Neutrino Interferometry for New Physics Search

TK, MPLA27(2012)1230024 IceCube, Nature Physics 14(2018)961 IceCube, To be published (2019)

Motivation

- String theory
- Loop quantum gravity
- Horava-Lifshitz gravity
- Lee-Wick theory
- Non-commutative field theory
- Supersymmetry
- etc

Physics

- Lorentz violating field
- Quantum foam
- Neutrino-dark matter coupling
- Neutrino-dark energy coupling
- Neutrino-torsion coupling
- Neutrino velocity \neq c
- Violation of equivalent principleetc

Find us on Facebook, "Institute of Physics Astroparticle Physics" https://www.facebook.com/IOPAPP

Teppei Katori Queen Mary University of London Quantum Sensors, King's College London, UK, March 6, 2019

Teppei Katori

18/12/13

IceCube, Nature Physics 14 (2018) 961 Mewes, Nature 560 (2018) 316 **Neutrino interferometry – Atmospheric neutrinos**



ueen Mary

University of London

Atmosphere Neutrino oscillation is a nature interferometer. Any extra interactions in the Lagrangian contribute the phase shift

For 20 TeV up-going atmospheric neutrinos (L~12700km), detectable phase shift by neutrino is $\overline{\psi}a^{\mu}\gamma_{\mu}\psi$, $a\sim 10^{-24} GeV$ If anomalous coupling with neutrinos in vacuum cause a phase shift in similar order, we can see it from spectrum distortion of atmospheric neutrinos



$$H \sim \frac{m^2}{2E} + \mathring{a}^{(3)} - E \cdot \mathring{c}^{(4)} + E^2 \cdot \mathring{a}^{(5)} - E^3 \cdot \mathring{c}^{(6)} \cdots$$

Neutrino interferometry – Atmospheric neutrinos

| dim. | method | type | sector | limits | ref. |
|------|-----------------------------------|------------------------------|----------|---|-----------|
| 3 | CMB polarization | astrophysical | photon | $\sim 10^{-43} \text{ GeV}$ | [6] |
| | He-Xe comagnetometer | tabletop | neutron | $\sim 10^{-34} { m GeV}$ | [10] |
| | torsion pendulum | tabletop | electron | $\sim 10^{-31}~{ m GeV}$ | [12] |
| | muon g-2 | accelerator | muon | $\sim 10^{-24} { m GeV}$ | [13] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $\begin{aligned} \text{Re}(\mathring{a}^{(3)}_{\mu\tau}) , \text{Im}(\mathring{a}^{(3)}_{\mu\tau}) &< 2.9 \times 10^{-24} \text{ GeV } (99\% \text{ C.L.}) \\ &< 2.0 \times 10^{-24} \text{ GeV } (90\% \text{ C.L.}) \end{aligned}$ | this work |
| 4 | GRB vacuum birefringence | astrophysical | photon | $\sim 10^{-38}$ | [7] |
| | Laser interferometer | LIGO | photon | $\sim 10^{-22}$ | [8] |
| | Sapphire cavity oscillator | tabletop | photon | $\sim 10^{-18}$ | [5] |
| | Ne-Rb-K comagnetometer | tabletop | neutron | $\sim 10^{-29}$ | [11] |
| | trapped Ca ⁺ ion | tabletop | electron | $\sim 10^{-19}$ | [14] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $ \operatorname{Re}(\overset{\circ}{c}{}^{(4)}_{\mu\tau}) , \operatorname{Im}(\overset{\circ}{c}{}^{(4)}_{\mu\tau}) \le 3.9 \times 10^{-28} (99\% \text{ C.L.}) \le 2.7 \times 10^{-28} (90\% \text{ C.L.})$ | this work |
| 5 | GRB vacuum birefringence | astrophysical | photon | $\sim 10^{-34} { m GeV^{-1}}$ | [7] |
| | ultra-high-energy cosmic ray | astrophysical | proton | $\sim 10^{-22}$ to $10^{-18} \text{ GeV}^{-1}$ | [9] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $ \operatorname{Re}(\mathring{a}_{\mu\tau}^{(5)}) , \operatorname{Im}(\mathring{a}_{\mu\tau}^{(5)}) < 2.3 \times 10^{-32} \text{ GeV}^{-1} (99\% \text{ C.L.}) < 1.5 \times 10^{-32} \text{ GeV}^{-1} (90\% \text{ C.L.})$ | this work |
| 6 | GRB vacuum birefringene | astrophysical | photon | $\sim 10^{-31} { m GeV}^{-2}$ | [7] |
| | ultra-high-energy cosmic ray | astrophysical | proton | $\sim 10^{-42}$ to 10^{-35} GeV ⁻² | [9] |
| | gravitational Cherenkov radiation | astrophysical | gravity | $\sim 10^{-31} { m GeV}^{-2}$ | [15] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $\frac{ \operatorname{Re}(\mathring{c}_{\mu\tau}^{(6)}) , \operatorname{Im}(\mathring{c}_{\mu\tau}^{(6)}) }{< 9.1 \times 10^{-37} \text{ GeV}^{-2} (99\% \text{ C.L.})}$ | this work |
| 7 | GRB vacuum birefringence | astrophysical | photon | $\sim 10^{-28} { m GeV^{-3}}$ | [7] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $ \operatorname{Re}(\mathring{a}_{\mu\tau}^{(7)}) , \operatorname{Im}(\mathring{a}_{\mu\tau}^{(7)}) < 8.3 \times 10^{-41} \text{ GeV}^{-3} (99\% \text{ C.L.}) < 3.6 \times 10^{-41} \text{ GeV}^{-3} (90\% \text{ C.L.})$ | this work |
| 8 | gravitational Cherenkov radiation | astrophysical | gravity | $\sim 10^{-46} { m GeV}^{-4}$ | [15] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $ \operatorname{Re}(\hat{c}_{\mu\tau}^{(8)}) , \operatorname{Im}(\hat{c}_{\mu\tau}^{(8)}) < 5.2 \times 10^{-45} \text{ GeV}^{-4} (99\% \text{ C.L.}) < 1.4 \times 10^{-45} \text{ GeV}^{-4} (90\% \text{ C.L.})$ | this work |

TABLE I: Comparison of attainable best limits of SME coefficients in various fields.



$$H \sim \frac{m^2}{2E} + \mathring{a}^{(3)} - E \cdot \mathring{c}^{(4)} + E^2 \cdot \mathring{a}^{(5)} - E^3 \cdot \mathring{c}^{(6)} \cdots$$

Neutrino interferometry – Atmospheric neutrinos

| dim. | method | type | sector | limits | ref. |
|------|-----------------------------------|------------------------------|----------|---|-----------|
| 3 | CMB polarization | astrophysical | photon | $\sim 10^{-43} \text{ GeV}$ | [6] |
| | He-Xe comagnetometer | tabletop | neutron | $\sim 10^{-34}~{ m GeV}$ | [10] |
| | torsion pendulum | tabletop | electron | $\sim 10^{-31}~{ m GeV}$ | [12] |
| | muon g-2 | accelerator | muon | $\sim 10^{-24}~{ m GeV}$ | [13] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $\begin{aligned} \text{Re}(\mathring{a}^{(3)}_{\mu\tau}) , \text{Im}(\mathring{a}^{(3)}_{\mu\tau}) &< 2.9 \times 10^{-24} \text{ GeV} (99\% \text{ C.L.}) \\ &< 2.0 \times 10^{-24} \text{ GeV} (90\% \text{ C.L.}) \end{aligned}$ | this work |
| 4 | GRB vacuum birefringence | astrophysical | photon | $\sim 10^{-38}$ | [7] |
| | Laser interferometer | LIGO | photon | $\sim 10^{-22}$ | [8] |
| | Sapphire cavity oscillator | tabletop | photon | $\sim 10^{-18}$ | [5] |
| | Ne-Rb-K comagnetometer | tabletop | neutron | $\sim 10^{-29}$ | [11] |
| | trapped Ca ⁺ ion | tabletop | electron | $\sim 10^{-19}$ | [14] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $ \operatorname{Re}(\hat{c}^{(4)}_{\mu\tau}) , \operatorname{Im}(\hat{c}^{(4)}_{\mu\tau}) \le 3.9 \times 10^{-28} (99\% \text{ C.L.}) \le 2.7 \times 10^{-28} (90\% \text{ C.L.})$ | this work |
| 5 | GRB vacuum birefringence | astrophysical | photon | $\sim 10^{-34} { m GeV^{-1}}$ | [7] |
| | ultra-high-energy cosmic ray | astrophysical | proton | $\sim 10^{-22}$ to $10^{-18} { m GeV^{-1}}$ | [9] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $\frac{ \operatorname{Re}(\hat{a}_{\mu\tau}^{(5)}) }{ \operatorname{Im}(\hat{a}_{\mu\tau}^{(5)}) } \leq 2.3 \times 10^{-32} \text{ GeV}^{-1} (99\% \text{ C.L.}) \\ < 1.5 \times 10^{-32} \text{ GeV}^{-1} (90\% \text{ C.L.})$ | this work |
| 6 | GRB vacuum birefringene | astrophysical | photon | $\sim 10^{-31} { m GeV^{-2}}$ | [7] |
| | ultra-high-energy cosmic ray | astrophysical | proton | $\sim 10^{-42}$ to 10^{-35} GeV ⁻² | [9] |
| | gravitational Cherenkov radiation | astrophysical | gravity | $\sim 10^{-31} { m GeV}^{-2}$ | [15] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $ \operatorname{Re} \left(\overset{\circ (6)}{c_{\mu\tau}} \right) , \operatorname{Im} \left(\overset{\circ (6)}{c_{\mu\tau}} \right) < 1.5 \times 10^{-36} \text{ GeV}^{-2} (99\% \text{ C.L.}) < 9.1 \times 10^{-37} \text{ GeV}^{-2} (90\% \text{ C.L.})$ | this work |
| 7 | GRB vacuum birefringence | astrophysical | photon | $\sim 10^{-28}~{ m GeV}^{-3}$ | [7] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $ \operatorname{Re}(\mathring{a}^{(7)}_{\mu\tau}) , \operatorname{Im}(\mathring{a}^{(7)}_{\mu\tau}) < 8.3 \times 10^{-41} \text{ GeV}^{-3} (99\% \text{ C.L.}) < 3.6 \times 10^{-41} \text{ GeV}^{-3} (90\% \text{ C.L.})$ | this work |
| 8 | gravitational Cherenkov radiation | astrophysical | gravity | $\sim 10^{-46} { m GeV^{-4}}$ | [15] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $\begin{aligned} \operatorname{Re}\left(\overset{\circ(8)}{c}_{\mu\tau}^{(8)}\right) , \operatorname{Im}\left(\overset{\circ(8)}{c}_{\mu\tau}^{(8)}\right) &< 5.2 \times 10^{-45} \text{ GeV}^{-4} (99\% \text{ C.L.}) \\ &< 1.4 \times 10^{-45} \text{ GeV}^{-4} (90\% \text{ C.L.}) \end{aligned}$ | this work |

TABLE I: Comparison of attainable best limits of SME coefficients in various fields.

Atomic experiments have strong limits on lower order couplings (renormalizable). Neutrino oscillations are doing good

$$H \sim \frac{m^2}{2E} + \mathring{a}^{(3)} - E \cdot \mathring{c}^{(4)} + E^2 \cdot \mathring{a}^{(5)} - E^3 \cdot \mathring{c}^{(6)} \cdots$$

Neutrino interferometry – Atmospheric neutrinos

| dim. | method | type | sector | limits | ref. |
|------|-----------------------------------|------------------------------|----------|--|-----------|
| 3 | CMB polarization | astrophysical | photon | $\sim 10^{-43}~{ m GeV}$ | [6] |
| | He-Xe comagnetometer | tabletop | neutron | $\sim 10^{-34}~{ m GeV}$ | [10] |
| | torsion pendulum | tabletop | electron | $\sim 10^{-31}~{ m GeV}$ | [12] |
| | muon g-2 | accelerator | muon | $\sim 10^{-24}~{ m GeV}$ | [13] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $\begin{aligned} \text{Re}(\mathring{a}^{(3)}_{\mu\tau}) , \text{Im}(\mathring{a}^{(3)}_{\mu\tau}) &< 2.9 \times 10^{-24} \text{ GeV } (99\% \text{ C.L.}) \\ &< 2.0 \times 10^{-24} \text{ GeV } (90\% \text{ C.L.}) \end{aligned}$ | this work |
| 4 | GRB vacuum birefringence | astrophysical | photon | $\sim 10^{-38}$ | [7] |
| | Laser interferometer | LIGO | photon | $\sim 10^{-22}$ | [8] |
| | Sapphire cavity oscillator | tabletop | photon | $\sim 10^{-18}$ | [5] |
| | Ne-Rb-K comagnetometer | tabletop | neutron | $\sim 10^{-29}$ | [11] |
| | trapped Ca ⁺ ion | tabletop | electron | $\sim 10^{-19}$ | [14] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $ \operatorname{Re}(\hat{c}^{\circ(4)}_{\mu\tau}) , \operatorname{Im}(\hat{c}^{\circ(4)}_{\mu\tau}) \le 3.9 \times 10^{-28} (99\% \text{ C.L.}) \le 2.7 \times 10^{-28} (90\% \text{ C.L.})$ | this work |
| 5 | GRB vacuum birefringence | astrophysical | photon | $\sim 10^{-34} { m GeV^{-1}}$ | [7] |
| | ultra-high-energy cosmic ray | astrophysical | proton | $\sim 10^{-22}$ to $10^{-18} { m GeV^{-1}}$ | [9] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $\frac{ \operatorname{Re}(\mathring{a}^{(5)}_{\mu\tau}) }{ \operatorname{Im}(\mathring{a}^{(5)}_{\mu\tau}) } \stackrel{< 2.3 \times 10^{-32} \text{ GeV}^{-1} (99\% \text{ C.L.})}{< 1.5 \times 10^{-32} \text{ GeV}^{-1} (90\% \text{ C.L.})}$ | this work |
| 6 | GRB vacuum birefringene | astrophysical | photon | $\sim 10^{-31} { m ~GeV^{-2}}$ | [7] |
| | ultra-high-energy cosmic ray | astrophysical | proton | $\sim 10^{-42}$ to 10^{-35} GeV ⁻² | [9] |
| | gravitational Cherenkov radiation | astrophysical | gravity | $\sim 10^{-31} { m GeV}^{-2}$ | [15] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $ \operatorname{Re}\left(\overset{\circ(6)}{c}_{\mu\tau}^{(6)}\right) , \operatorname{Im}\left(\overset{\circ(6)}{c}_{\mu\tau}^{(6)}\right) < 1.5 \times 10^{-36} \text{ GeV}^{-2} (99\% \text{ C.L.}) < 9.1 \times 10^{-37} \text{ GeV}^{-2} (90\% \text{ C.L.})$ | this work |
| 7 | GRB vacuum birefringence | astrophysical | photon | $\sim 10^{-28} \text{ GeV}^{-3}$ | [7] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $\begin{aligned} \operatorname{Re}(\mathring{a}^{(7)}_{\mu\tau}) , \operatorname{Im}(\mathring{a}^{(7)}_{\mu\tau}) &< 8.3 \times 10^{-41} \text{ GeV}^{-3} \text{ (99\% C.L.)} \\ &< 3.6 \times 10^{-41} \text{ GeV}^{-3} \text{ (90\% C.L.)} \end{aligned}$ | this work |
| 8 | gravitational Cherenkov radiation | astrophysical | gravity | $\sim 10^{-46} { m GeV}^{-4}$ | [15] |
| | neutrino oscillation | $\operatorname{atmospheric}$ | neutrino | $ \operatorname{Re}(\hat{c}_{\mu\tau}^{(8)}) , \operatorname{Im}(\hat{c}_{\mu\tau}^{(8)}) < 5.2 \times 10^{-45} \text{ GeV}^{-4} (99\% \text{ C.L.}) < 1.4 \times 10^{-45} \text{ GeV}^{-4} (90\% \text{ C.L.})$ | this work |

TABLE I: Comparison of attainable best limits of SME coefficients in various fields.

Astrophysical observations have strong limits on higher order couplings non-renormalizable). Neutrino oscillations are better tools there

$$H \sim \frac{m^2}{2E} + \mathring{a}^{(3)} - E \cdot \mathring{c}^{(4)} + E^2 \cdot \mathring{a}^{(5)} - E^3 \cdot \mathring{c}^{(6)} \cdots$$

Argüelles, TK, Salvado, PRL115(2015)161303 Diaz et al, PRD85(2013)096005, Ellis et al, PLB789(2018)352 Neutrino interferometry – Astrophysical high-energy neutrinos

Combination of longer baseline and higher energy makes extra-terrestrial neutrino to be the most sensitive source of fundamental physics.

Vacuum Cherenkov radiation and Blazar neutrino ToF can limit new physics of neutrino.

Neutrino mixing properties of astrophysical neutrinos can push this limit further.

- In principle, we can do it without knowing neutrino flavours at the production
- This is the most sensitive, as long as we assume new physics cause neutrino mixing





Neutrino interferometry - Astrophysical high-energy neutrinos

Any new physics can end up in the effective Hamiltonian

$$h_{eff} = \frac{1}{2E} U^{\dagger} M^2 U + \sum_{n} \left(\frac{E}{\Lambda_n}\right)^n \tilde{U}_n^{\dagger} O_n \tilde{U}_n = V^{\dagger} \Delta V$$

neutrino oscillation formula

$$P_{\alpha \to \beta}(L) = 1 - 4\sum_{i>j} \operatorname{Re}(V_{\alpha i}^* V_{\beta i}^* V_{\alpha j} V_{\beta j}) \sin^2\left(\frac{\Delta_{ij}}{2}L\right) + 2\sum_{i>j} \operatorname{Im}(V_{\alpha i}^* V_{\beta i}^* V_{\alpha j} V_{\beta j}) \sin\left(\Delta_{ij}L\right)$$

neutrino mixing formula

$$P_{\alpha \to \beta}(L \to \infty) \sim 1 - 2\sum_{i>j} \operatorname{Re}(V_{\alpha i}^* V_{\beta i}^* V_{\alpha j} V_{\beta j}) = \sum_i |V_{\alpha i}|^2 |V_{\beta i}|^2$$

Information of small contamination of new physics appears on neutrino flavour mixings

At high energy, neutrino mass term is suppressed



Neutrino interferometry - Astrophysical high-energy neutrinos



Neutrino interferometry - Astrophysical high-energy neutrinos



Neutrino interferometry - Astrophysical high-energy neutrinos



New IceCube data (2018)

Blazar neutrino

- IC170922A and TXS 0506+056
- Optical coincidence
- Clustering from this direction



IceCube, Science361(2018)147



0.0 HESE with ternary topology ID 1.0 best fit: 0.35 : 0.45 : 0.2 Sensitivity, E^{-2.9} spectrum 0,2 1:1:1 flavor composition 0.8 Fraction of _____ WORK IN PROGRESS 68 × 0 °.~ 0.2 .68 % 4.0 .95 % 0.0 *S*. 9.0 0.0 0 Fraction of ν_{e}

New flavour ratio measurement

- Likelihood is very shallow and fit confuse between ν_e and ν_τ

- New flavour ratio result has some power to distinguish between ν_{e} and ν_{τ}



11

Argüelles (IceCube), Neutrino 2018

University of London

IceCube preliminary

Fraction of "

Fraction or 2

IceCube flavor ratio test for new physics (to be published)

We start to exclude possible new physics in Planck scale signal region

- This moment, we can exclude only 2 scenarios
- dimnesion-3 vacuum operator limit ($\bar{\psi}a^{\mu}\gamma_{\mu}\psi, \bar{\psi}b^{\mu}\gamma_{\mu}\gamma_{5}\psi$) ~ 10⁻²⁵ GeV
- dimnesion-4 vacuum operator limit ($\bar{\psi}c^{\mu\nu}\gamma_{\mu}\partial_{\nu}\psi, \bar{\psi}d^{\mu\nu}\gamma_{\mu}\gamma_{5}\partial_{\nu}\psi$) ~ 10⁻³³
- dimension-6 vacuum operator limit ~10⁻⁴⁰ GeV⁻²



12

Conclusion

Neutrino interferometry is a powerful technique to look for new physics if new physics couple with neutrinos and they cause neutrino mixings.

Spectrum distortion of atmospheric neutrino is used to look for new physics.

Astrophysical neutrino mixing sensitivity reaches to naïve expectation of Planck scale. However, in this moment, the sensitivity is limited. We need more statistics and better particle identification algorithm to find new physics.



Thank you for your attention!

Neutrino interferometry - Astrophysical high-energy neutrinos

Combination of longer baseline and higher energy makes extra-terrestrial neutrino to be the most sensitive source of fundamental physics.



New physics limits and projected sensitivity

University of London