

# QE kinematics

1. Introduction
2. MiniBooNE phase space study
3. T2K phase space study
4. Conclusion

Teppei Katori

Queen Mary University of London

Valencia Neutrino Interaction T2K meeting (VANISH)

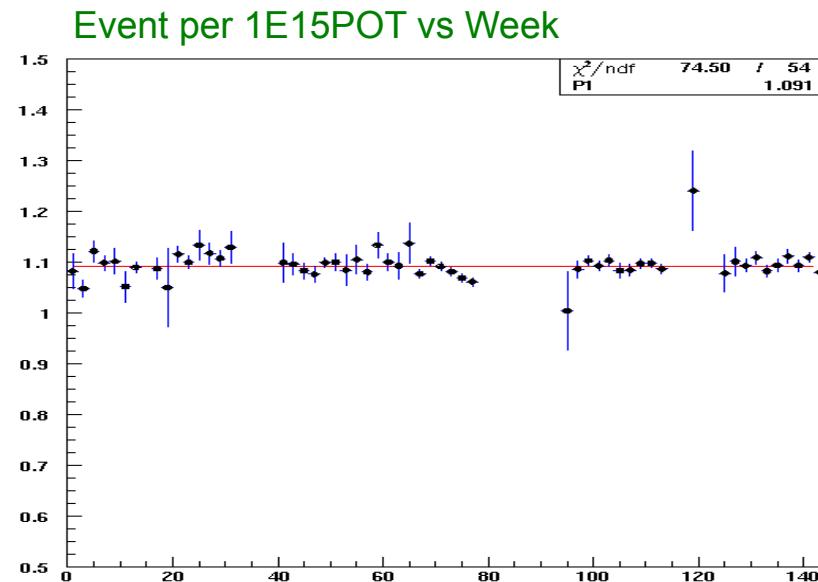
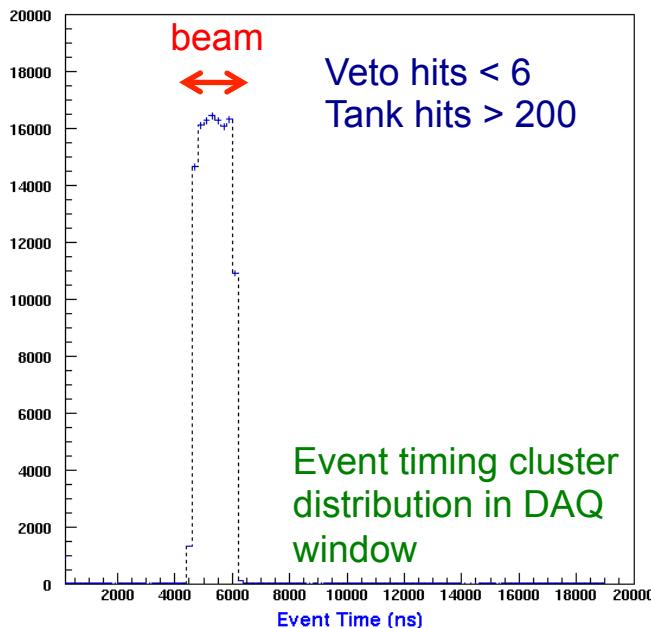
University of Valencia, Valencia, Spain, Apr. 3, 2014

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# 1. Introduction

## MiniBooNE detector

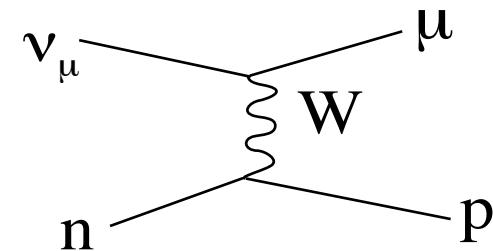
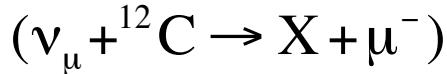
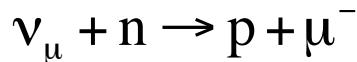
- 800 MeV (700 MeV) neutrinos from Booster neutrino beamline at Fermilab.
- 12m diameter (10m fiducial) spherical Cherenkov detector filled with mineral oil ( $\text{CH}_2$ ).
- CC interaction is isolated by veto hits <6 and tank hits > 200.
- Ring pattern is used for particle ID.
- Roughly 1 neutrino interaction per minute for  $5\text{E}16$  proton per hour (neutrino mode).



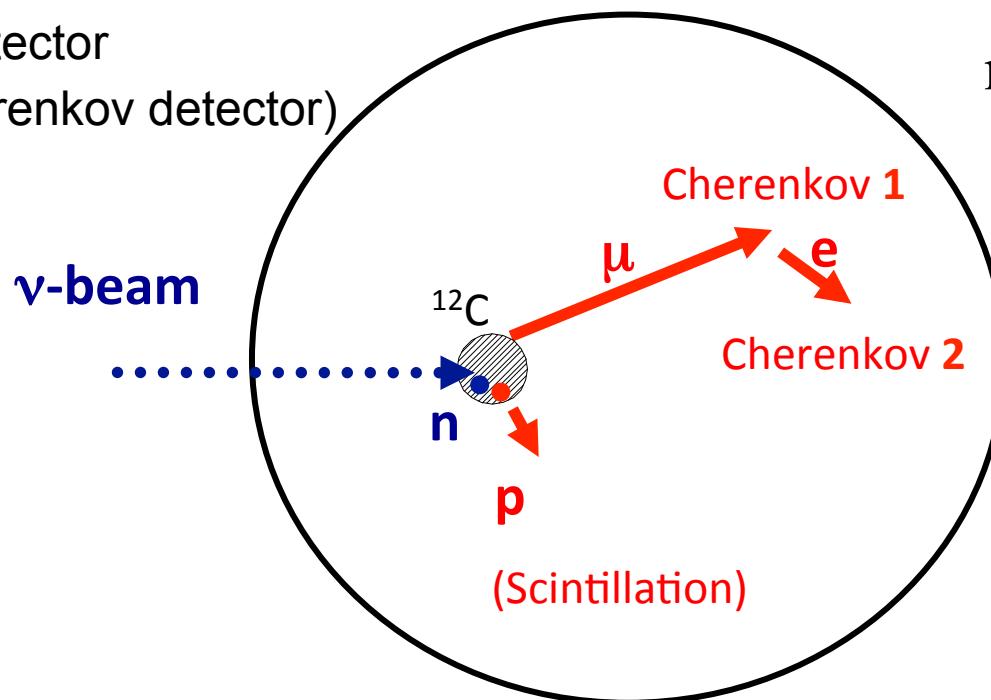
# 1. Introduction

## MiniBooNE CCQE definition

$\nu_\mu$  charged current quasi-elastic ( $\nu_\mu$ CCQE) interaction is an important channel for the neutrino oscillation physics and the most abundant (~40%) interaction type in MiniBooNE detector



MiniBooNE detector  
(spherical Cherenkov detector)

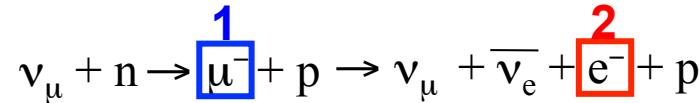


muon like Cherenkov light and subsequent decayed electron (Michel electron) like Cherenkov light are the signal of CCQE event

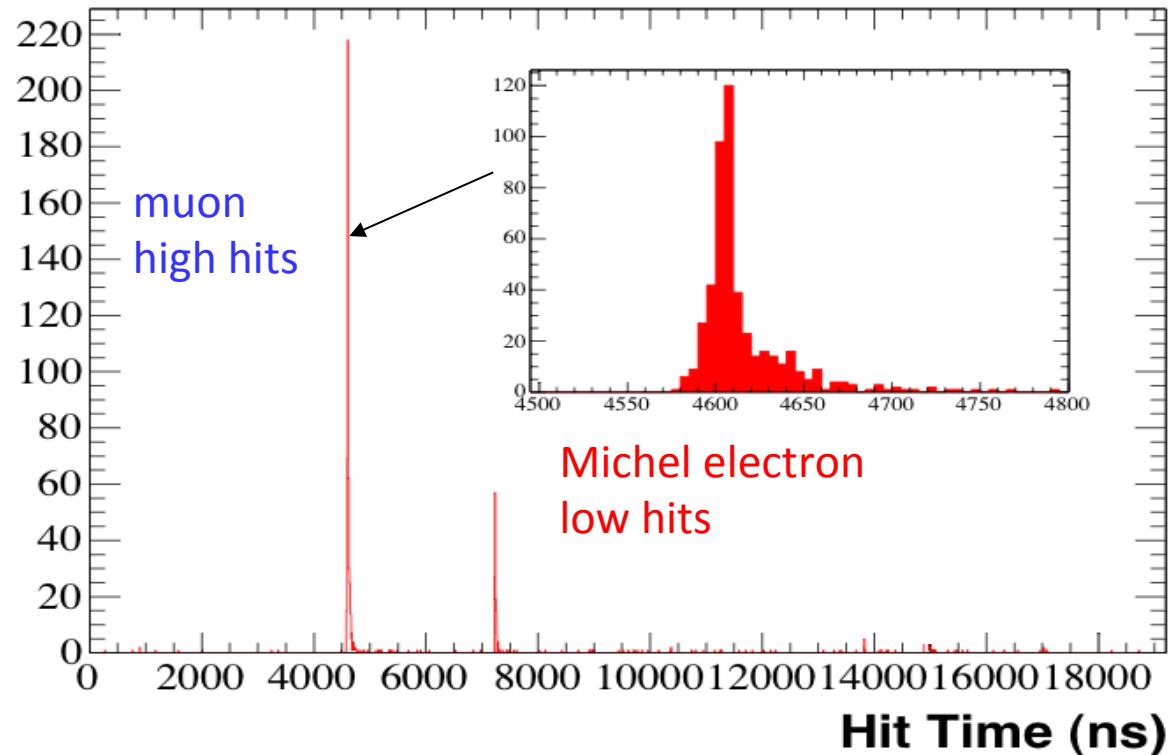
# 1. Introduction

## MiniBooNE CCQE definition

$\nu_\mu$  CCQE interactions ( $\nu + n \rightarrow \mu + p$ ) has characteristic two “subevent” structure from muon decay



27% efficiency  
77% purity  
146,070 events with  
5.58E20 POT



# 1. Introduction

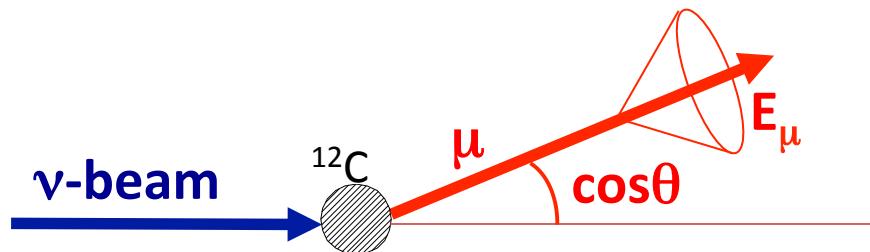
## MiniBooNE CCQE observables

All kinematics are specified from 2 observables, muon energy  $E_\mu$  and muon scattering angle  $\theta_\mu$

Energy of the neutrino  $E_\nu^{\text{QE}}$  and 4-momentum transfer  $Q^2_{\text{QE}}$  can be reconstructed by these 2 observables, under the assumption of CCQE interaction with bound neutron at rest (“QE assumption”)

$$E_\nu^{\text{QE}} = \frac{ME_\mu - 0.5m_\mu^2}{M - E_\mu + p_\mu \cos\theta_\mu}$$

$$Q^2_{\text{QE}} = -m_\mu^2 + 2E_\nu^{\text{QE}}(E_\mu - p_\mu \cos\theta_\mu)$$

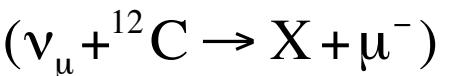


# 1. Overview of MiniBooNE xs measurements

CCQE (CC0 $\pi$ )

PRD81(2010)092005

FERIMILAB-THESIS-2008-64



Signal definition: 1  $\mu$  + 0  $\pi$  + N protons

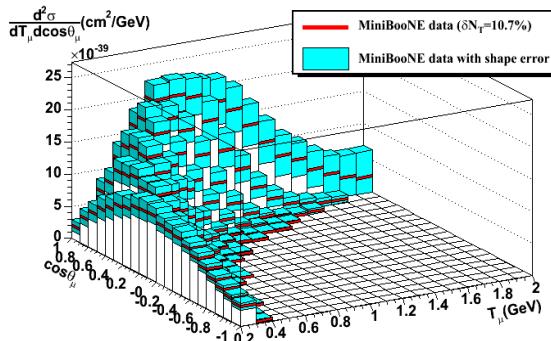
## Why we measure

- Test CCQE models

## Why MiniBooNE measure

- Largest sample (~40%) to test detector efficiency, veto efficiency, event uniformity, timing, etc
- Best sample to study  $\nu_e$ CCQE kinematics (=oscillation signal)
- $\nu_\mu$ CCQE to constraint  $\nu_e$  from  $\mu$ -decay in oscillation sample

It is important to measure CCQE!

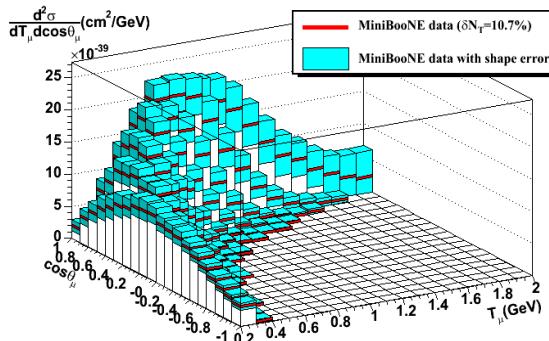


MiniBooNE collaboration,  
PRD81(2010)092005

# 1. Overview of MiniBooNE xs measurements

## MiniBooNE CCQE basics

- high statistics
- high purity
- error dominated by flux (as usual)



MiniBooNE collaboration,  
PRD81(2010)092005

integrated protons on target		$5.58 \times 10^{20}$
energy-integrated $\nu_\mu$ flux		$2.88 \times 10^{11} \nu_\mu/\text{cm}^2$
CCQE candidate events		146070
CCQE efficiency ( $R < 550$ cm)		26.6%
background channel	events	fraction
NCE	45	<0.1%
CC1 $\pi^+$	26866	18.4%
CC1 $\pi^0$	3762	2.6%
NC1 $\pi^\pm$	535	0.4%
NC1 $\pi^0$	43	<0.1%
other $\nu_\mu$	328	0.2%
all non- $\nu_\mu$	1977	1.4%
total background	33556	23.0%

source	normalization error (%)
neutrino flux prediction	8.66
background cross sections	4.32
detector model	4.60
kinematic unfolding procedure	0.60
statistics	0.26
total	10.7

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## 2. Neutrino experiment

Experiment measure the interaction rate  $R$ ,

$$R \sim \int \Phi \times \sigma \times \varepsilon$$

- $\Phi$  : neutrino flux
- $\sigma$  : cross section
- $\varepsilon$  : efficiency

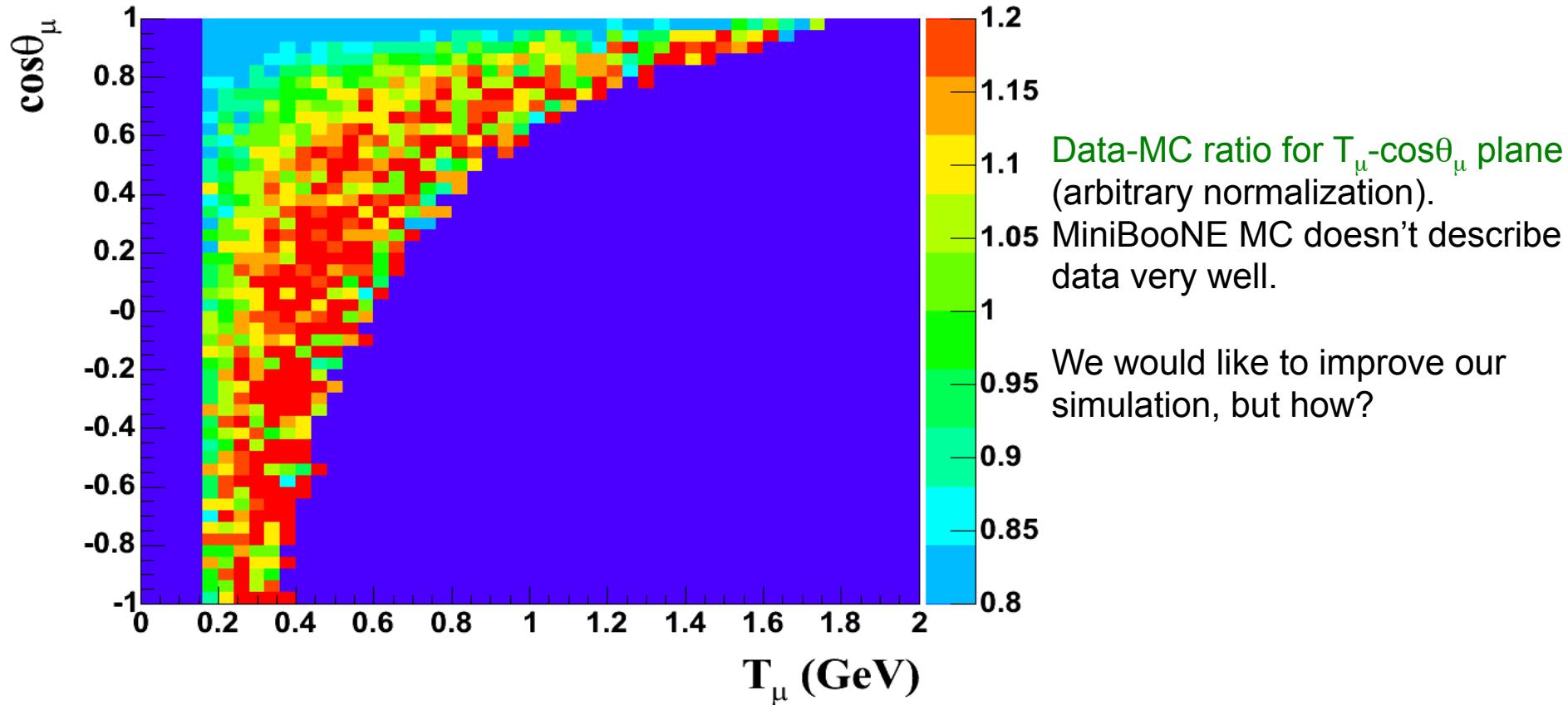
When do you see data-MC disagreement, how to interpret the result?

## 2. MiniBooNE phase space

CCQE kinematic space ( $T_\mu$ - $\cos\theta_\mu$  plane) in MiniBooNE

Since observables are muon energy ( $T_\mu$ ) and angle ( $\cos\theta_\mu$ ), these 2 variables completely specify the kinematic space.

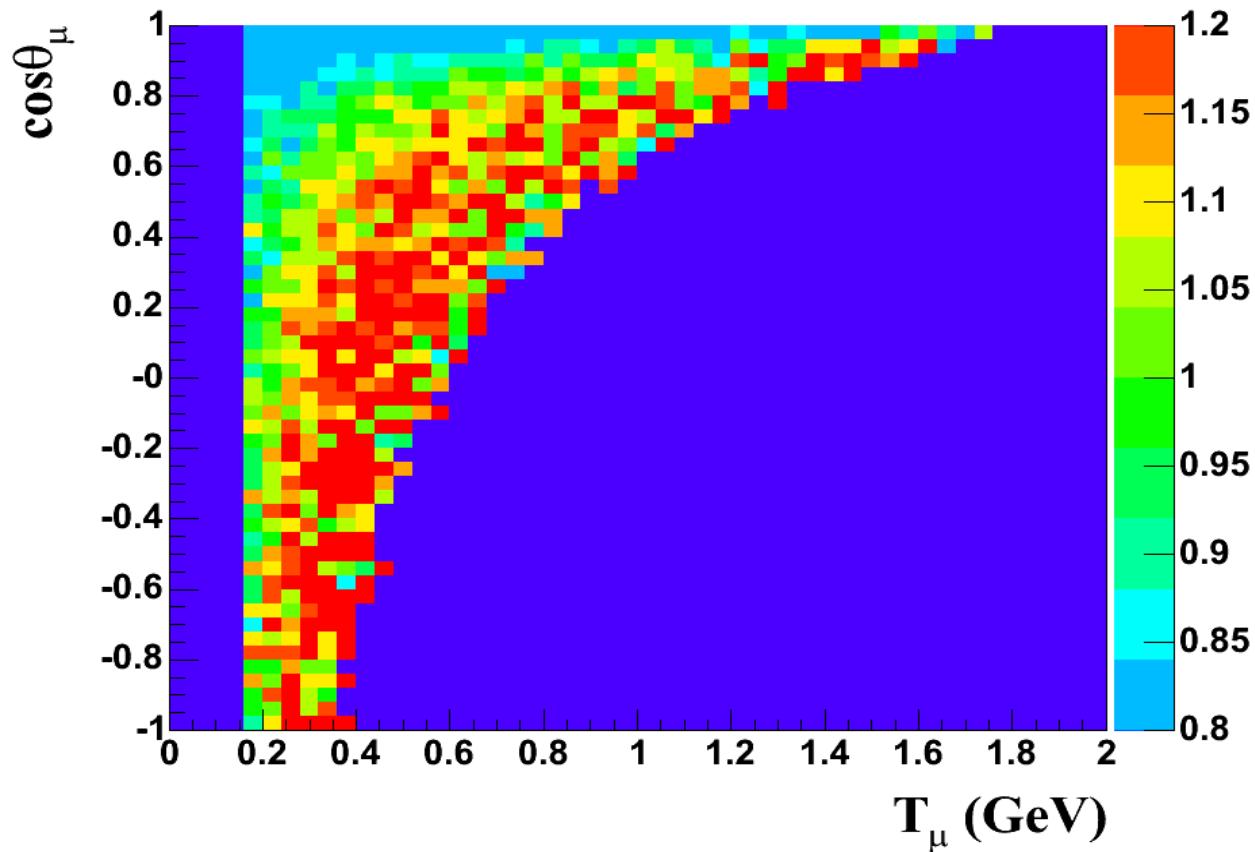
$$\frac{d\sigma^2}{dEd\Omega} \sim \frac{d\sigma^2}{dEd(\cos\vartheta)}$$



## 2. MiniBooNE phase space

Without knowing flux, you cannot modify cross section model

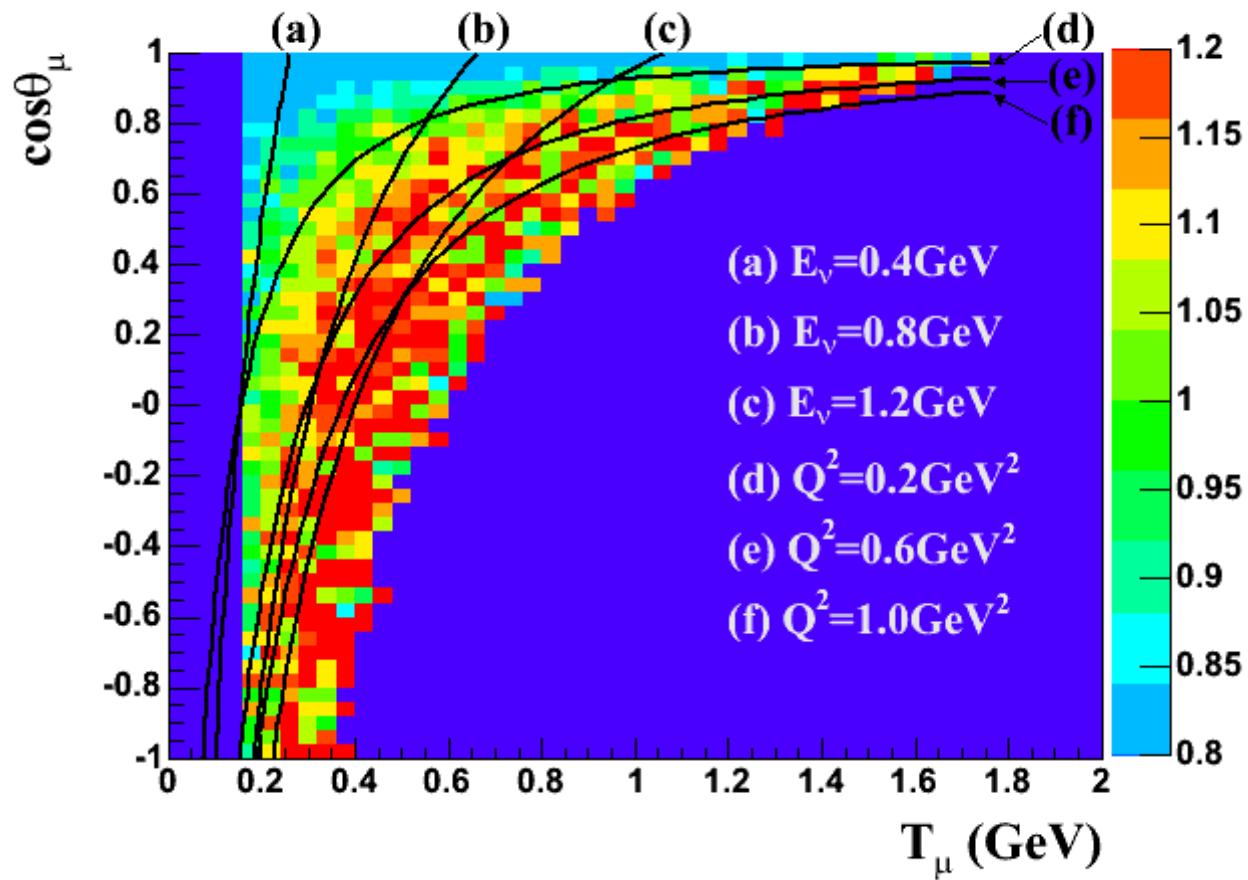
$$R \sim \int \Phi \times \sigma$$



## 2. MiniBooNE phase space

Without knowing flux, you cannot modify cross section model

$$R(E_\nu, Q^2) \sim \int \Phi(E_\nu) \times \sigma(Q^2)$$



The data-MC disagreement follows equal  $Q^2$ -lines, not equal  $E_\nu$ -lines.

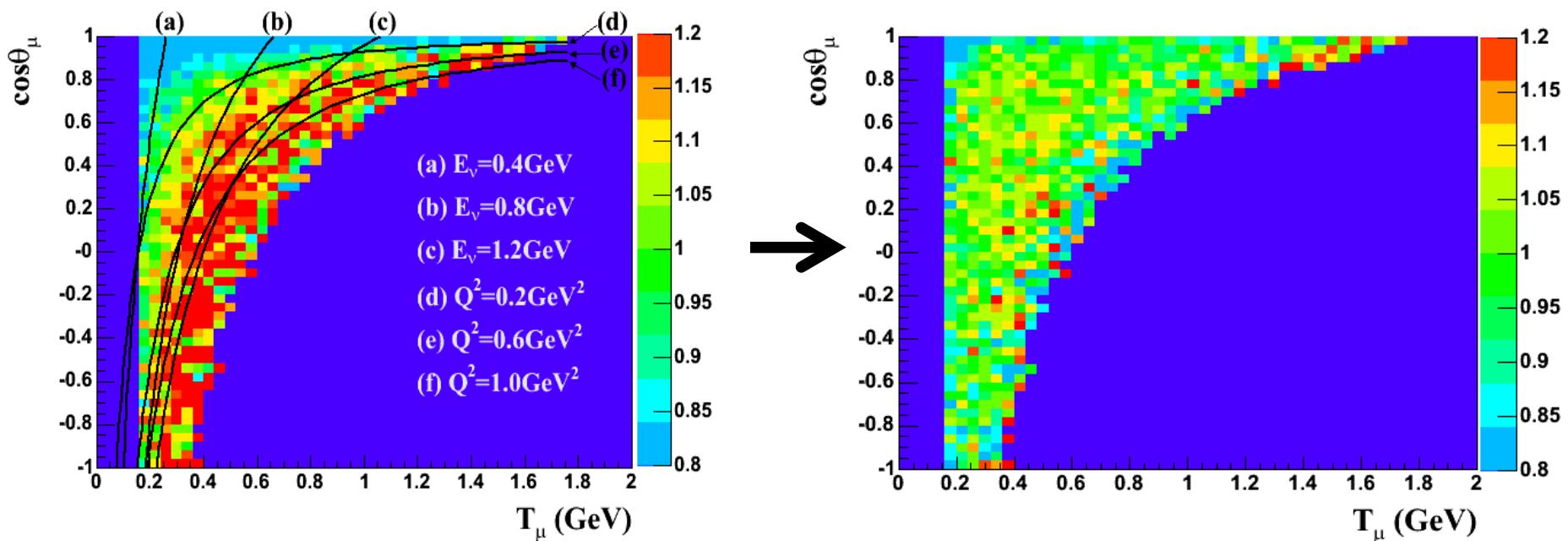
→ Something wrong in cross section model, not flux model.

## 2. MiniBooNE phase space

Without knowing flux, you cannot modify cross section model

$$R(E_\nu, Q^2) \sim \int \Phi(E_\nu) \times \sigma(Q^2)$$

After tuning cross section parameters, data and MC agree.



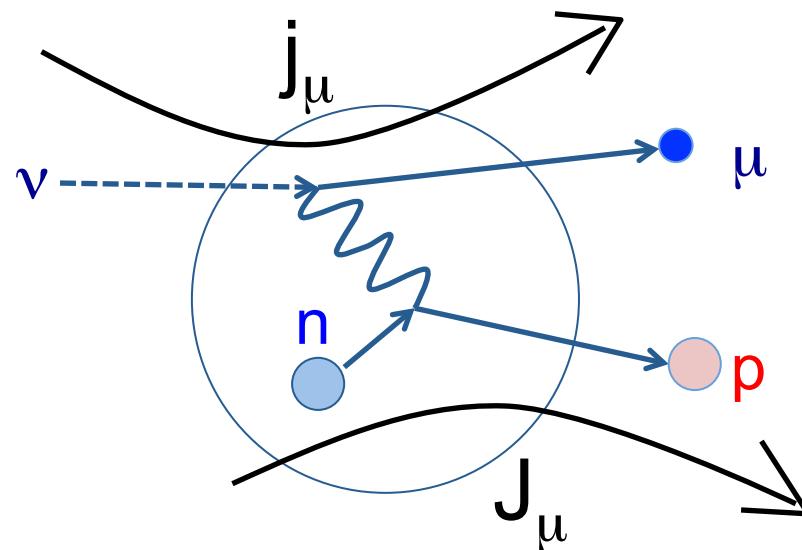
## 2. Cross section formula

### Cross section formula

- product of Leptonic and hadronic tensor

$$d\sigma = \frac{\overline{|M^2|}}{F} dLips, \quad dLips = \frac{dQ^2}{16\pi M E_\nu}$$

$$\overline{|M^2|} = \frac{G_F^2}{4} \sum_{\text{spin}} J_\mu J^\dagger j^\mu j^{\nu\dagger} = \frac{G_F^2 \cos^2 \theta_c}{4} W_{\mu\nu} L^{\mu\nu}$$



Question:  
→ how to model  $W_{\mu\nu}$ ?

## 2. Smith-Moniz formalism

Nucleus is described by the collection of incoherent nucleons.

$$(W_{\mu\nu})_{ab} = \int_{E_{lo}}^{E_{hi}} f(\vec{k}, \vec{q}, w) T_{\mu\nu} dE : \text{hadronic tensor}$$

$f(\vec{k}, \vec{q}, w)$  : nucleon phase space distribution

$T_{\mu\nu} = T_{\mu\nu} (F_1, F_2, F_A, F_P)$  : nucleon form factors

$F_A(Q^2) = g_A / (1 + Q^2/M_A^2)^2$  : Axial vector form factor

$E_{hi}$  : the highest energy state of nucleon

$E_{lo}$  : the lowest energy state of nucleon

Although Smith-Moniz formalism offers variety of choice, one can solve this equation analytically if the nucleon space is simple.



[Home](#) » Dr. Ernest Moniz

ABOUT US

DR. ERNEST MONIZ - SECRETARY OF ENERGY



## 2. Relativistic Fermi Gas (RFG) model

Nucleus is described by the collection of incoherent **Fermi gas particles**.

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$E_{hi}$  : the highest energy state of nucleon =  $\sqrt{(p_F^2 + M^2)}$

$E_{lo}$  : the lowest energy state of nucleon =  $\kappa \left( \sqrt{(p_F^2 + M^2)} - \omega + E_B \right)$

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MiniBooNE tuned following 2 parameters using  $Q^2$  distribution by least  $\chi^2$  fit;

$M_A$  = effective axial mass

$\kappa$  = effective Pauli blocking parameter

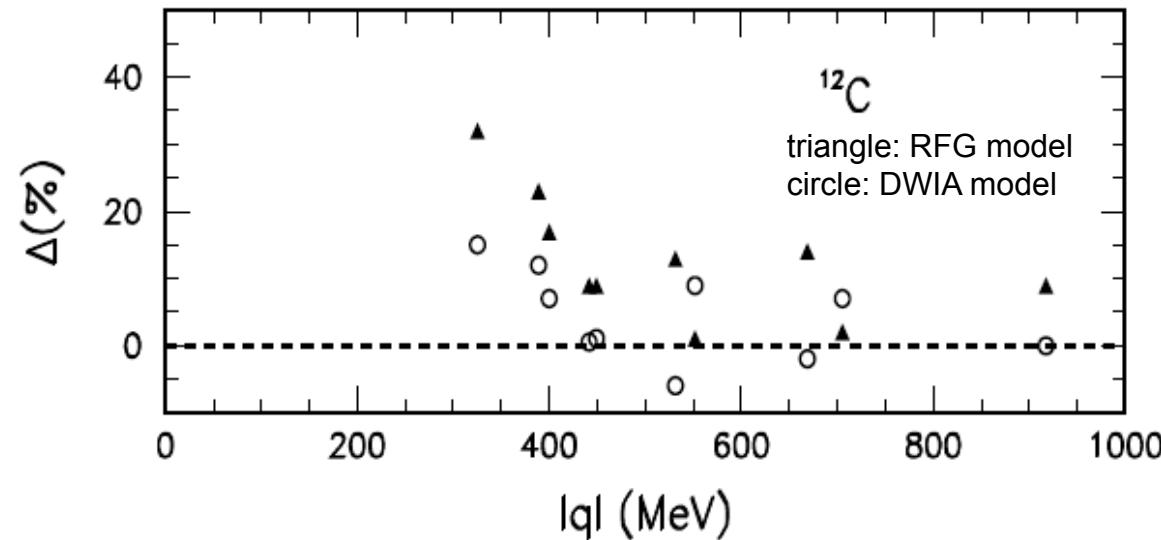
## 2. Relativistic Fermi Gas (RFG) model

### Relativistic Fermi Gas (RFG) Model

Nucleus is described by the collection of incoherent Fermi gas particles. All details come from hadronic tensor.

In low  $|q|$ , The RFG model systematically over predicts cross section for electron scattering experiments at low  $|q|$  ( $\sim$ low  $Q^2$ )

Data and predicted xs difference for  $^{12}\text{C}$

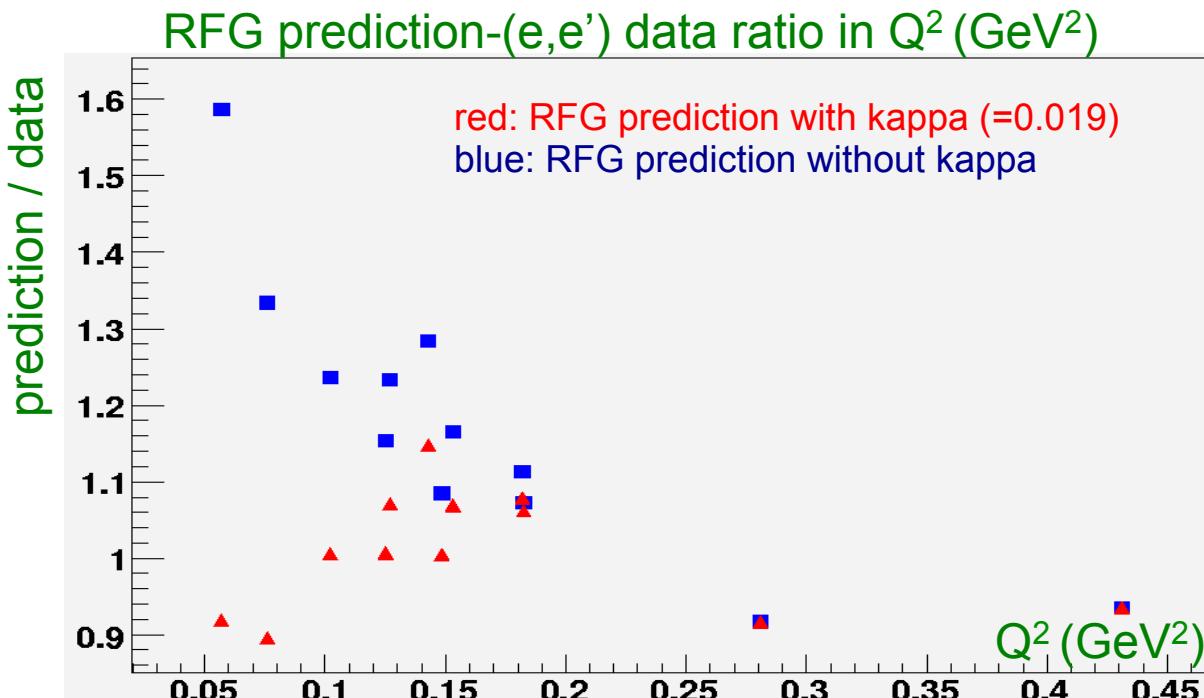


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## 2. Summary

MiniBooNE studied kinematic space to disentangle flux and cross section effect.

Multi-dimentional distribution is very powerful to test underlining models (or find mistake)

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### 3. MiniBooNE phase space study

Violation of impulse approximation

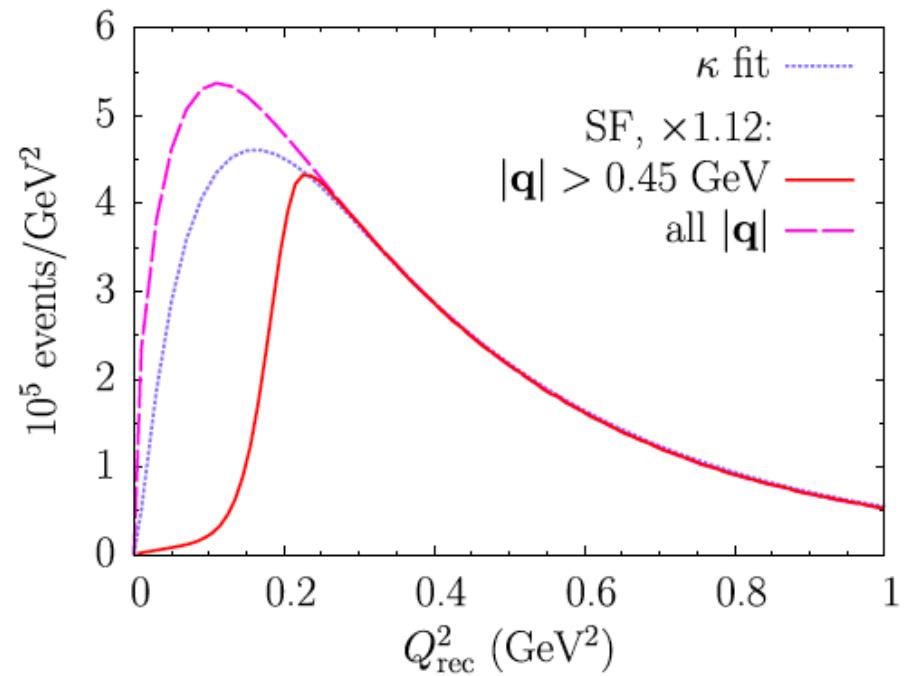
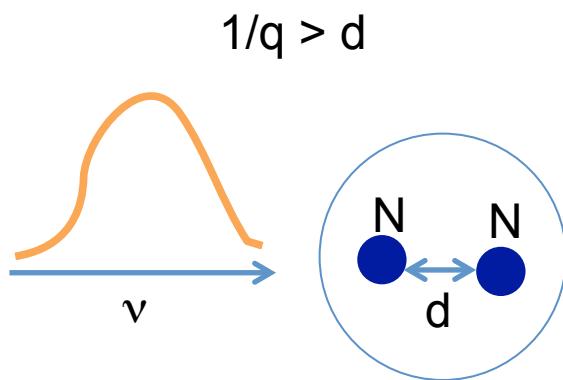


FIG. 6 (color online). Comparison of the MiniBooNE parametrization of the data (dotted line), labeled as the  $\kappa$  fit, to the spectral function calculation (dashed line). The solid line depicts the contribution to the latter from the region where the IA is expected to be valid. The SF results are multiplied by a factor 1.12 to make them match the  $\kappa$  fit.

### 3. MiniBooNE phase space study

Violation of impulse approximation

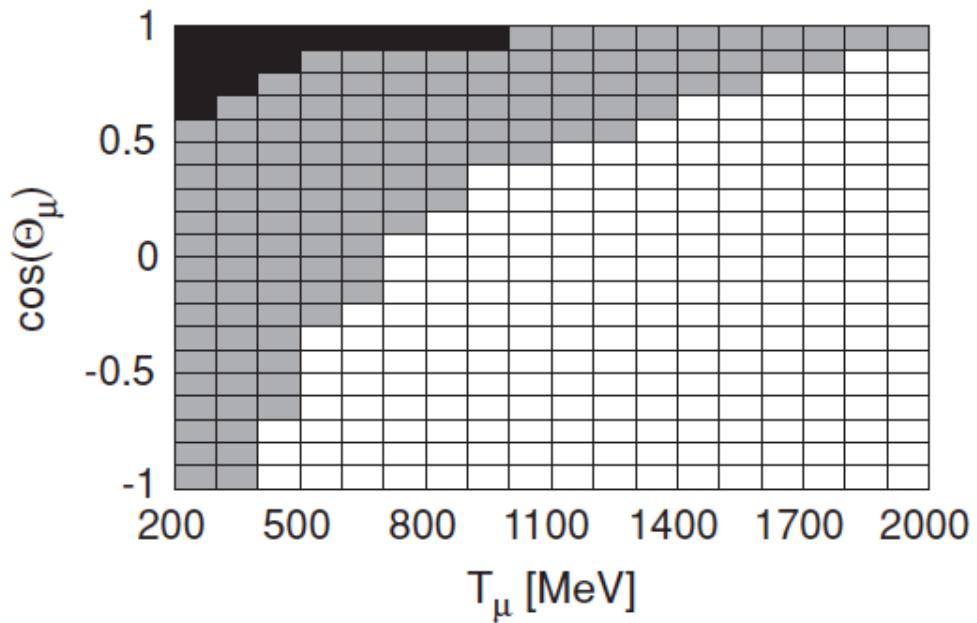
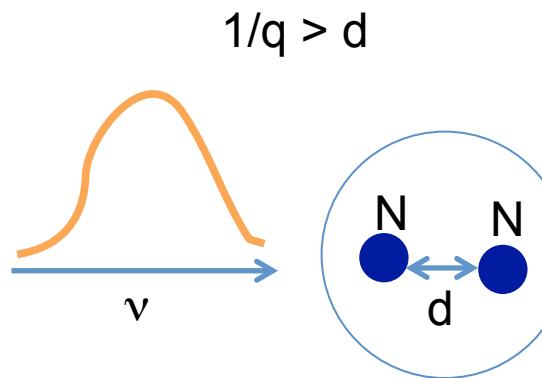
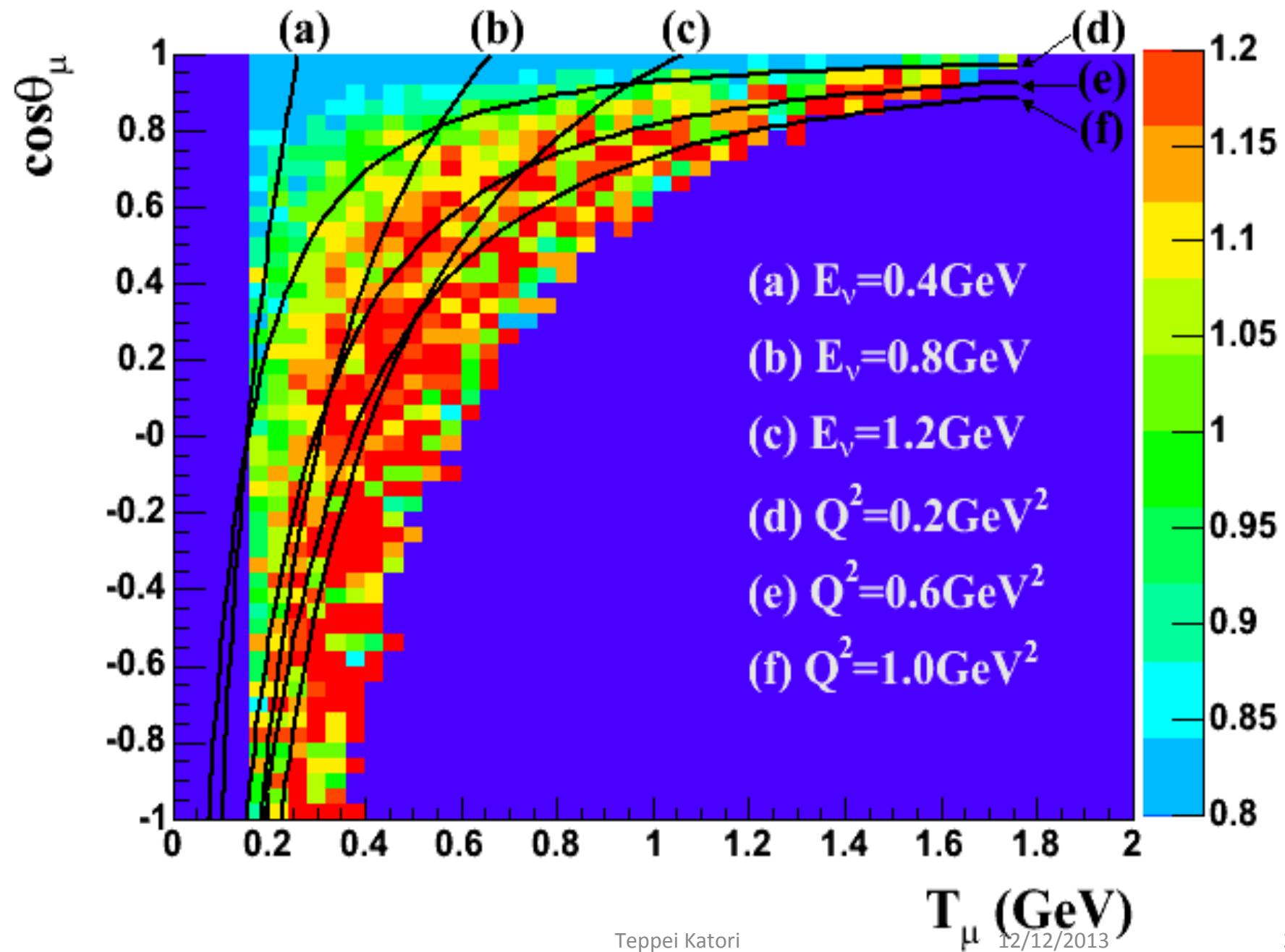
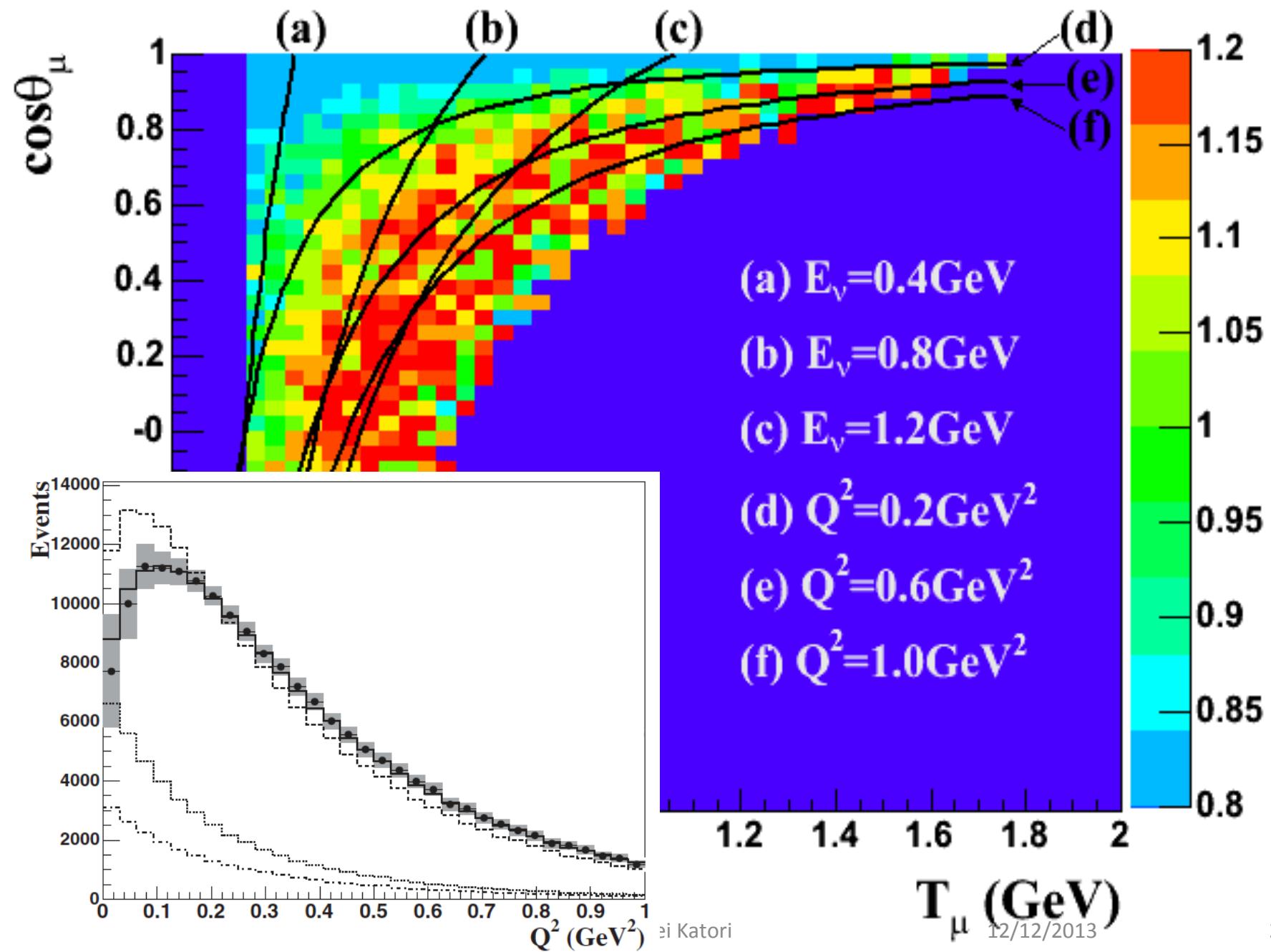
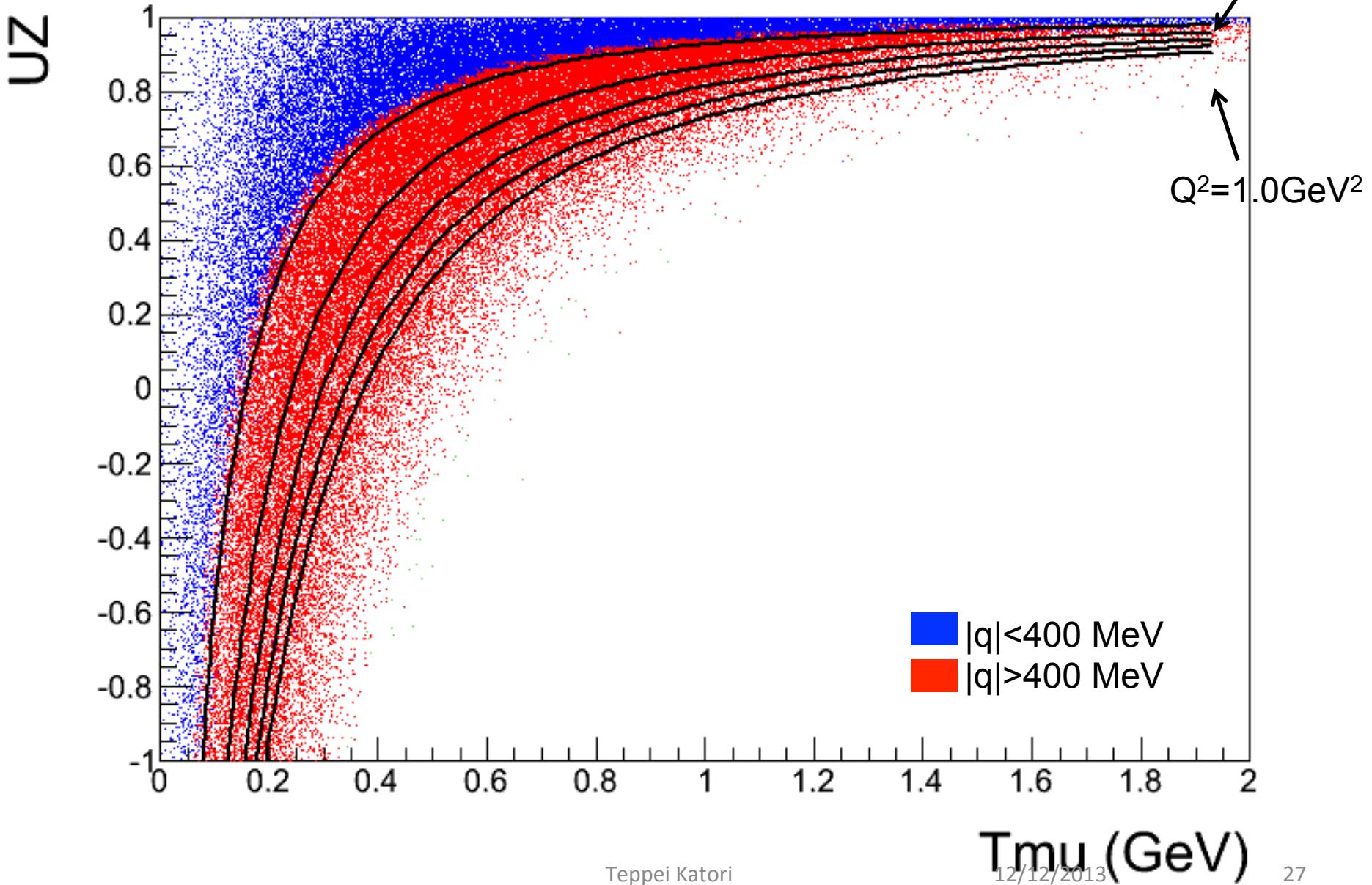


FIG. 3. Bins excluded from the fitting procedure for  $q_{\text{cut}} = 400$  MeV/ $c$  are shown in black. Bins with a nonzero cross section measured by MiniBooNE are shown in gray.

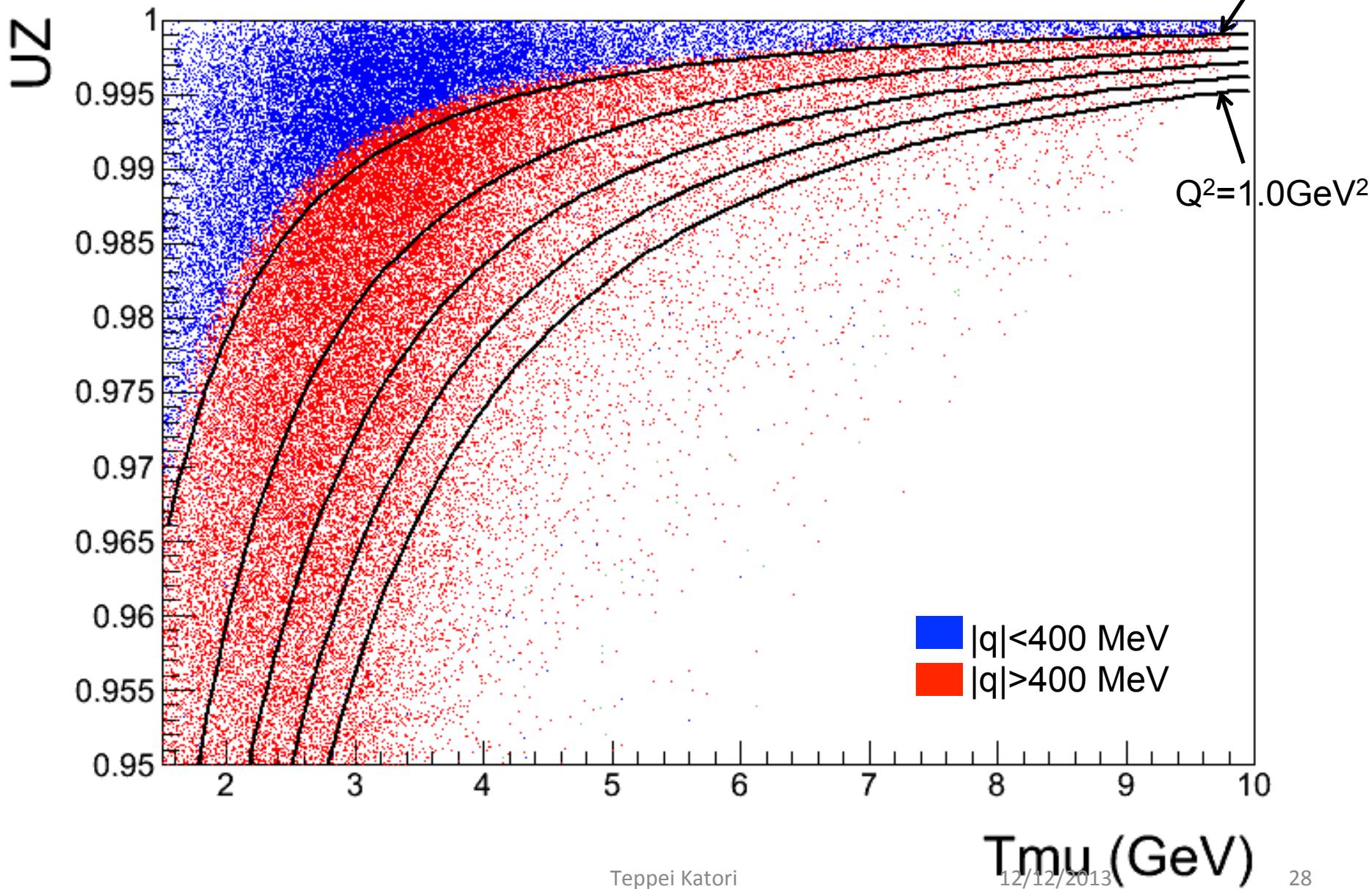




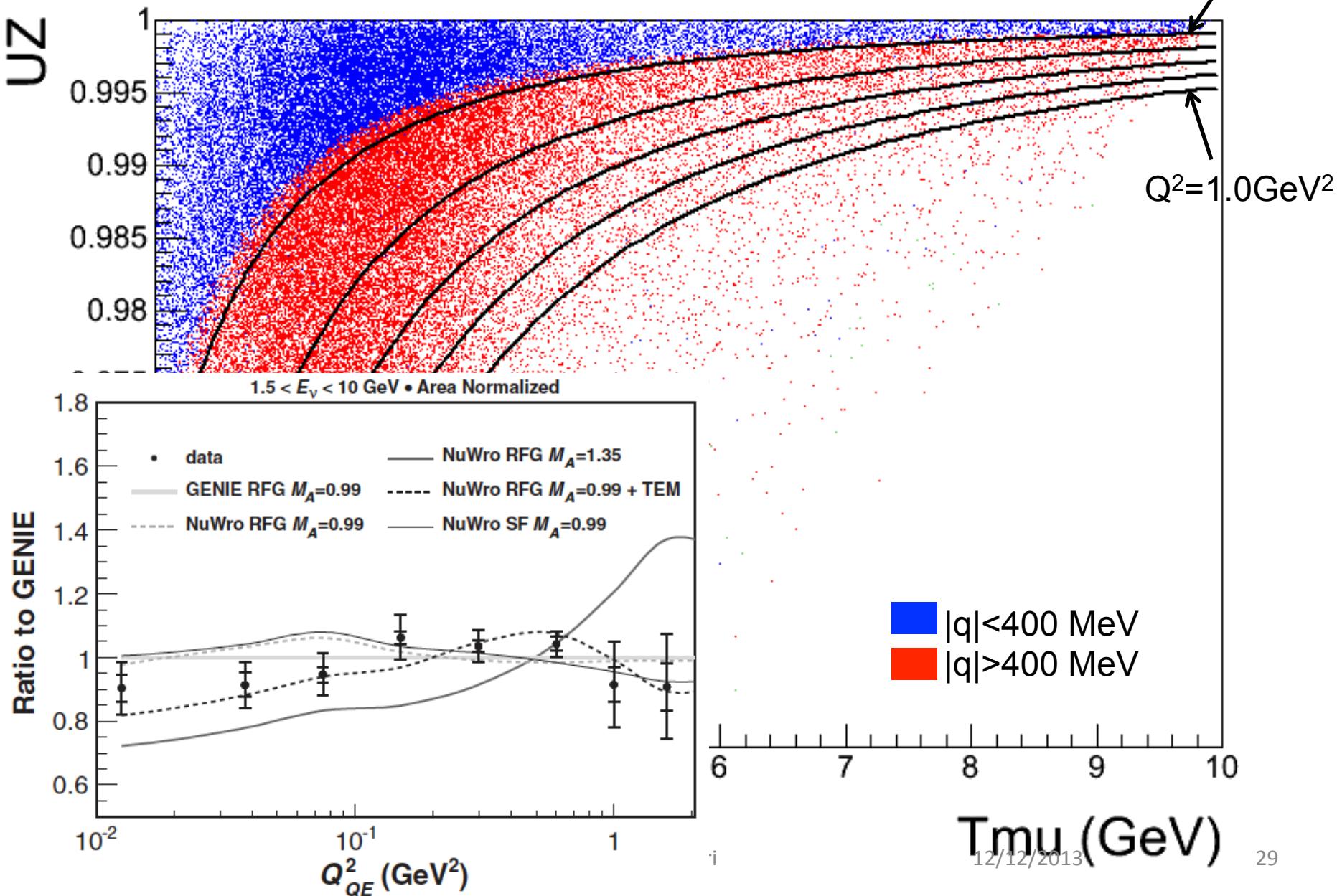
### 3. MiniBooNE phase space study



### 3. MINERvA phase space study



### 3. MINERvA phase space study

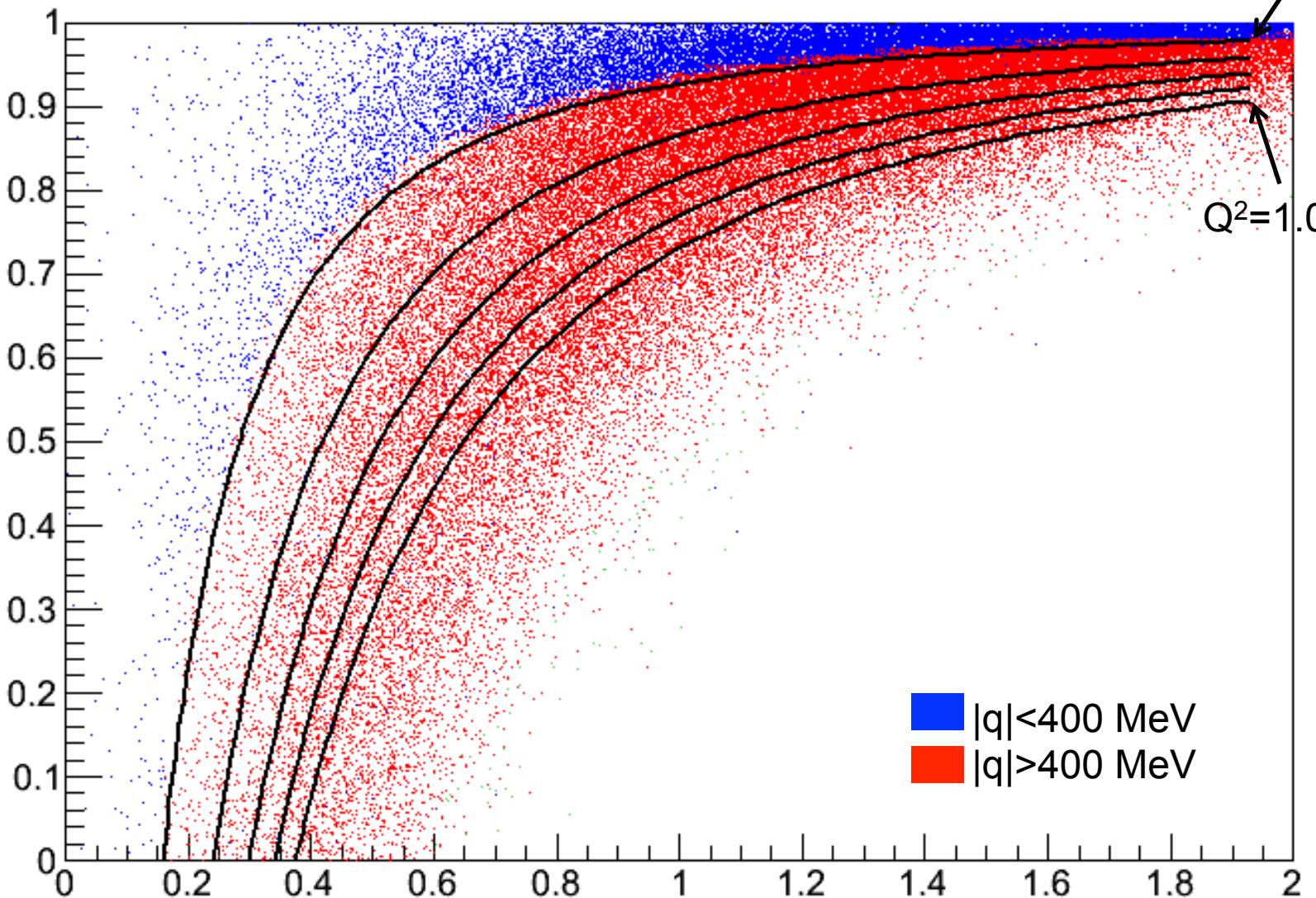


### 3. T2K phase space study

$Q^2=0.2\text{GeV}^2$

$Q^2=1.0\text{GeV}^2$

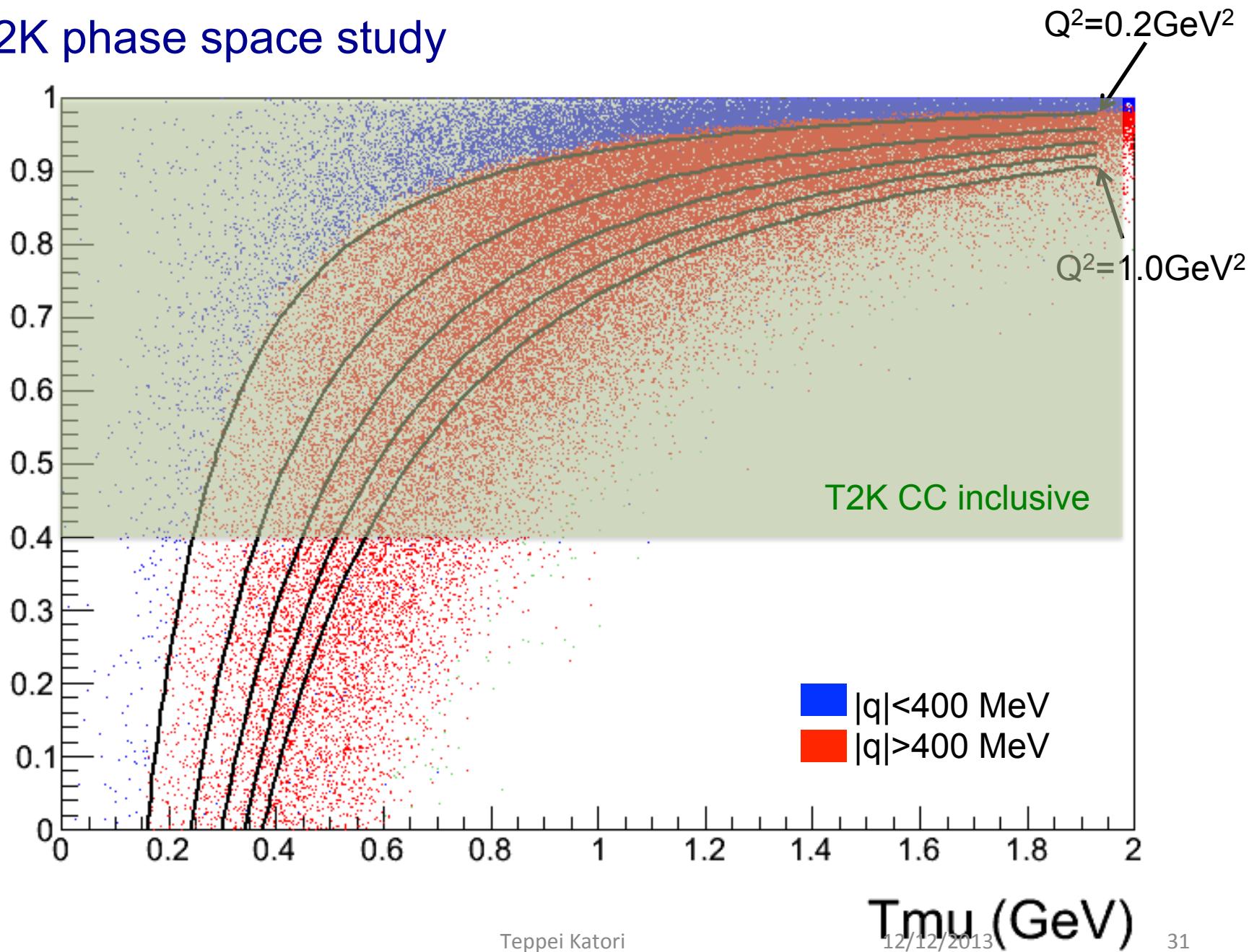
UZ



$|q| < 400 \text{ MeV}$   
 $|q| > 400 \text{ MeV}$

### 3. T2K phase space study

UZ



### 3. T2K phase space study

UZ

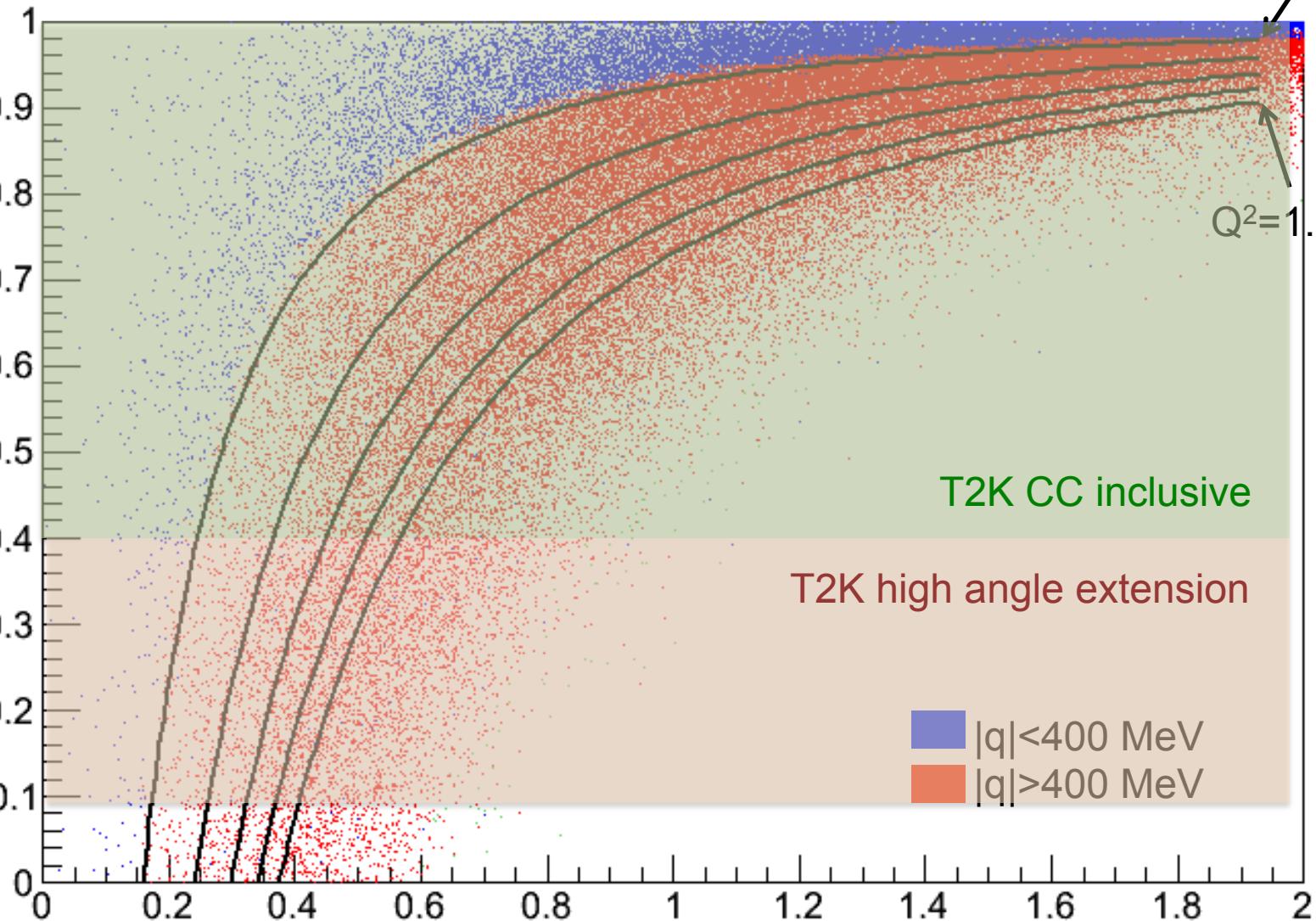
$Q^2=0.2\text{GeV}^2$

$Q^2=1.0\text{GeV}^2$

T2K CC inclusive

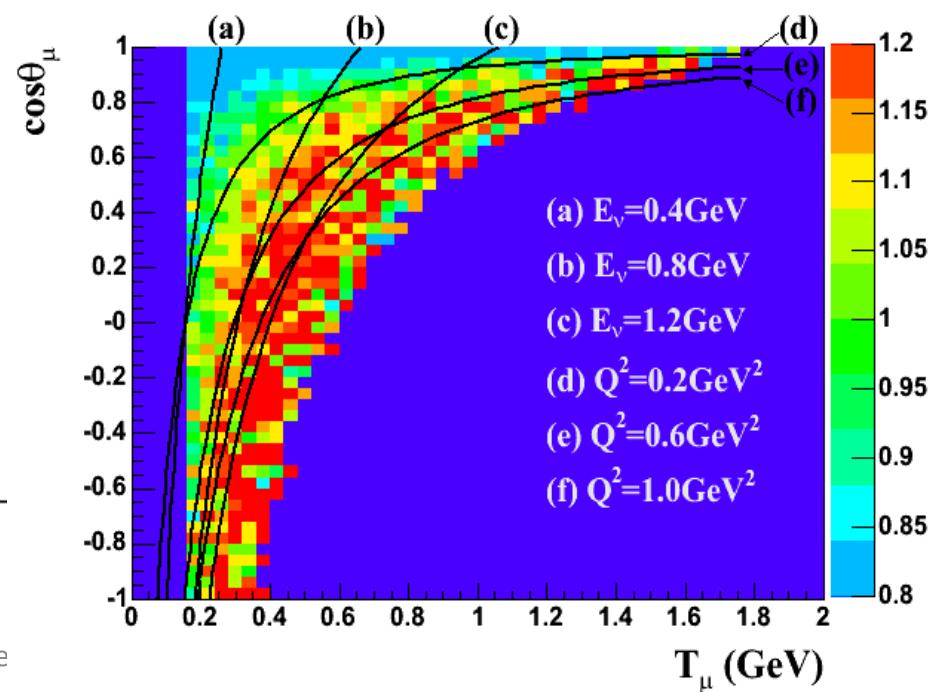
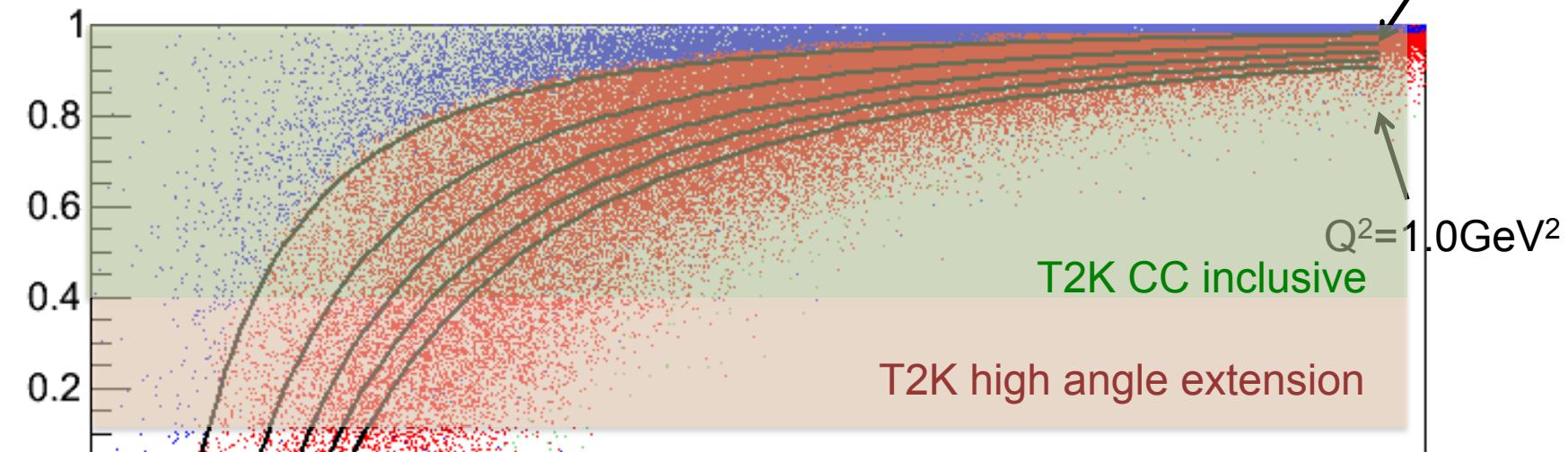
T2K high angle extension

| $q$ |<400 MeV  
| $q$ |>400 MeV

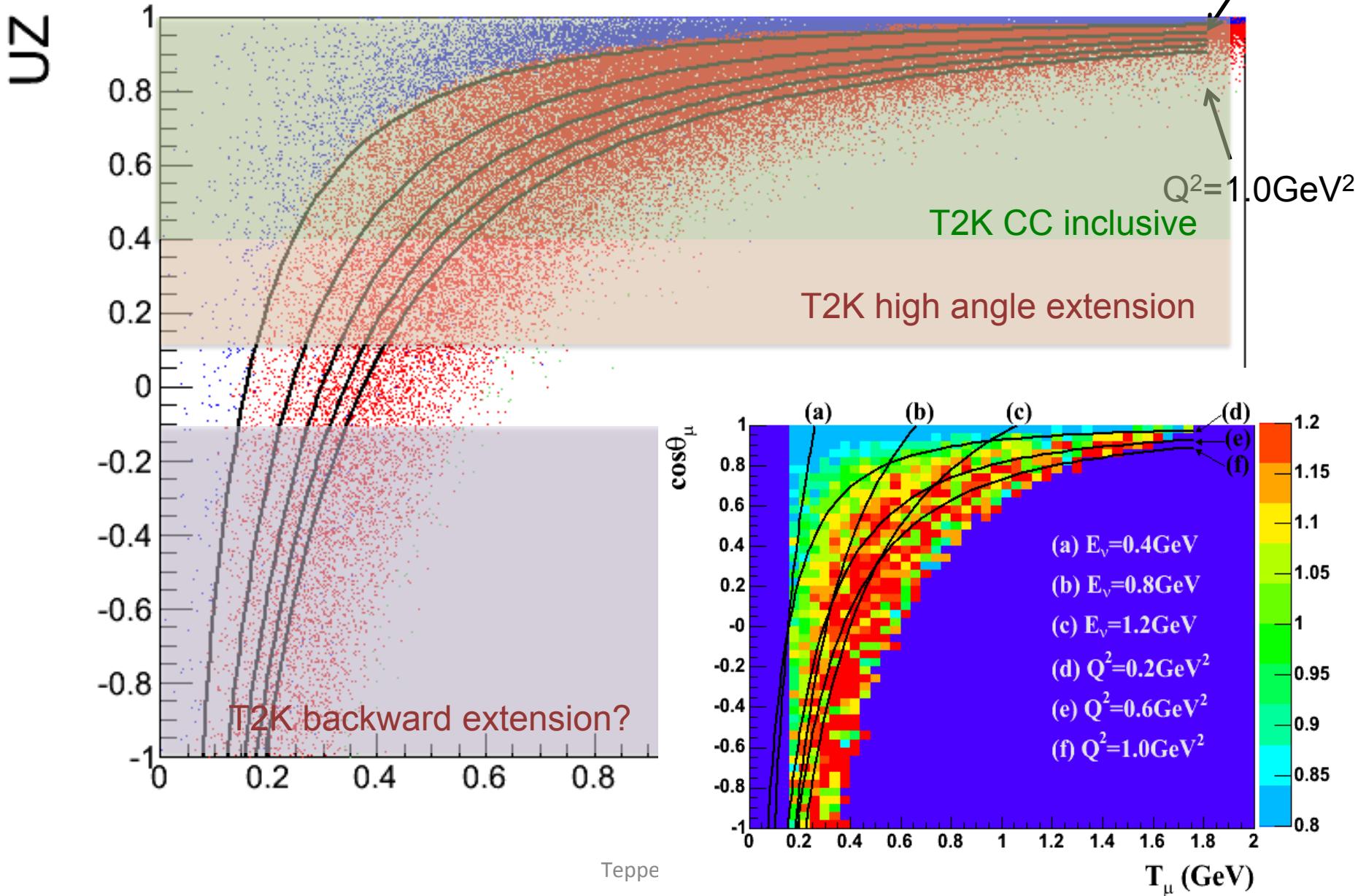


### 3. T2K phase space study

UZ



### 3. T2K phase space study



## 4. Conclusion

Theoretically MB/MINERvA/T2K phase space overlap each other. However Fermi motion and unknown effect can make  $Q^2$  distribution strange. So measuring wider muon kinematic space seems to me the best test to check MB result.

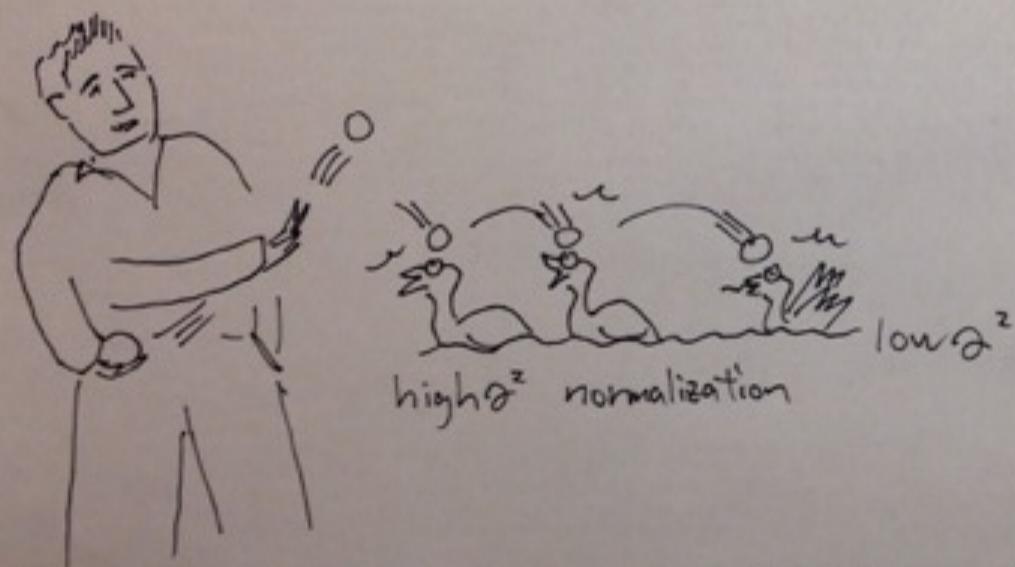
Low  $Q^2$  region ( $Q^2 < 0.2 \text{ GeV}^2$ ) corresponds to low 3 momentum transfer, where impulse approximation breaks down (including SF?).

Theoretical models working in low  $Q^2$  region is essential (Nieves, Martini).

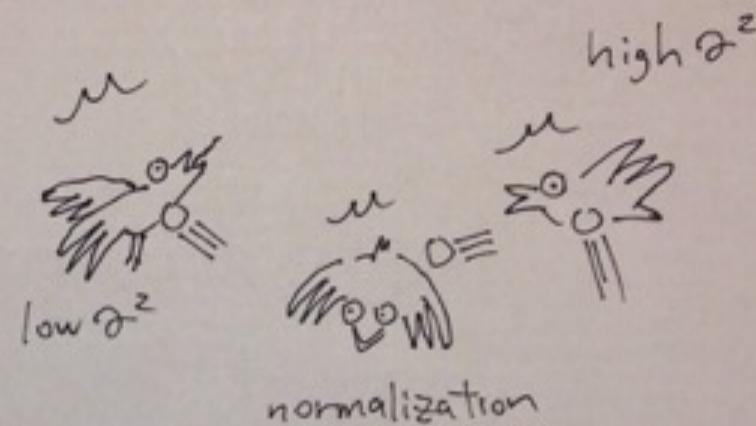
QE+2p-2h+RPA kills three birds with one stone

- 1<sup>st</sup> bird = high Q<sup>2</sup> problem
- 2<sup>nd</sup> bird = normalization
- 3<sup>rd</sup> bird = low Q<sup>2</sup> problem

Juan Nieves



$QE + 2p - 2h + RPA$  kills  
three birds with one stone



Tepper K.  
12/12/13

### 3. MiniBooNE phase space study

#### Violation of Sachs interpretation

- Dipole form factor is motivated by spherically symmetric exponential charge distribution.

$$\rho(r) = \rho_0 \exp(-Mr) \xleftrightarrow{Fourier} G(|q|^2) \sim \frac{1}{\left(1 + \frac{|q|^2}{M^2}\right)^2}.$$

This interpretation works only up to  $|q| < M \sim 1 \text{ GeV}$ .

→ deviation from dipole form factor is necessary at high  $Q^2$ ?

