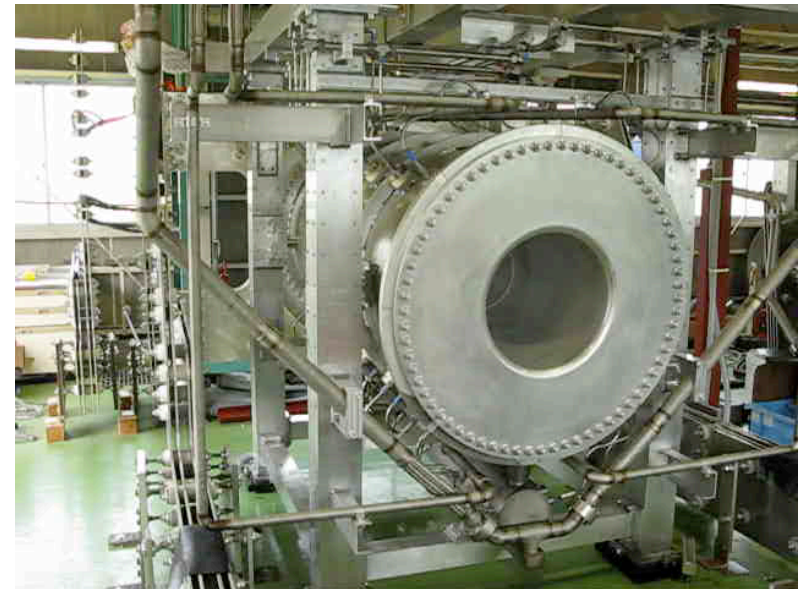


2. Neutrino beamline

Secondary beamline

- Protons collide the graphite target (in the Horn 1) to produce mesons, and these mesons decay in the decay volume to produce neutrinos (decay-in-flight).
- In **neutrino mode**, 3 magnetic horns focus positive mesons and defocus negative mesons to produce neutrino beam (flux $\sim \times 17$). In **antineutrino mode**, horn current is reversed to focus negative mesons.

Horn 3 test (250 kA, ~ 1.7 T)



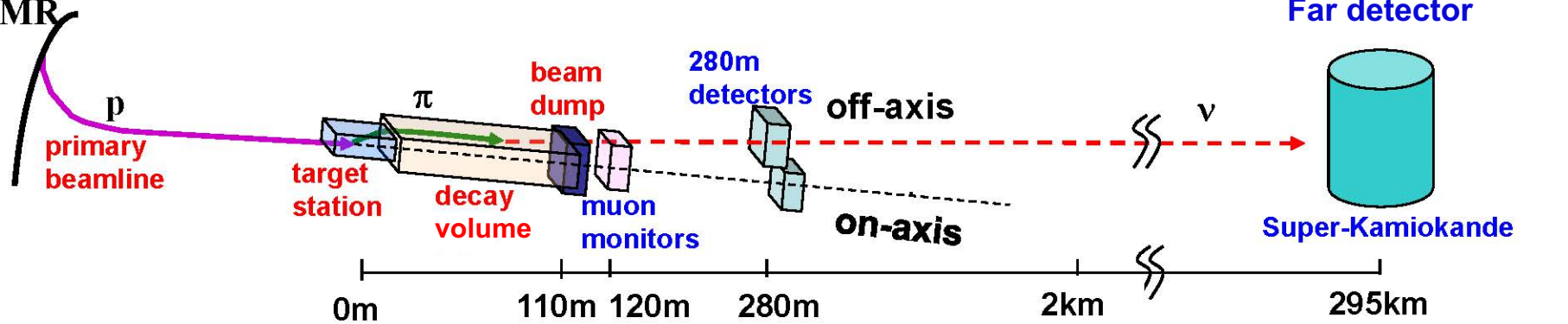
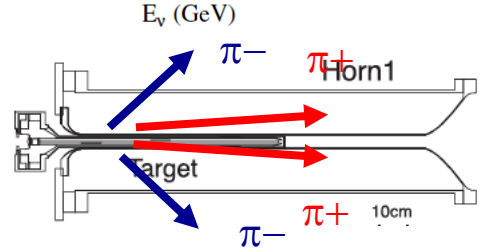
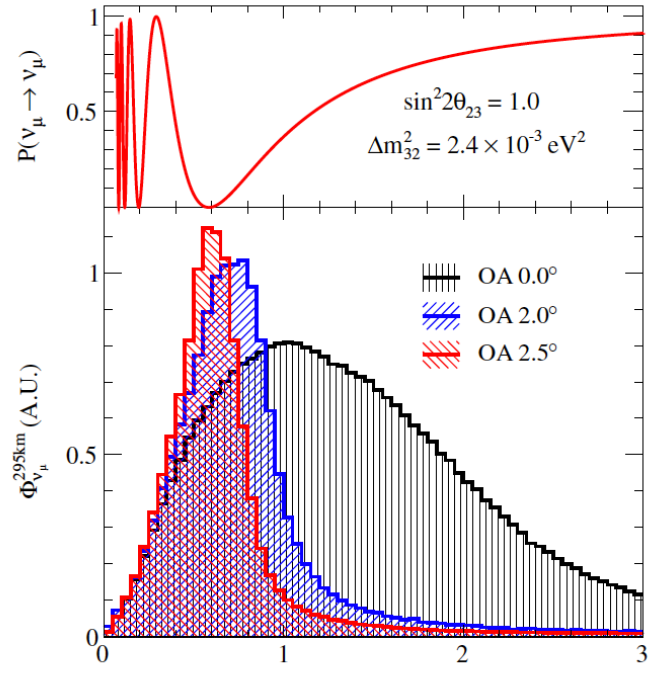
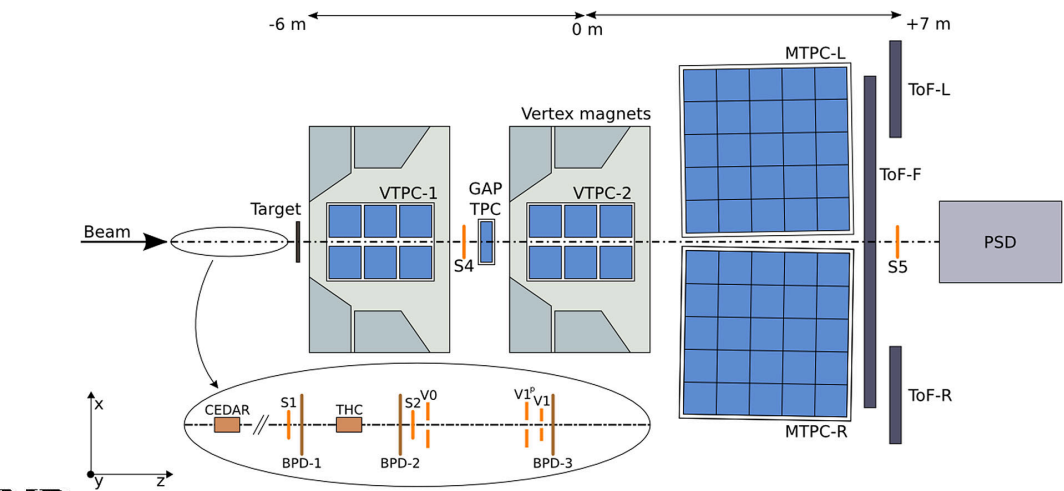
2. Neutrino beamline

Off-axis beam

- 2.5° off-axis to make ~0.6 GeV narrow band beam

CERN NA61/SHINE

- Hadron production at the target is simulated with the data from the hadron measurement



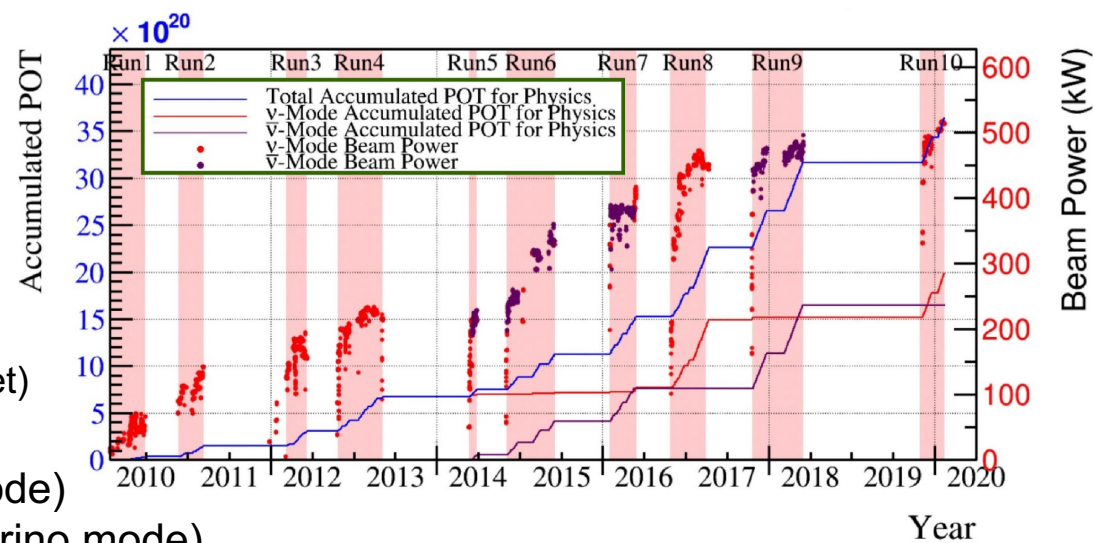
2. Neutrino beamline

2009 – 2018 data

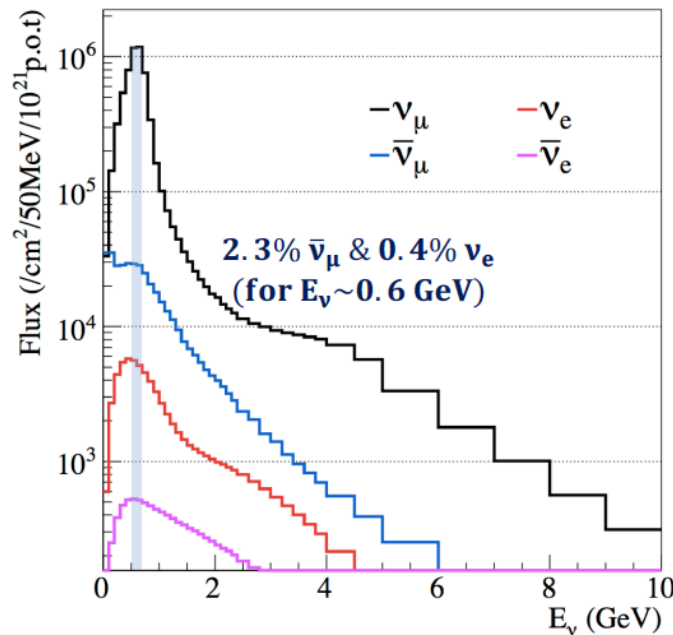
- Neutrino mode, $1.49\text{E}21$ POT
- Antineutrino mode, $1.64\text{E}21$ POT
(POT=protons on target)

Neutrino flux prediction

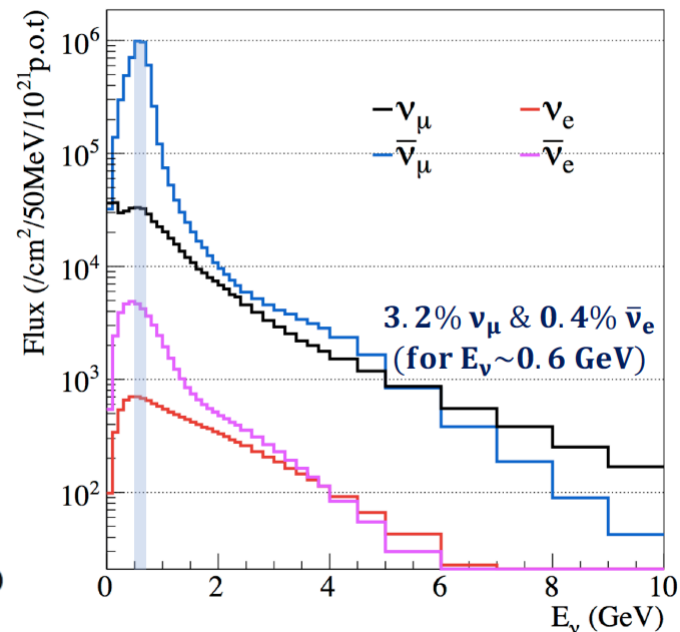
- muon neutrino dominant (neutrino mode)
- muon antineutrino dominant (antineutrino mode)
- ~9% error at the flux peak
- replica target NA61/SHINE data can reduce error to ~5%



Neutrino Mode Flux at SK



Antineutrino Mode Flux at SK



1. Neutrino oscillation experiments

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3. Neutrino detector

4. Neutrino interaction physics

5. Oscillation result

6. SK-Gd, T2K-Upgrade, Hyper-Kamiokande

7. Conclusion

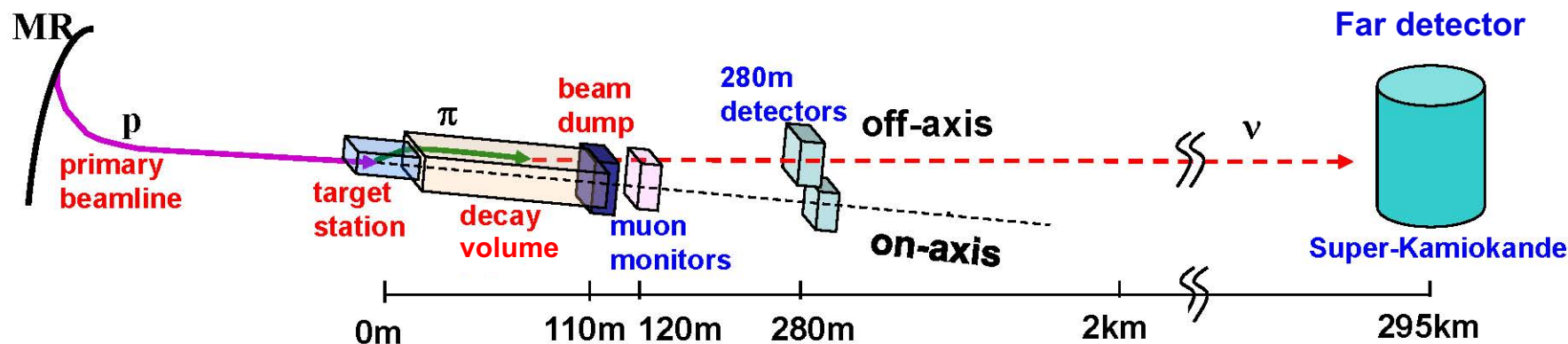
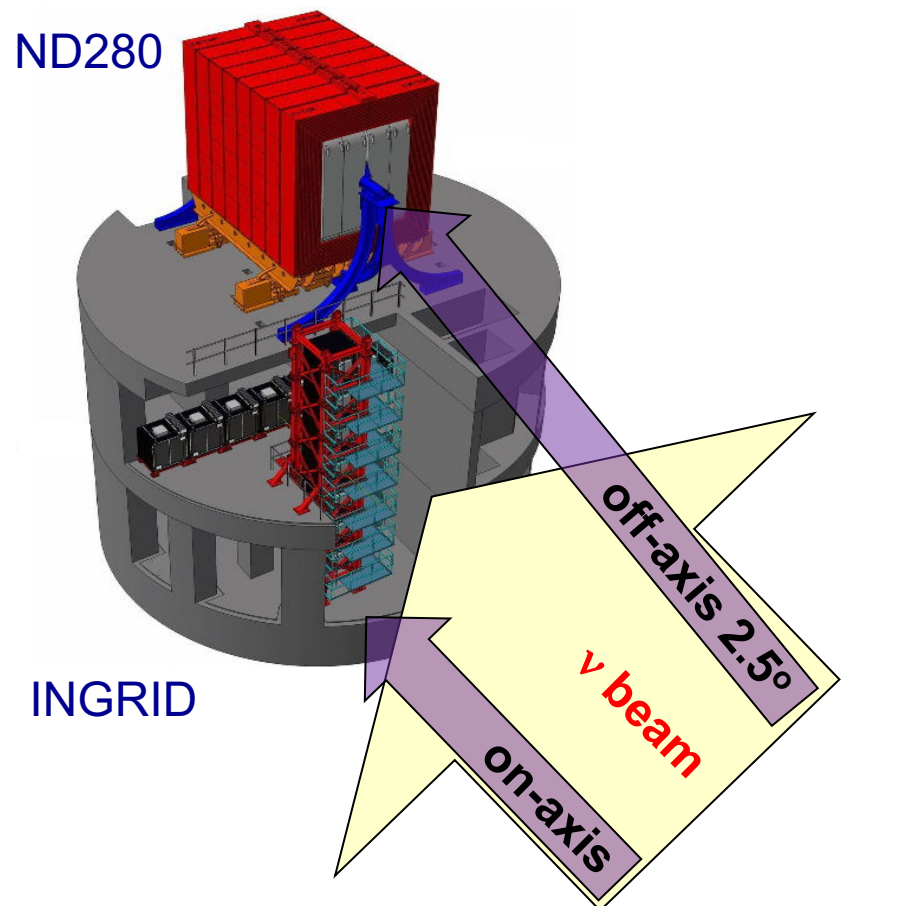
3. Near detectors

INGRID

- on-axis near detector
- Mainly for neutrino flux monitoring

ND280

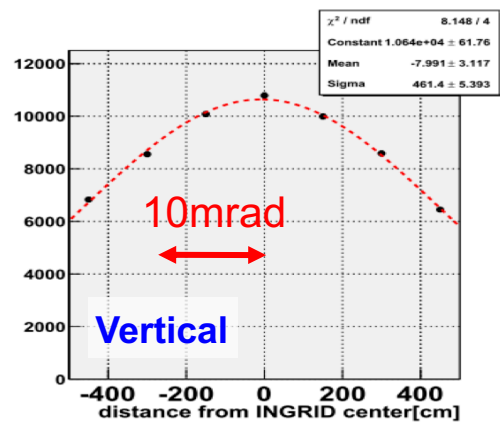
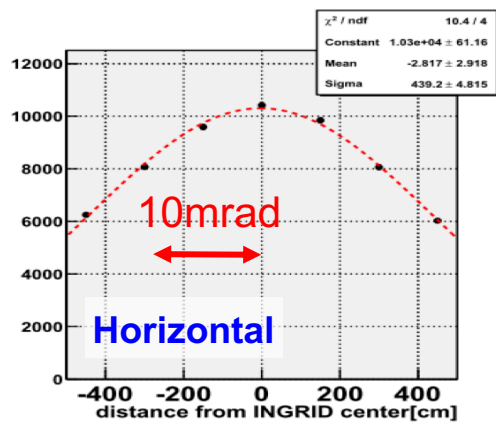
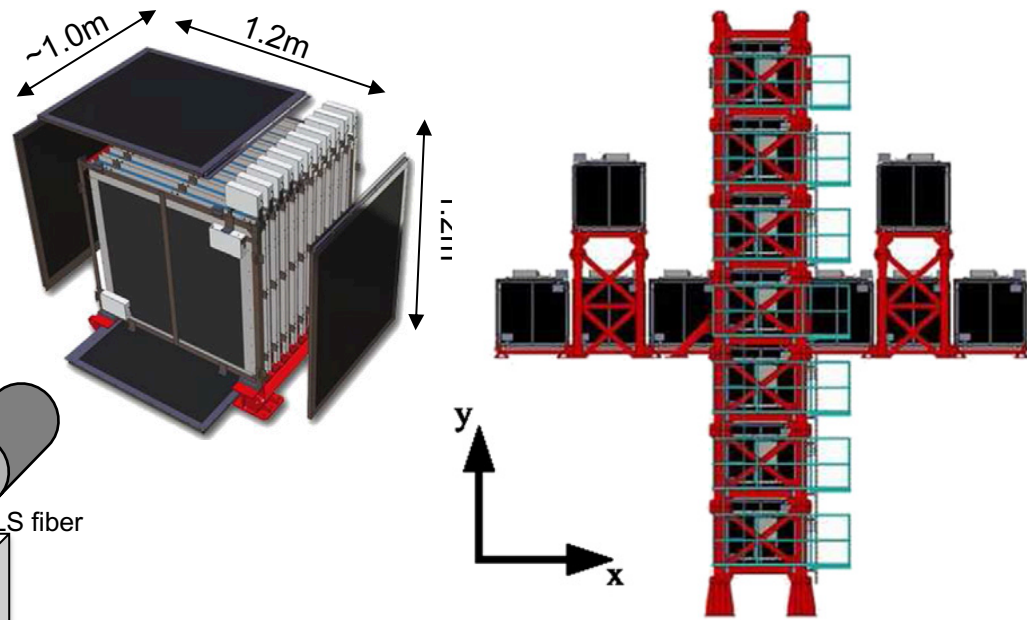
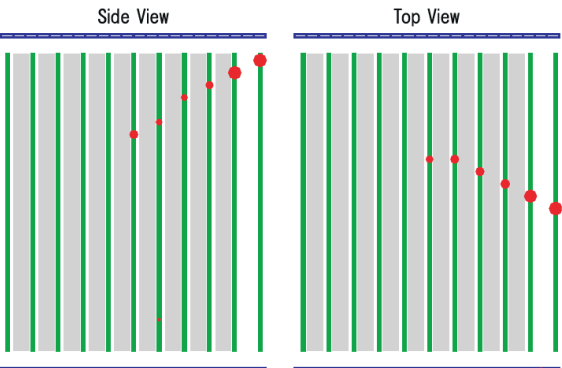
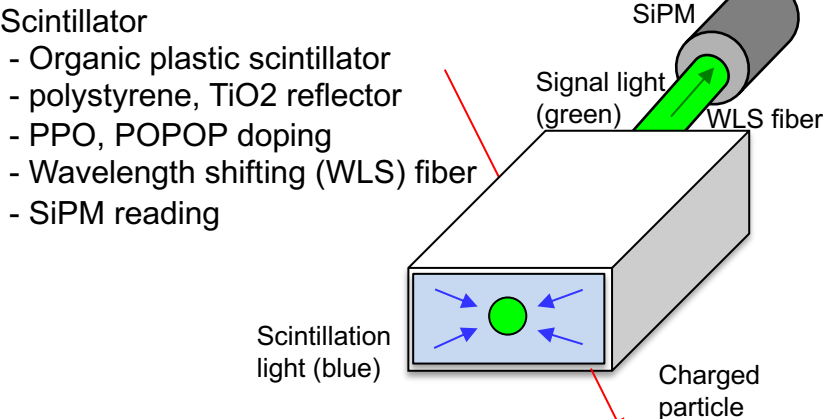
- off-axis near detector
- Data are used to constrain various systematics



3. On-axis detector

INGRID

- An array of 16 modules
- Scintillator-iron tracker
- nominal accuracy ~ 0.1 mrad



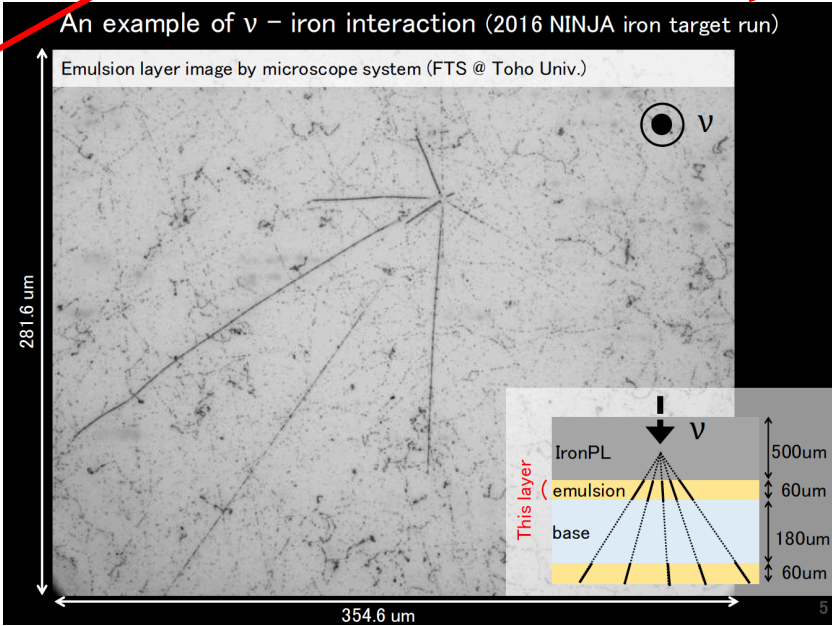
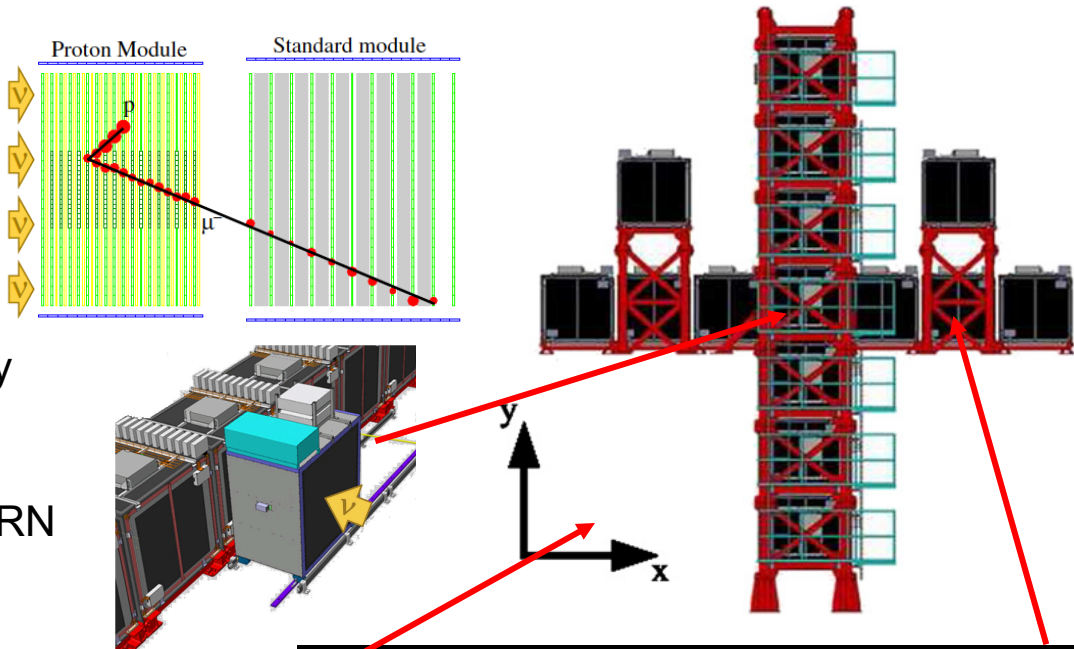
3. On-axis detectors

Proton module
- Fully active module

WAGASCI
- Water target 3-d scintillator array

BabyMIND
- Magnetized tracker made at CERN

NINJA
- Emulsion neutrino detector



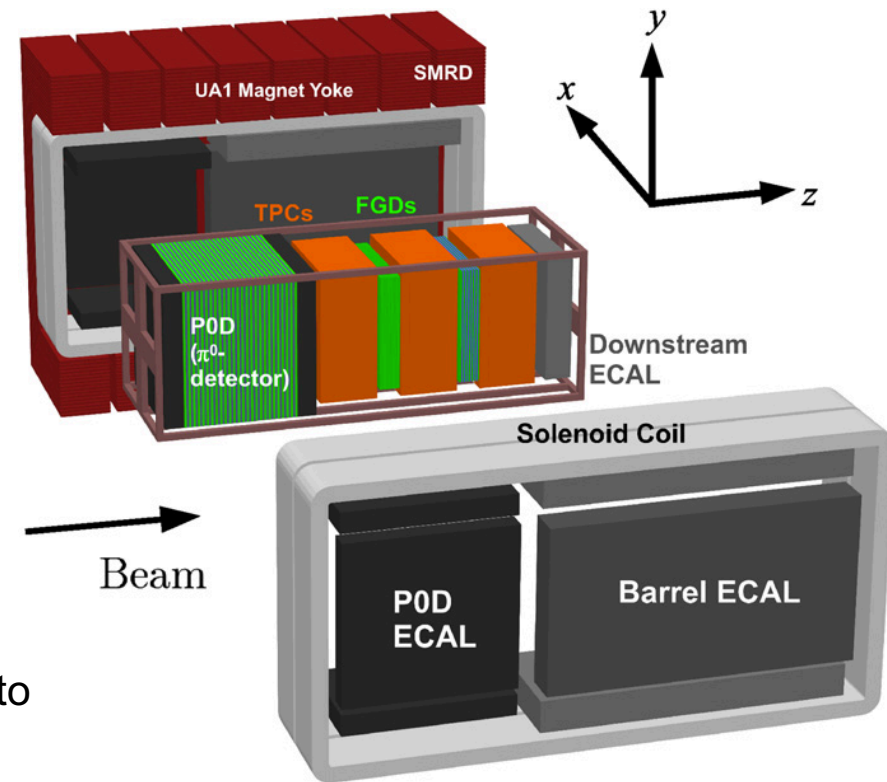
3. Off-axis detector

ND280

- P0D: Water-scintillator tracker
- FGD: Fully active scintillator tracker
- TPC: Ar gas TPC
- ECal: Lead-scintillator calorimeter
- SMRD: Iron-scintillator tracker
- UA1 magnet

Near detector data

- 14 samples are used for the oscillation analysis to constrain flux and cross-section systematic errors



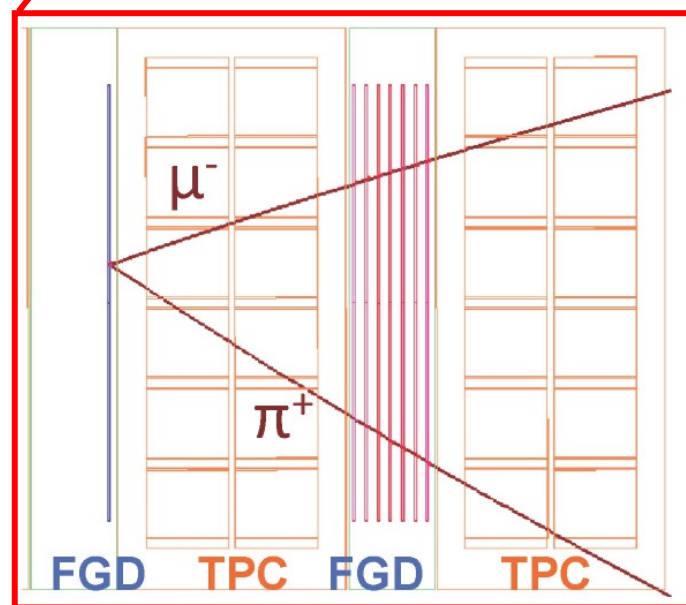
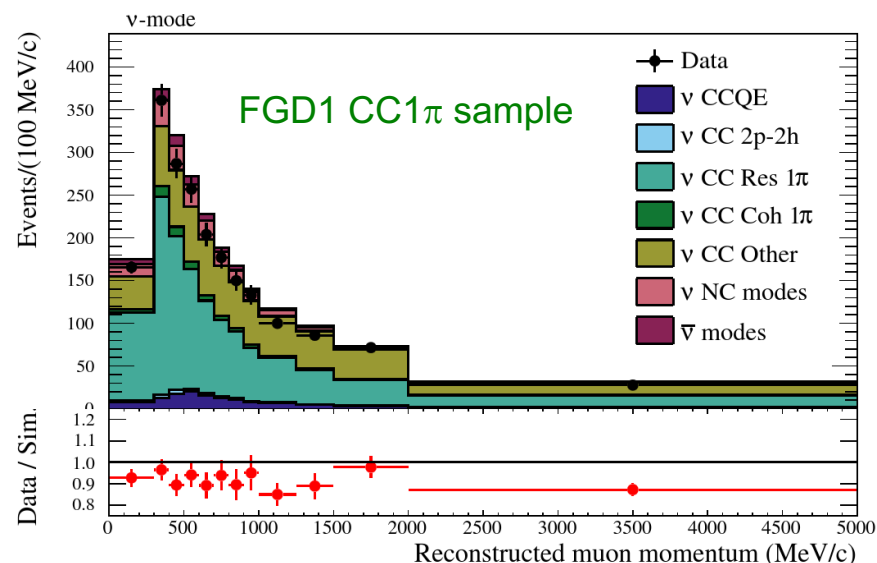
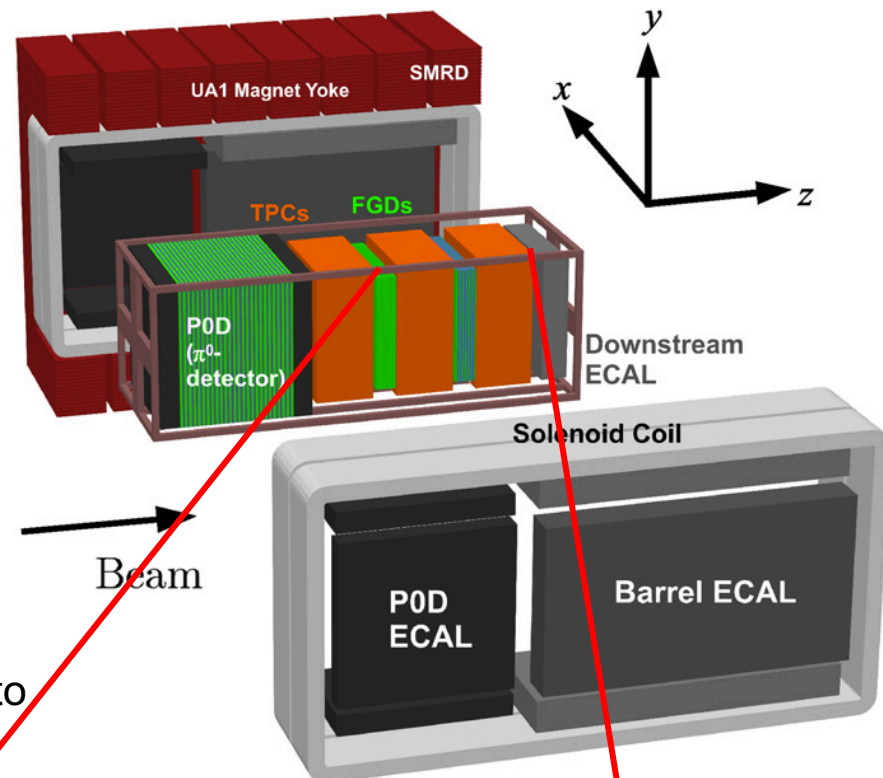
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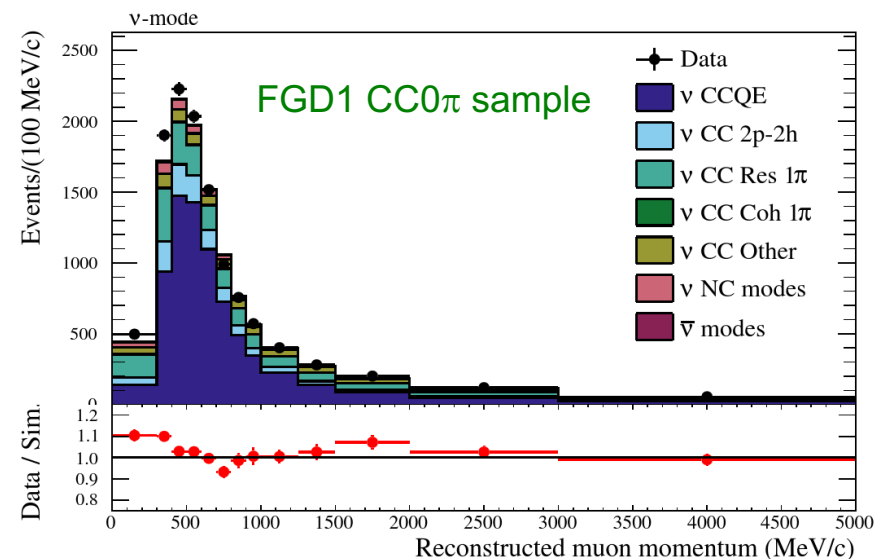
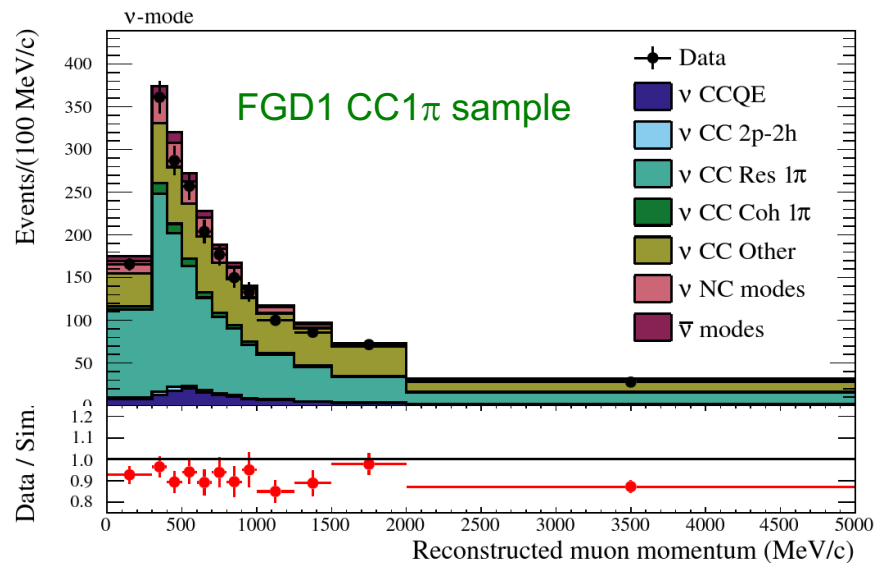
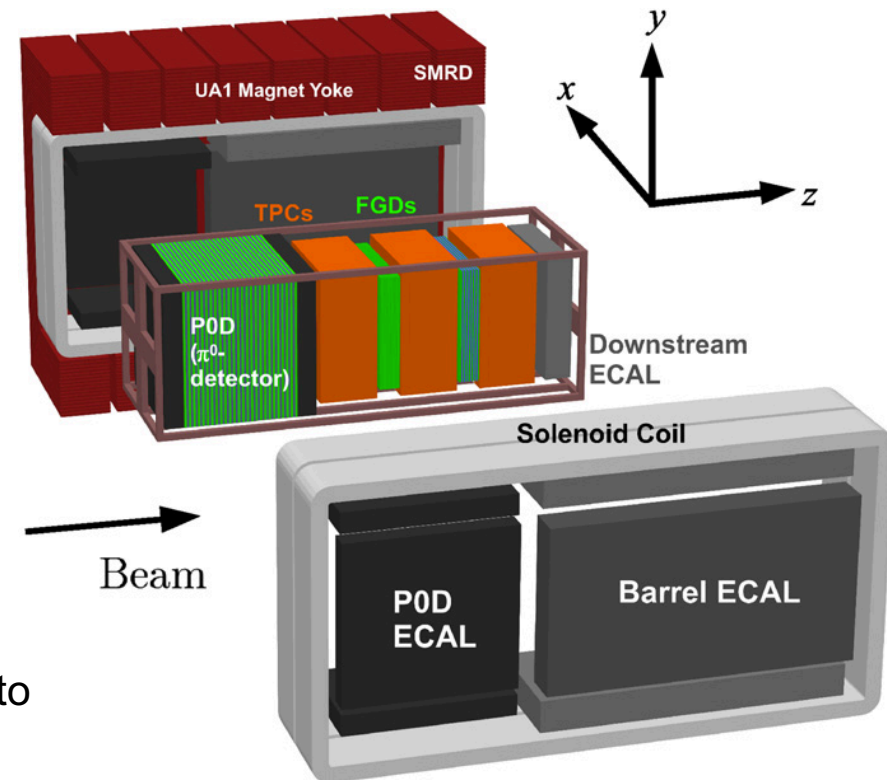
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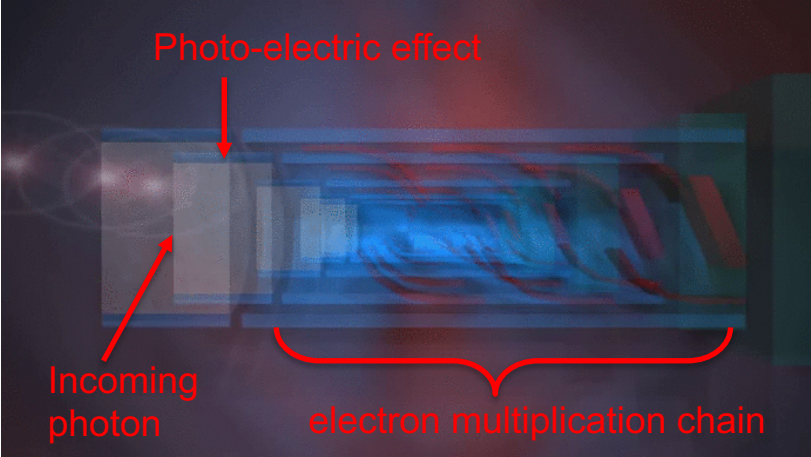


3. Far detector

Super-Kamiokande

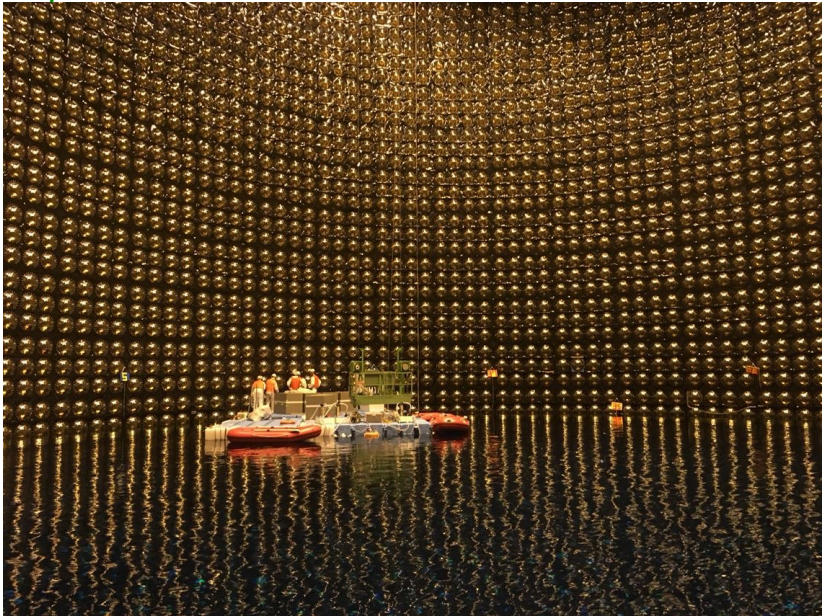
- 50 kton water Cherenkov detector
- 2015 Nobel prize
- 11,146 20-inch PMTs (inner detector)
- 1,885 8-inch PMTs (outer detector)

Photo-multiplier tube (PMT)



20-inch PMT is quite big...

Super-K inner detector



Super-K outer detector



OD PMT unit
- 8-inch PMT
- wave-length
shifting plate

White Tyvek
reflector

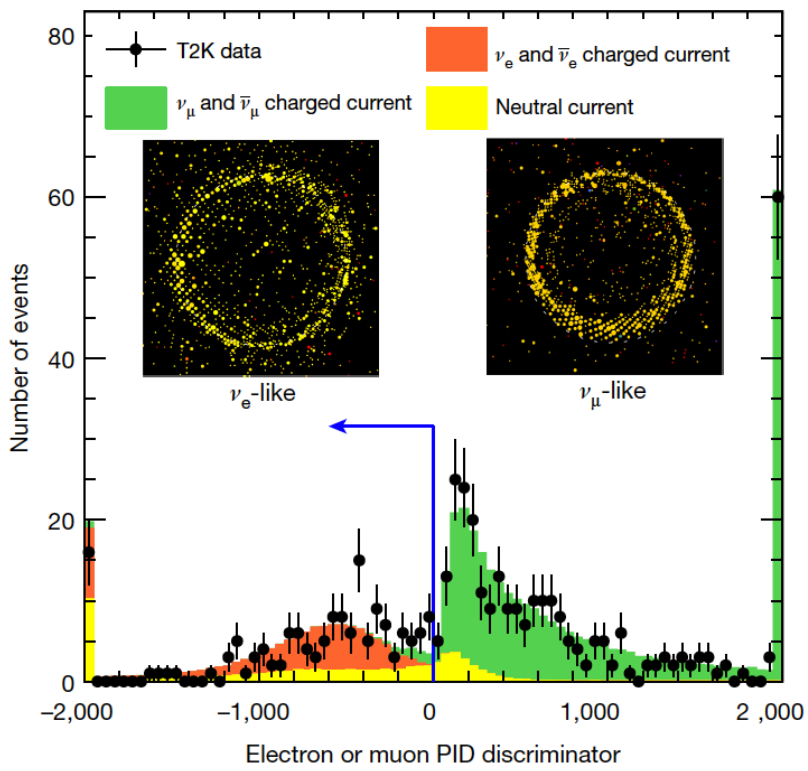
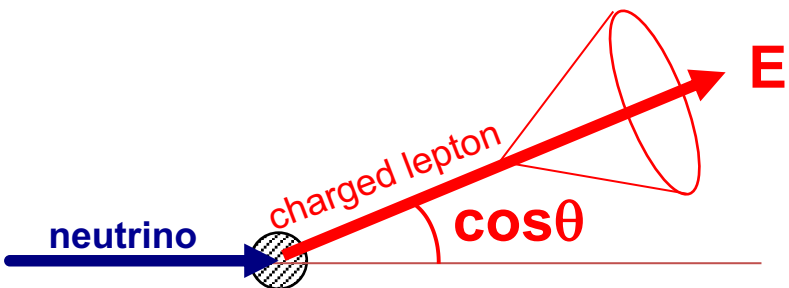
katori@fnal.gov

3. Far detector

Event reconstruction

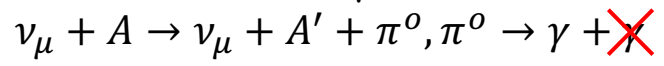
- From measured time and charge information from all PMTs, particle identification (PID) and kinematics are reconstructed
- From reconstructed charged lepton kinematics, neutrino energy is reconstructed

$$E_{\nu}^{QE} = \frac{ME - 0.5m_l^2}{M - E + p\cos\theta}$$



$\nu_e(\bar{\nu}_e)$ measurement has 2 major backgrounds

1. Intrinsic background
 $\nu_e(\bar{\nu}_e)$ contamination in the beam (~0.5%)
2. misID background
Gamma rays counted as electron (positron). Majority of them are from neutral current π^0 production where one of γ is undetected



1. Neutrino oscillation experiments

2. Neutrino beam

3. Neutrino detector

4. Neutrino interaction physics

5. Oscillation result

6. SK-Gd, T2K-Upgrade, Hyper-Kamiokande

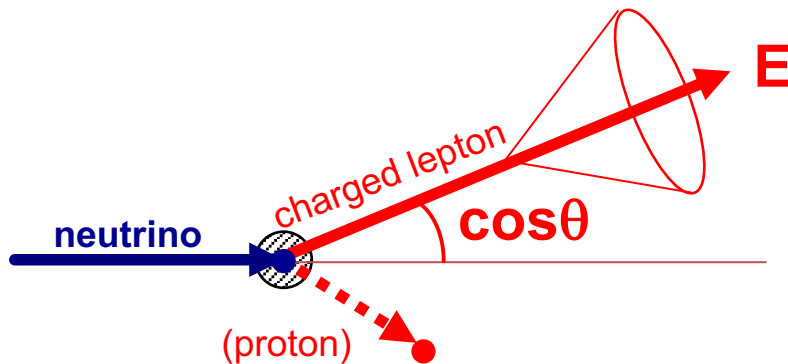
7. Conclusion

4. Charged current quasi-elastic (CCQE) scattering

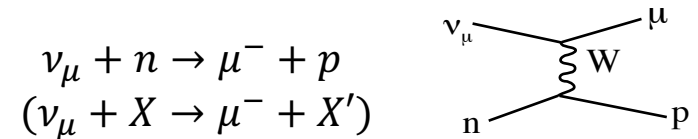
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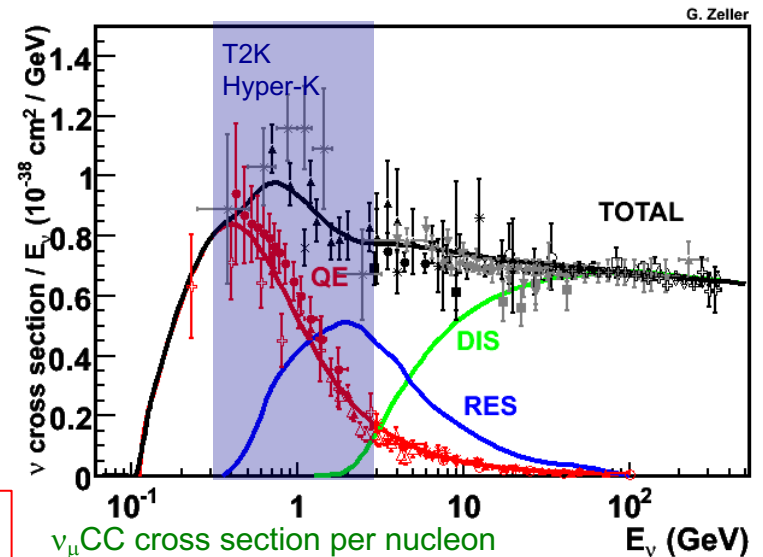
CCQE is the most abundant interaction at ~ 1 GeV.



Neutrino energy is reconstructed from the observed lepton kinematics

“QE assumption”

1. assuming neutron at rest
2. assuming interaction is CCQE (2-body kinematics)



All neutrino cross-section channels (including CCQE) have large error

4. CCQE puzzle

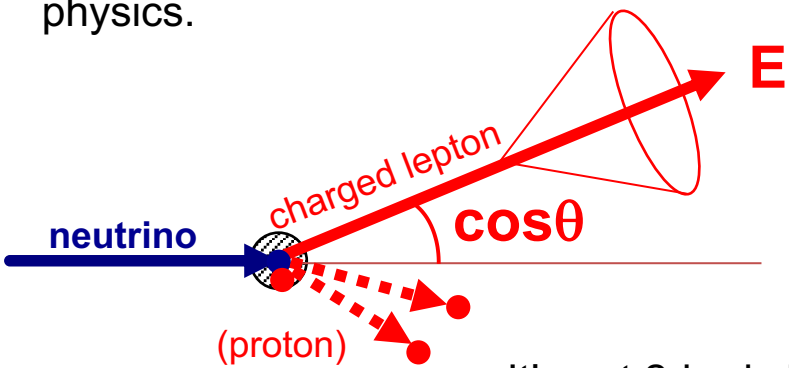
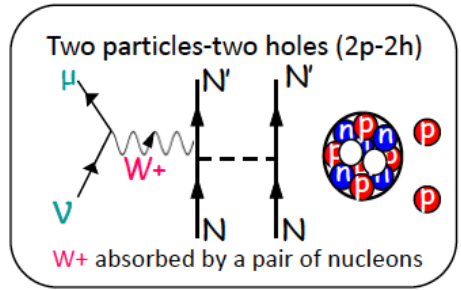
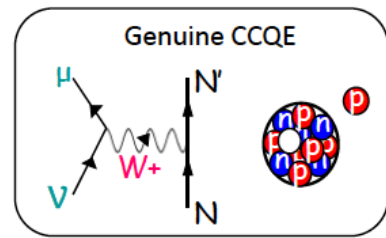
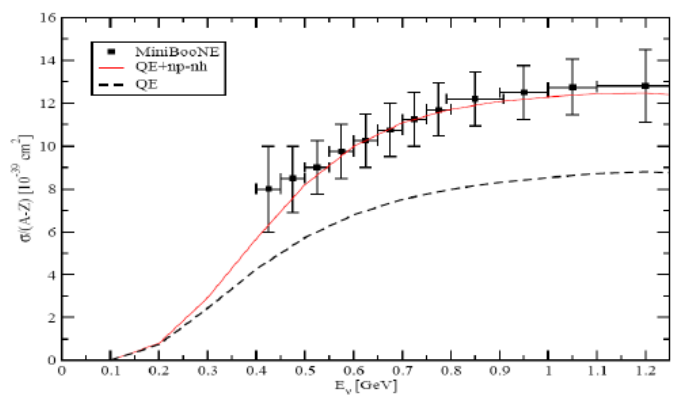
Nuclear correlations

- Martini et al pointed out that neutrino interactions around 1 GeV can be modified ~30% by correlated nucleons (2p2h, 2-body current, meson exchange current, etc)

A large community effort (both theorists and experimentalists) to understand the role of nucleon correlations in neutrino interaction physics.

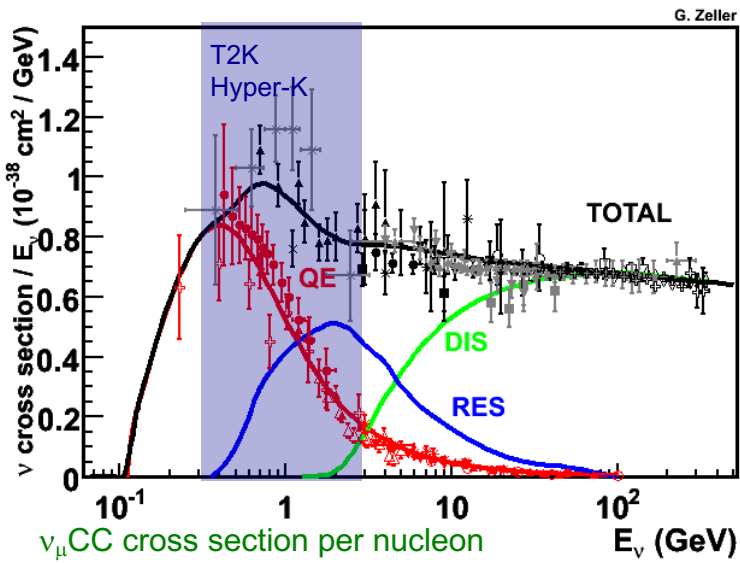
An explanation of this puzzle

Inclusion of the multinucleon emission channel (np-nh)



It's not 2-body kinematics

$$E_{\nu}^{QE} \neq \frac{ME - 0.5m_l^2}{M - E + p \cos \theta}$$



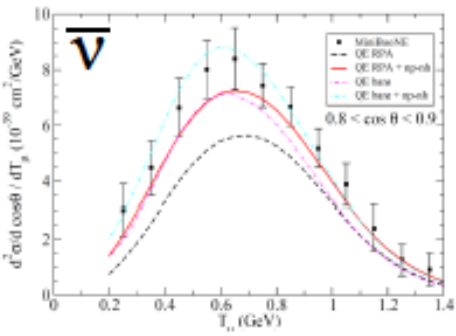
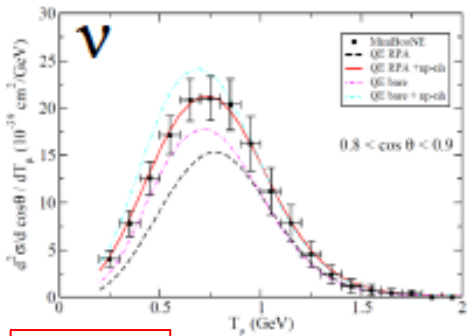
4. CCQE puzzle (2019)

Advanced nuclear models can reproduce MiniBooNE CCQE-like data, but there are large systematics errors on nuclear parameters.

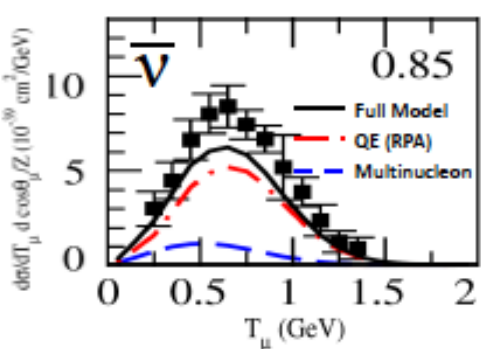
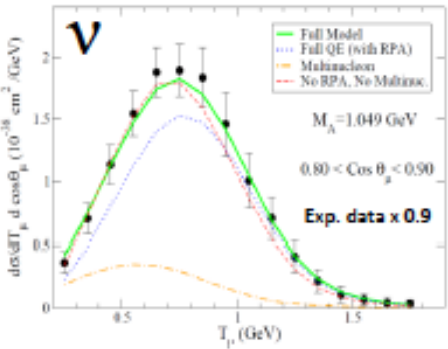
- Martini – RPA+2p2h
- Nieves – Valencia 2p2h model
- SuSA – Superscaling+MEC
- Giusti – Relativistic Green’s function
- Butkevich – RDWIA+MEC

We use Valencia 2p2h model for our simulation

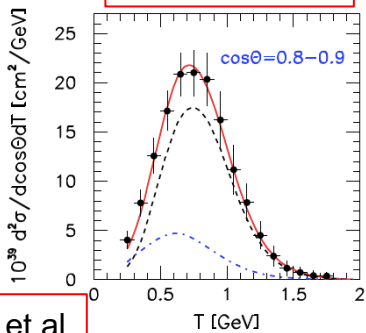
Martini et al



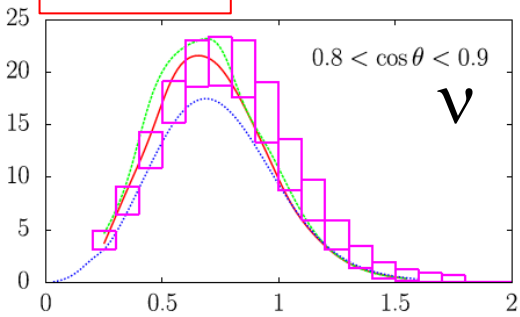
Valencia



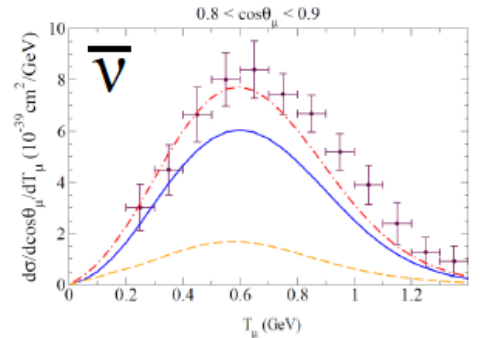
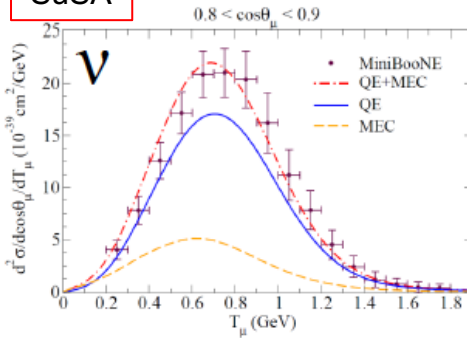
Butkevich et al



Giusti et al



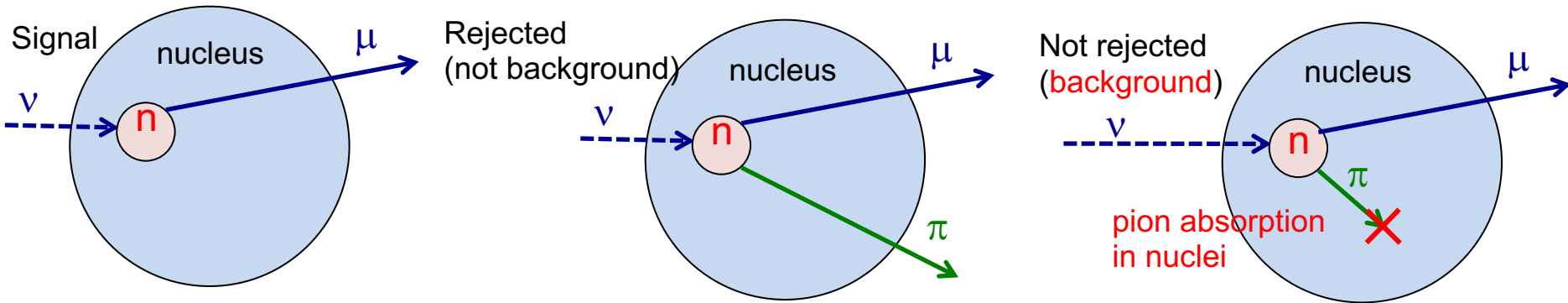
SuSA



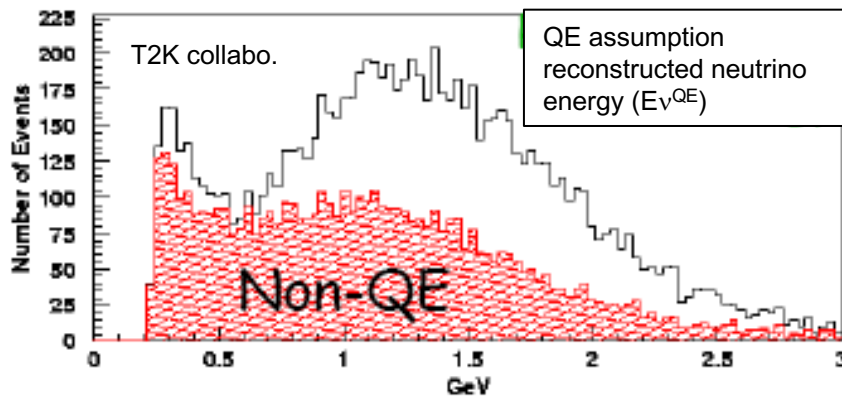
4. Neutrino-induced single pion production

Baryon resonant pion production + final state interaction (FSI)

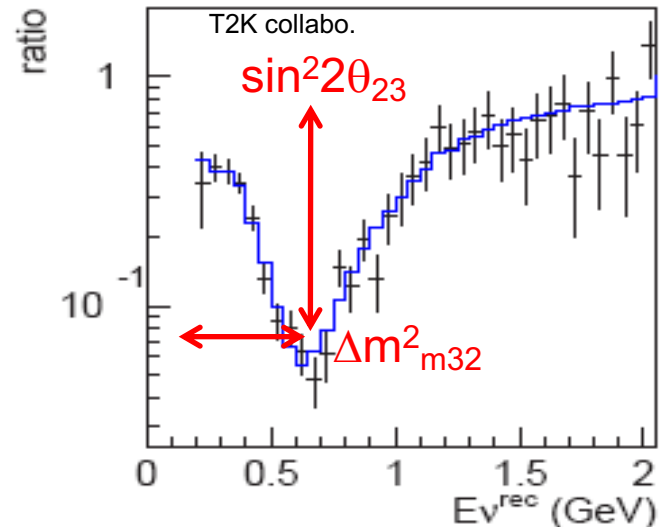
- Neutrino induced pion productions have large errors
- Final state interaction of hadrons have large errors



muon neutrino disappearance simulation



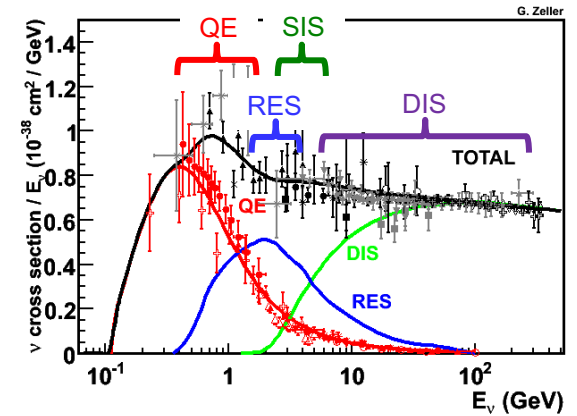
katori@fnal.gov



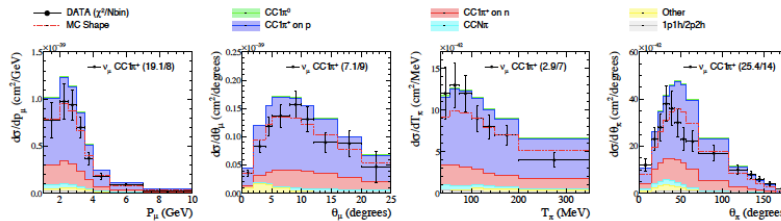
4. Pion puzzle (2019)

MINERvA simultaneous fit for 4 different data set

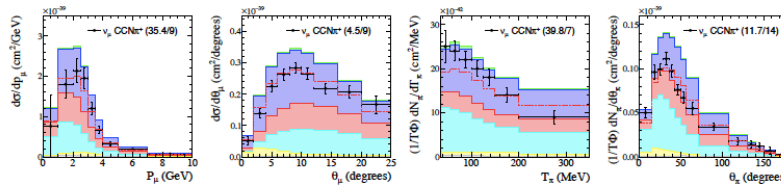
- Most advanced study in this community
- Not conclusive on baryon resonance and FSI models



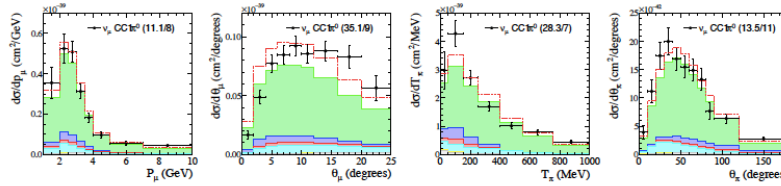
$\nu_\mu \text{CC}1\pi^+$



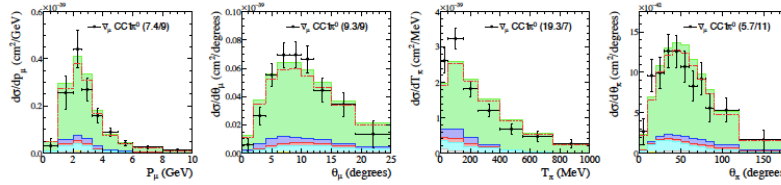
$\nu_\mu \text{CCN}\pi^+$



$\nu \text{CC}1\pi^0$

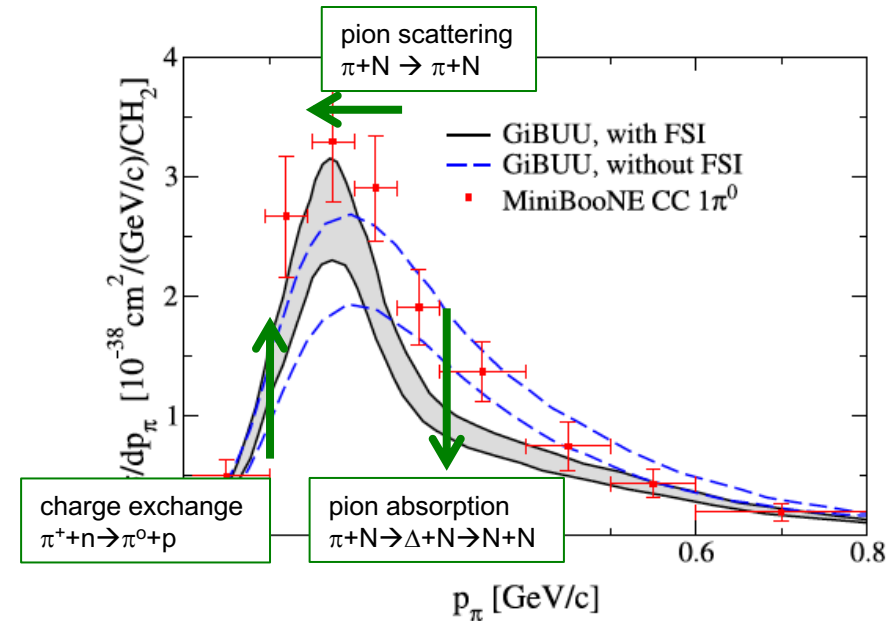


$\bar{\nu} \text{CC}1\pi^0$



GiBUU vs. MiniBooNE CC\$\pi^0\$ data

- You need to simulate both CC\$\pi^0\$ and CC\$\pi^\pm\$, and FSI including inelastic scattering, charge exchange, pion absorption



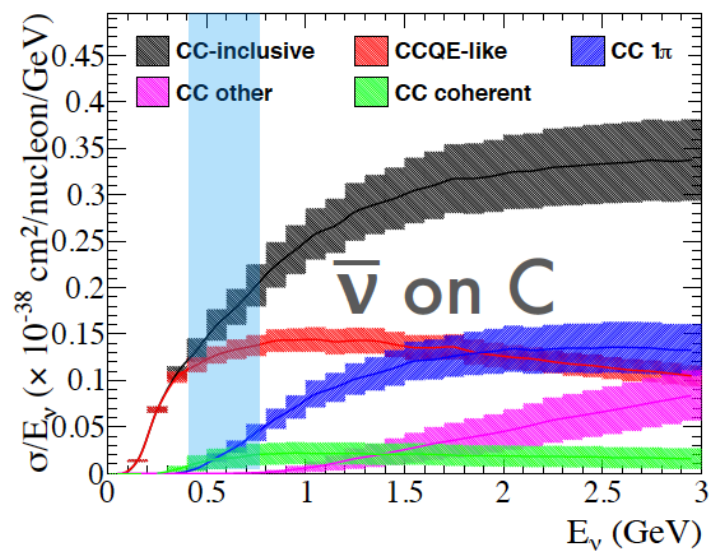
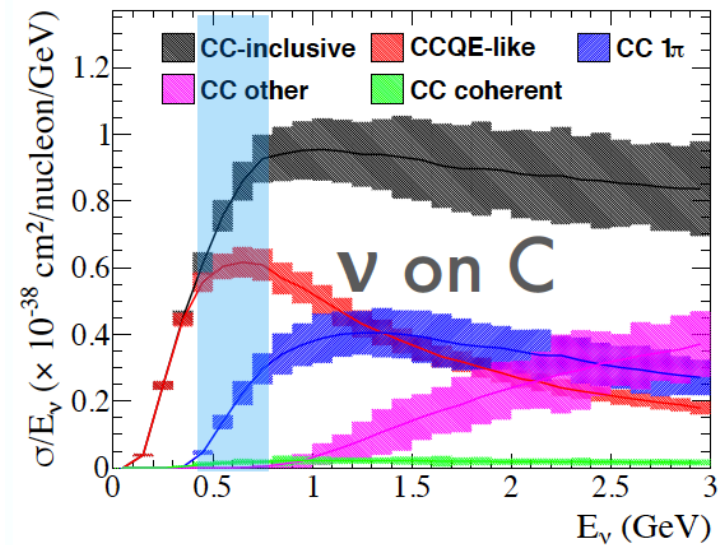
4. Neutrino interaction physics, external data constraints

We accept large systematic errors on neutrino interaction models

We need to constrain these errors **internally**, using the data from the ND280 near detector data

PDG (2019)
Section 42. Monte Carlo Neutrino Generators
Section 50. Neutrino Cross Section Measurements

NuSTEC (<https://nustec.fnal.gov/>)
New theory-experiment collaboration to promote neutrino interaction physics



1. Neutrino oscillation experiments

2. Neutrino beam

3. Neutrino detector

4. Neutrino interaction physics

5. Oscillation result

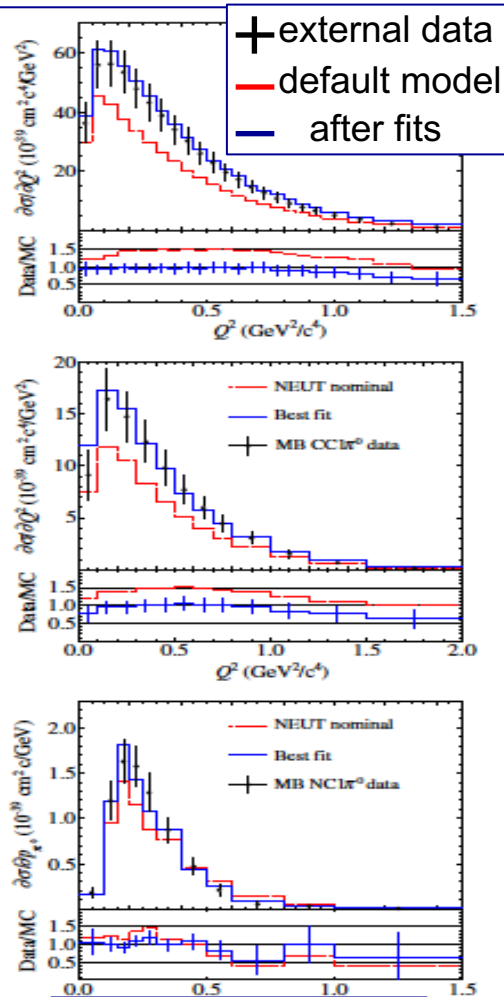
6. SK-Gd, T2K-Upgrade, Hyper-Kamiokande

7. Conclusion

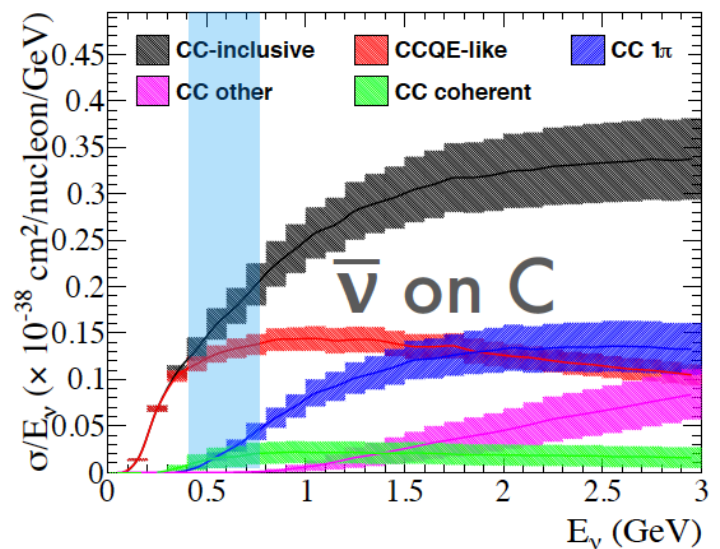
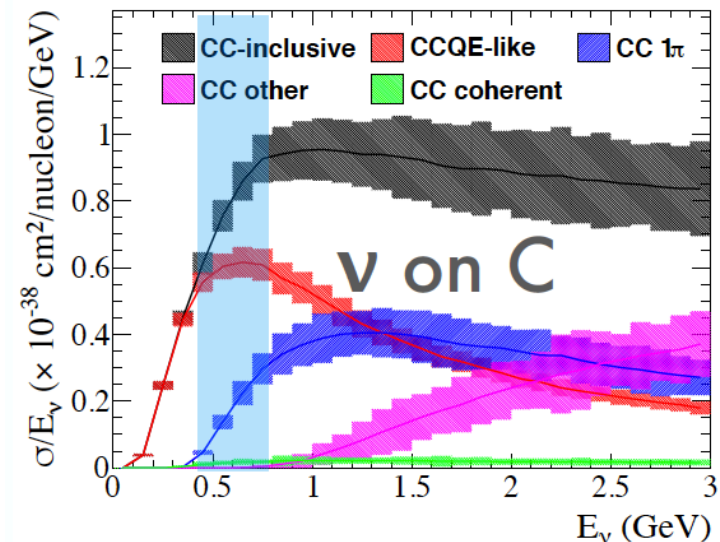
5. T2K oscillation results

External constraint

MiniBooNE, MINERvA, SciBooNE
K2K, MINOS, Bubble chambers



External data give
initial guess of cross-
section systematics



External data fit

5. T2K oscillation results

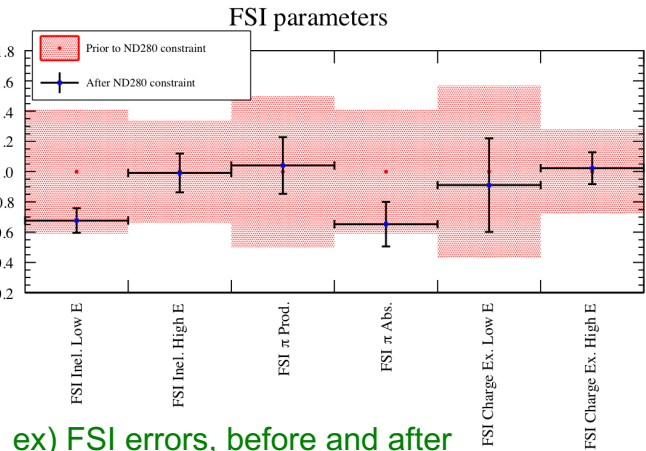
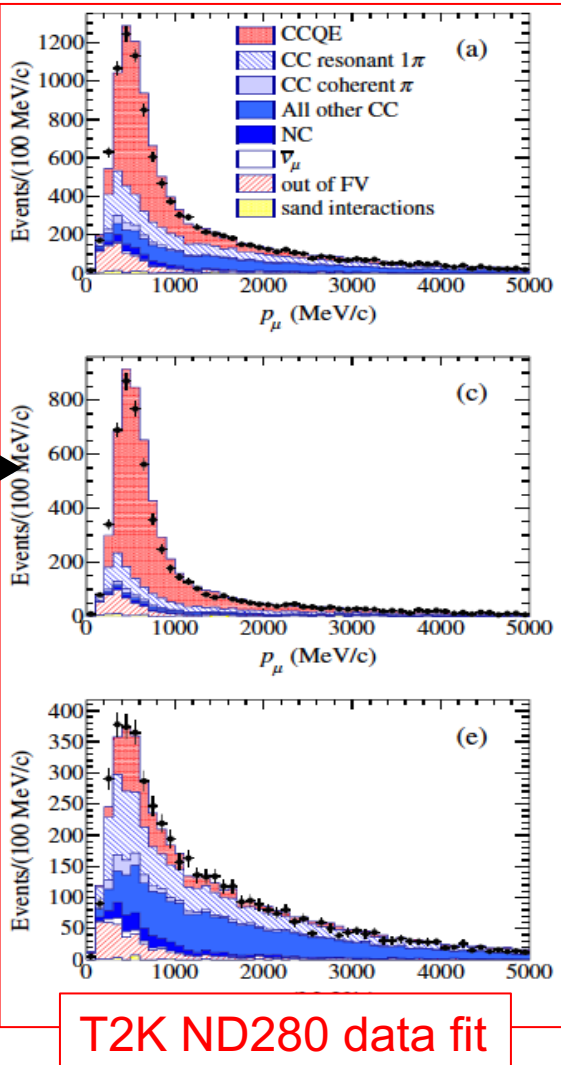
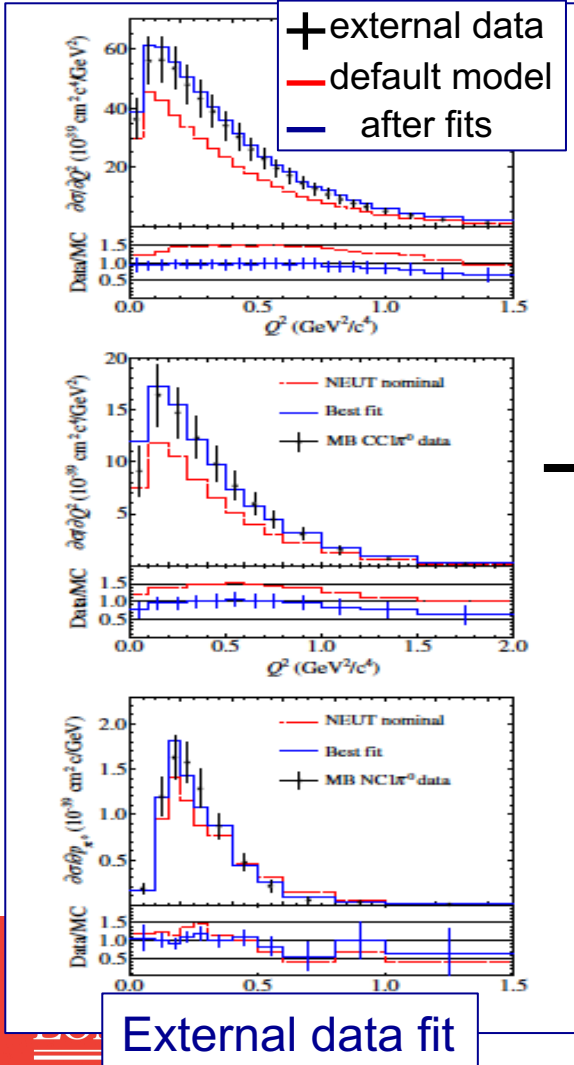
External constraint
MiniBooNE, MINERvA, SciBooNE
K2K, MINOS, Bubble chambers

Internal constraint
Near detector
oscillation non-sensitive channels

Internal data can constrain
systematic errors for the event rate
(flux x cross-section)

SuperK sample systematic error

sample	Without ND280	With ND280
ν μ -like ring	14.6%	5.1%
ν e-like ring	16.9%	8.8%
$\bar{\nu}$ μ -like ring	12.5%	4.5%
$\bar{\nu}$ e-like ring	14.4%	7.1%



ex) FSI errors, before and after
internal constraints

5. T2K oscillation results

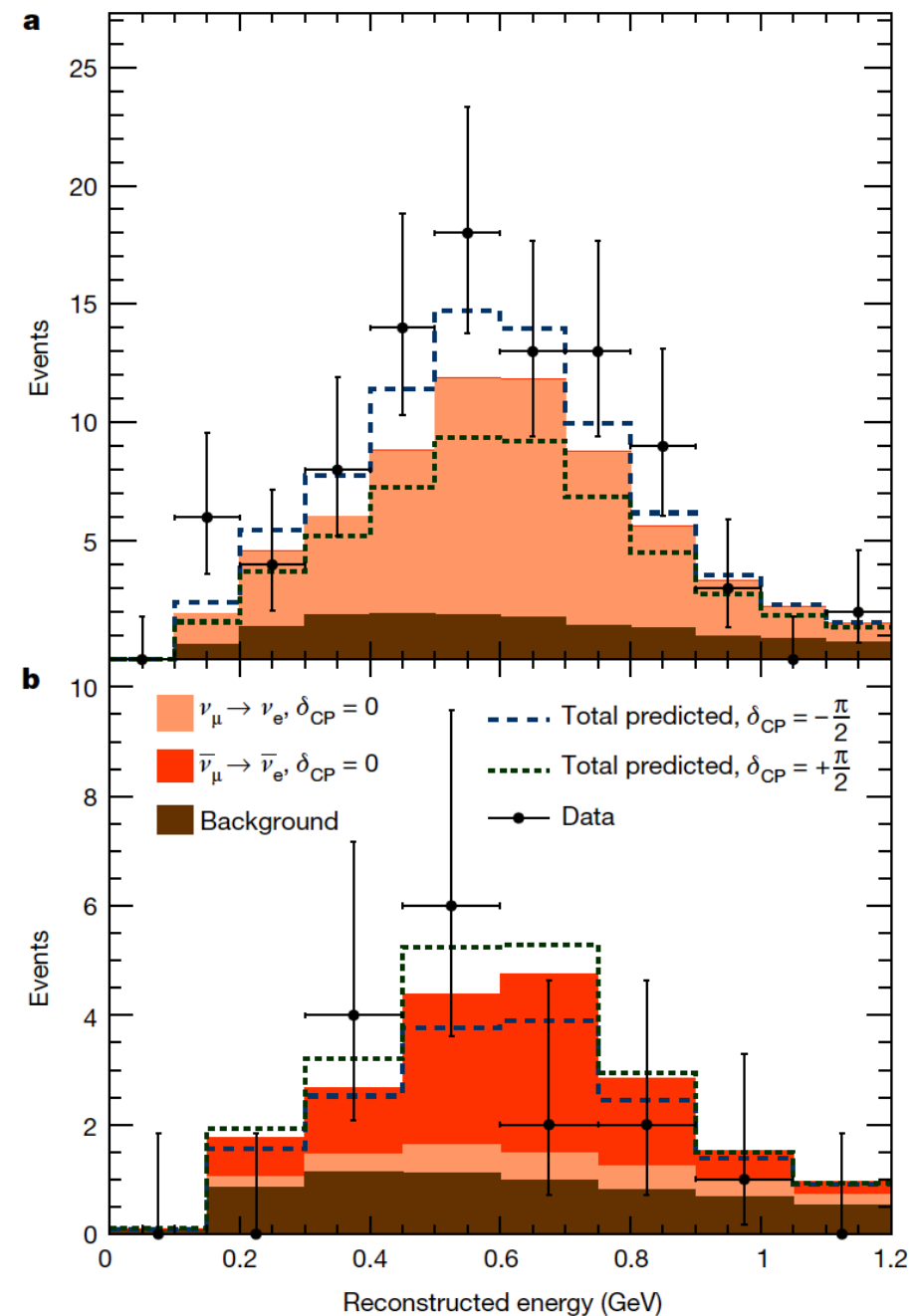
SuperK data prefer a model with negative CP violation angle ($\sim -\pi/2$)

- Enhancement of $P(\nu_\mu \rightarrow \nu_e)$
- Suppression of $P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)$

c	1e0de ν -mode	1e0de $\bar{\nu}$ -mode	1e1de ν -mode
$\nu_\mu \rightarrow \nu_e$	59.0	3.0	5.4
$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$	0.4	7.5	0.0
Background	13.8	6.4	1.5
Total predicted	73.2	16.9	6.9
Systematic uncertainty	8.8%	7.1%	18.4%
Data	75	15	15

2009 – 2018 data

- Neutrino mode, 1.49E21 POT
- Antineutrino mode, 1.64E21 POT



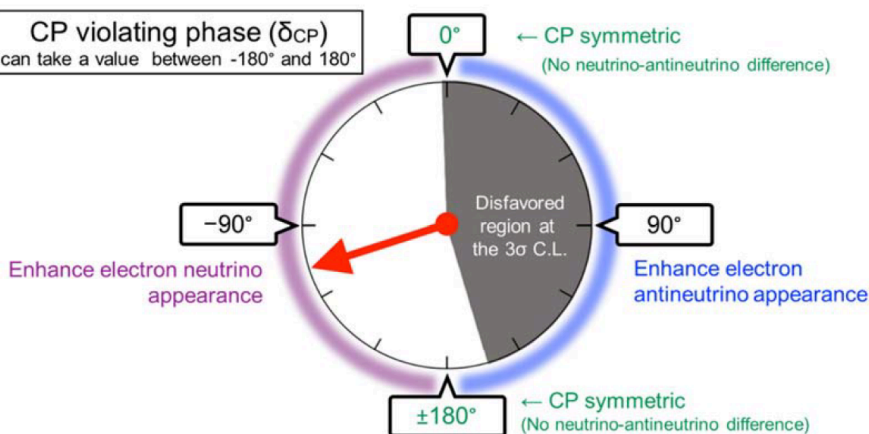
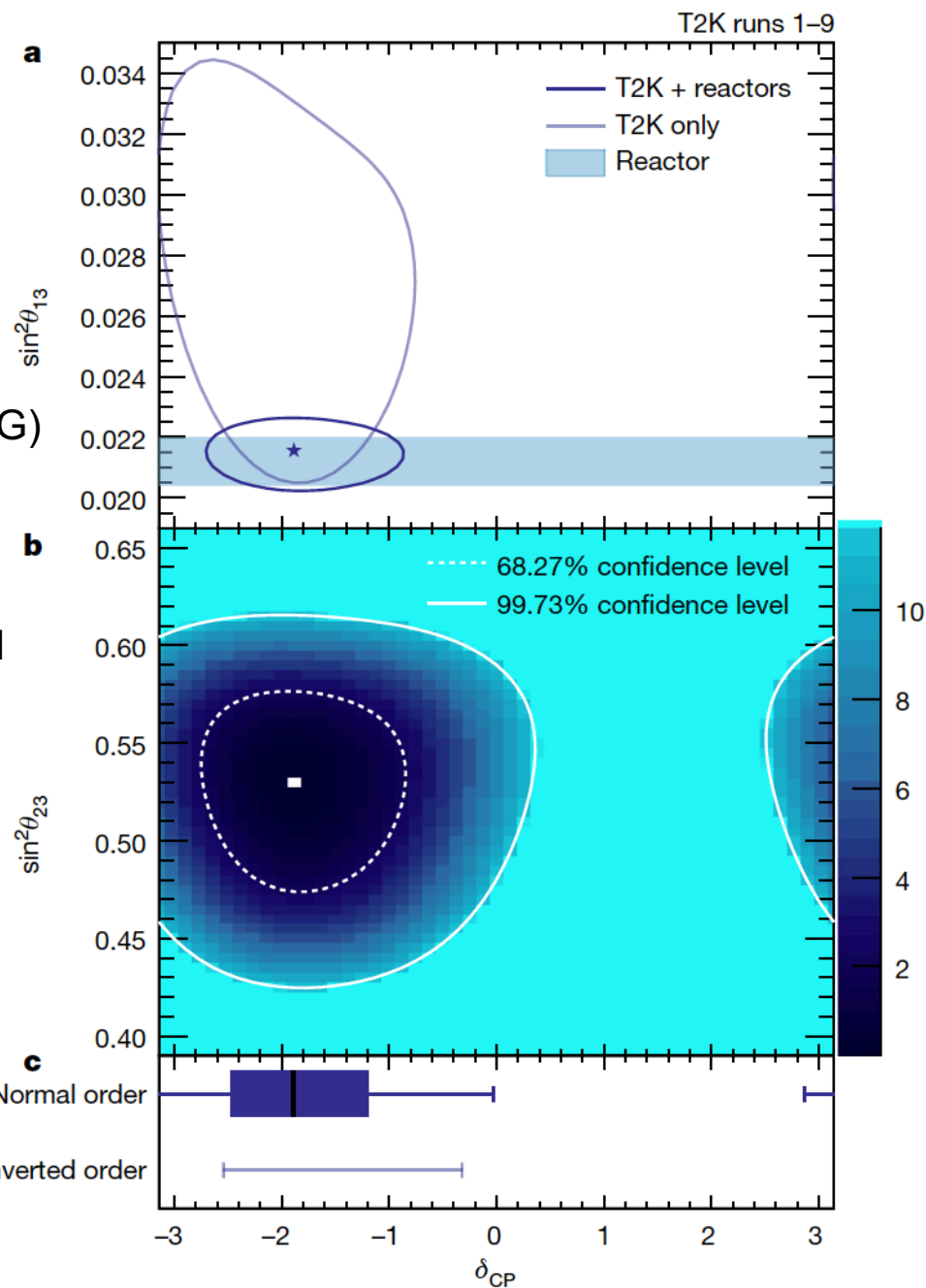
5. T2K oscillation results

All oscillation parameters are fit by assuming normal or inverted mass ordering.

- δ_{CP} , $\sin^2\theta_{23}$, Δm^2_{32} : flat prior
- $\sin^2\theta_{12}$, $\sin^2\theta_{13}$, Δm^2_{21} : external constraint (PDG)

Now the 3σ contour is closed, more data or new generation experiments can find the right value from here (Note, zero CP violation is not rejected with 3σ).

Normal ordering is favoured with 89% posterior probability.



1. Neutrino oscillation experiments

2. Neutrino beam

3. Neutrino detector

4. Neutrino interaction physics

5. Oscillation result

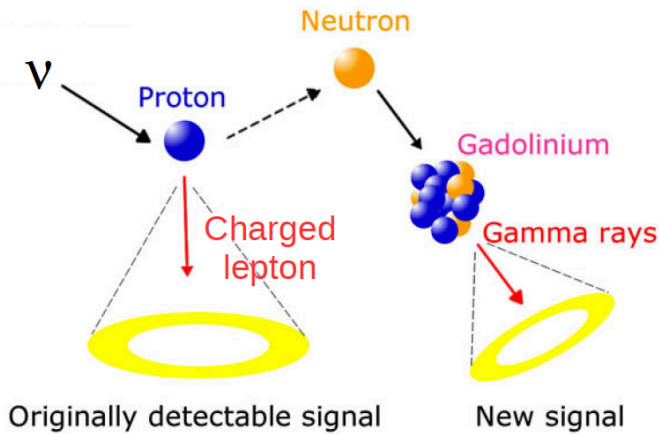
6. SK-Gd, T2K-Upgrade, Hyper-Kamiokande

7. Conclusion

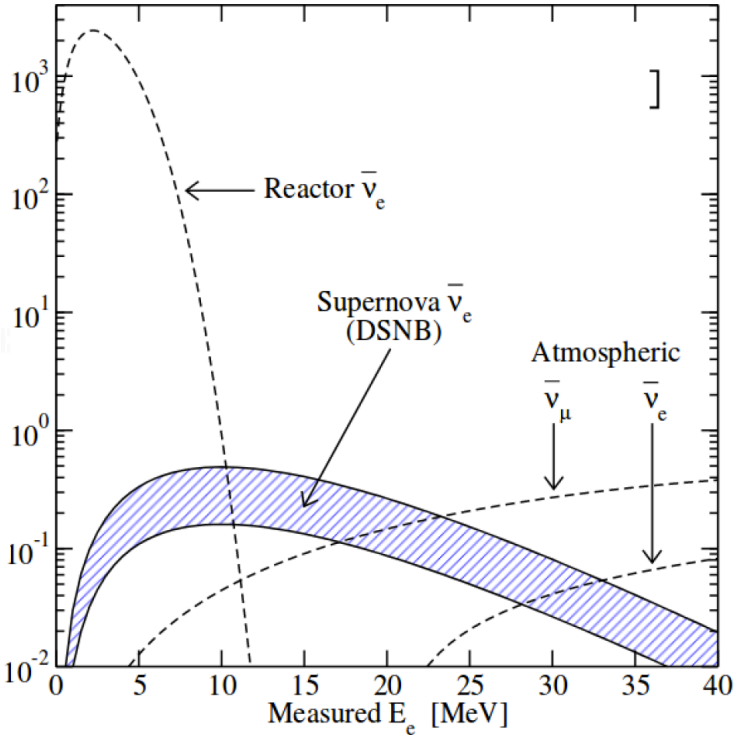
6. SK-Gd

SuperK is planned to be doped with 0.1% of Gd

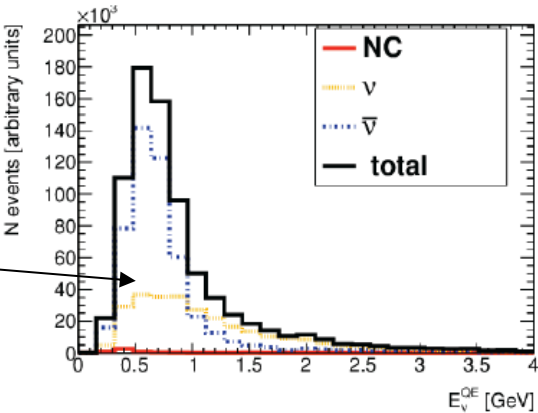
This improves neutron tagging efficiency to be ~90%, making SK-Gd to be visible for DSNB (diffused supernova neutrino background).



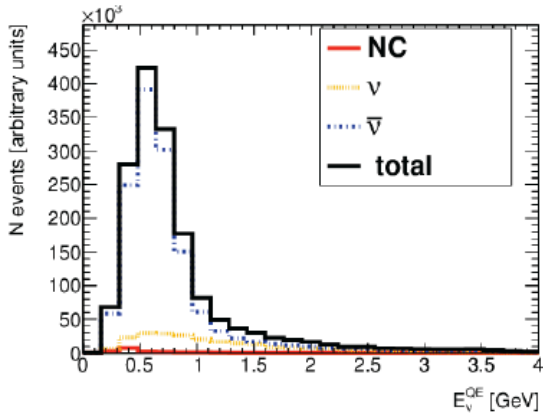
Antineutrino beam has large background of neutrinos (wrong-sign background)



Neutron tagging is also useful to suppress background for antineutrino oscillation measurements.

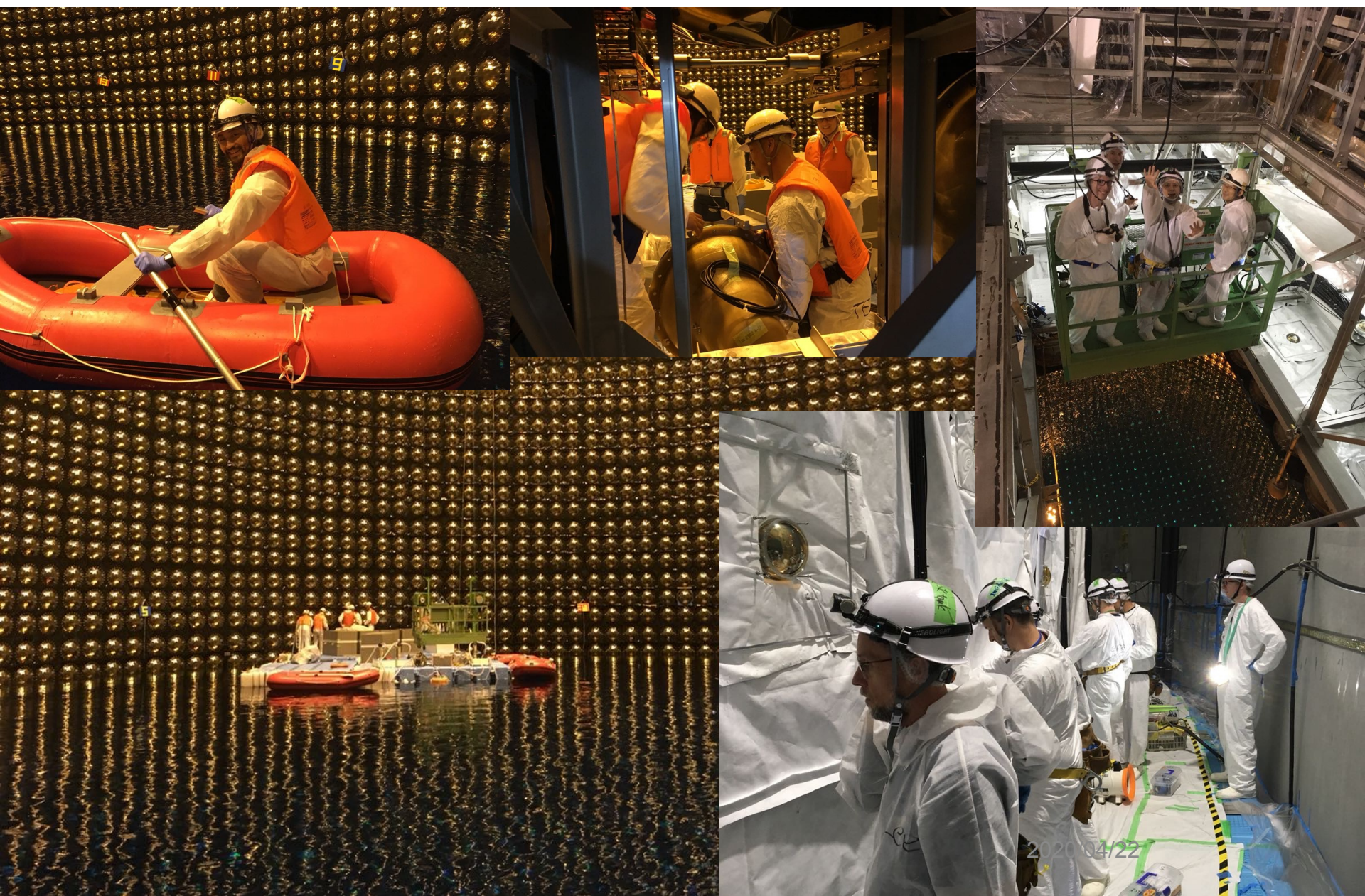


(b) $N_{\text{neutrons}} = 0$



(c) $N_{\text{neutrons}} \geq 1$

Super-Kamiokande detector refurbishment 2018



6. T2K upgrade

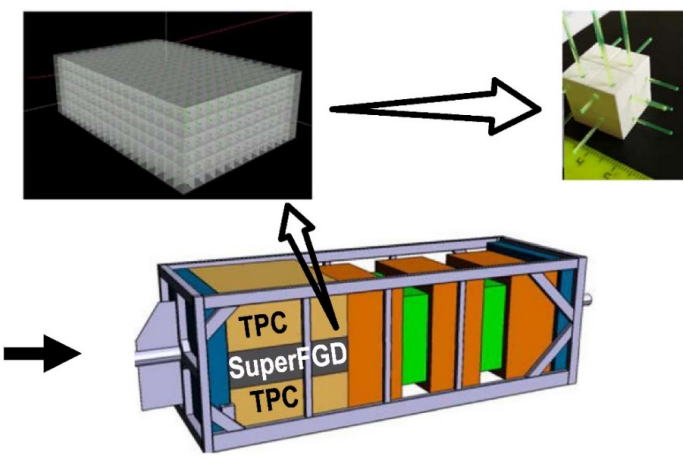
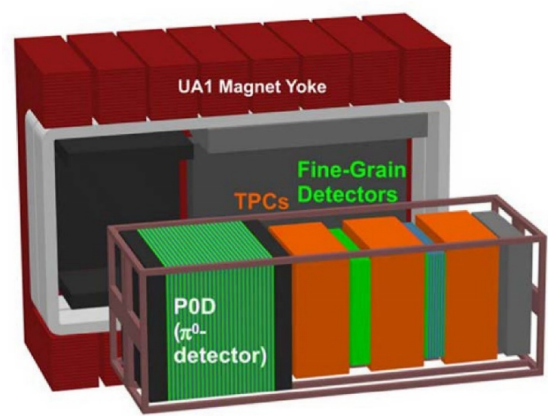
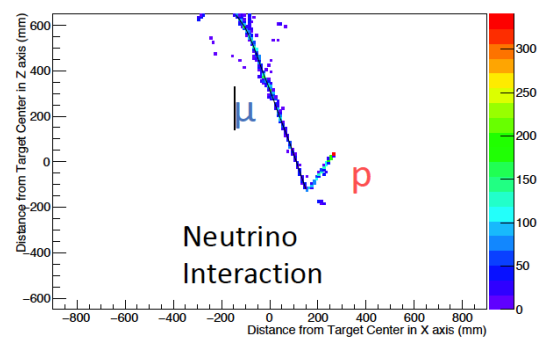
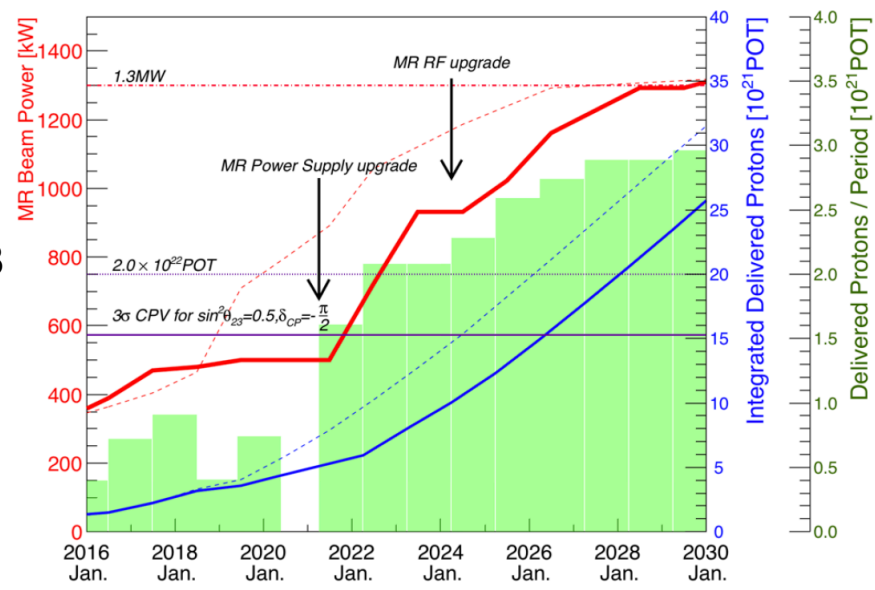
J-PARC neutrino beam upgrade

- Current beam power is ~ 500kW
- Beam power will be increased to 1.3MW in ~2028
- 2.48 sec period → 1.16 sec period
- 2.6E14 ppp → 3.2E14 ppp

ND280 detector upgrade

- POD is replaced with new detectors
- SuperFGD: 3-d scintillator tracker
- High pressure gas TPC: high-angle event measurement

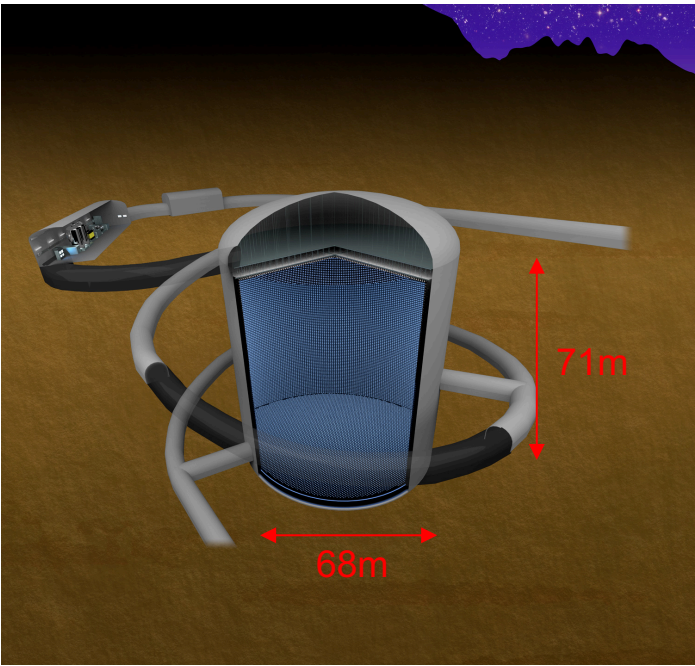
T2K-II Target POT (Protons-On-Target)



6. Hyper-Kamiokande

260 kton water Cherenkov tank

- ~x8.4 fiducial volume of SuperK
- Construction starts in this year!
- MeV to TeV physics
- solar, atmospheric, beam neutrinos
- proton decay, new physics search



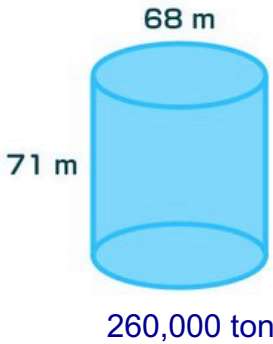
HiggsTan
<https://higgstan.com/>



Kamiokande
1983 start



Super-Kamiokande
1996 start



Hyper-Kamiokande
2020 construction start
2027 data taking (plan)

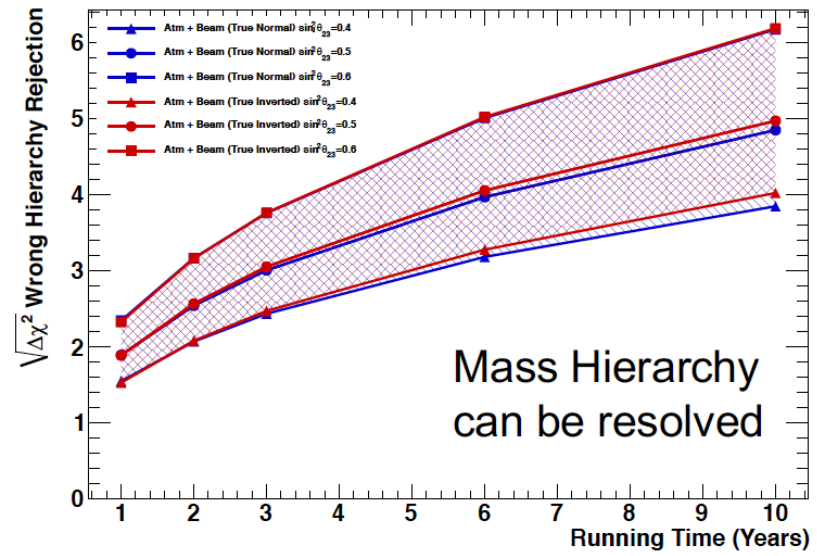
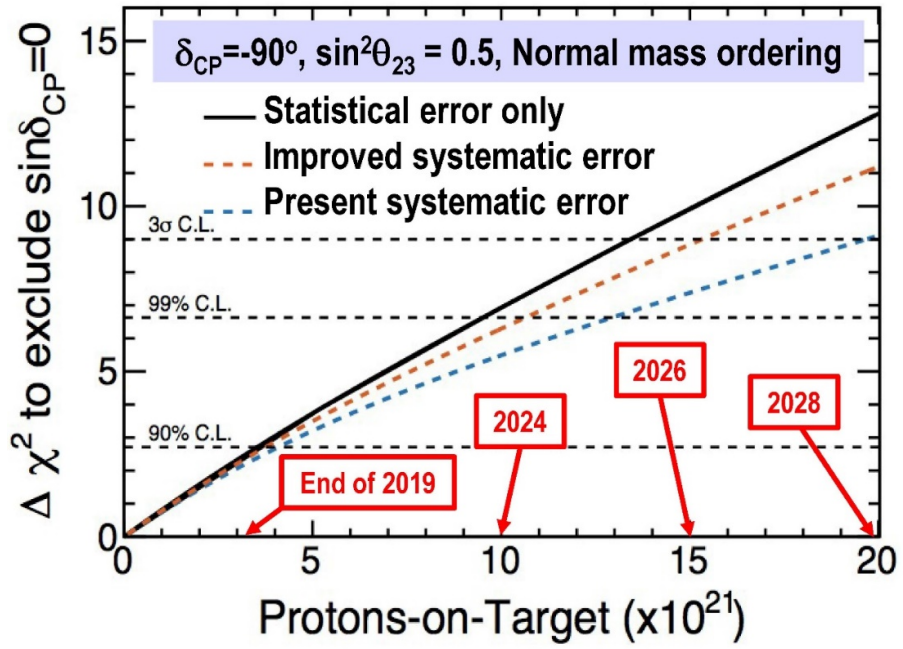
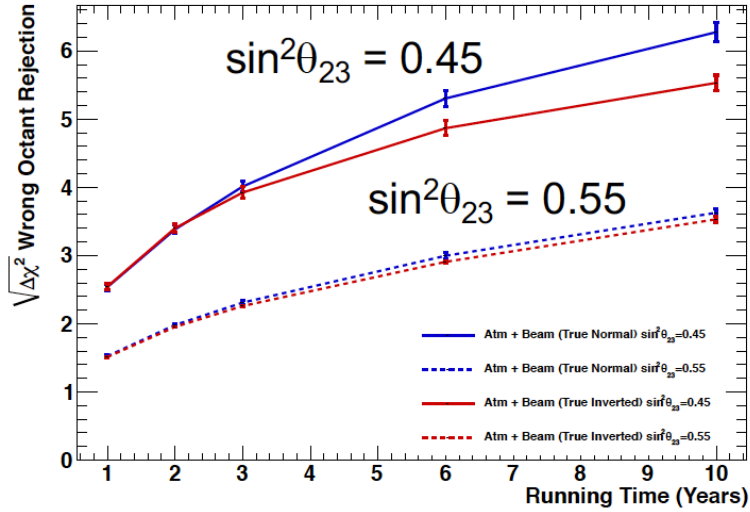


6. Hyper-Kamiokande

Oscillation program

- 5σ δ_{CP} measurement depends on systematic error reduction
- mass ordering is determined through atmospheric neutrinos
- θ_{23} will be measured

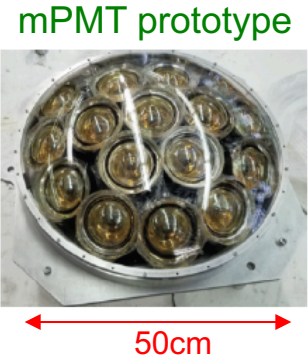
HyperK will find all oscillation parameters!



6. Hyper-Kamiokande

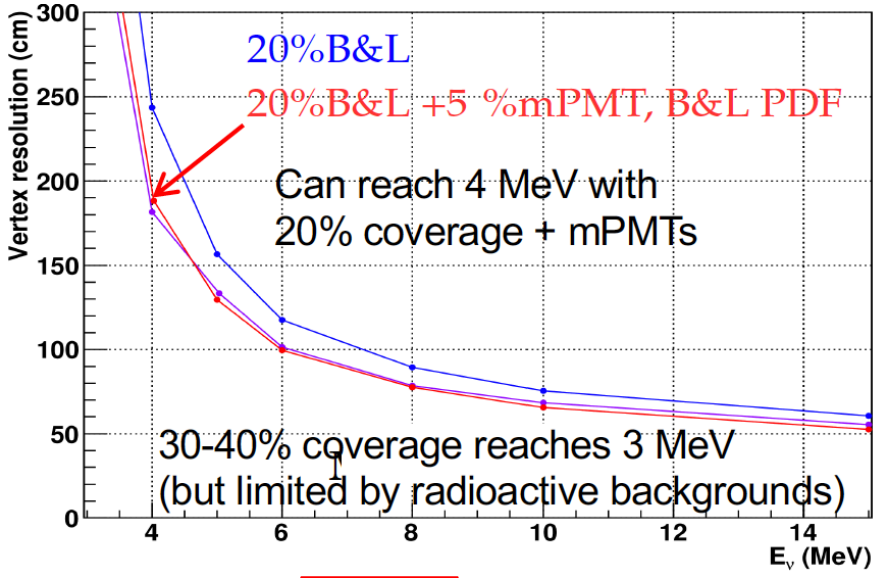
Solar neutrino

- Need low PMT dark rate
- mPMT helps a lot



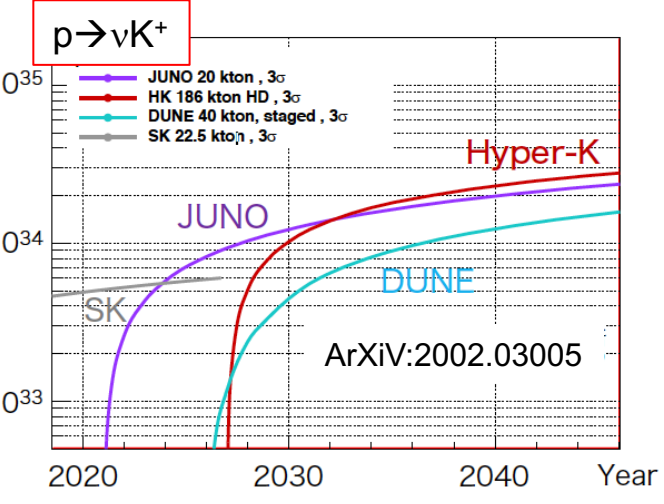
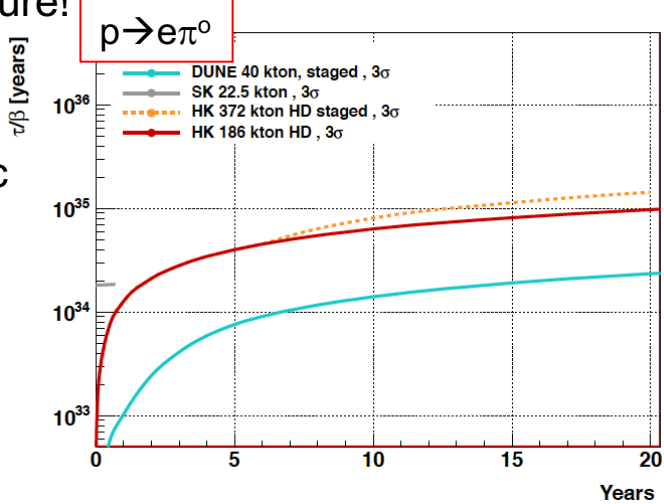
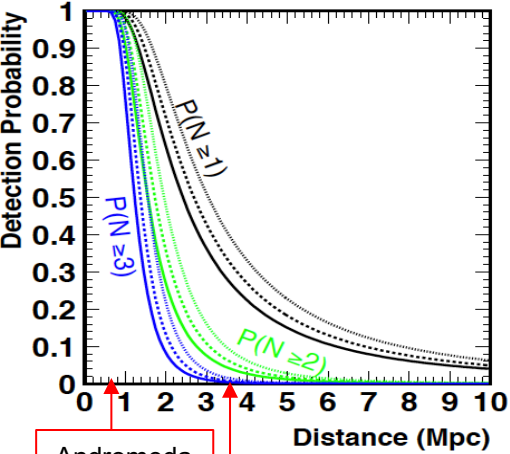
Proton decay

- 3σ sensitivity in 20yrs
- $p \rightarrow e\pi^0$, $\tau \sim 1E35yr$
- $p \rightarrow \nu K^+$, $\tau \sim 2E34yr$
- Many channels to measure!



Supernova neutrinos

- Search extends to ~Mpc



Future of T2K, SK-Gd, T2K-upgrade, and Hyper-Kamiokande are very bright!



Conclusion

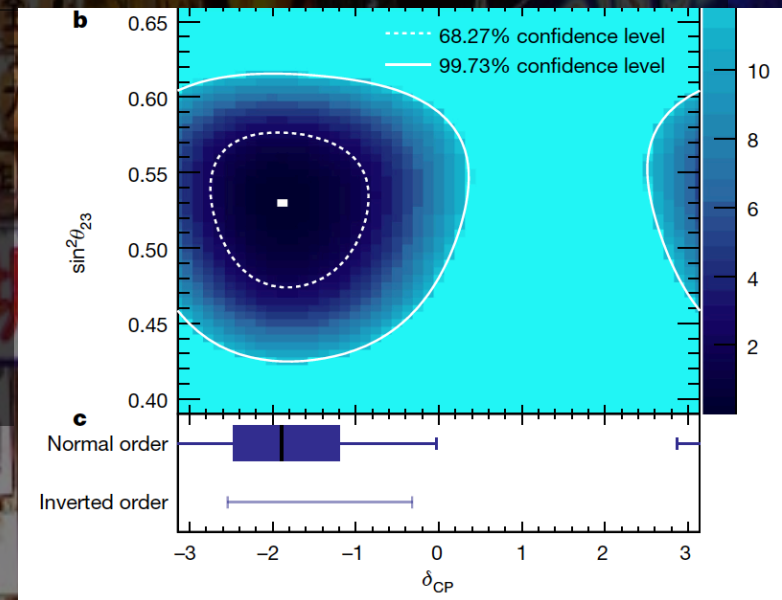
T2K is the second generation long-baseline neutrino oscillation experiment in Japan

Neutrinos from the J-PARC neutrino beam are measured by the Super-Kamiokande detector

2009-2018 data shows asymmetric oscillations, and neutrino oscillation is enhanced, and antineutrino oscillation is suppressed. This can be interpreted as negative CP violation phase.

$\delta\text{CP}=0$ is rejected more than 2σ , and 3σ interval is $[-3.41, -0.03]$ (normal ordering), and $[-2.54, -0.32]$ (inverted ordering)

The future plans include the SK-Gd, T2K-Upgrade, and Hyper-Kamiokande, and all of them are ongoing projects!



The T2K collaboration



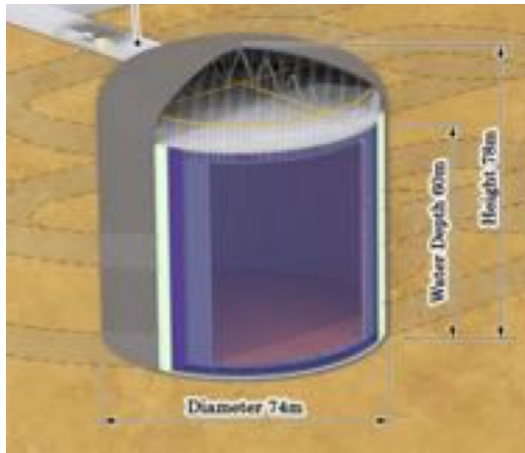
Thank you for your attention!

Backup

1. Next goal of high energy physics

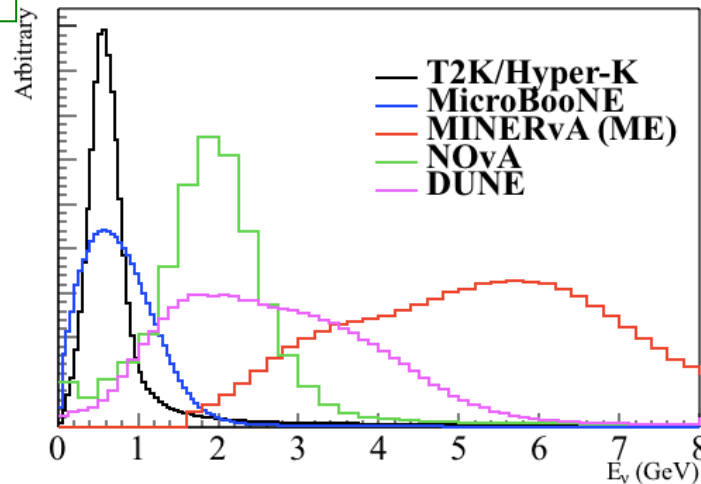
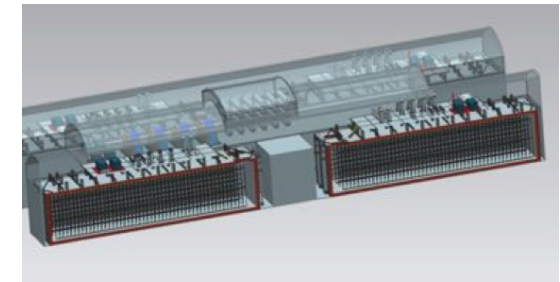
Hyper-Kamiokande (Japan)

- Water target
- Narrow band 0.6 GeV
- Low spatial resolution
- High time resolution



DUNE (USA)

- Argon target
- wide band 1-4 GeV
- High spatial resolution
- Low time resolution



Low energy beam (~1 GeV)

- shorter baseline (lower flux reduction)
- lower neutrino production
- lower interaction rate
- kinematic energy reconstruction

High energy beam (~few GeV)

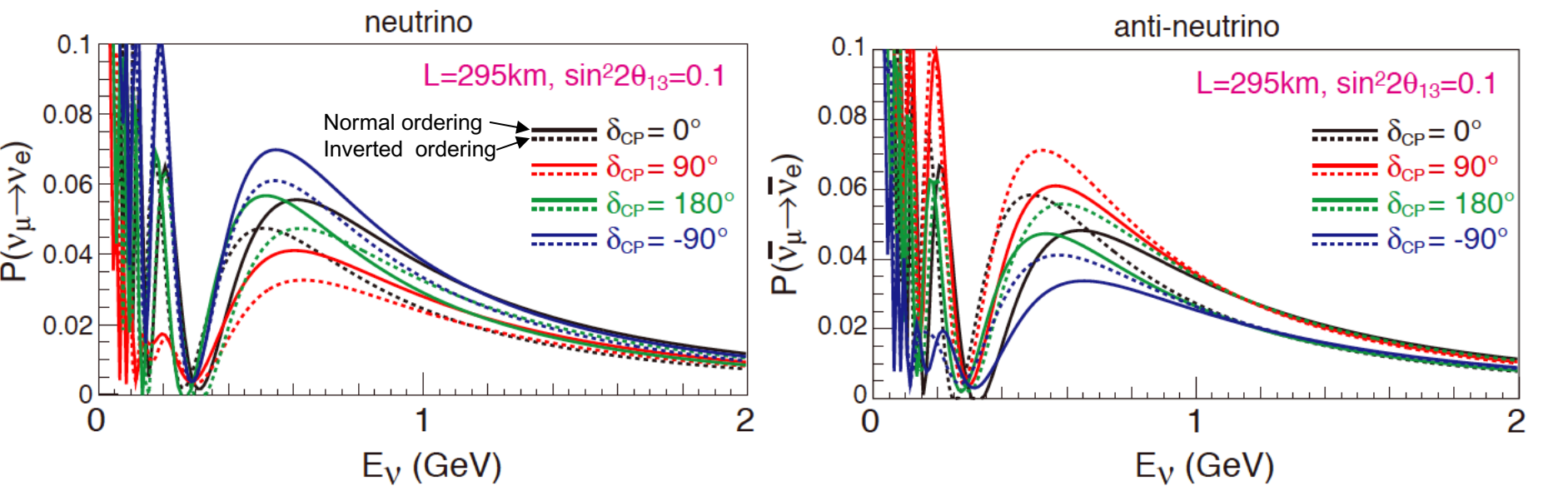
- longer baseline (higher flux reduction)
- higher neutrino production
- higher interaction
- calorimetric energy reconstruction

$$P_{\mu \rightarrow e}(L/E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L(km)}{E(GeV)} \right)$$

1. Neutrino oscillations

$$\begin{aligned} P(\nu_\mu \rightarrow \nu_e) = & 4c_{13}^2 s_{13}^2 s_{23}^2 \cdot \sin^2 \Delta_{31} \\ & + 8c_{13}^2 s_{12} s_{13} s_{23} (c_{12} c_{23} \cos \delta_{CP} - s_{12} s_{13} s_{23}) \cdot \cos \Delta_{32} \cdot \sin \Delta_{31} \cdot \sin \Delta_{21} \\ & - 8c_{13}^2 c_{12} c_{23} s_{12} s_{13} s_{23} \sin \delta_{CP} \cdot \sin \Delta_{32} \cdot \sin \Delta_{31} \cdot \sin \Delta_{21} \\ & + 4s_{12}^2 c_{13}^2 (c_{12}^2 c_{23}^2 + s_{12}^2 s_{23}^2 s_{13}^2 - 2c_{12} c_{23} s_{12} s_{23} s_{13} \cos \delta_{CP}) \cdot \sin^2 \Delta_{21} \\ & - 8c_{13}^2 s_{13}^2 s_{23}^2 \cdot \frac{aL}{4E_\nu} (1 - 2s_{13}^2) \cdot \cos \Delta_{32} \cdot \sin \Delta_{31} \\ & + 8c_{13}^2 s_{13}^2 s_{23}^2 \frac{a}{\Delta m_{31}^2} (1 - 2s_{13}^2) \cdot \sin^2 \Delta_{31}, \end{aligned}$$

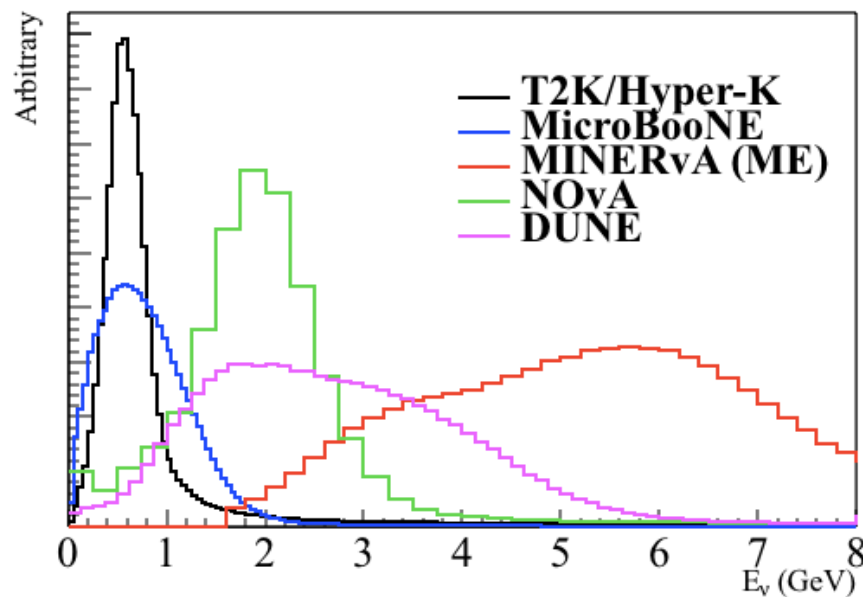
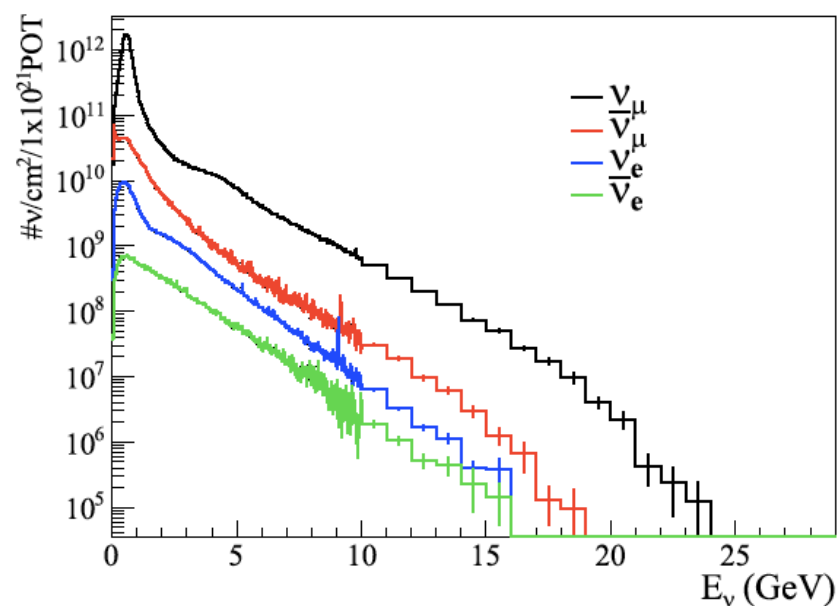
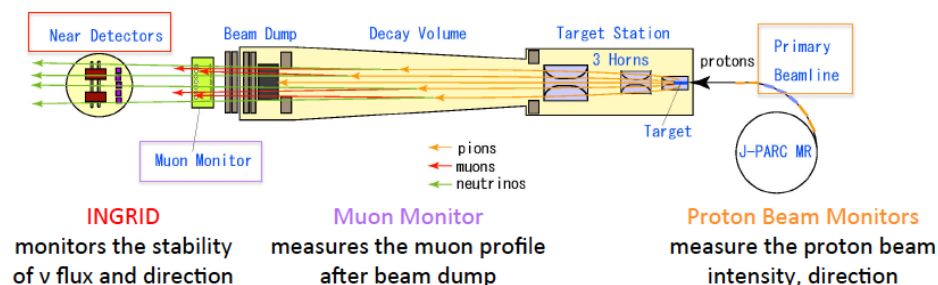
Normal ordering has a higher oscillation probability than inverted ordering for neutrino oscillation



2. Typical neutrino beams for oscillation experiments

e.g.) J-PARC neutrino beam (T2K)

- pion decay-in-flight (high flux)
- off-axis beam (narrow band)
- but has components up to ~ 10 GeV
- typical beam 1-10 GeV
- $\sim 5\%$ normalization error (best case)



2. LINAC → RCS → MR → extraction

LINAC

- 1st stage
- 400 MeV

RCS (Rapid Cycling Synchrotron)

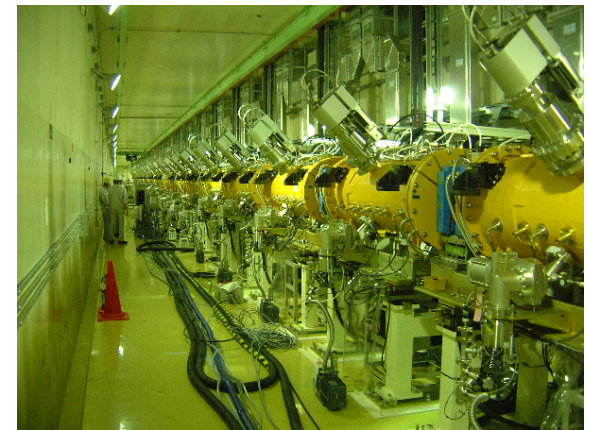
- 2nd stage
- 3 GeV, 25 Hz
- 4 beam pulse are extracted with 40ms intervals

Main Ring

- 3rd stage
- 30 GeV
- 8 bunches in 1 spill, 2.48 sec interval

Extraction beam line

- Superconducting bending magnet



3. Off-axis detector

ND280

- POD: Water-scintillator tracker
- FGD: Fully active scintillator tracker
- TPC: Ar gas TPC
- ECal: Lead-scintillator calorimeter
- SMRD: Iron-scintillator tracker
- UA1 magnet

Near detector data

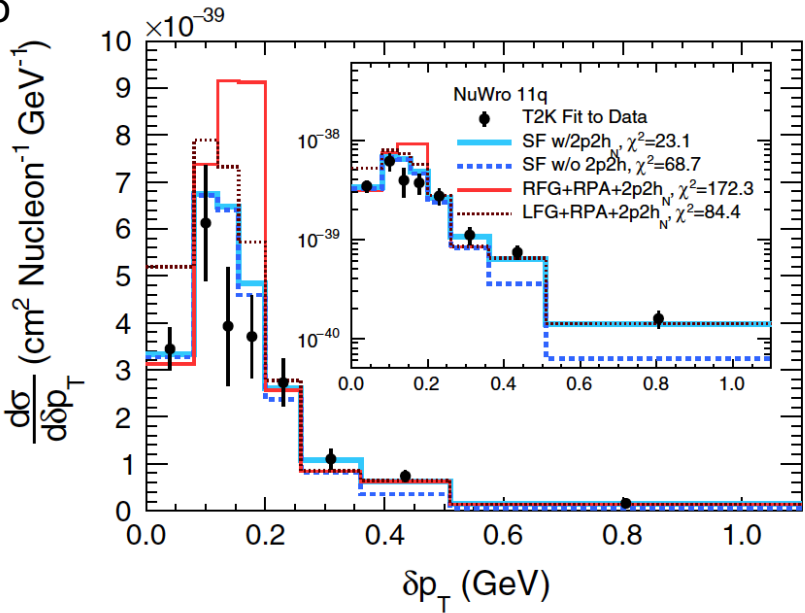
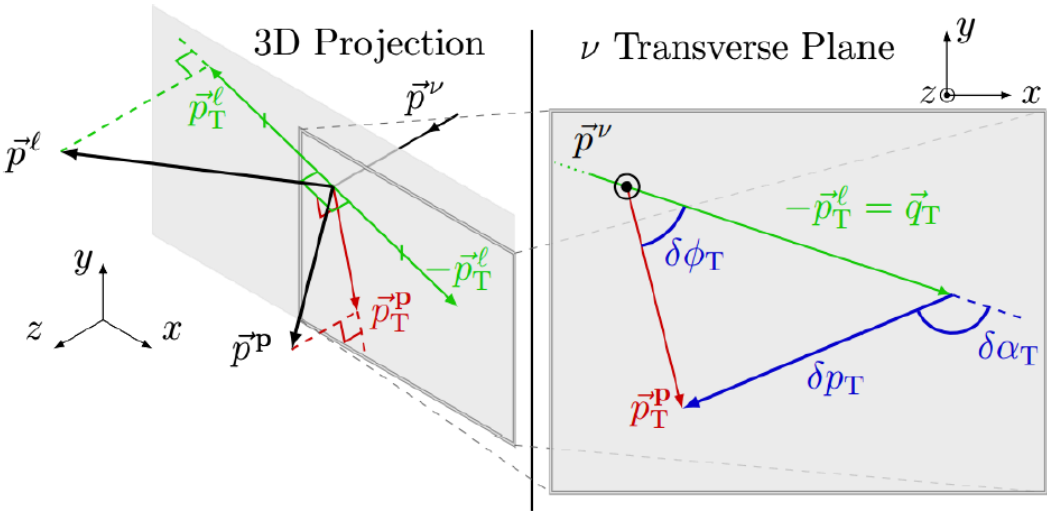
- 14 samples are used for the oscillation analysis to constrain flux and cross-section systematic errors

Neutrino cross section measurements

- many results are published
- neutrino interaction models are studied

ex) $1\mu+1p$ measurement

- Data prefer advanced nuclear models

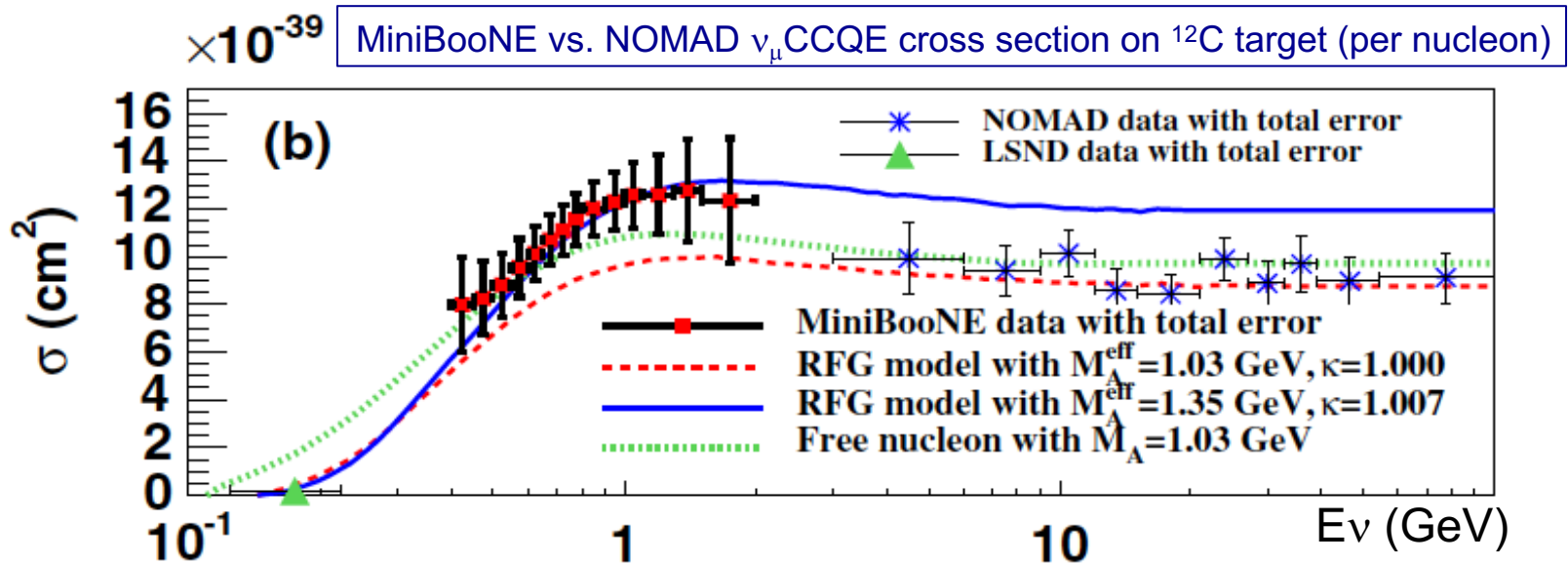


4. Charged Current Quasi-Elastic scattering (CCQE)

CCQE puzzle

1. low Q^2 suppression $\rightarrow$ Low forward efficiency? (detector?)
2. high Q^2 enhancement $\rightarrow$ Axial mass > 1.0 GeV? (physics?)
3. large normalization $\rightarrow$ Beam simulation is wrong? (flux?)

CCQE interaction on nuclear targets are precisely measured by electron scattering
 - Lepton universality = precise prediction for neutrino CCQE cross-section...?



4. Summary of CCQE for oscillation physics

Community is converged: the origin of CCQE puzzle is multi-nucleon correlation

This moment...

New nuclear models available in simulations do not **quantitatively** describe T2K and MINERvA CCQE $\mu+p$ data

large M_A error $\rightarrow$ large 2p2h error

It is crucial to have correct CCQE, 2p2h, pion production models to understand data simultaneously. Otherwise M_A error stays around 20-30%.

We have good theorists who make models, and good experimentalists who measure data, but we are still lacking people between them.

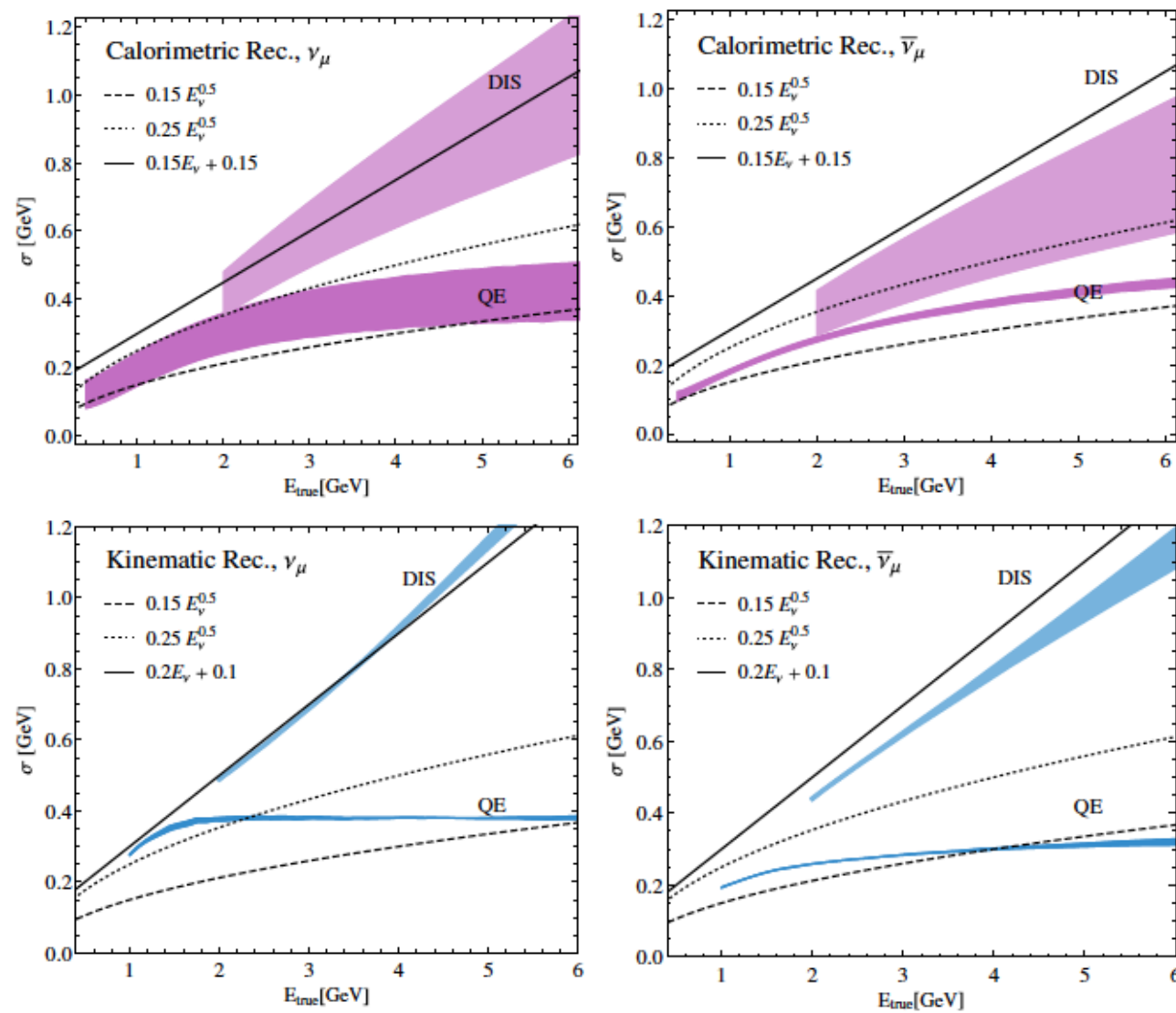


4. Kinematic E reconstruction vs calorimetric E reconstruction

Calorimetric energy reconstruction suffers invisible hadrons (=neutrons)

It largely depends on neutrino interaction and hadron simulation

- multiplicity
 - kinematics
 - nuclear effect
 - re-scattering
 - charge exchange
 - baryonic resonance
 - nucleon correlation
- etc



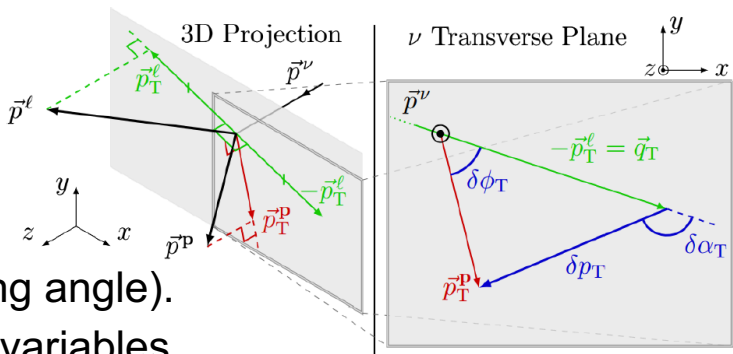
4. Neutrino hadron measurement for nuclear correlations

We want to constrain nuclear model from neutrino data

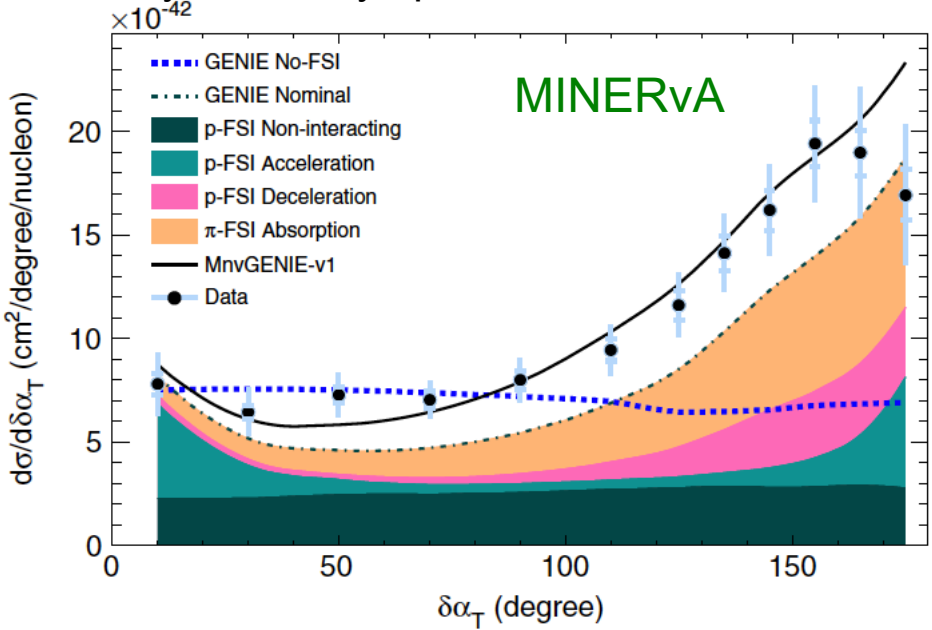
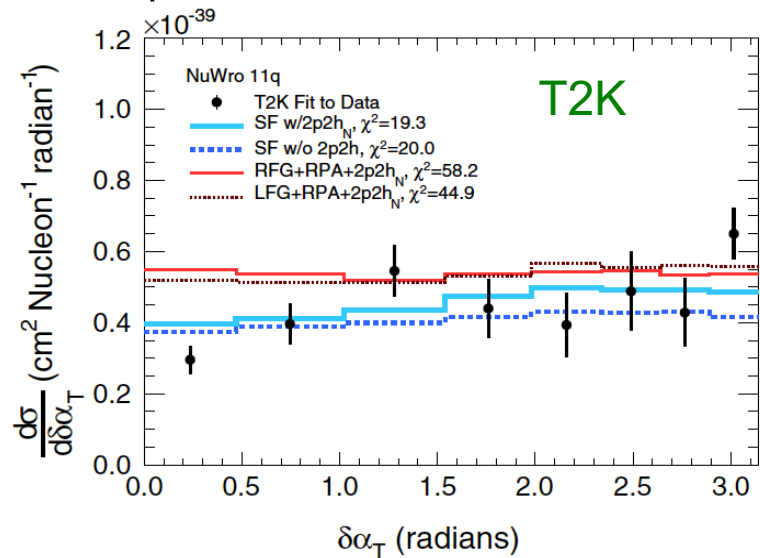
- Final state hadron measurement is the key

1 muon + 1 proton sample

- 5 dof (mu E and cosθ, proton E and cosθ, mu-p opening angle).
- Low statistics, and these are converted to 3 kinematic variables.

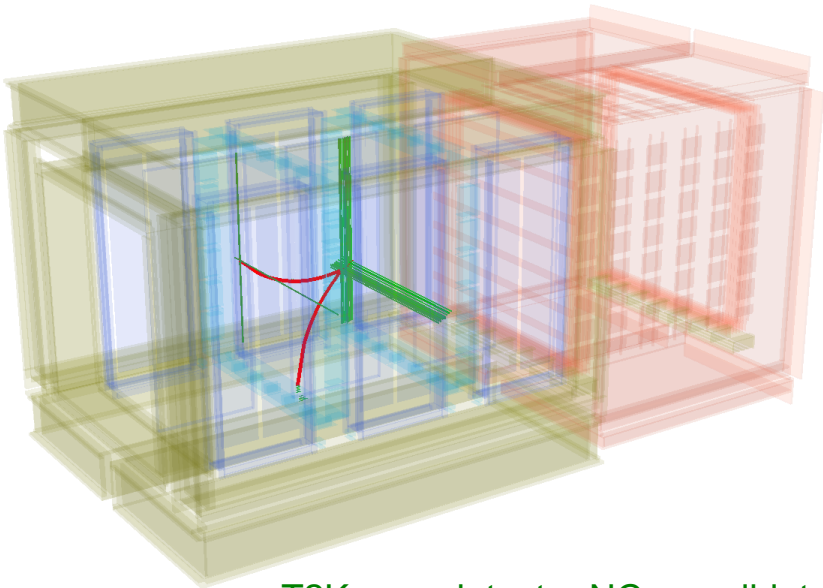
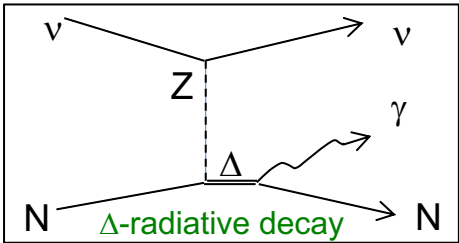


Data prefer advanced nuclear models, but it's not easy to identify 2p2h model

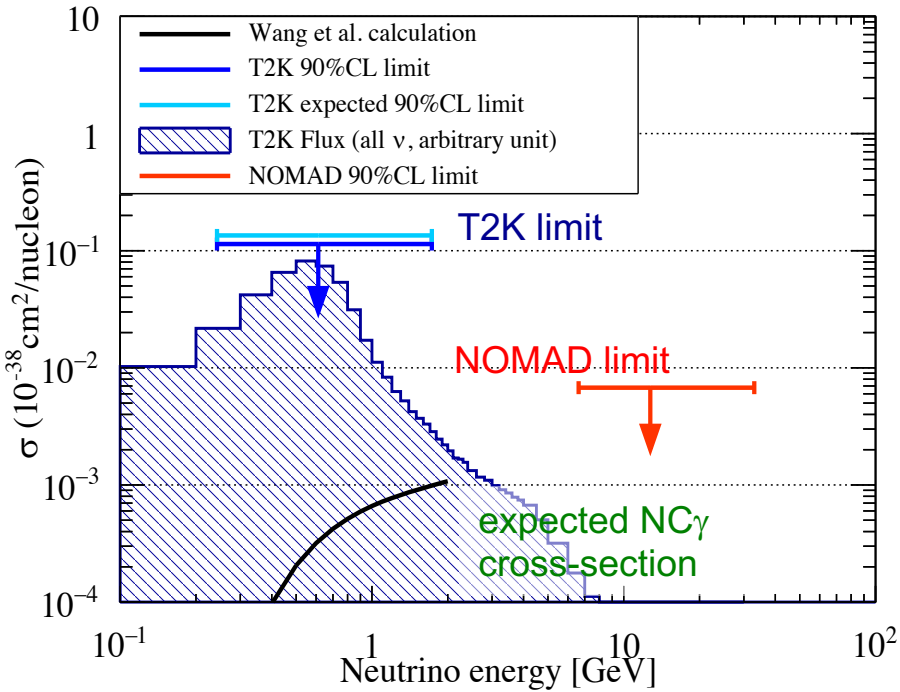


4. T2K Neutrino NC single photon production (NC_γ)

Neutrino induced NC single photon production (NC_γ) process is not experimentally identified. NC_γ is misID background for every electron-neutrino appearance oscillation experiment. T2K and NOMAD set limits on this process, but $\sim x3$ higher cross-section can explain all MiniBooNE excess.



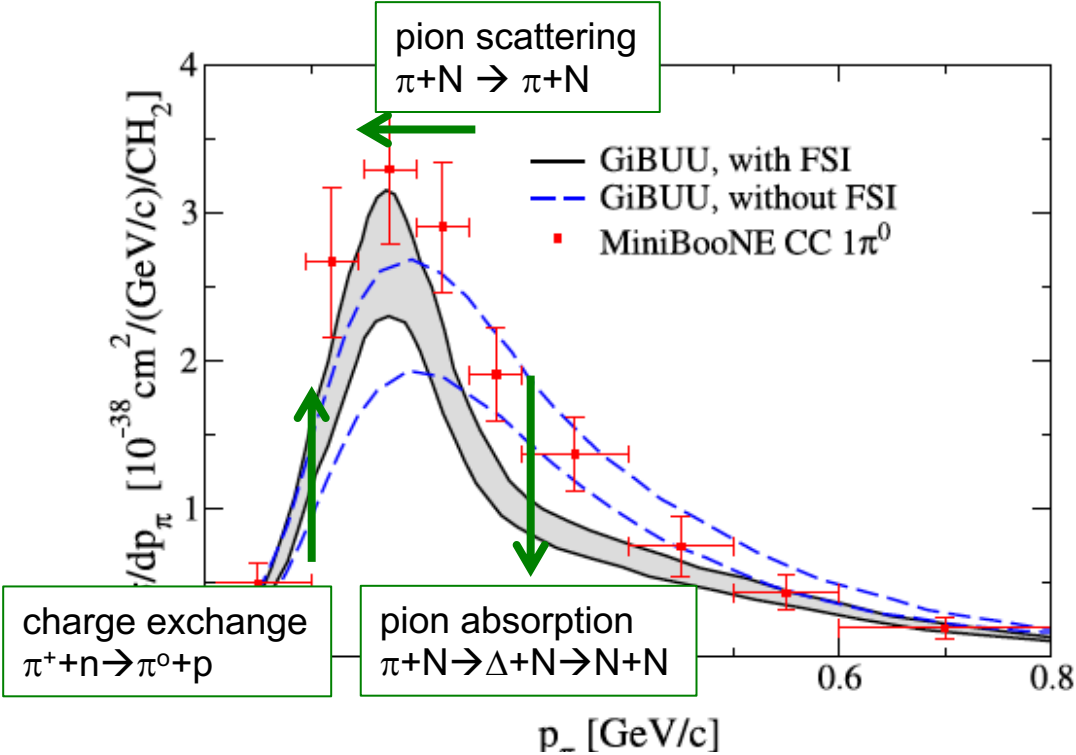
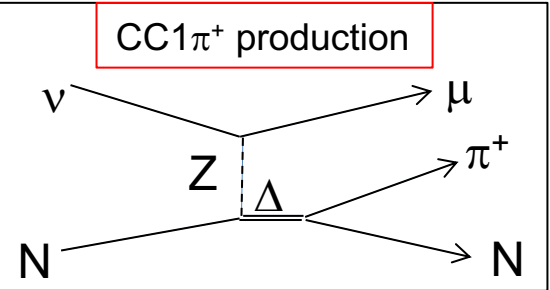
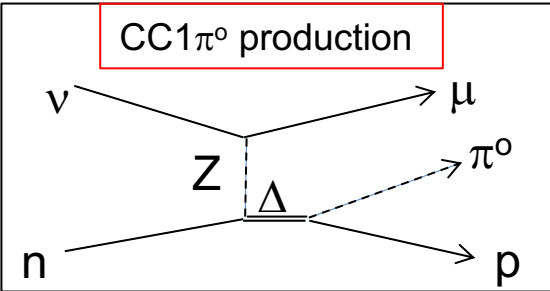
T2K near detector NC_γ candidate



4. Pion puzzle (2019)

Pion production + final state interaction (FSI)

- Neutrino induced pion productions have large errors
- Final state interaction of hadrons have large errors



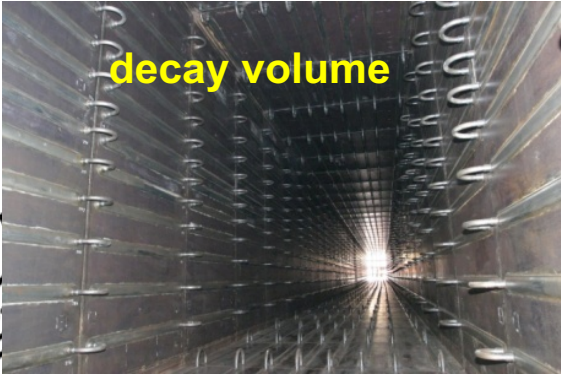
ex) Giessen BUU transport model

- Developed for heavy ion collision, and now used to calculate final state interactions of pions in nuclear media

Neutrino beam line and components



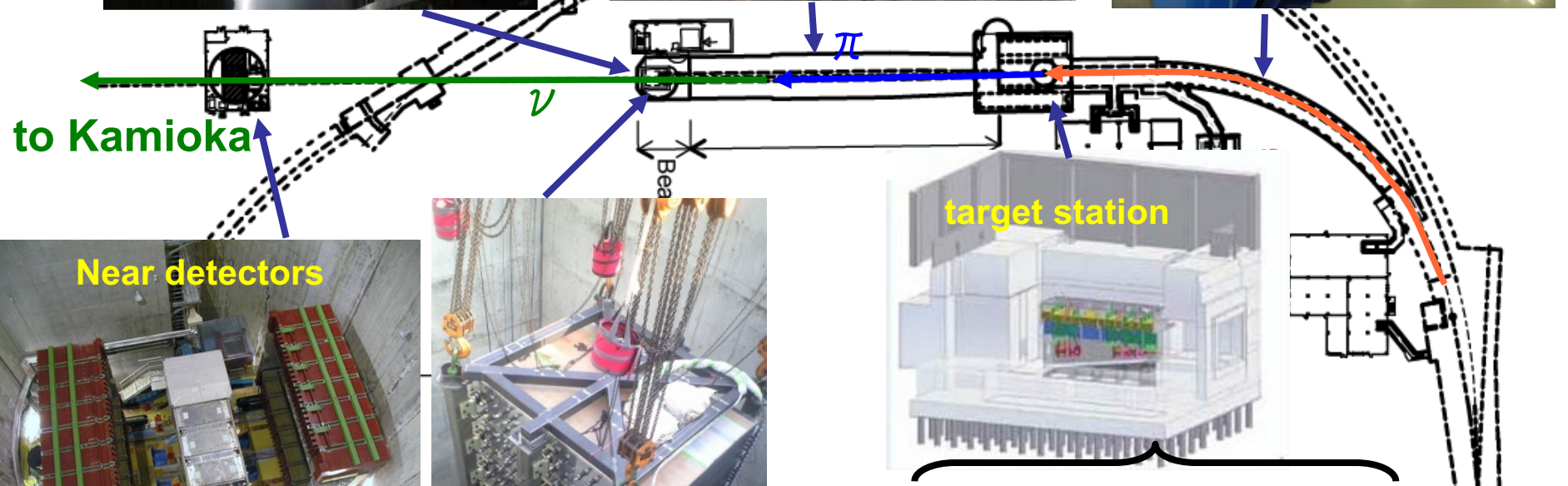
muon monitors



decay volume



superconducting magnet



to Kamioka

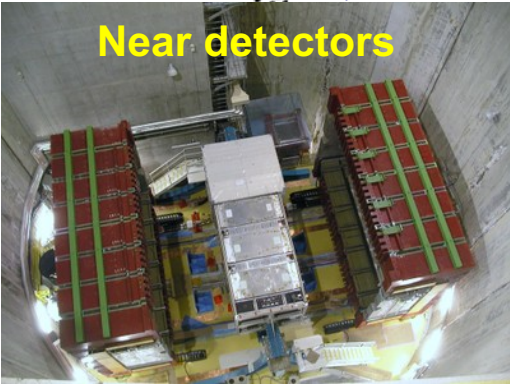
ν

π

Bea

target station

The diagram illustrates the neutrino beam line from the target station to the detectors. A red line represents the beam path, which is deflected by a superconducting magnet. A blue line represents the decay volume where pions (π) decay into neutrinos (ν). The beam then travels through a beam dump and is directed towards the detectors. The distance from the target station to the detectors is labeled 'Bea'.



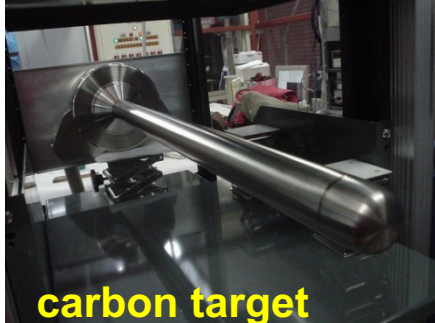
Near detectors



beam dump



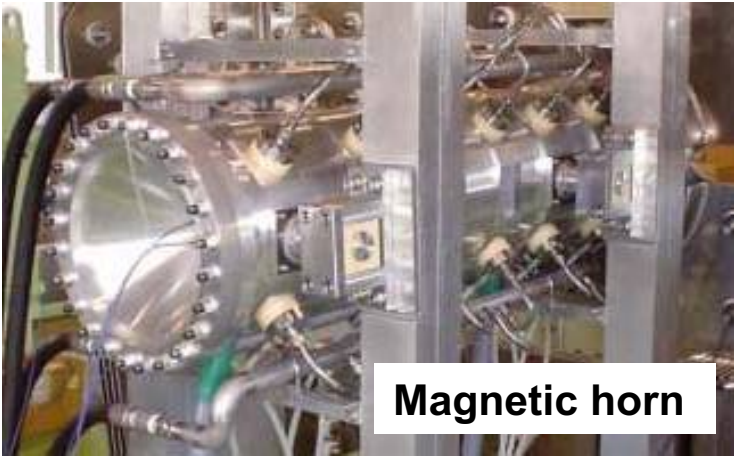
magnetic horn



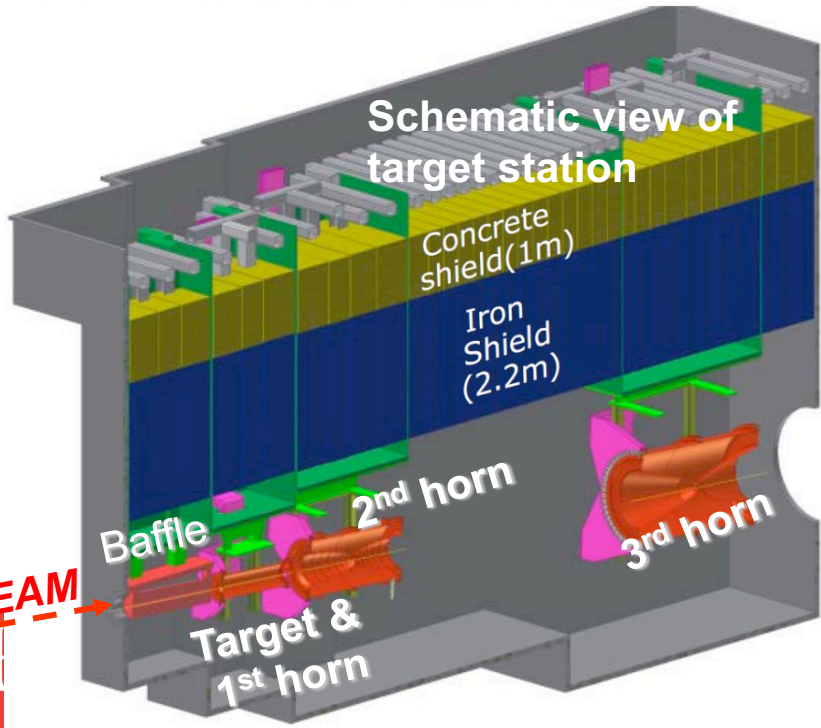
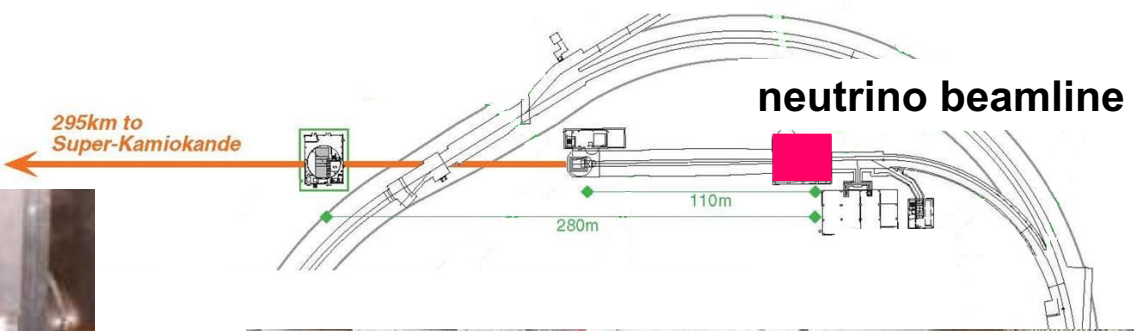
carbon target



Target Station



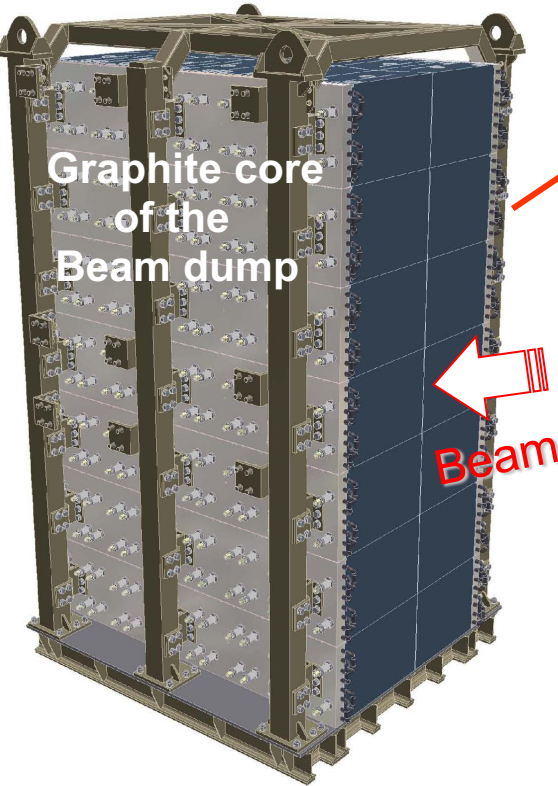
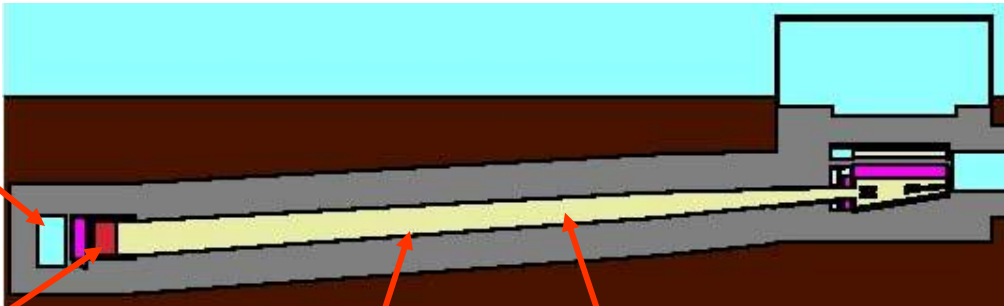
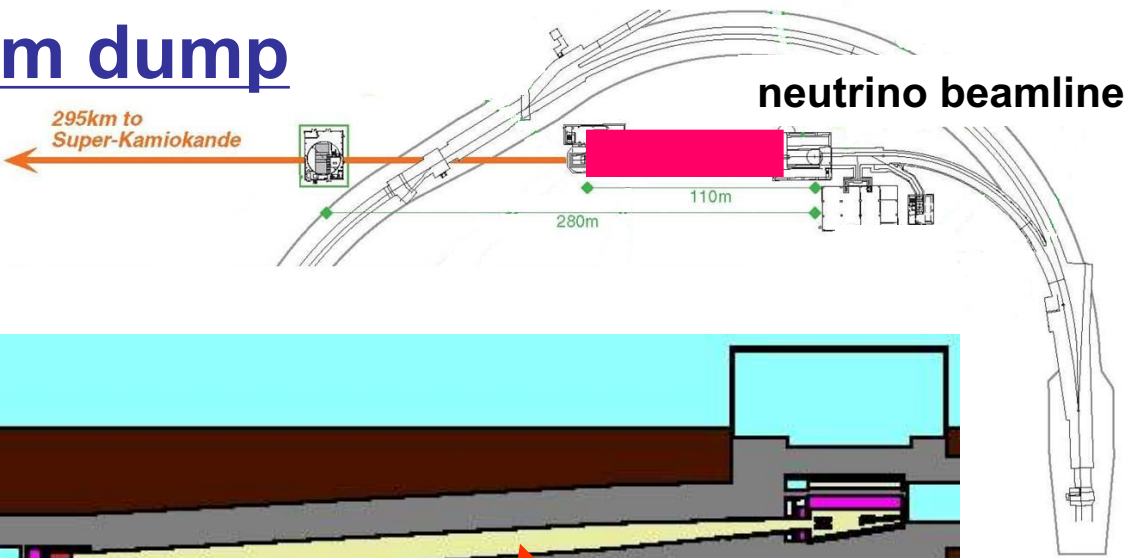
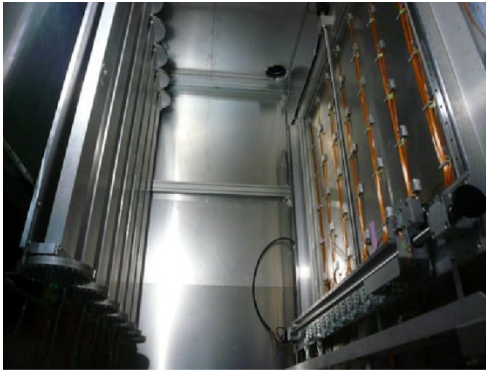
Magnetic horn



Civil construction of target station

Decay pipe and Beam dump

Muon monitors in Muon Pit

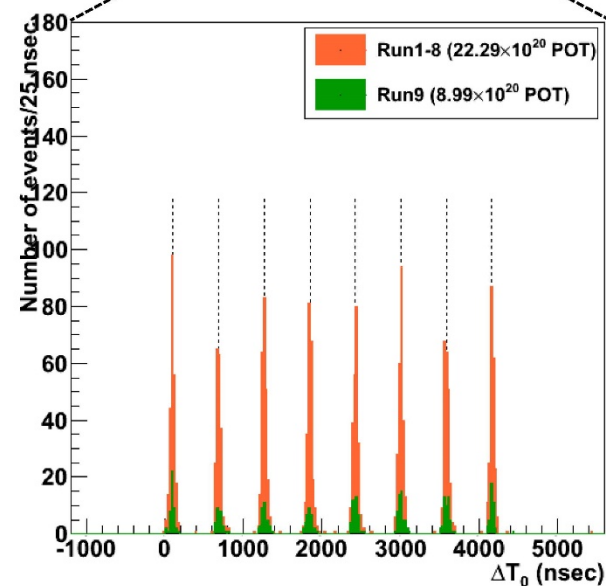
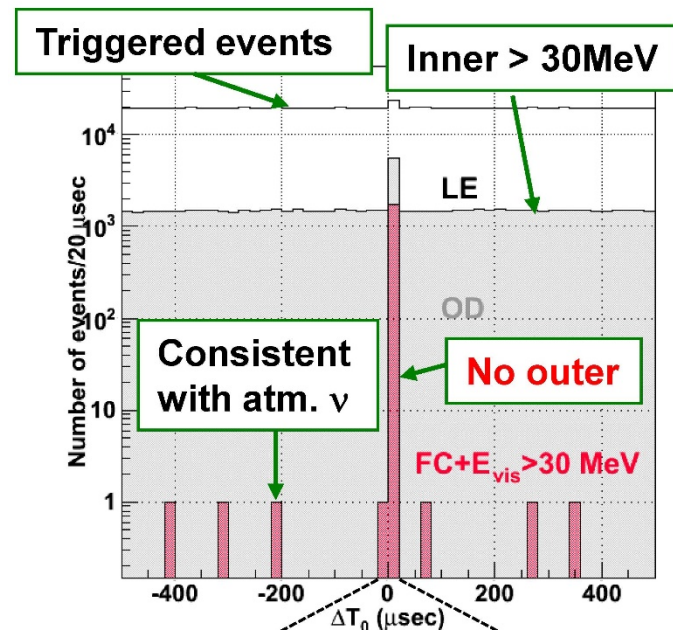


Super-Kamiokande Event Selection

- Data recorded between Jan. 2010 and May 2018 are used. It corresponds to $1.49 \times 10^{21} (\nu) + 1.64 \times 10^{21} (\bar{\nu})$.

Event Selection Criteria

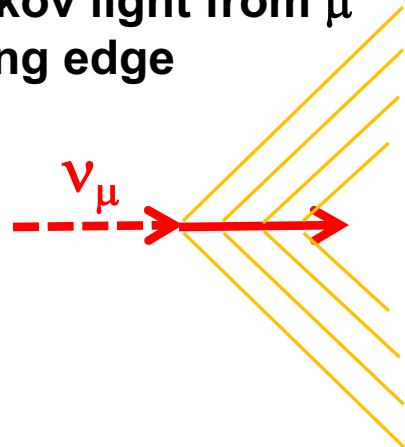
1. Total energy deposit in the inner detector is larger than **30 MeV** equivalent.
2. No outer detector activity
3. The event time agrees with **$\sim 5 \mu\text{sec}$** beam period in 2.48 sec accelerator cycle. (8 bunch structure can be found.)
4. 1 Ring events
→ μ/e particle identification is applied



μ/e identification in Super-Kamiokande

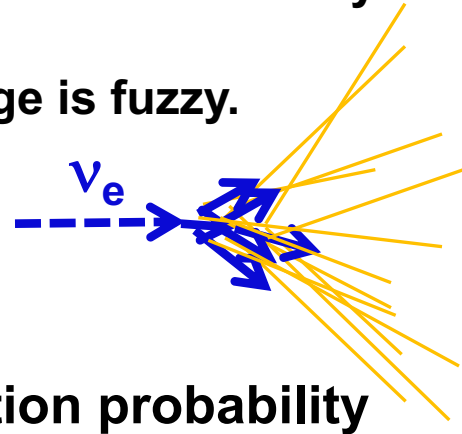
$\nu_\mu \rightarrow \mu$

Only direct Cherenkov light from μ
Clear Cherenkov ring edge

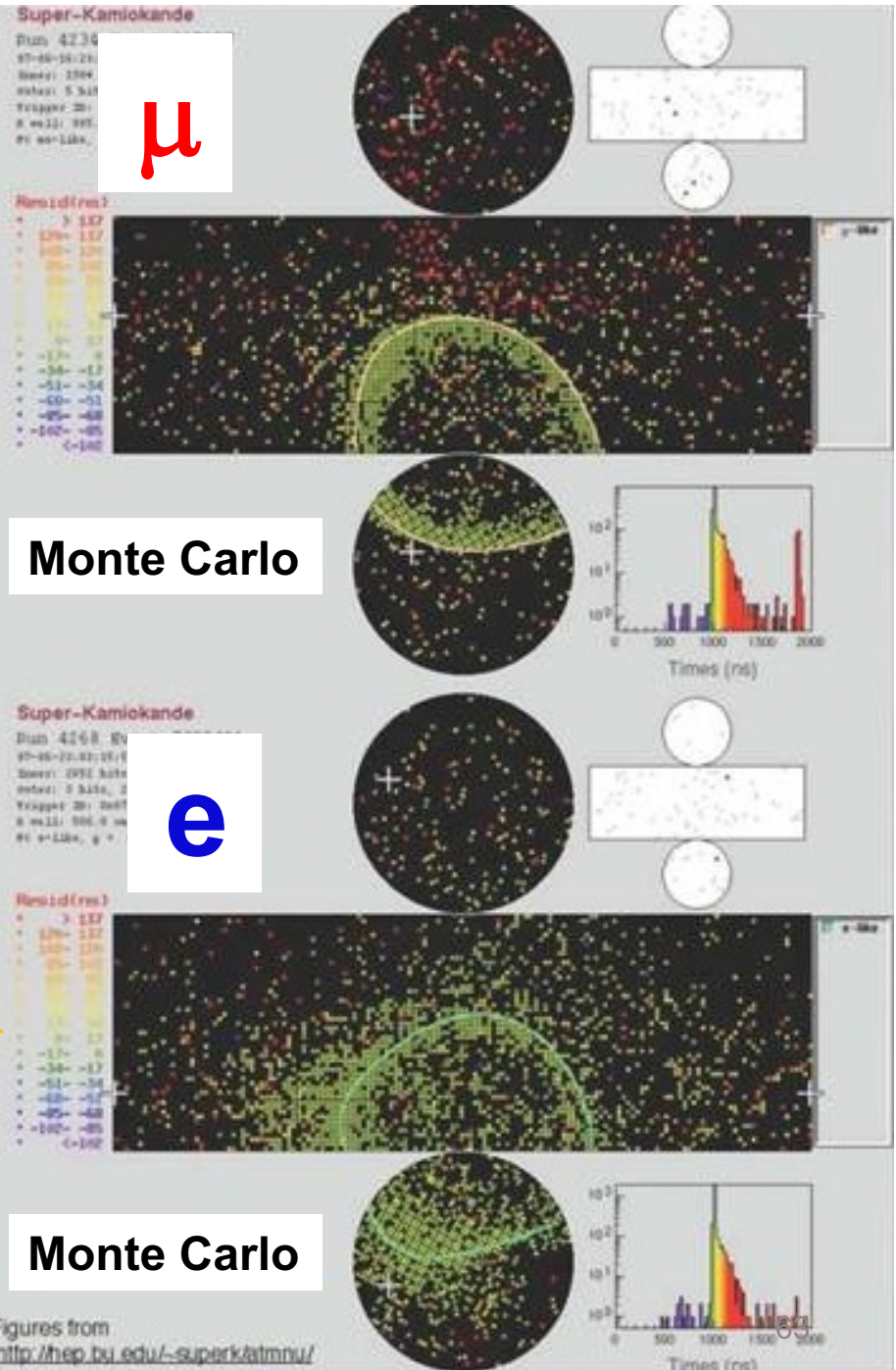


$\nu_e \rightarrow e$

Cherenkov light from e-m shower.
Electrons and positrons are heavily scattered.
Cherenkov ring edge is fuzzy.



• μ/e misidentification probability is less than **1 %**.



Event Selection

Examine Particle ID of 1 ring events

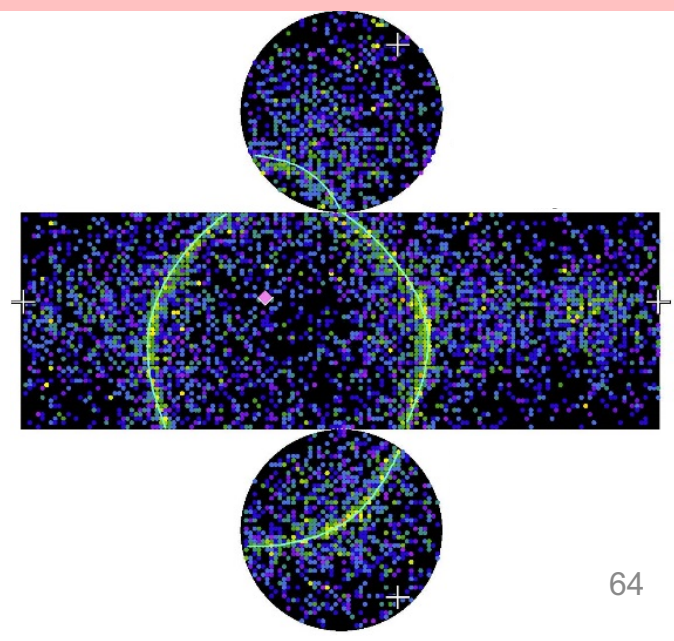
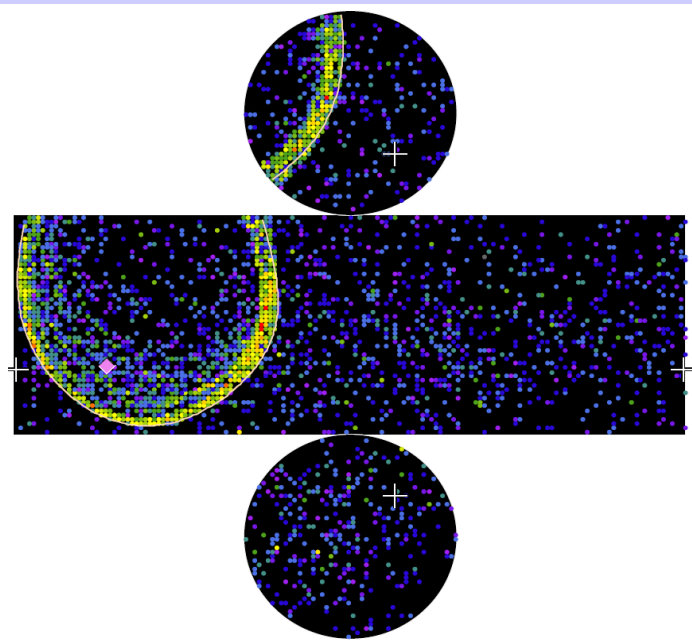
ν_μ selection

- μ -like PID
- $p_\mu > 200 \text{ MeV/c}$
- Michel electron 1 or 0

ν_e selection

- e-like PID
- $p_e > 100 \text{ MeV/c}$
- $E_{\text{rec}} < 1250 \text{ MeV}$
- π^0 rejection

π^0 rejection :
Forced 2nd ring is
assumed. Invariant
mass and
likelihood for π^0
are examined.



Results of ν_μ and $\bar{\nu}_\mu$ disappearance

- **Disappearance** of muon neutrino events as well as a **distortion** of neutrino energy spectrum is obvious for both ν_μ and $\bar{\nu}_\mu$.
- Oscillation parameters for anti-neutrinos well agree with the parameters for neutrinos within statistical errors.
- These results are updates of past publications.

PRL 112, 181801(2014),
PRD 91, 072010(2015)
for 6.57×10^{20} POT, **120** ν_μ data

PRL 116, 181801(2016)
for 4.01×10^{20} POT, **34** $\bar{\nu}_\mu$ data

