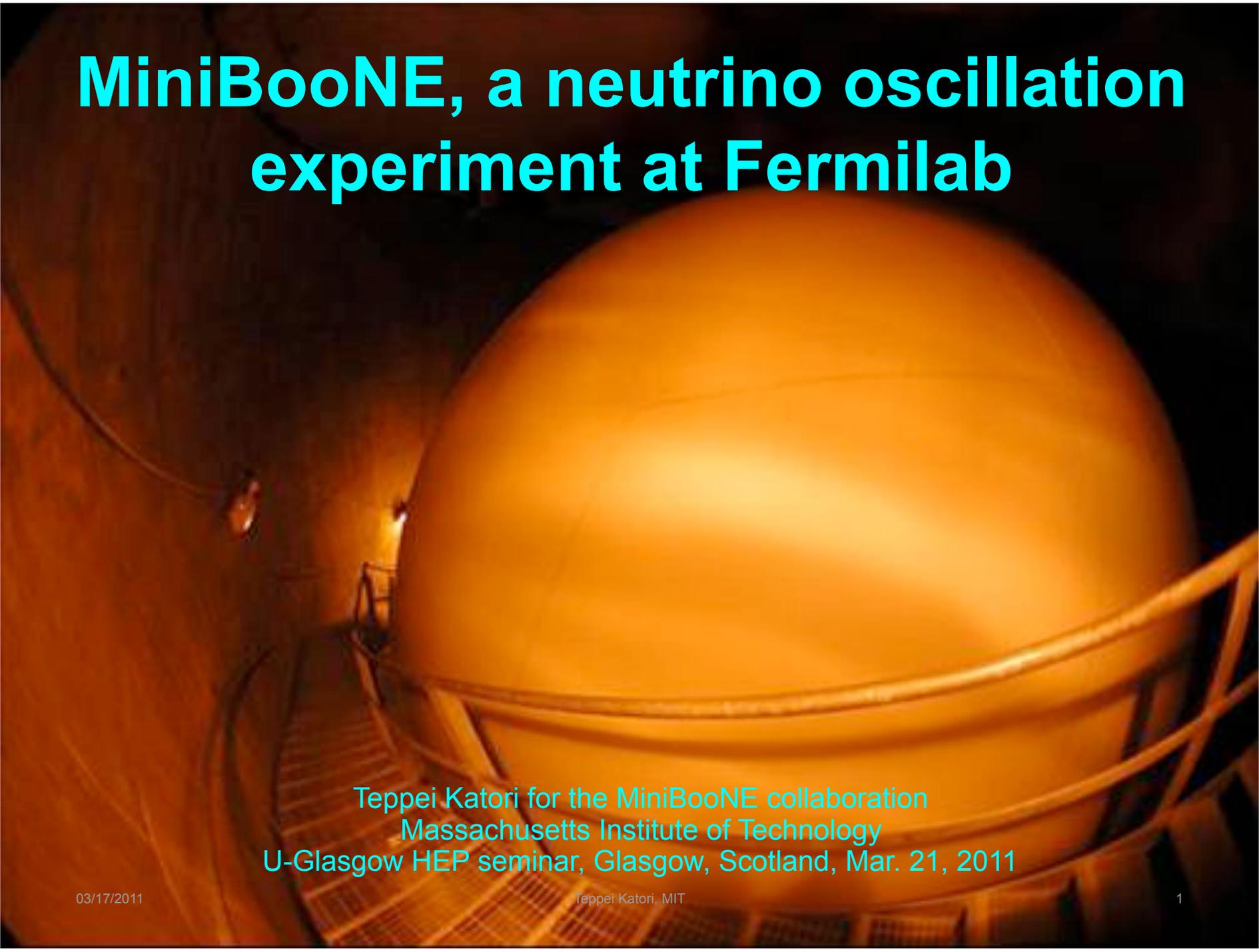


MiniBooNE, a neutrino oscillation experiment at Fermilab



Teppei Katori for the MiniBooNE collaboration
Massachusetts Institute of Technology
U-Glasgow HEP seminar, Glasgow, Scotland, Mar. 21, 2011

MiniBooNE, a neutrino oscillation experiment at Fermilab

Outline

1. Introduction
2. Neutrino beam
3. Events in the detector
4. Cross section model
5. Neutrino oscillation result
6. Anti-neutrino oscillation result
7. Outlook

- 1. Introduction**
- 2. Neutrino beam**
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1. Neutrino oscillation

The neutrino weak eigenstate is described by neutrino Hamiltonian eigenstates, ν_1 , ν_2 , and ν_3 and their mixing matrix elements.

$$|\nu_e\rangle = \sum_{i=1}^3 U_{ei} |\nu_i\rangle$$

The time evolution of neutrino weak eigenstate is written by Hamiltonian mixing matrix elements and eigenvalues of ν_1 , ν_2 , and ν_3 .

$$|\nu_e(t)\rangle = \sum_{i=1}^3 U_{ei} e^{-i\lambda_i t} |\nu_i\rangle$$

Then the transition probability from weak eigenstate ν_μ to ν_e is (no CP violation)

$$P_{\mu \rightarrow e}(t) = \left| \langle \nu_e(t) | \nu_\mu \rangle \right|^2 = -4 \sum_{i>j} (U_{\mu i} U_{\mu j} U_{ei} U_{ej}) \sin^2 \left(\frac{\Delta_{ij}}{2} t \right)$$

1. Neutrino oscillation

In the vacuum, 2 neutrino state effective Hamiltonian has a form,

$$H_{eff} \rightarrow \begin{pmatrix} \frac{m_{ee}^2}{2E} & \frac{m_{e\mu}^2}{2E} \\ \frac{m_{e\mu}^2}{2E} & \frac{m_{\mu\mu}^2}{2E} \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \frac{m_1^2}{2E} & 0 \\ 0 & \frac{m_2^2}{2E} \end{pmatrix} \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix}$$

Therefore, 2 massive neutrino oscillation model is

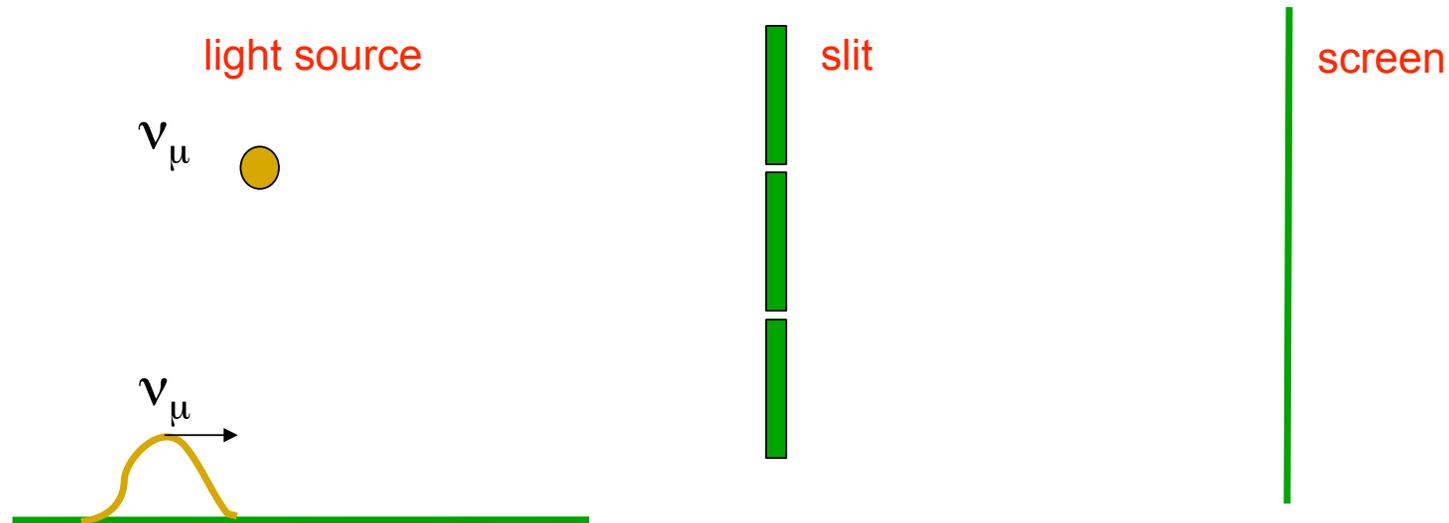
$$P_{\mu \rightarrow e}(t) = \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2}{4E} t \right)$$

Or, conventional form

$$P_{\mu \rightarrow e}(L/E) = \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2 (eV^2) \frac{L(m)}{E(MeV)} \right)$$

1. Neutrino oscillation

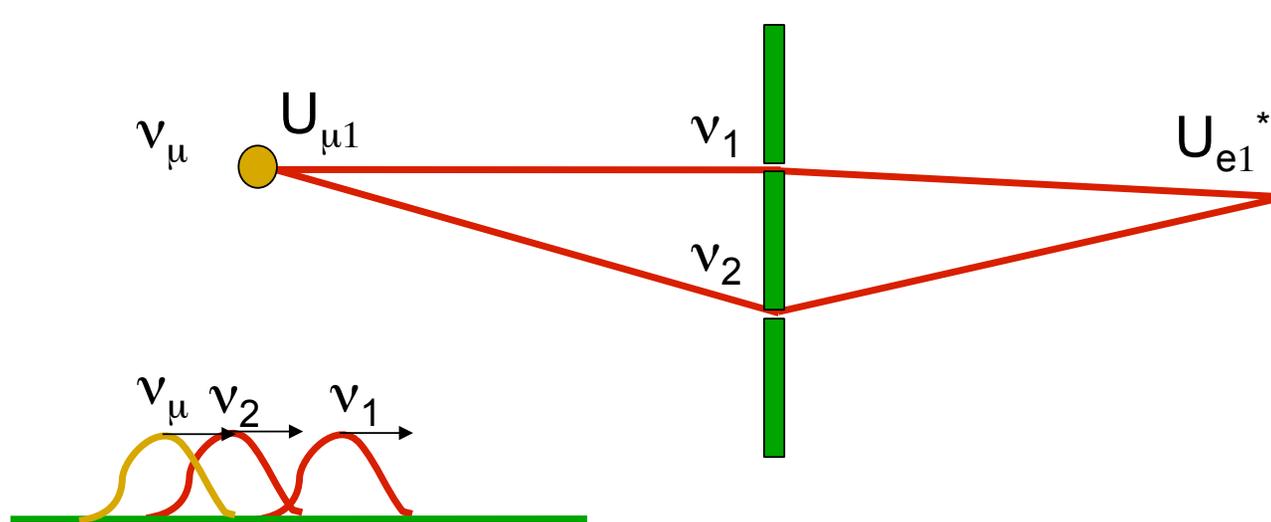
Neutrino oscillation is an interference experiment (cf. double slit experiment)



If 2 light paths have different length, they have different phase rotation, and they make interference pattern on the screen.

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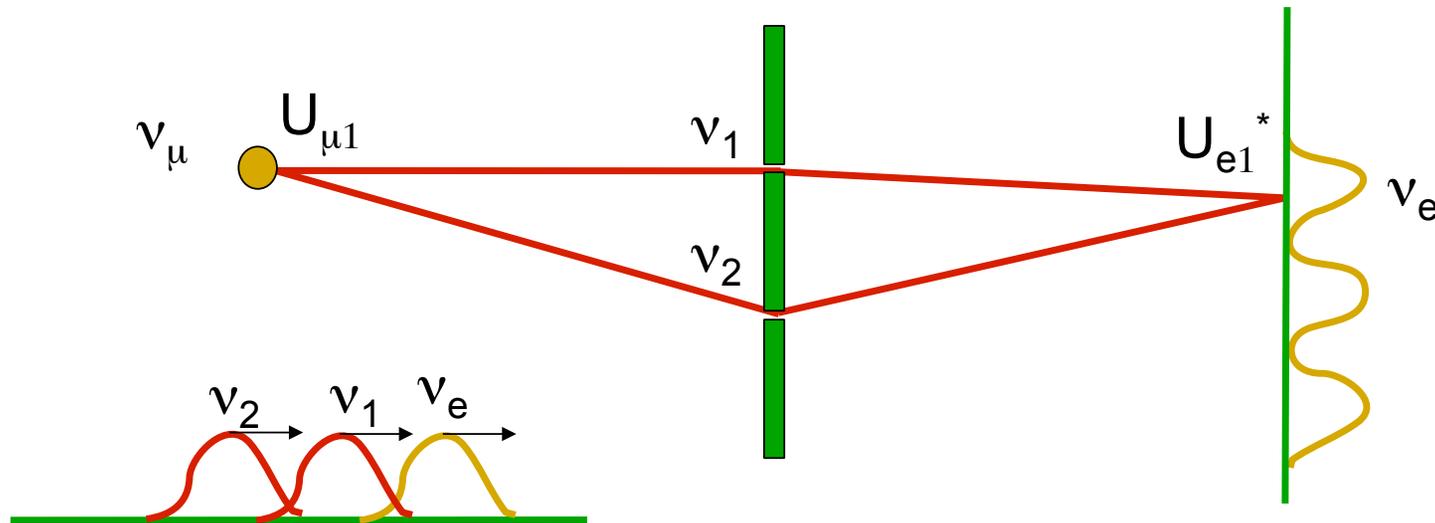


If 2 light paths have different length, they have different phase rotation, and they make interference pattern on the screen.

For massive neutrino model, if ν_2 is heavier than ν_1 , they have different group velocities hence different phase rotation, thus the superposition of those 2 wave packet no longer makes same state

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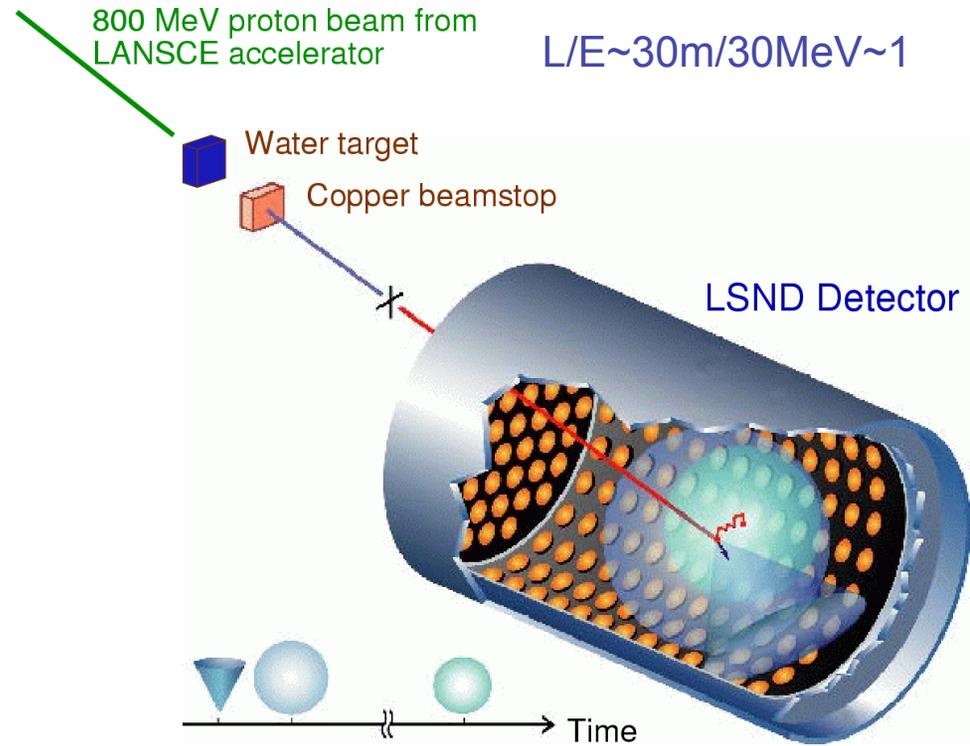
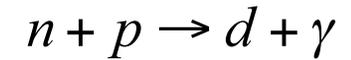
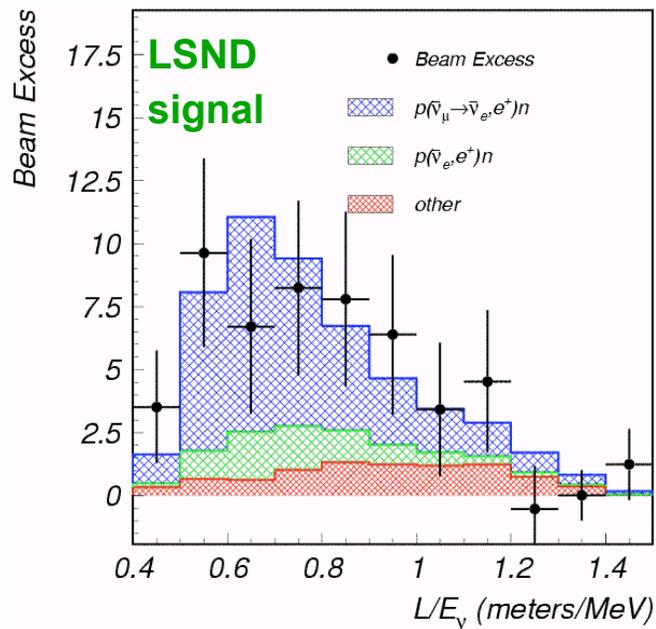
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For massive neutrino model, if ν_2 is heavier than ν_1 , they have different group velocities hence different phase rotation, thus the superposition of those 2 wave packet no longer makes same state

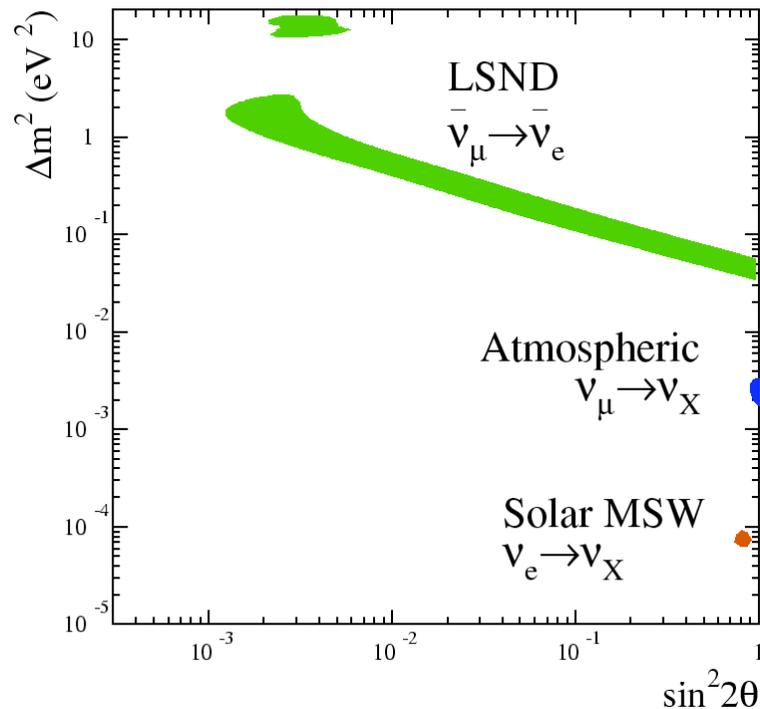
1. LSND experiment

LSND experiment at Los Alamos observed excess of anti-electron neutrino events in the anti-muon neutrino beam.

$87.9 \pm 22.4 \pm 6.0$ ($3.8.\sigma$)



1. LSND experiment

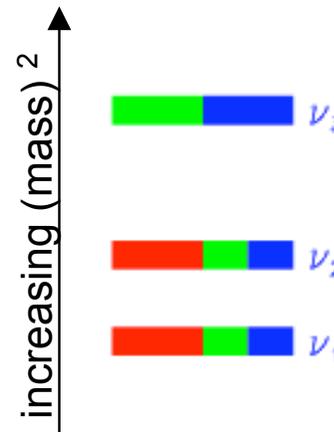


3 types of neutrino oscillations are found:

- LSND neutrino oscillation: $\Delta m^2 \sim 1 \text{eV}^2$
- Atmospheric neutrino oscillation: $\Delta m^2 \sim 10^{-3} \text{eV}^2$
- Solar neutrino oscillation : $\Delta m^2 \sim 10^{-5} \text{eV}^2$

But we cannot have so many Δm^2 !

$$\Delta m_{13}^2 \neq \Delta m_{12}^2 + \Delta m_{23}^2$$



We need to test LSND signal

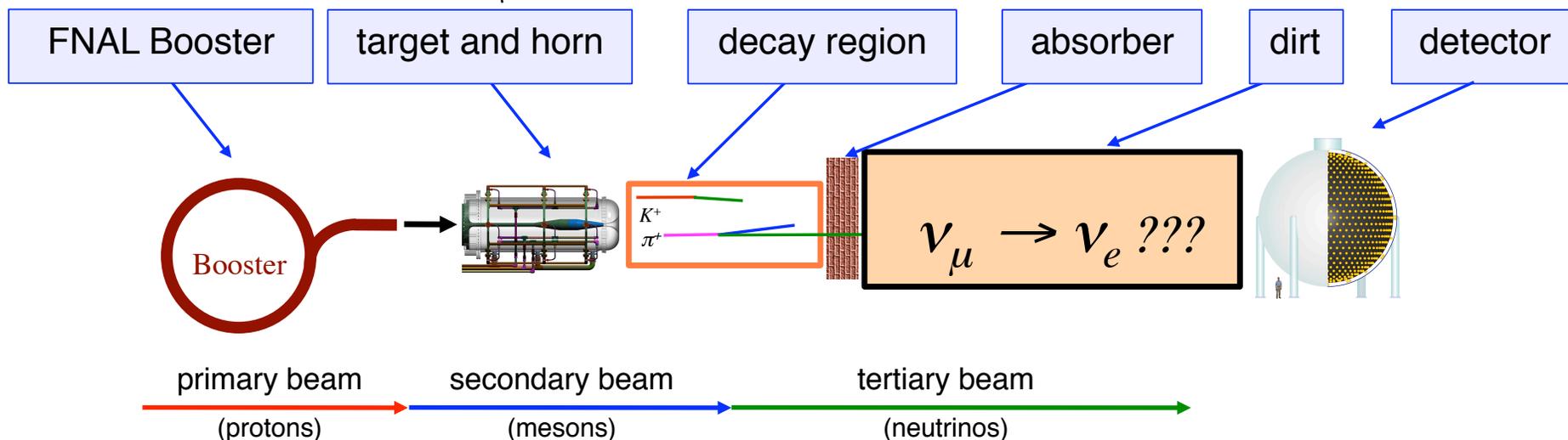
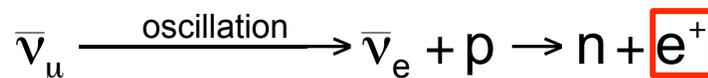
MiniBooNE experiment is designed to have same $L/E \sim 500\text{m}/500\text{MeV} \sim 1$ to test LSND $\Delta m^2 \sim 1 \text{eV}^2$

1. MiniBooNE experiment

Keep L/E same with LSND, while changing systematics, energy & event signature;

$$P(\nu_\mu \rightarrow \nu_e) = \sin^2 2\theta \sin^2(1.27 \Delta m^2 L/E)$$

MiniBooNE is looking for **the single isolated electron like events**, which is the signature of ν_e events



MiniBooNE has;

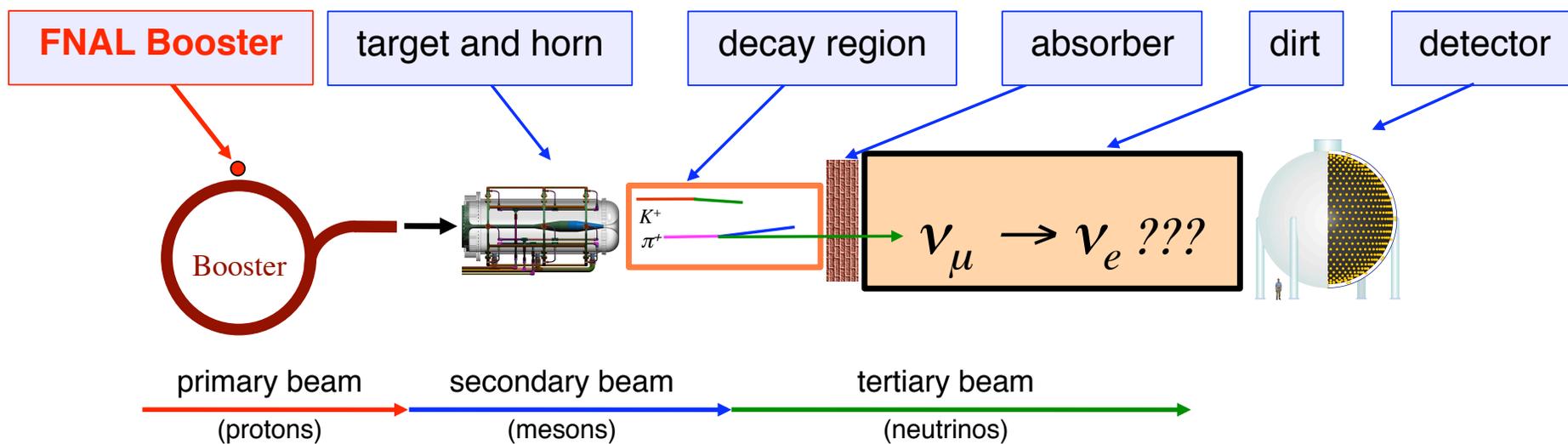
- higher energy (~500 MeV) than LSND (~30 MeV)
- longer baseline (~500 m) than LSND (~30 m)

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2. Neutrino beam

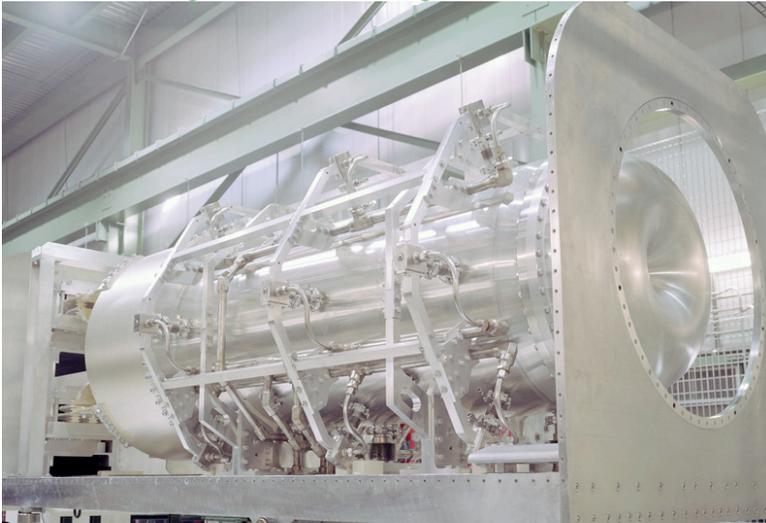


MiniBooNE extracts beam
from the 8 GeV Booster



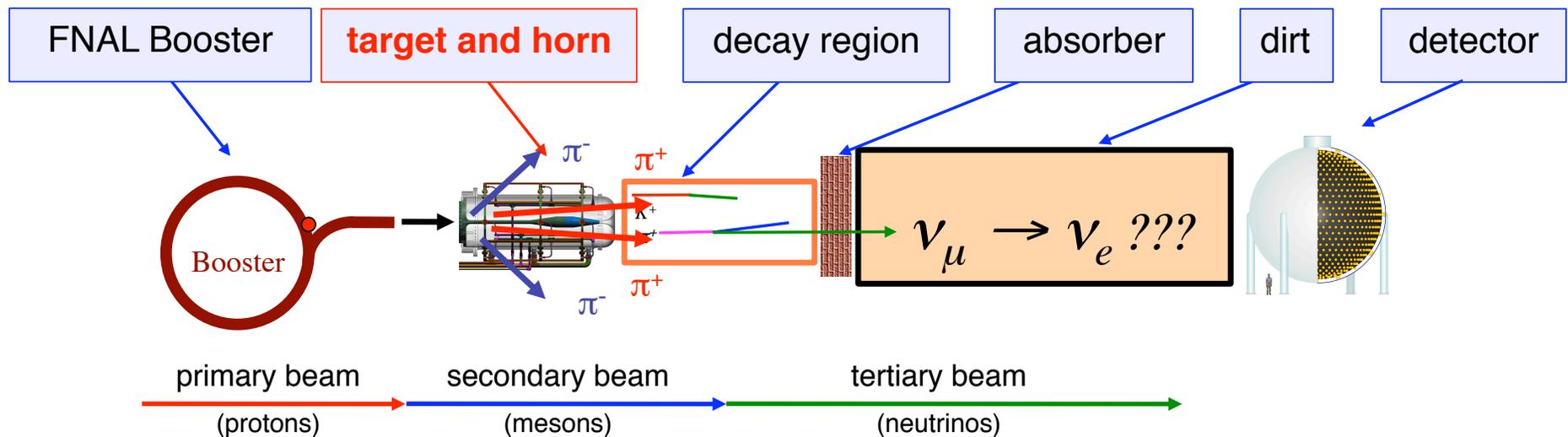
2. Neutrino beam

Magnetic focusing horn



8GeV protons are delivered to
a 1.7λ Be target

within a magnetic horn
(2.5 kV, 174 kA) that
increases the flux by $\times 6$



2. Neutrino beam

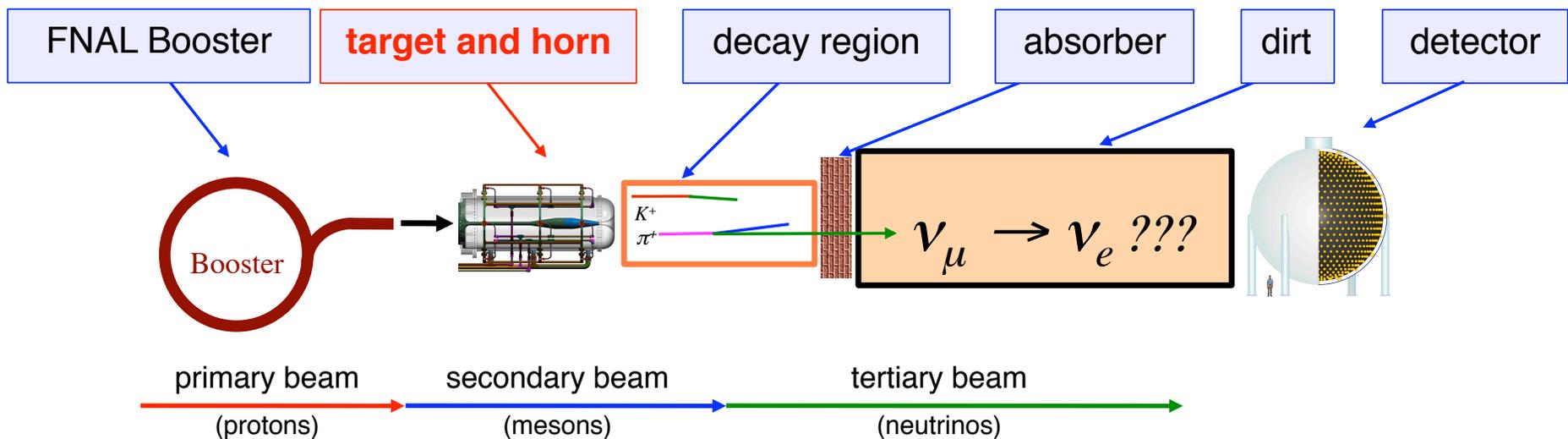
HARP experiment (CERN)



Modeling of meson production is based on the measurement done by HARP collaboration.

- Identical, but 5% λ Beryllium target
- 8.9 GeV/c proton beam momentum

HARP collaboration,
Eur.Phys.J.C52(2007)29



2. Neutrino beam

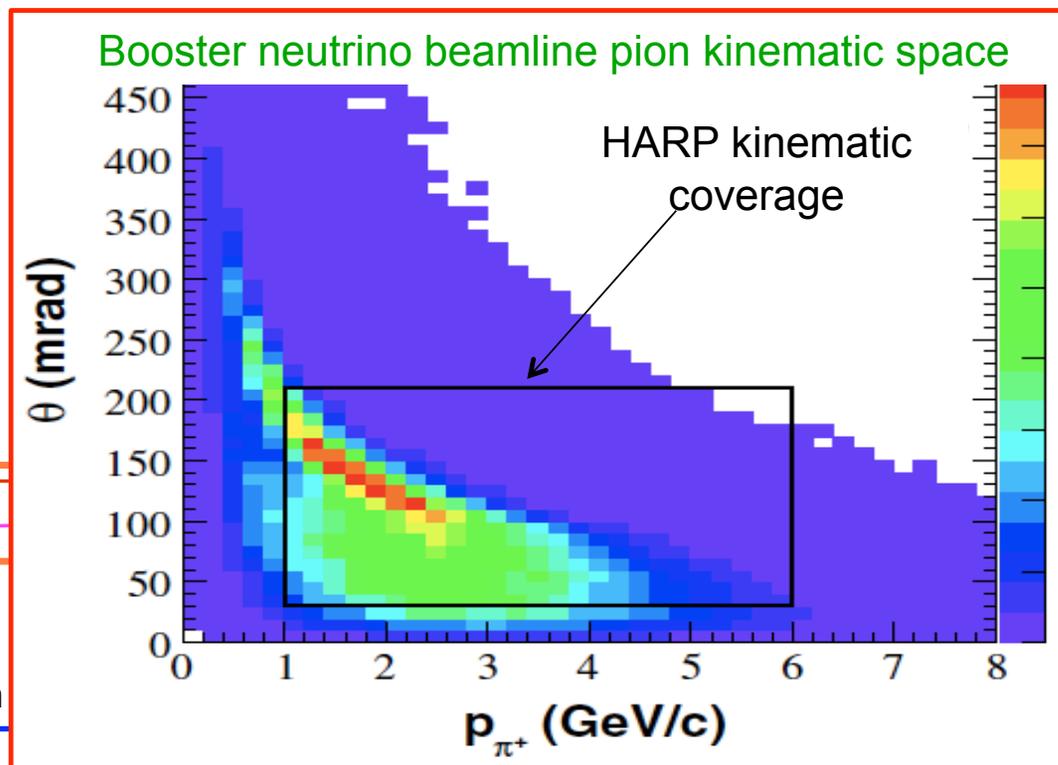
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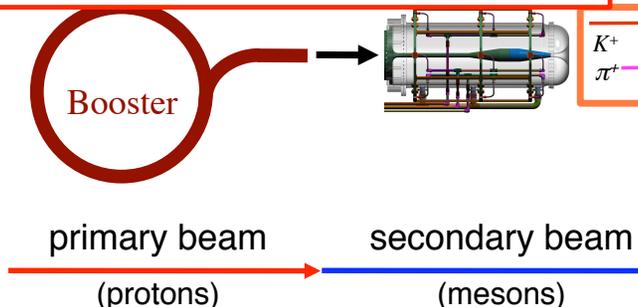
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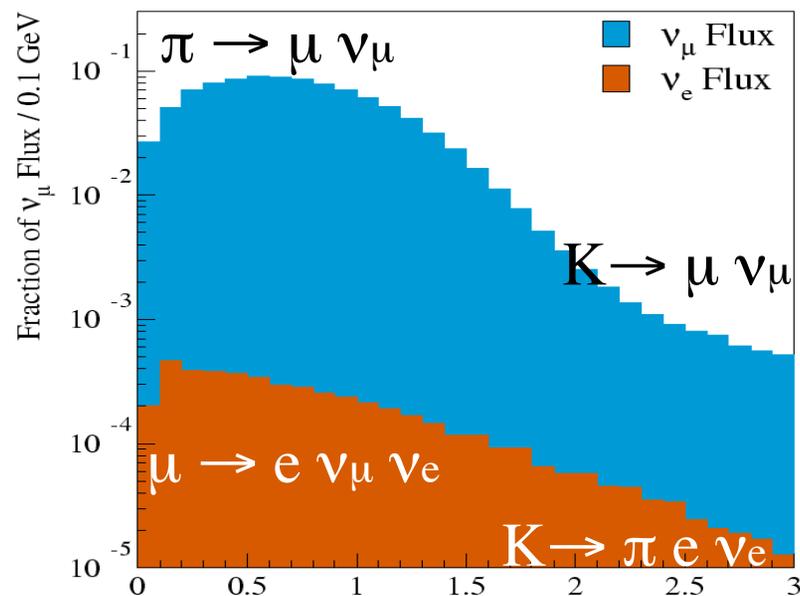
- Identical, but 5% λ Beryllium target
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Majority of pions create neutrinos in MiniBooNE are directly measured by HARP (>80%)



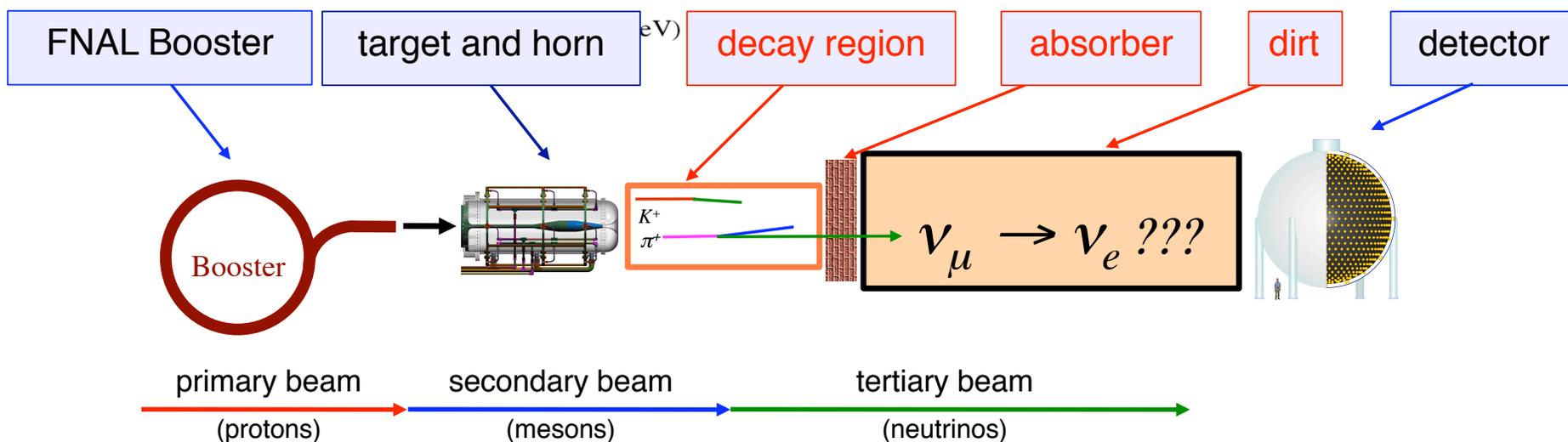
2. Neutrino beam



Neutrino flux from simulation by GEANT4

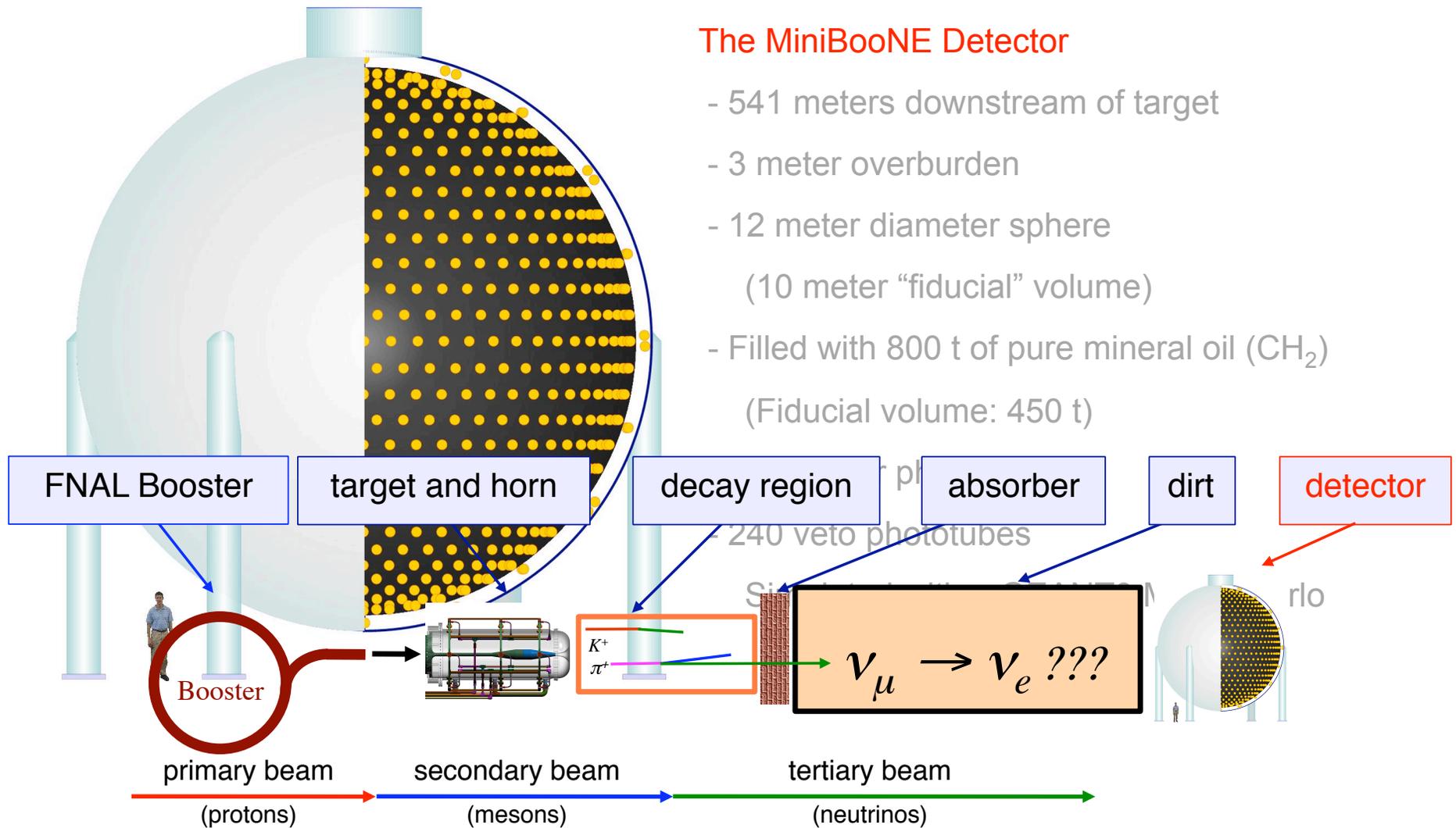
MiniBooNE is the ν_e (anti ν_e) appearance oscillation experiment, so we need to know the distribution of beam origin ν_e and anti ν_e (intrinsic ν_e)

	neutrino mode	antineutrino mode
intrinsic ν_e contamination	0.6%	0.6%
intrinsic ν_e from μ decay	49%	55%
intrinsic ν_e from K decay	47%	41%
others	4%	4%
wrong sign fraction	6%	16%



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3. Events in the Detector



03/17/2011

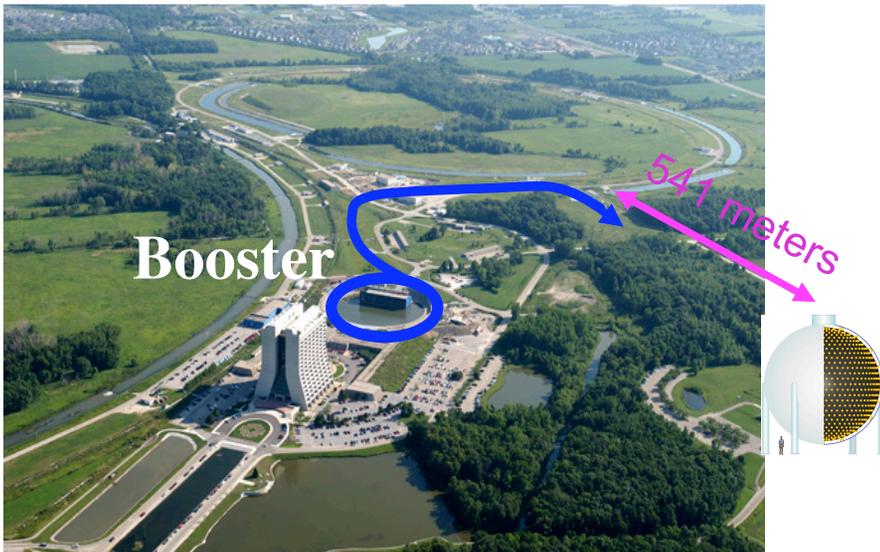
Teppei Katori, MIT

3. Events in the Detector

The MiniBooNE Detector

- 541 meters downstream of target
- 3 meter overburden
- 12 meter diameter sphere
(10 meter “fiducial” volume)
- Filled with 800 t of pure mineral oil (CH_2)
(Fiducial volume: 450 t)
- 1280 inner phototubes,
- 240 veto phototubes

Simulated with a GEANT3 Monte Carlo



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• **Muons**

3. Events in the Detector

MiniBooNE collaboration,
NIM.A599(2009)28

– **Sharp, clear rings**

• **Long, straight tracks**

• **Electrons**

– Scattered rings

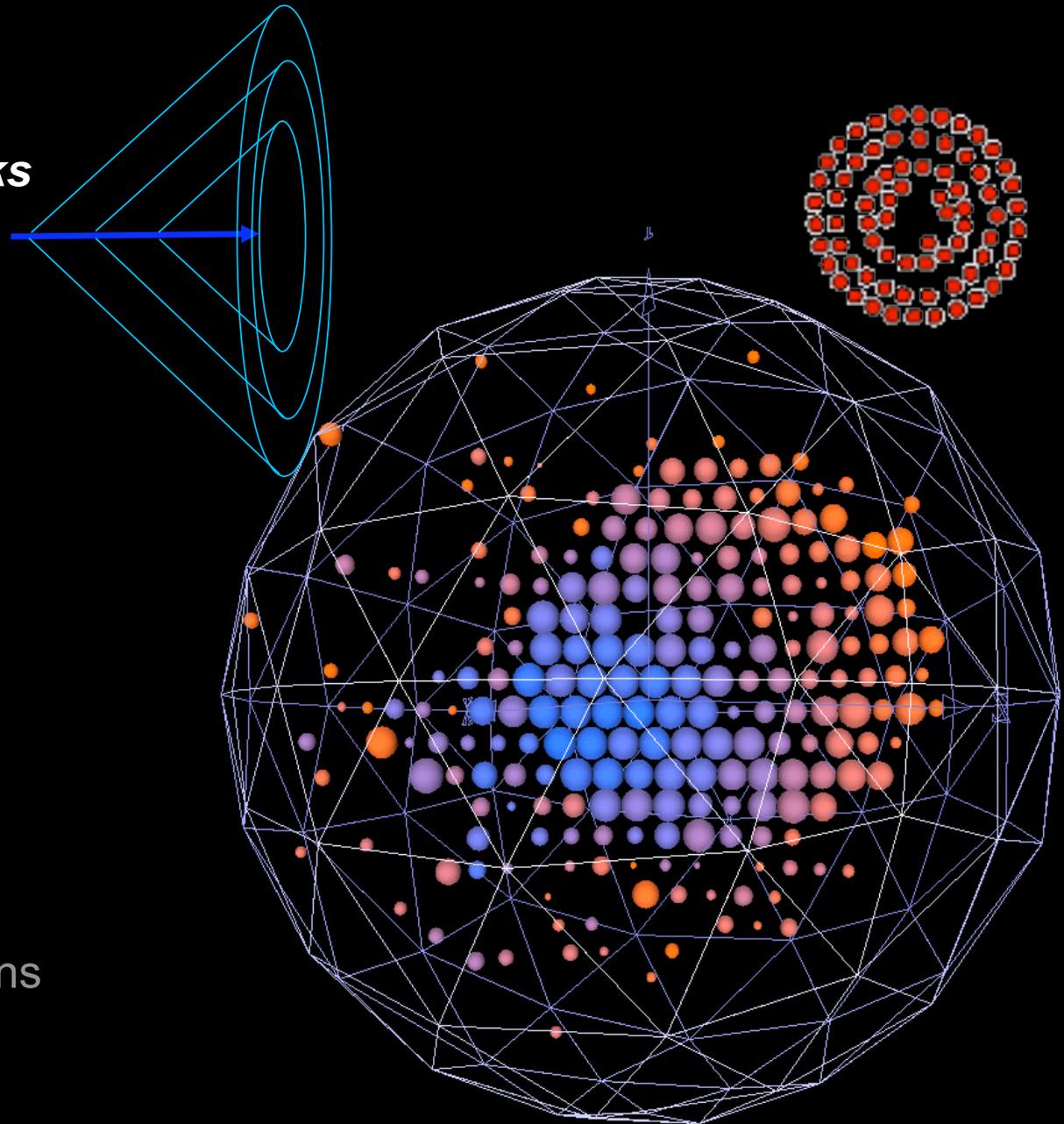
• Multiple scattering

• Radiative processes

• **Neutral Pions**

– Double rings

• Decays to two photons



•Muons

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MiniBooNE collaboration,
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• Long, straight tracks

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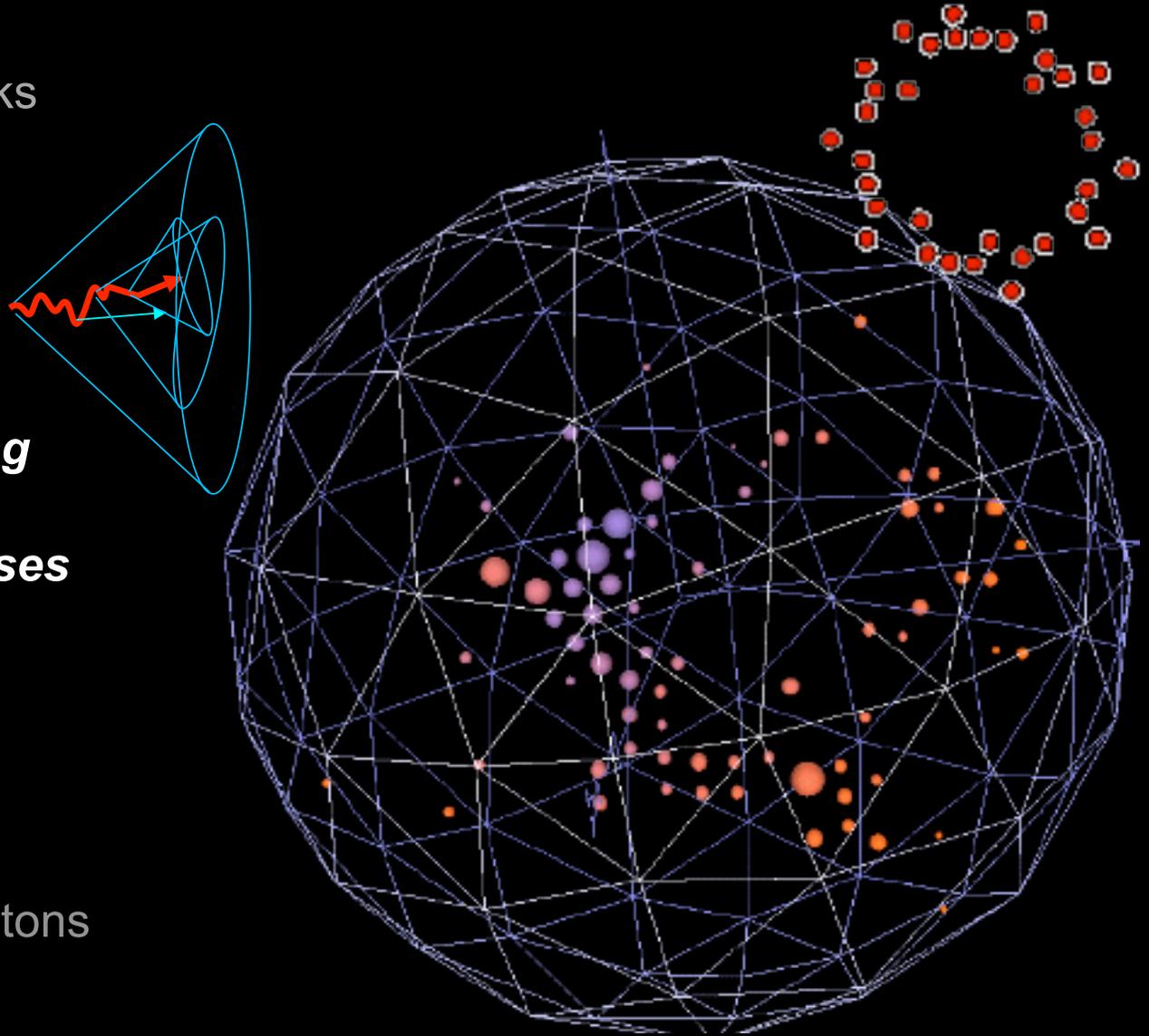
• *Multiple scattering*

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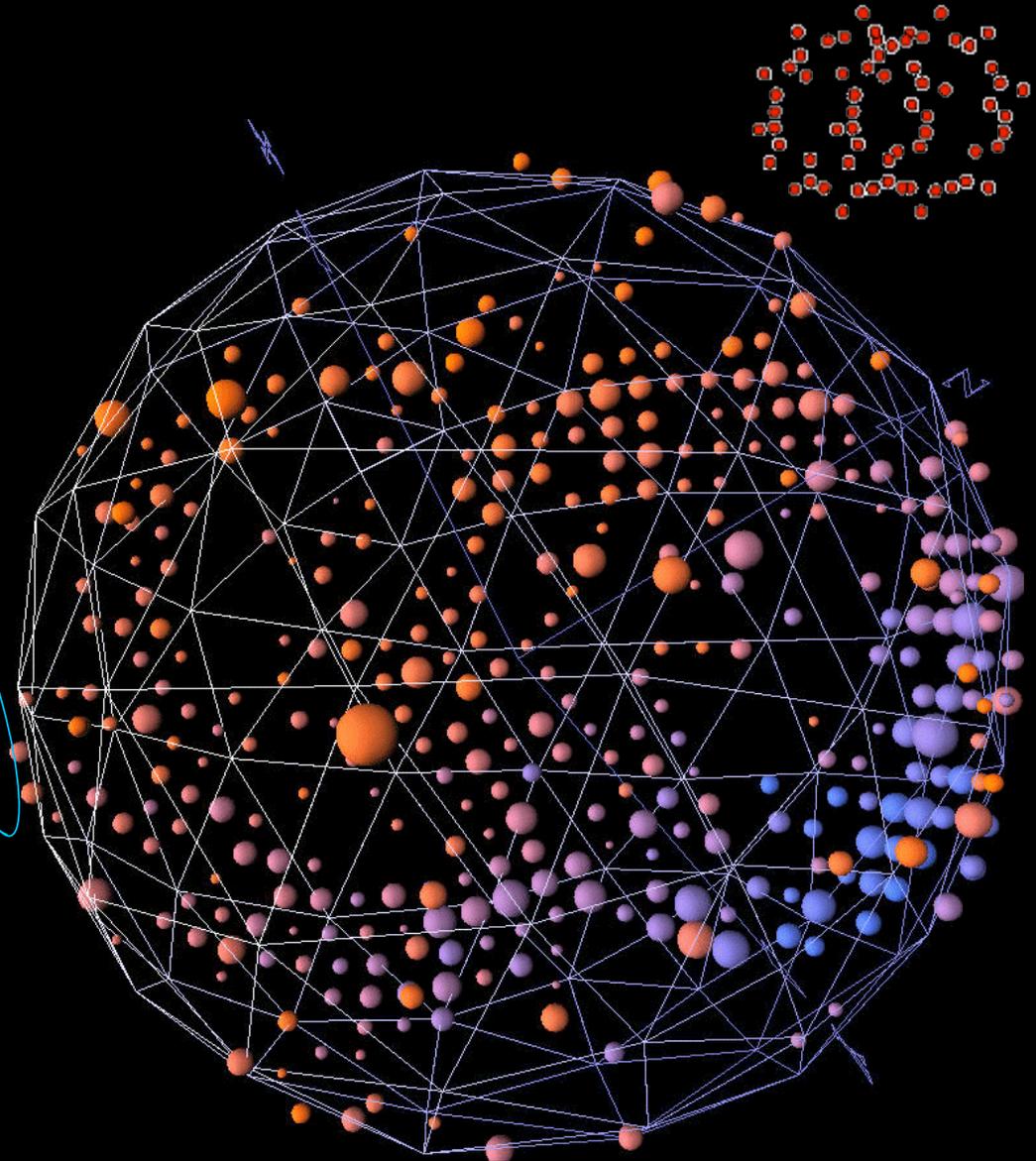
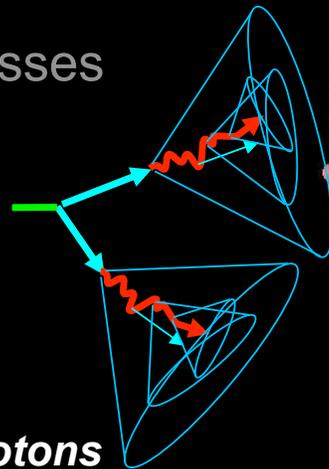
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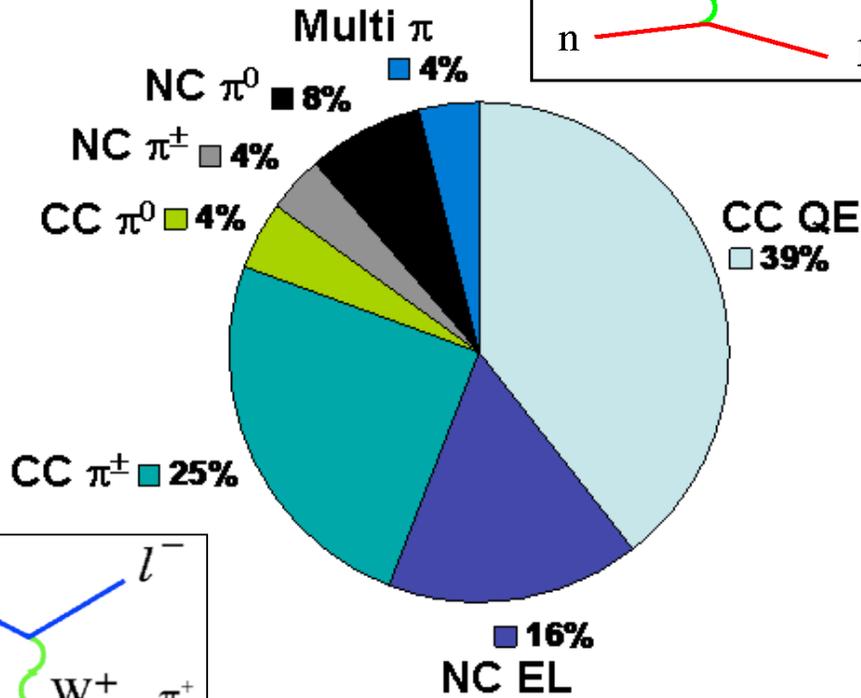
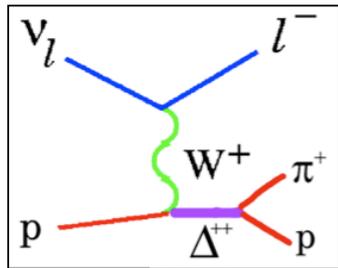
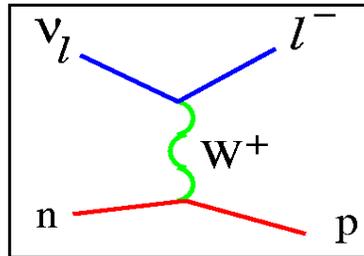
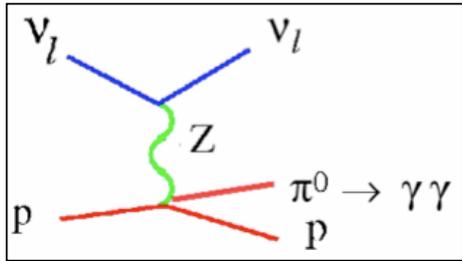
– **Double rings**

• **Decays to two photons**



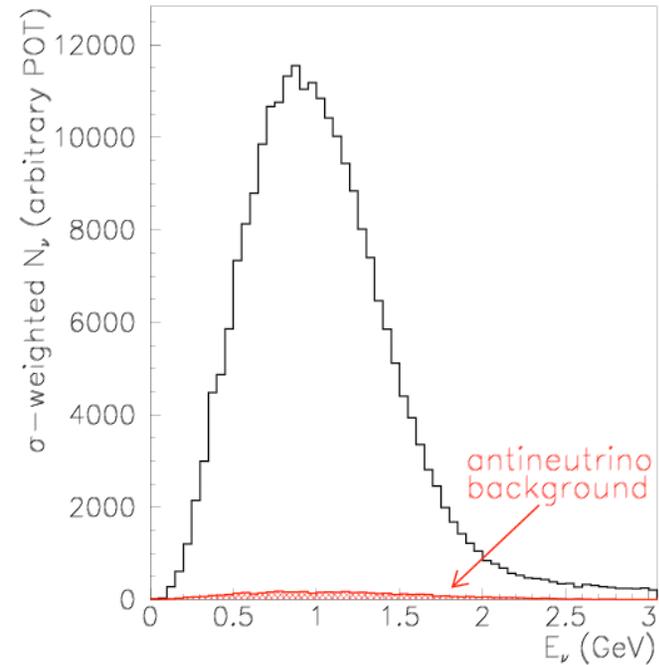
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4. Cross section model



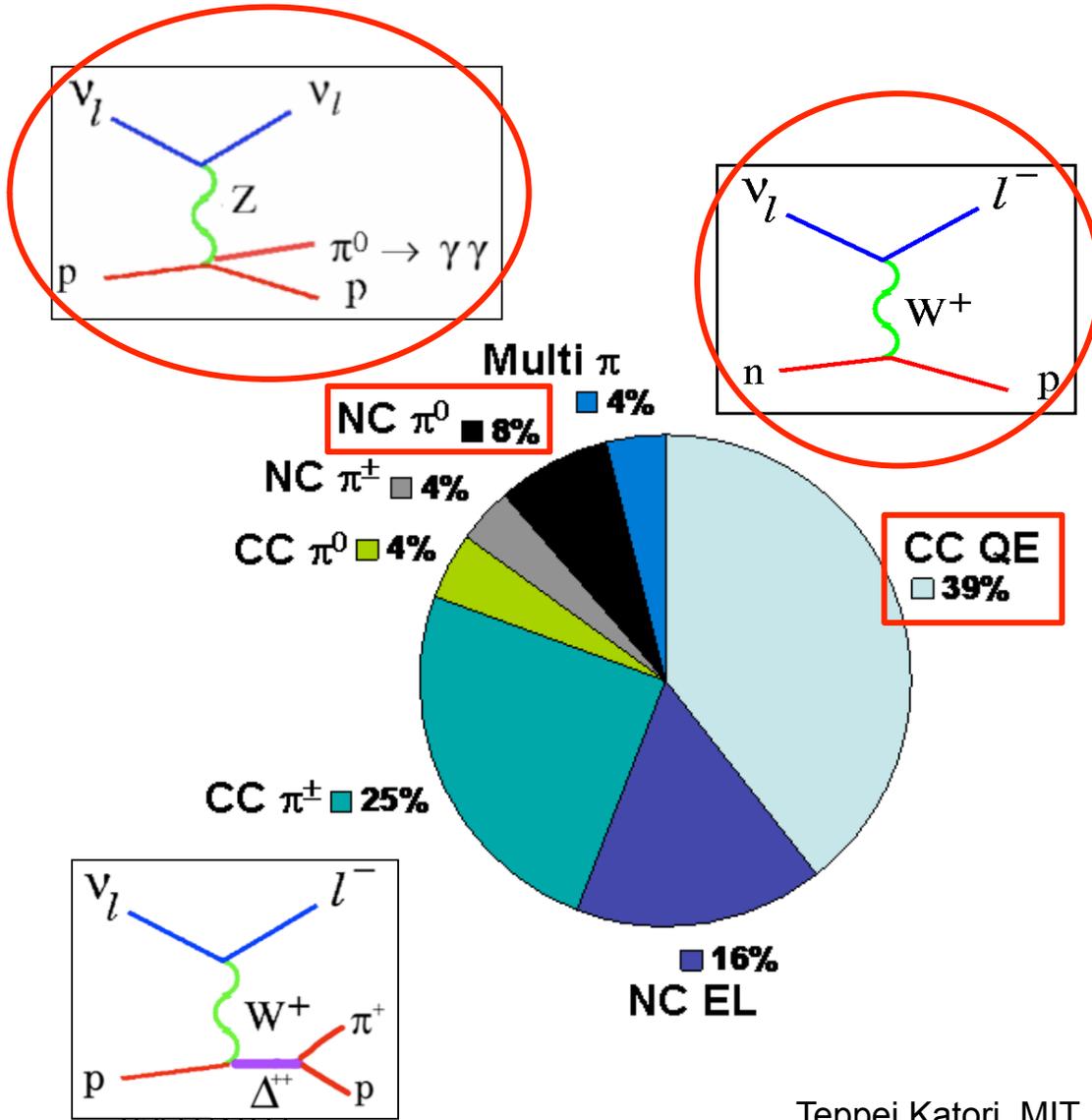
Predicted event rates before cuts
(NUANCE Monte Carlo)

Casper, Nucl.Phys.Proc.Suppl.112(2002)161



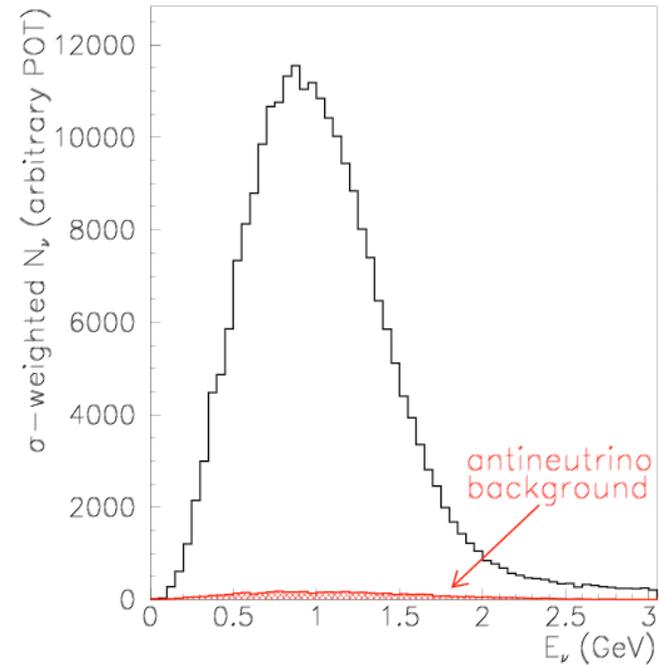
Event neutrino energy (GeV)

4. Cross section model



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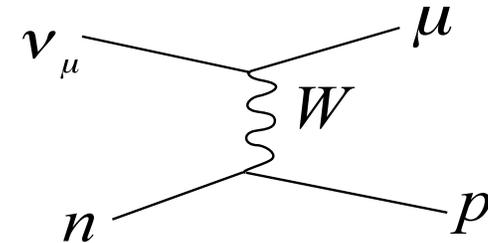
4. CCQE event measurement

CCQE (Charged Current Quasi-Elastic) event

ν_μ charged current quasi-elastic (ν_μ CCQE) interaction is the most abundant (~40%) and the fundamental interaction in MiniBooNE detector

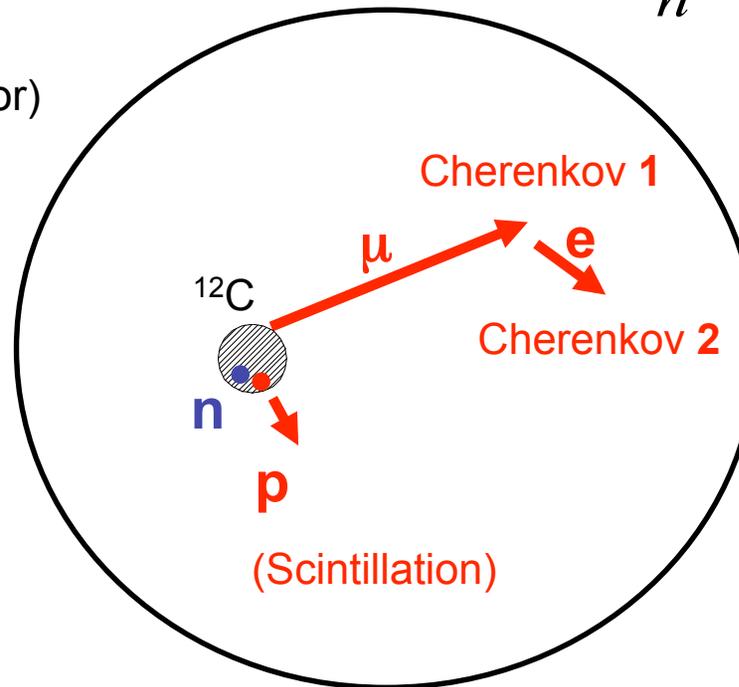
$$\nu_\mu + n \rightarrow p + \mu^-$$

$$(\nu_\mu + {}^{12}\text{C} \rightarrow X + \mu^-)$$



MiniBooNE detector
(spherical Cherenkov detector)

ν -beam



muon like Cherenkov light and subsequent decayed electron (Michel electron) like Cherenkov light are the signal of CCQE event

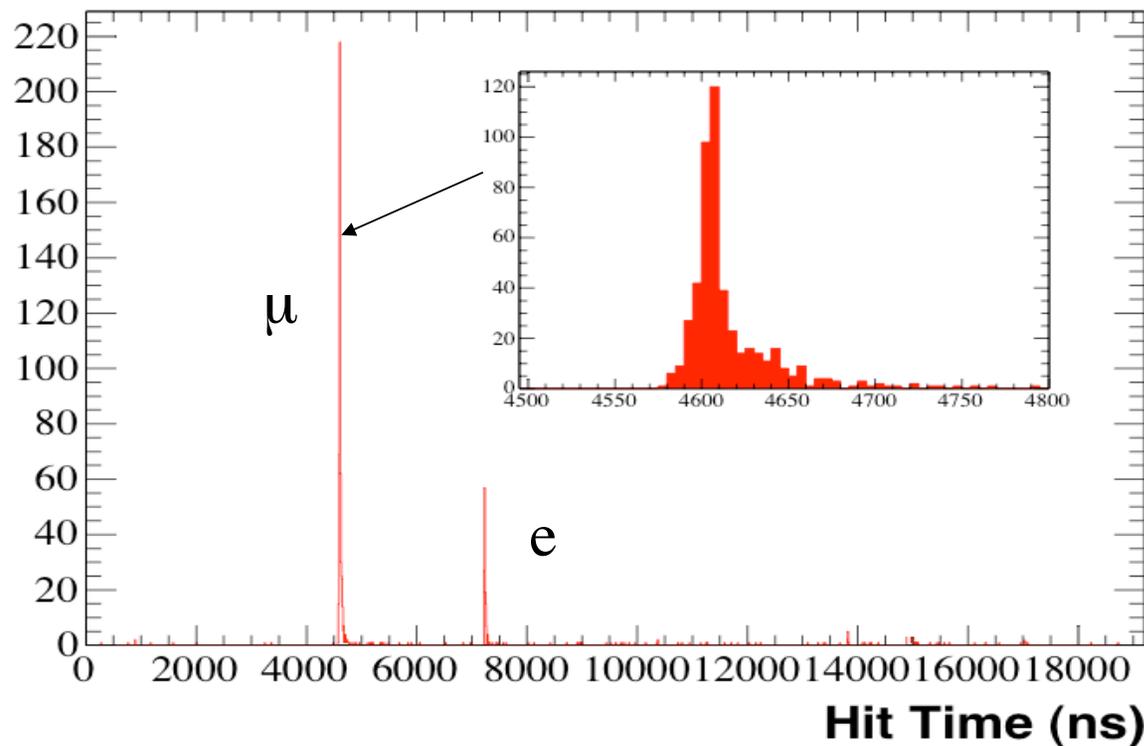
4. CCQE event measurement

19.2 μs beam trigger window with the 1.6 μs spill

Multiple hits within a ~ 100 ns window form “subevents”

ν_μ CCQE interactions ($\nu+n \rightarrow \mu+p$) with characteristic two “subevent” structure from stopped $\mu \rightarrow \nu_\mu \nu_e e$

Number of tank hits for CCQE event



4. CCQE event measurement

All kinematics are specified from 2 observables, muon energy E_μ and muon scattering angle θ_μ

Energy of the neutrino E_ν^{QE} and 4-momentum transfer Q_2^{QE} can be reconstructed by these 2 observables, under the assumption of CCQE interaction with bound neutron at rest (“QE assumption”). CCQE is the signal channel of ν_e candidate.

$$E_\nu^{\text{QE}} = \frac{2(M - E_B)E_\mu - (E_B^2 - 2ME_B + m_\mu^2 + \Delta M^2)}{2[(M - E_B) - E_\mu + p_\mu \cos \theta_\mu]}$$

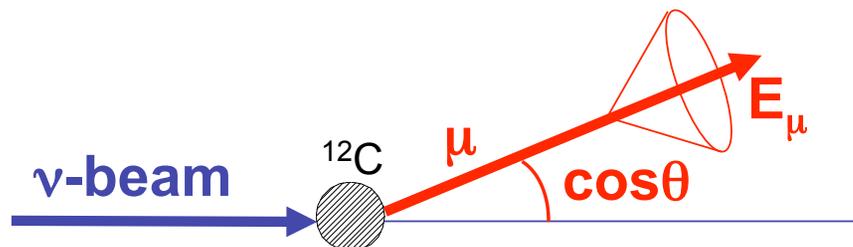
$$Q_{\text{QE}}^2 = -m_\mu^2 + 2E_\nu^{\text{QE}}(E_\mu - p_\mu \cos \theta_\mu)$$

$$\nu_\mu + n \rightarrow p + \mu^-$$

$$(\nu_\mu + {}^{12}\text{C} \rightarrow X + \mu^-)$$

$$\nu_e + n \rightarrow p + e^-$$

$$(\nu_e + {}^{12}\text{C} \rightarrow X + e^-)$$



4. Relativistic Fermi Gas (RFG) model

Relativistic Fermi Gas (RFG) Model

Carbon is described by the collection of incoherent Fermi gas particles. All details come from hadronic tensor.

$$(W_{\mu\nu})_{ab} = \int_{E_{lo}}^{E_{hi}} f(\vec{k}, \vec{q}, w) T_{\mu\nu} dE : \text{hadronic tensor}$$

$f(\vec{k}, \vec{q}, w)$: nucleon phase space density function

$T_{\mu\nu} = T_{\mu\nu}(F_1, F_2, F_A, F_P)$: nucleon tensor

$F_A(Q^2) = g_A / (1 + Q^2 / M_A^2)^2$: Axial form factor

E_{hi} : the highest energy state of nucleon = $\sqrt{(p_F^2 + M^2)}$

E_{lo} : the lowest energy state of nucleon = $\kappa \left(\sqrt{(p_F^2 + M^2)} - \omega + E_B \right)$

We tuned following 2 parameters using Q^2 distribution by least χ^2 fit;

M_A = effective axial mass

κ = Pauli blocking parameter

4. CCQE cross section model tuning

The data-MC agreement in Q^2 (4-momentum transfer) is not good
We tuned nuclear parameters in Relativistic Fermi Gas model

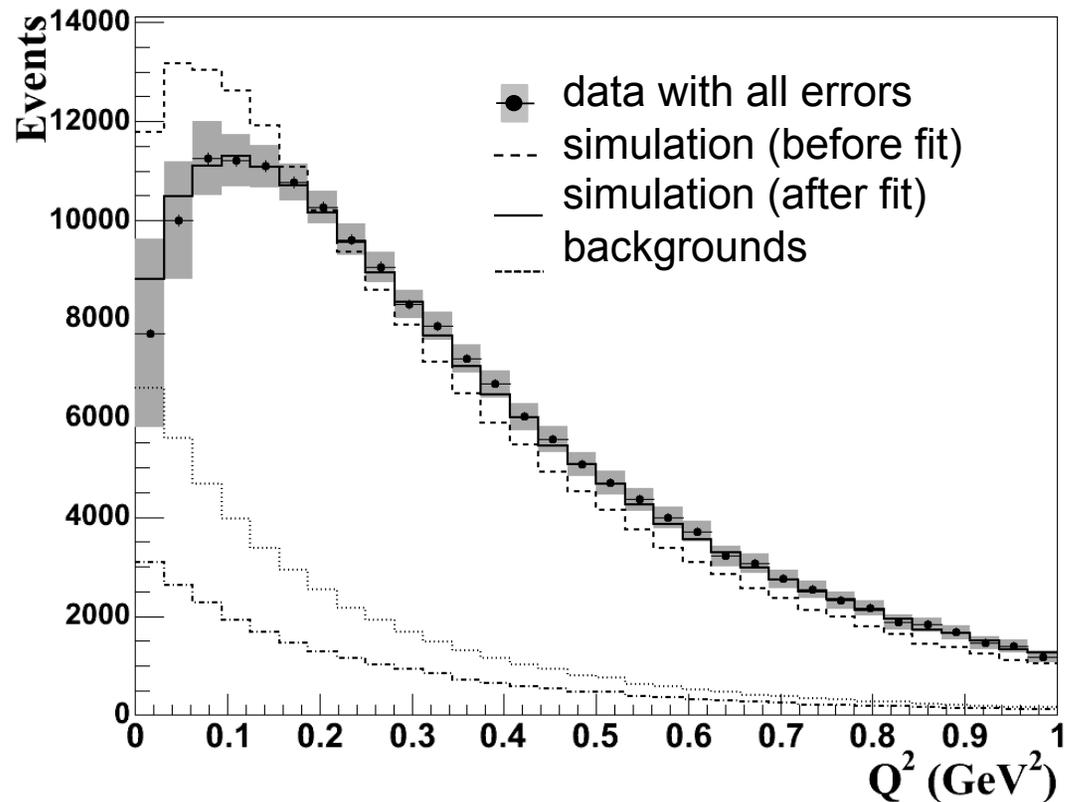
Q^2 fits to MB ν_μ CCQE data using the nuclear parameters:

M_A^{eff} - effective axial mass
 κ - Pauli Blocking parameter

Relativistic Fermi Gas Model with tuned parameters describes ν_μ CCQE data well

This improved nuclear model is used in ν_e CCQE model, too.

Q^2 distribution before and after fitting

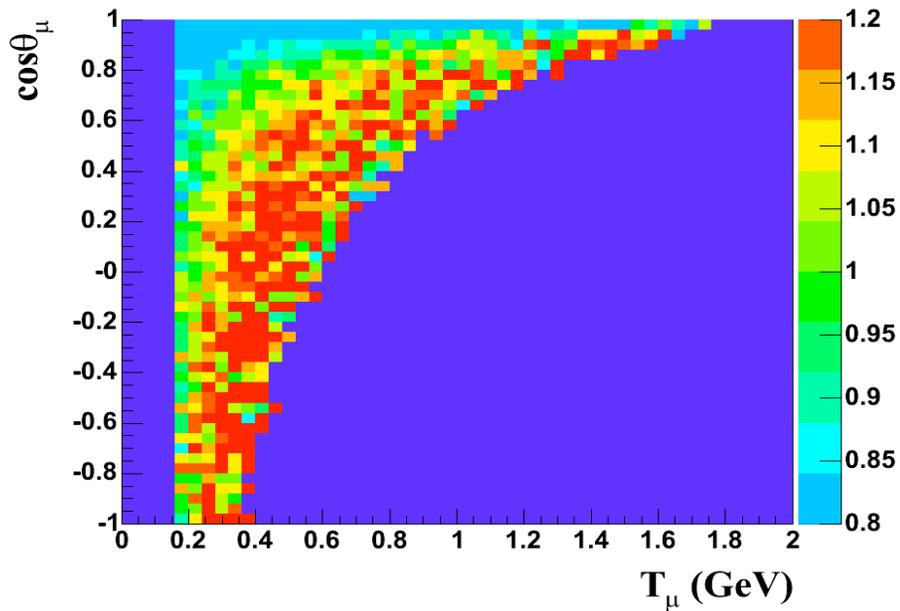


4. CCQE cross section model tuning

Without knowing flux perfectly, we cannot modify cross section model

$$R(\text{interaction}) \propto \int (\text{flux}) \times (\text{xs})$$

Data-MC ratio for T_μ - $\cos\theta_\mu$ plane, before tuning

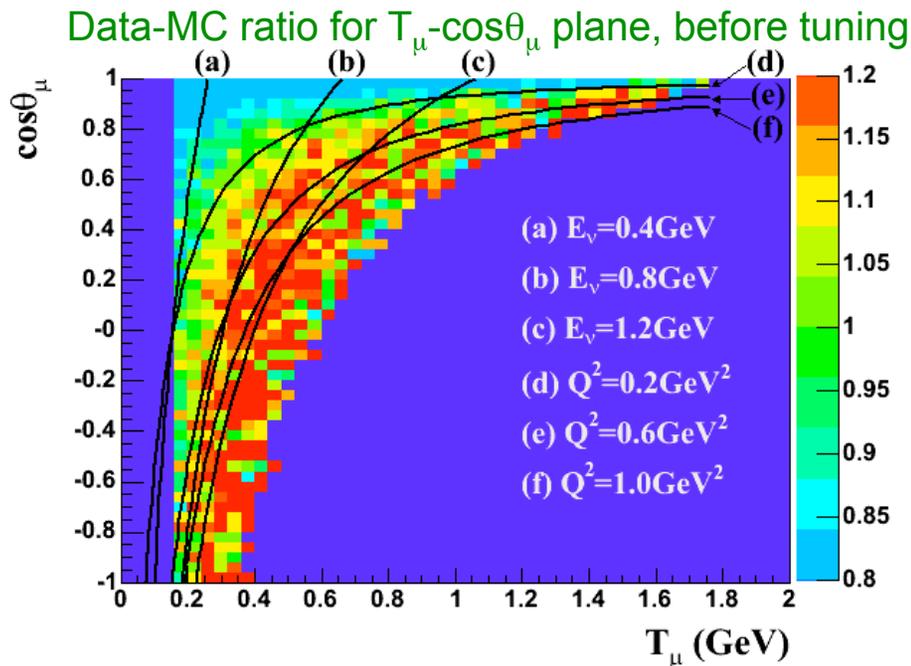


4. CCQE cross section model tuning

Without knowing flux perfectly, we cannot modify cross section model

$$R(\text{interaction}[E_\nu, Q^2]) \propto \int (\text{flux}[E_\nu]) \times (\text{xs}[Q^2])$$

Data-MC mismatching follows Q^2 lines, not E_ν lines, therefore we can see the problem is not the flux prediction, but the cross section model

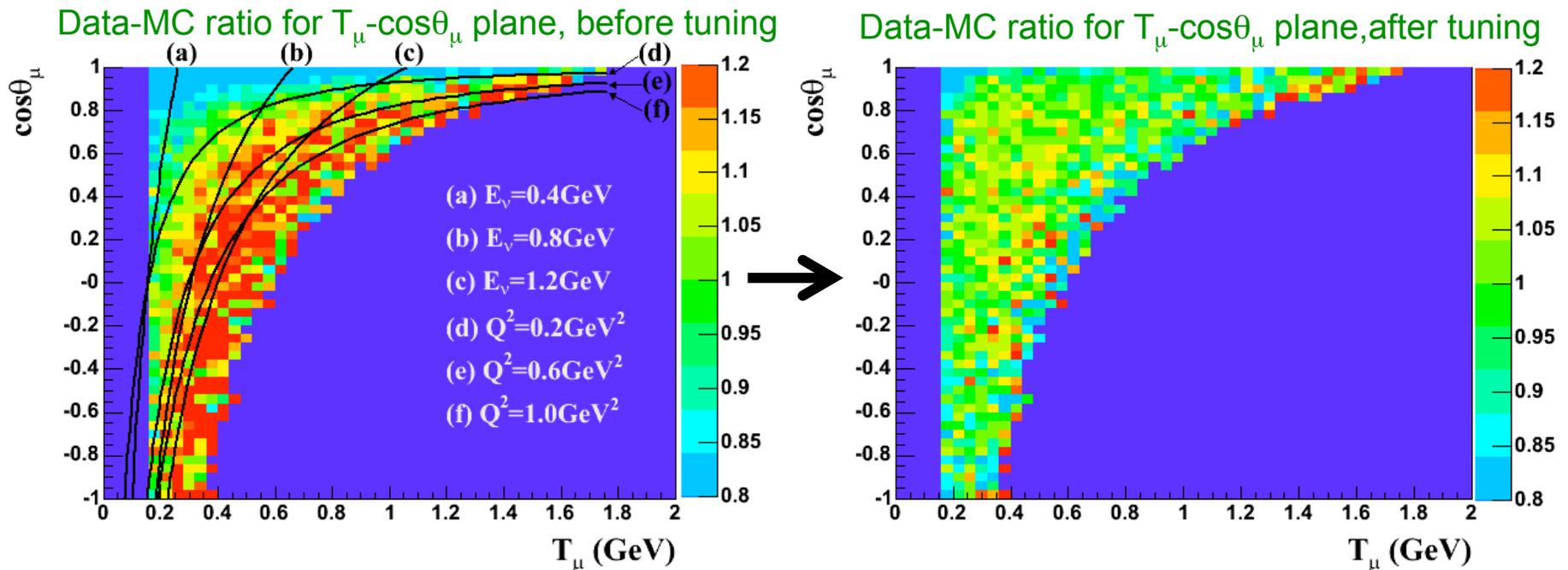


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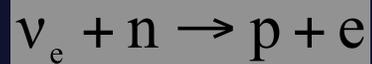
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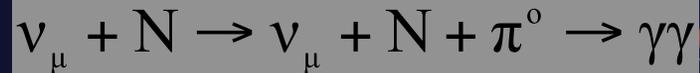
4. NC π^0 rate tuning

NC π^0 (neutral current π^0 production)

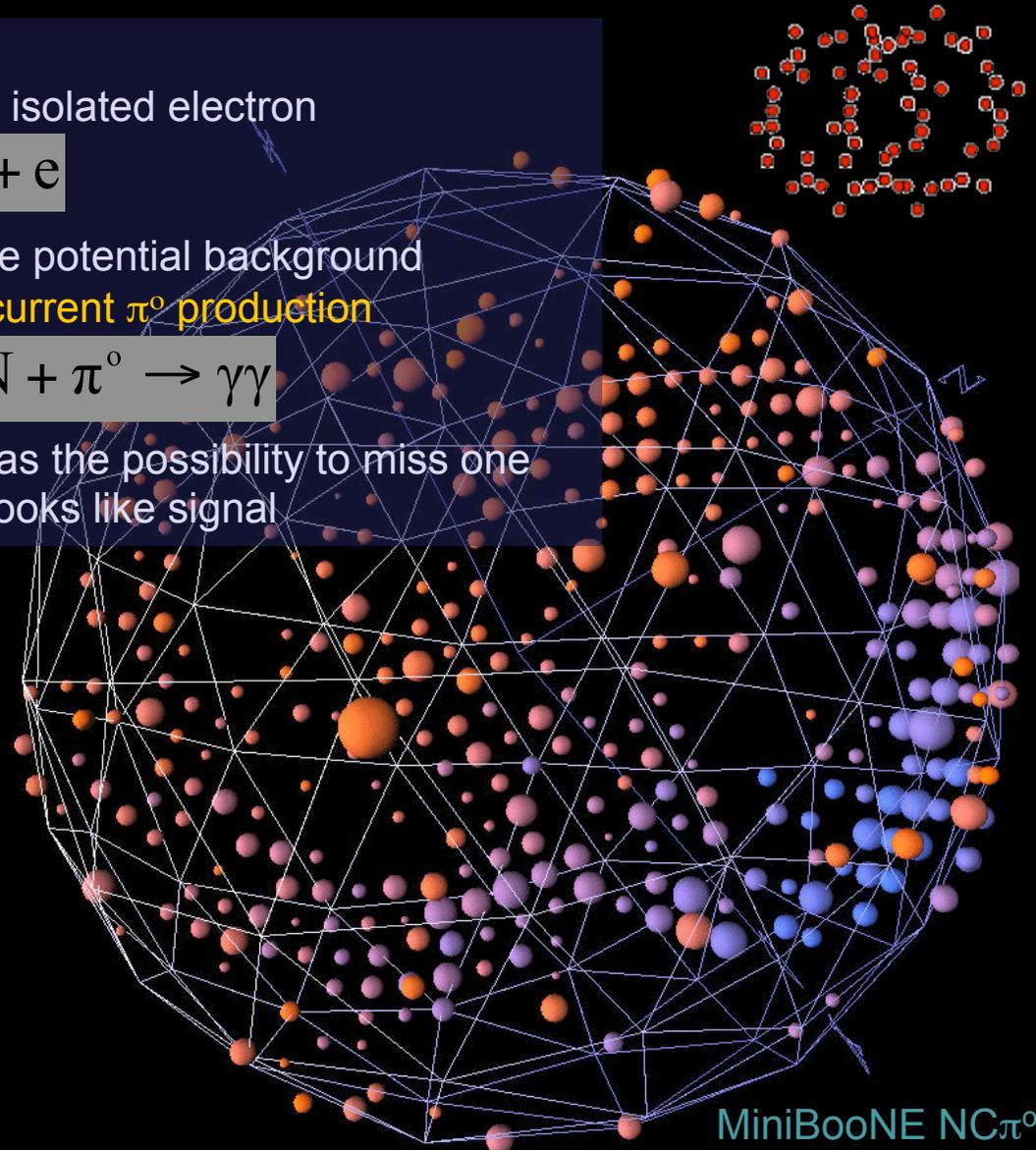
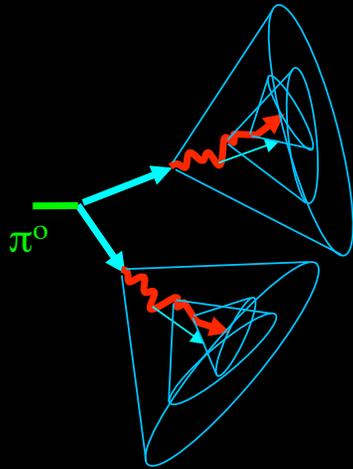
The signal of ν_e candidate is a single isolated electron



- single electromagnetic shower is the potential background
- the notable background is **Neutral current π^0 production**



Because of kinematics, one always has the possibility to miss one gamma ray, and hence this reaction looks like signal

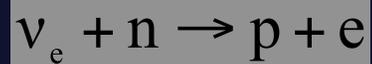


MiniBooNE NC π^0
candidate

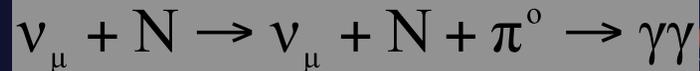
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NC π^0 (neutral current π^0 production)

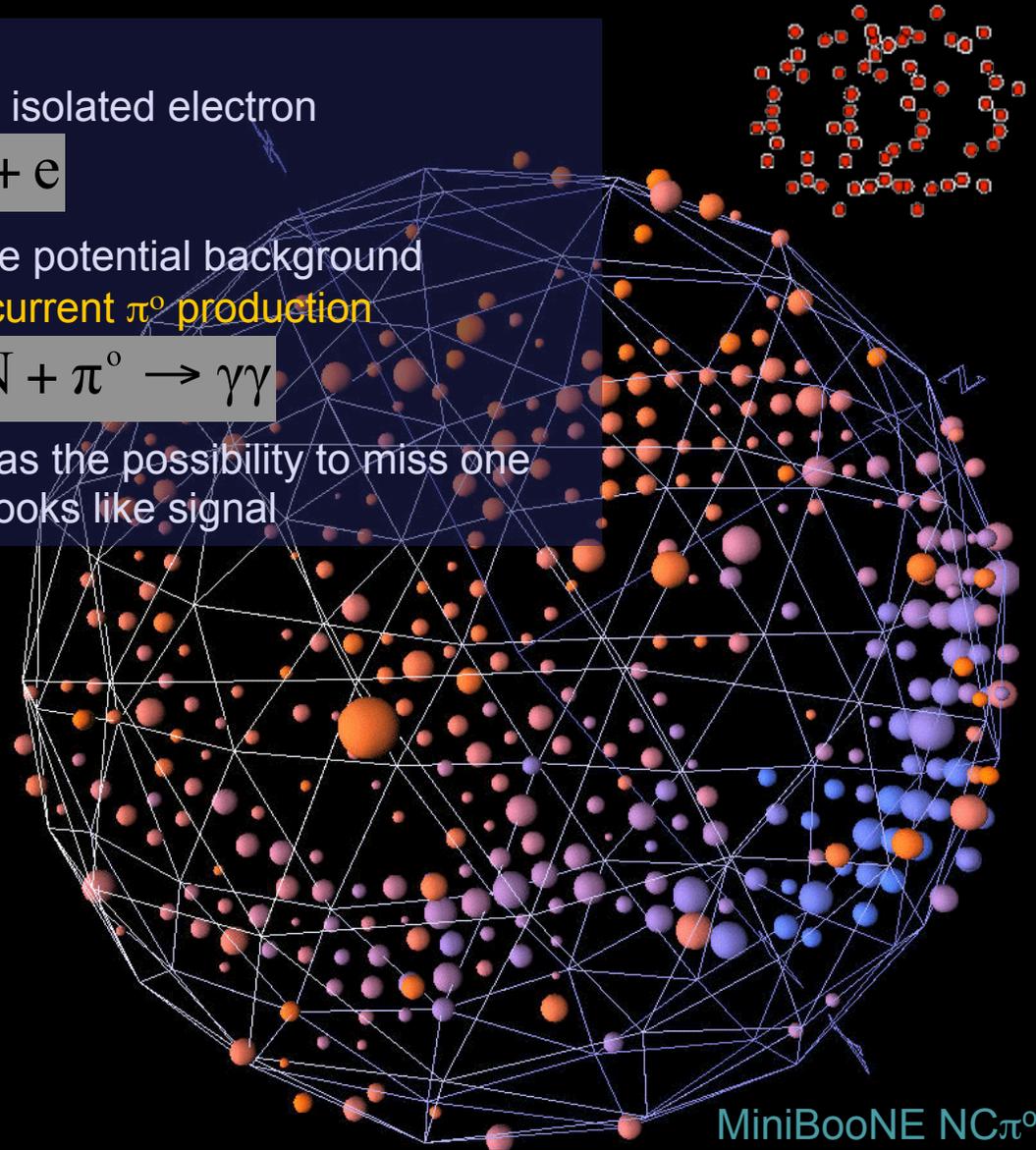
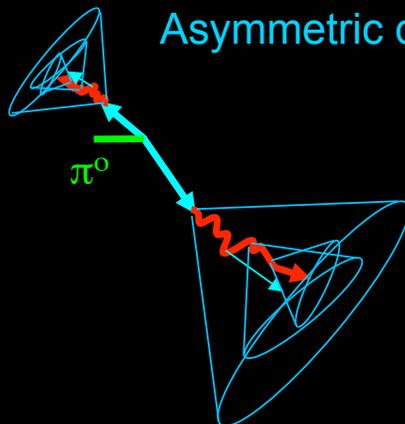
The signal of ν_e candidate is a single isolated electron



- single electromagnetic shower is the potential background
- the notable background is **Neutral current π^0 production**

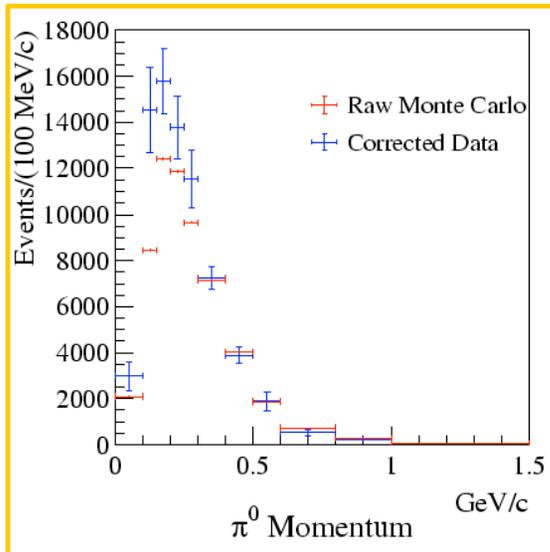


Because of kinematics, one always has the possibility to miss one gamma ray, and hence this reaction looks like signal

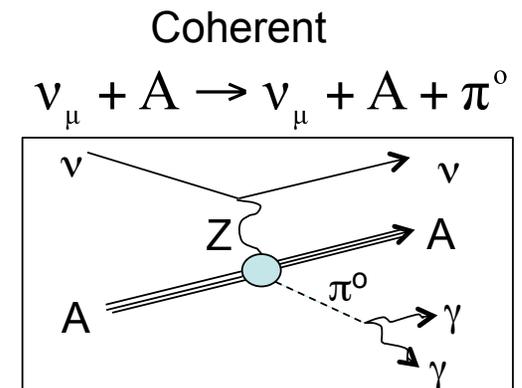
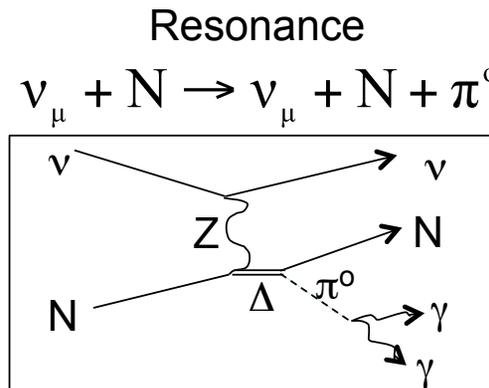
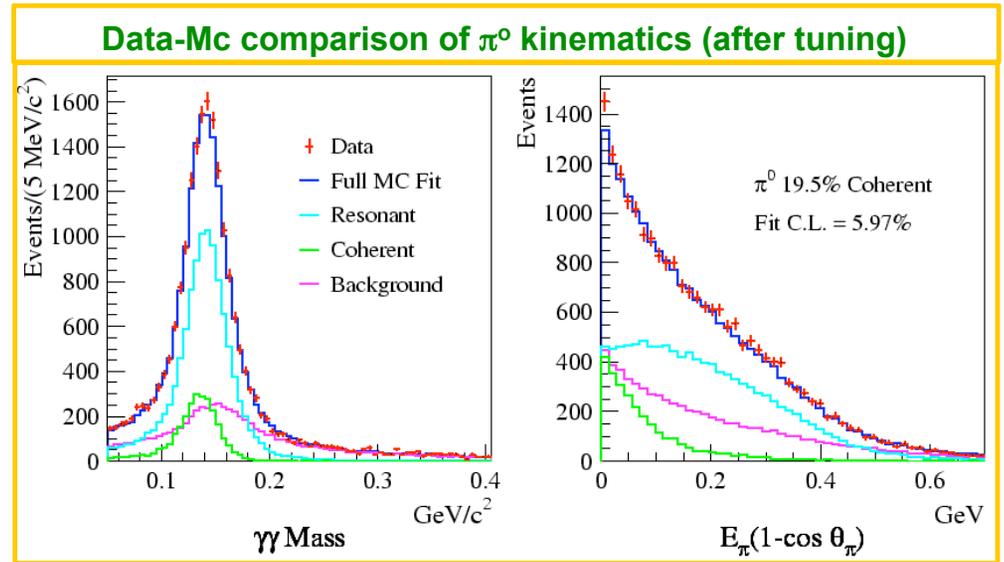


4. NC π^0 rate tuning

We tuned NC π^0 rate from our NC π^0 measurement. Since loss of gamma ray is pure kinematic effect, after tuning we have a precise prediction for intrinsic NC π^0 background for ν_e appearance search.



MiniBooNE collaboration
PLB664(2008)41



4. MiniBooNE cross section results

NuInt09, May18-22, 2009, Sitges, Spain

All talks proceedings are available on online (open access),
<http://proceedings.aip.org/proceedings/confproceed/1189.jsp>



1. charged current quasielastic (CCQE) cross section measurement
by Teppei Katori, [PRD81\(2010\)092005](#)
2. neutral current elastic (NCE) cross section measurement
by Denis Perevalov, [PRD82\(2010\)092005](#)
3. neutral current π^0 production (NC π^0) cross section measurement (ν and anti- ν)
by Colin Anderson, [PRD81\(2010\)013005](#)
4. charged current single pion production (CC π^+) cross section measurement
by Mike Wilking, [arXiv:1011.3572](#)
5. charged current single π^0 production (CC π^0) cross section measurement
by Bob Nelson, [arXiv:1010.3264](#)
6. CC π^+ /CCQE cross section ratio measurement
by Steve Linden, [PRL103\(2009\)081801](#)
7. anti- ν CCQE measurement
by Joe Grange, [arXiv:1102.1964](#)

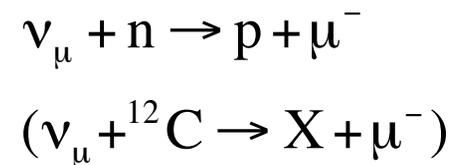
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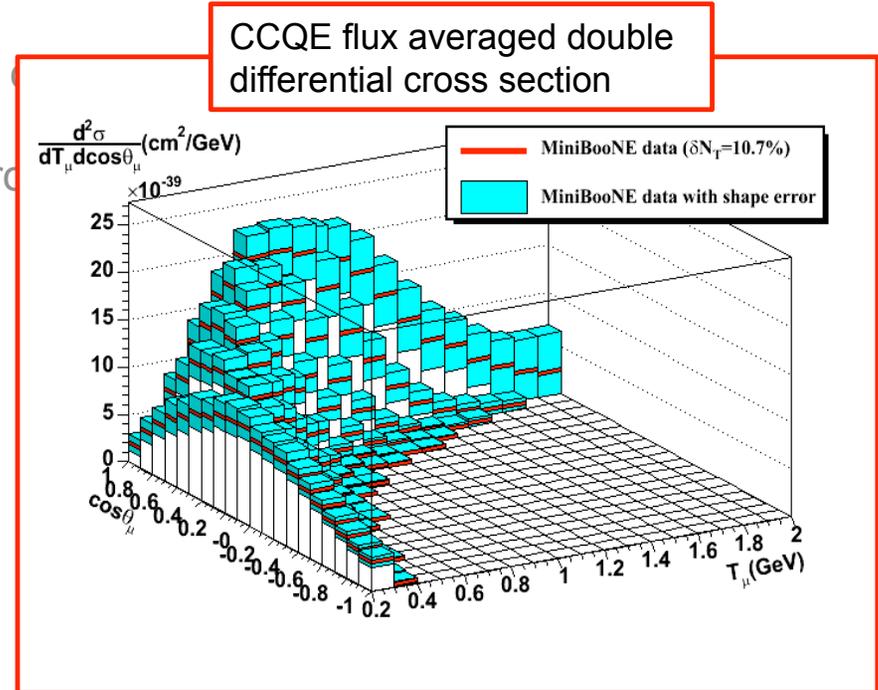
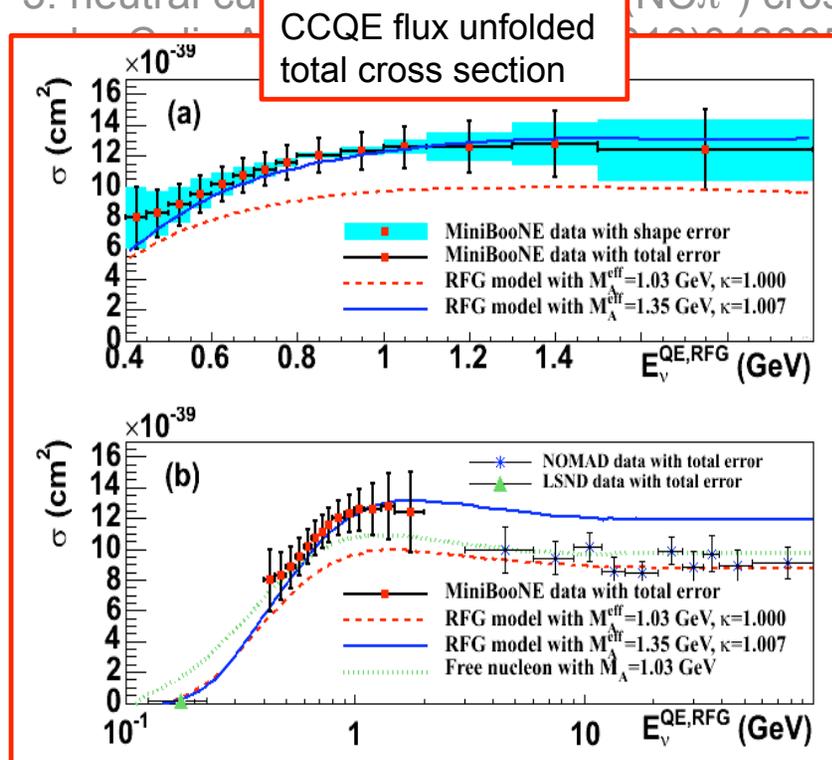
All talks proceedings are available on online (open access),
<http://proceedings.aip.org/proceedings/confproceed/1189.jsp>



Main results are flux averaged differential cross section, they are function of measured variables (e.g., muon angle, pion energy, etc), not inferred variables (e.g., neutrino energy, momentum transfer)



3. neutral current π^0 production (NC π^0) cross section measurement (ν and anti- ν)



4. MiniBooNE cross section results

NuInt11, Mar. 07-11, 2011, Dehradun, India

In NuInt11, MiniBooNE presented 2 new anti-neutrino measurement results.



1. charged current quasielastic (CCQE) cross section measurement
by Teppei Katori, [PRD81\(2010\)092005](#)

2. neutral current elastic (NCE) cross section measurement
by Denis Perevalov, [PRD82\(2010\)092005](#)

3. neutral current π^0 production (NC π^0) cross section measurement (ν
by Colin Anderson, [PRD81\(2010\)013005](#)

4. charged current single pion production (CC π^+) cross section measurement
by Mike Wilking, [arXiv:1011.3572](#)

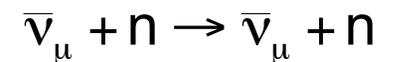
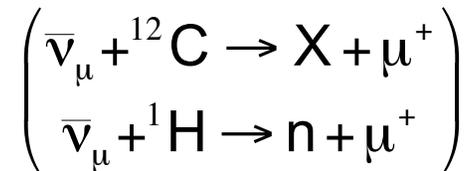
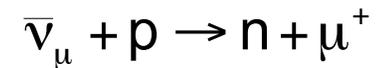
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6. CC π^+ /CCQE cross section ratio measurement
by Steve Linden, [PRL103\(2009\)081801](#)

7. anti- ν CCQE measurement
by Joe Grange, [arXiv:1102.1964](#)

8. **new anti- ν CCQE measurement**
by Joe Grange

9. **anti- ν NCE measurement**
by Ranjan Dharmapalan



Anti-neutrino interaction measurement is the best place to test neutrino interaction models developed in neutrino mode.

1. Introduction
2. Neutrino beam
3. Events in the detector
4. Cross section model
- 5. Neutrino oscillation result**
6. Anti-neutrino oscillation result
7. Outlook

5. Oscillation analysis background summary

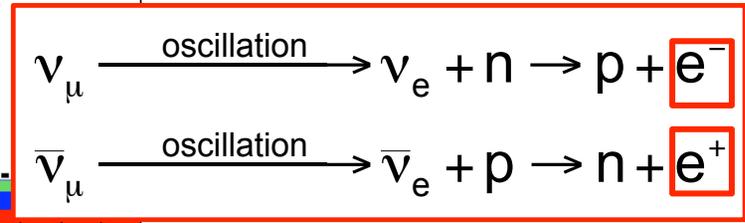
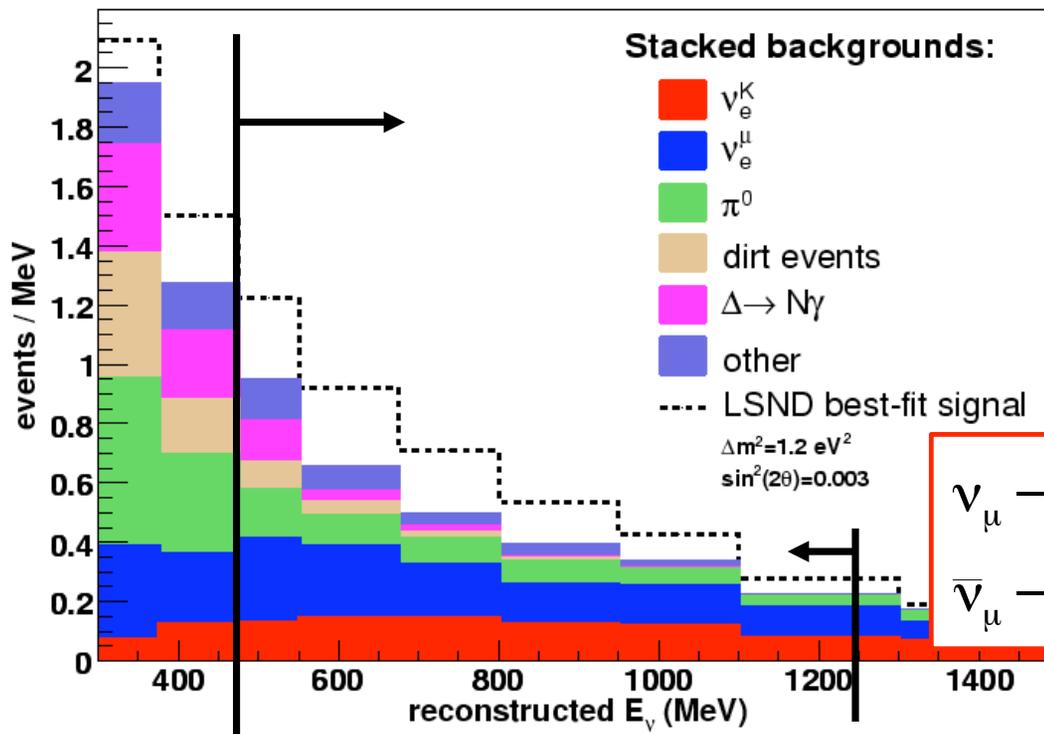
Oscillation analysis summary

- Oscillation analysis uses $475\text{MeV} < E < 1250\text{MeV}$

475 MeV – 1250 MeV

ν_e^K	94
ν_e^μ	132
π^0	62
dirt	17
$\Delta \rightarrow N \gamma$	20
other	33
<hr/>	
total	358

LSND best-fit $\nu_\mu \rightarrow \nu_e$ 126



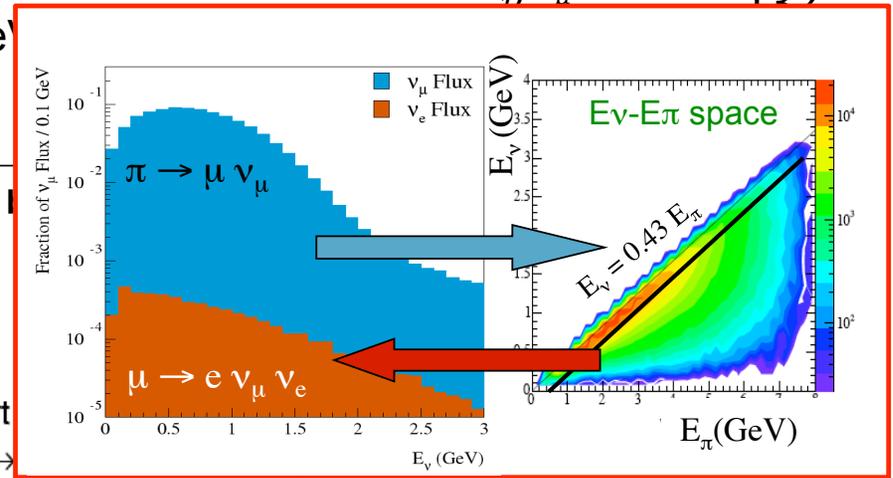
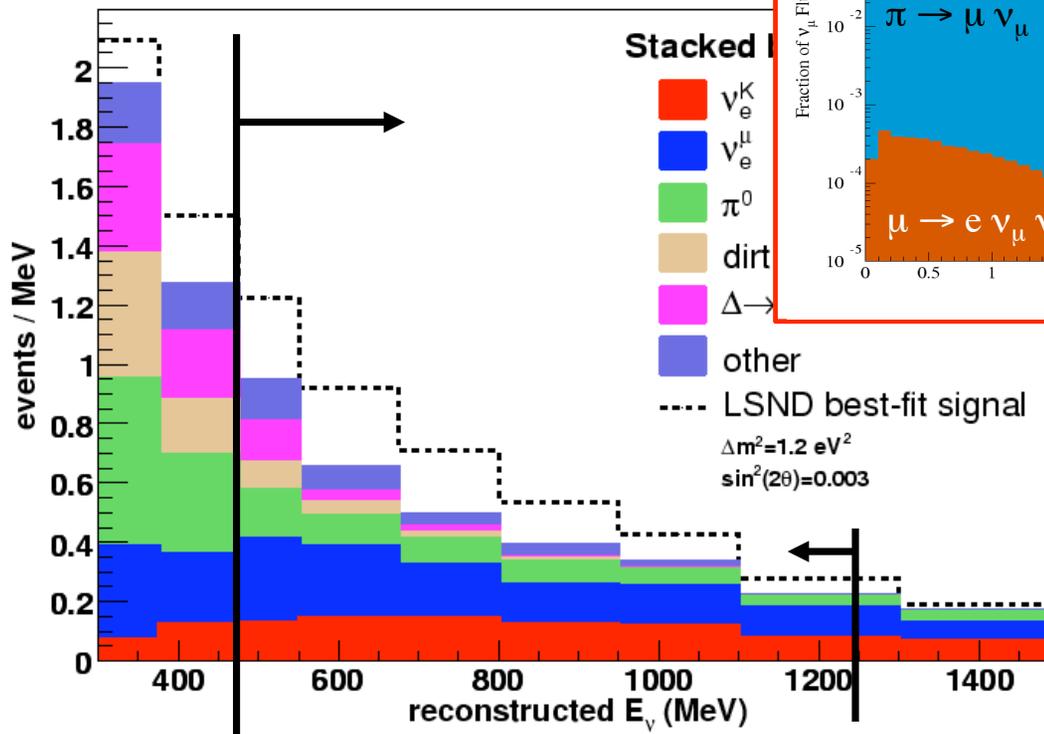
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475 MeV – 1250 MeV

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ν_μ	122

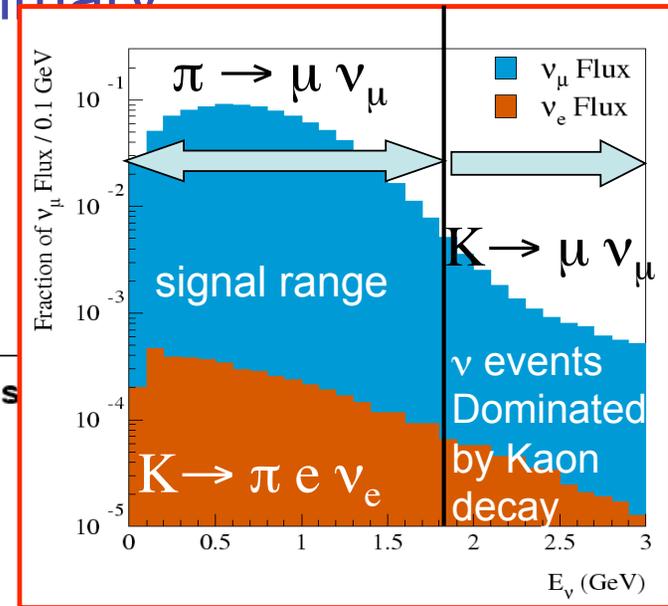
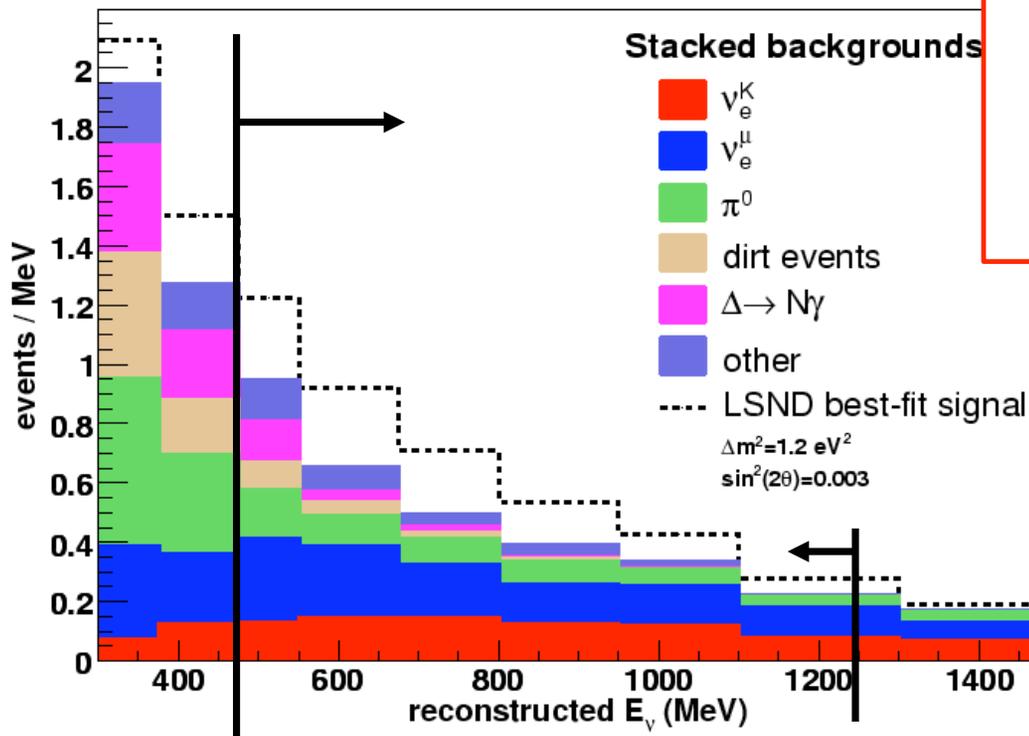


ν_e from μ decay is constrained from ν_μ CCQE measurement

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ν_e from μ decay is constrained from ν_μ CCQE measurement

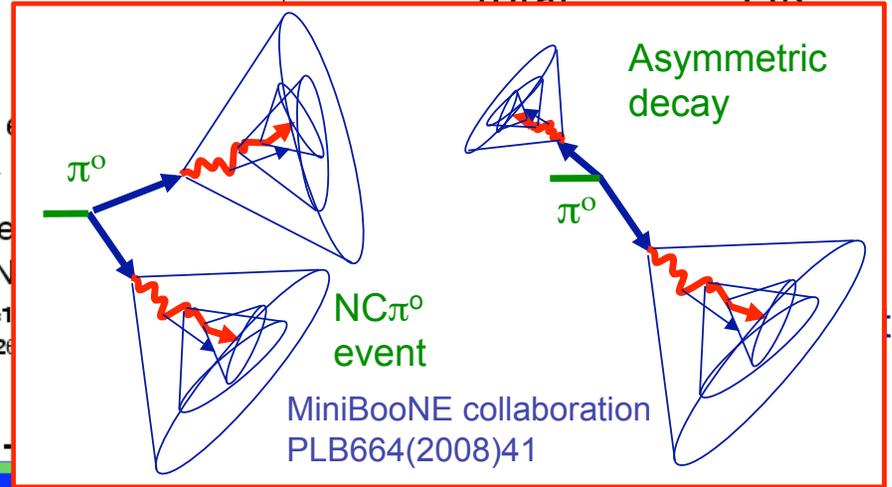
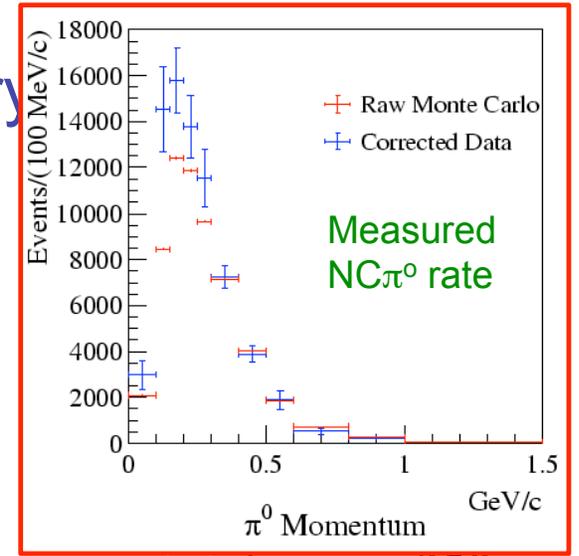
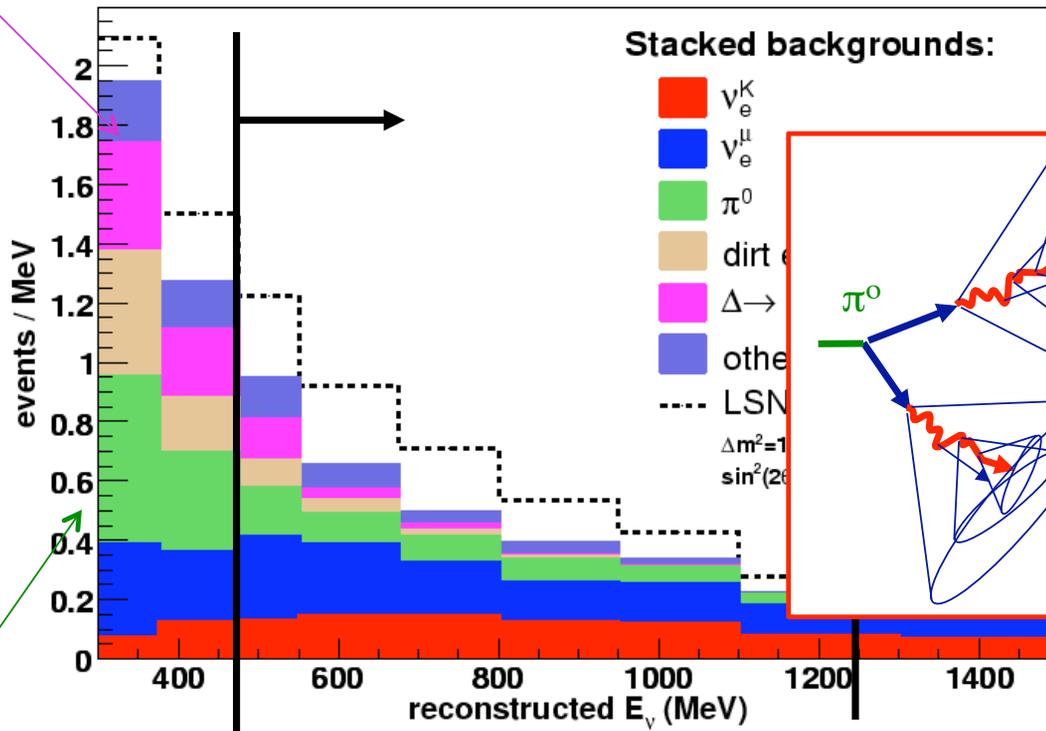
ν_e from K decay is constrained from high energy ν_e event measurement

5. Oscillation analysis background summary

Oscillation analysis summary

- Oscillation analysis uses $475\text{MeV} < E < 1250\text{MeV}$

Δ resonance rate is constrained from measured $\text{CC}\pi^0$ rate



Asymmetric π^0 decay is constrained from measured $\text{CC}\pi^0$ rate ($\pi^0 \rightarrow \gamma$)

ν_e from K decay is constrained from high energy ν_e event measurement

5. Oscillation analysis background summary

475 MeV - 1250 MeV

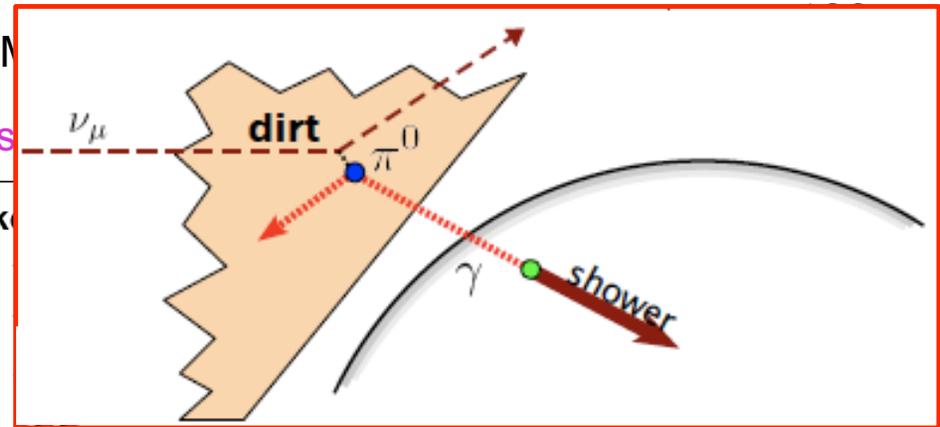
ν_e^K

94

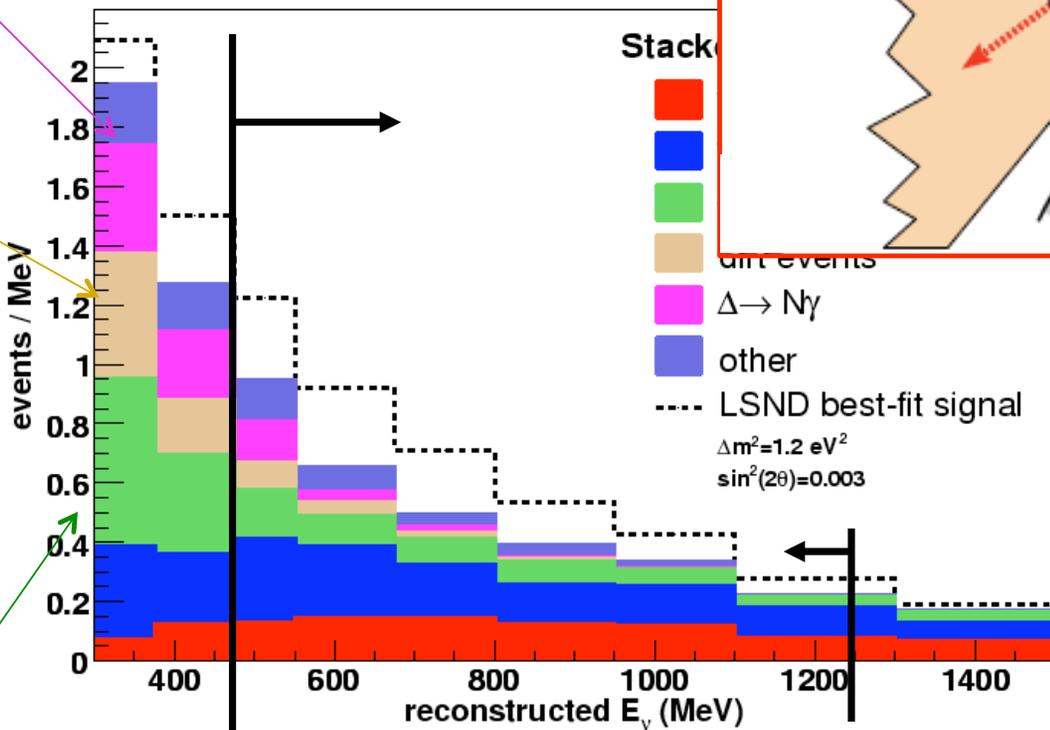
Oscillation analysis summary

- Oscillation analysis uses $475\text{MeV} < E < 1250\text{MeV}$

Δ resonance rate is constrained from meas



dirt rate is measured from dirt data sample



ν_e from μ decay is constrained from ν_μ CCQE measurement

ν_e from K decay is constrained from high energy ν_e event measurement

Asymmetric π^0 decay is constrained from measured CC π^0 rate ($\pi^0 \rightarrow \gamma$)

5. Oscillation analysis background summary

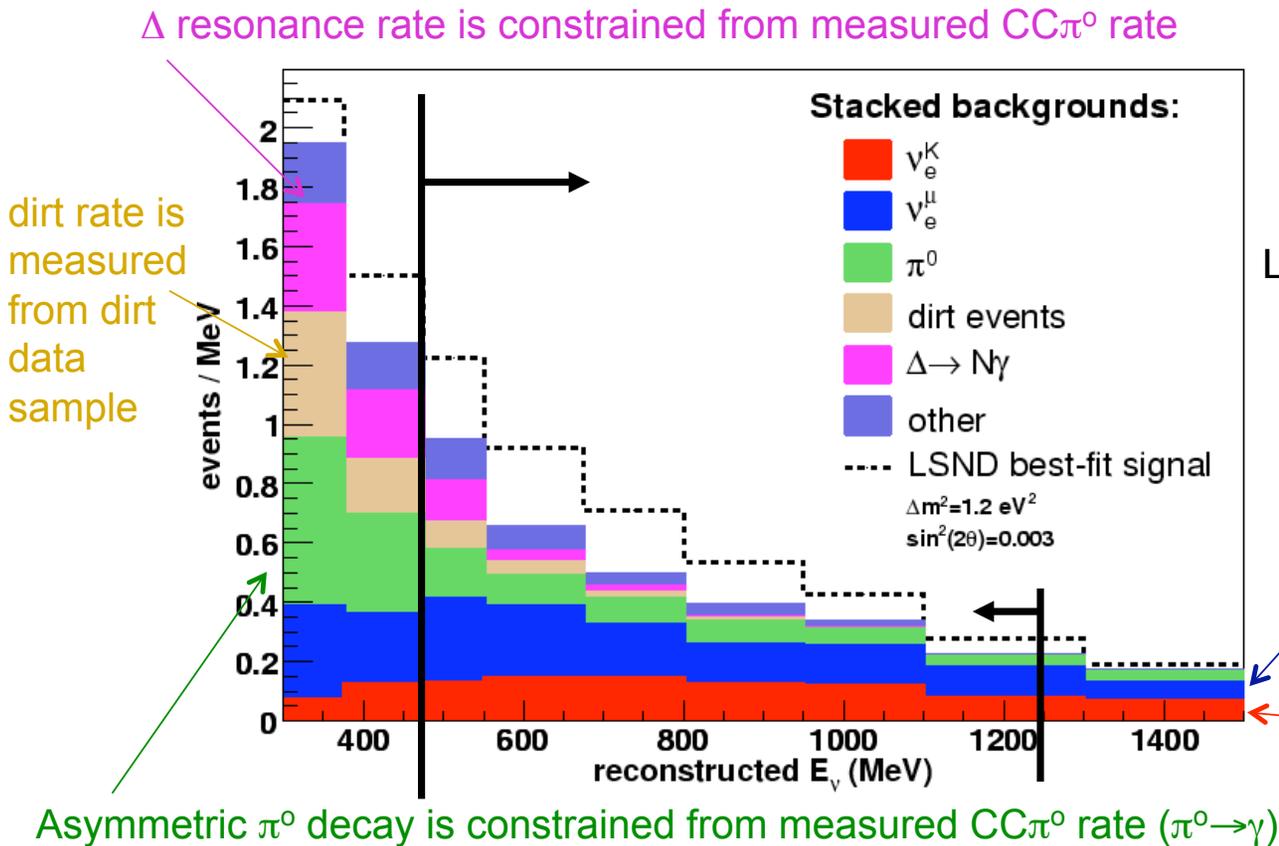
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other	33
<hr/>	
total	358

LSND best-fit $\nu_\mu \rightarrow \nu_e$ 126



03/17/2011

All backgrounds are measured in other data sample and their errors are constrained!

5. Error analysis - Multisim

Input error matrix
keep all correlation
of systematics

"multisim"
nonlinear error propagation

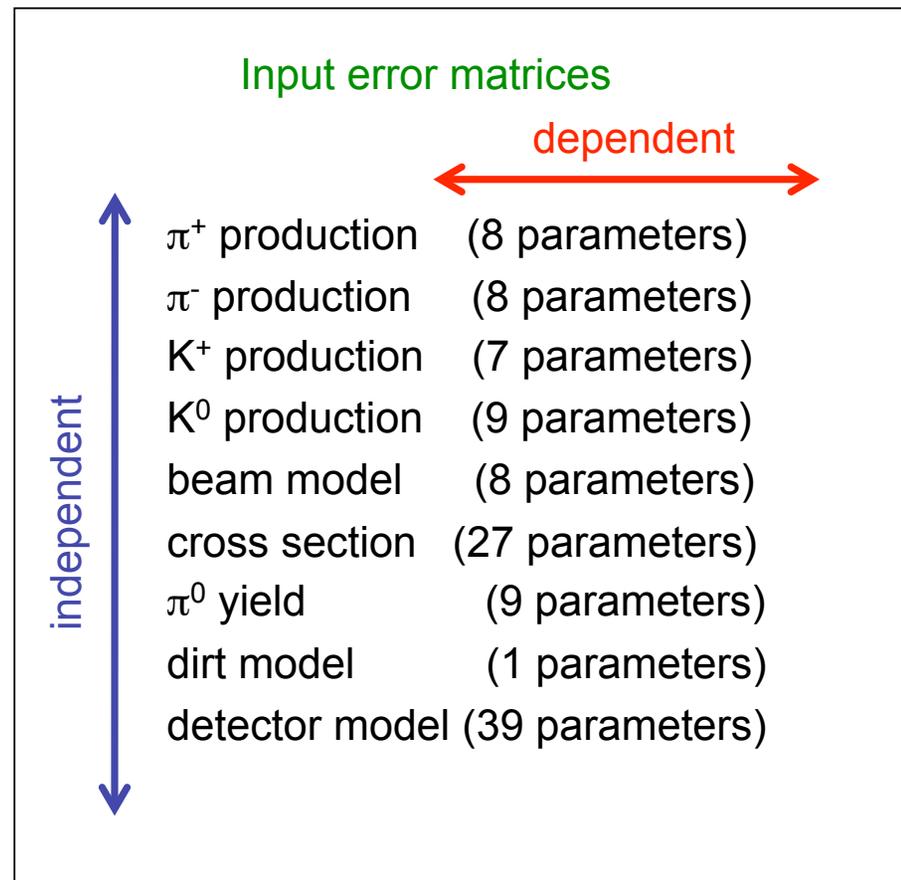
Output error matrix
keep all correlation
of E_ν^{QE} bins

Multi-simulation (Multisim) method

many fake experiments (~ 1000) with different parameter set give the variation of correlated systematic errors for each independent error matrix

The total error matrix is the sum of all independent error matrix

The total error matrix is used for oscillation fit to extract the best fit Δm^2 and $\sin^2 2\theta$.



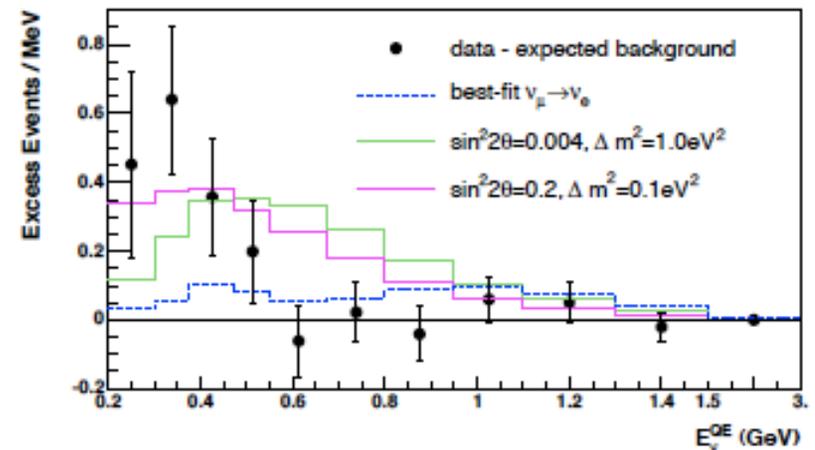
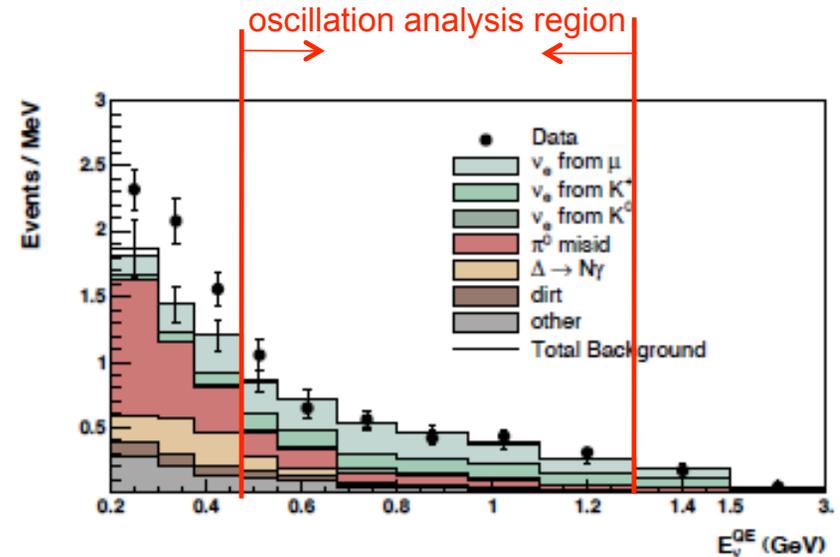
5. Oscillation analysis result

ν_e appearance oscillation result

- The best fit result shows no sign of an excess in the analysis region (where the LSND signal is expected from 1 sterile neutrino interpretation). However, there is 3.4σ excess in low energy region (300-475MeV)

- The shape of low energy excess is not described by any of two neutrino massive oscillation models

- many models are suggested to explain low energy excess, and can test those exotic models through antineutrino oscillation data

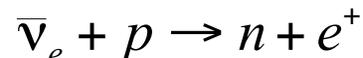
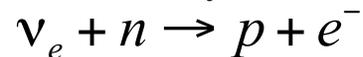


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5. Neutrino oscillation result
- 6. Anti-neutrino oscillation result**
7. Outlook

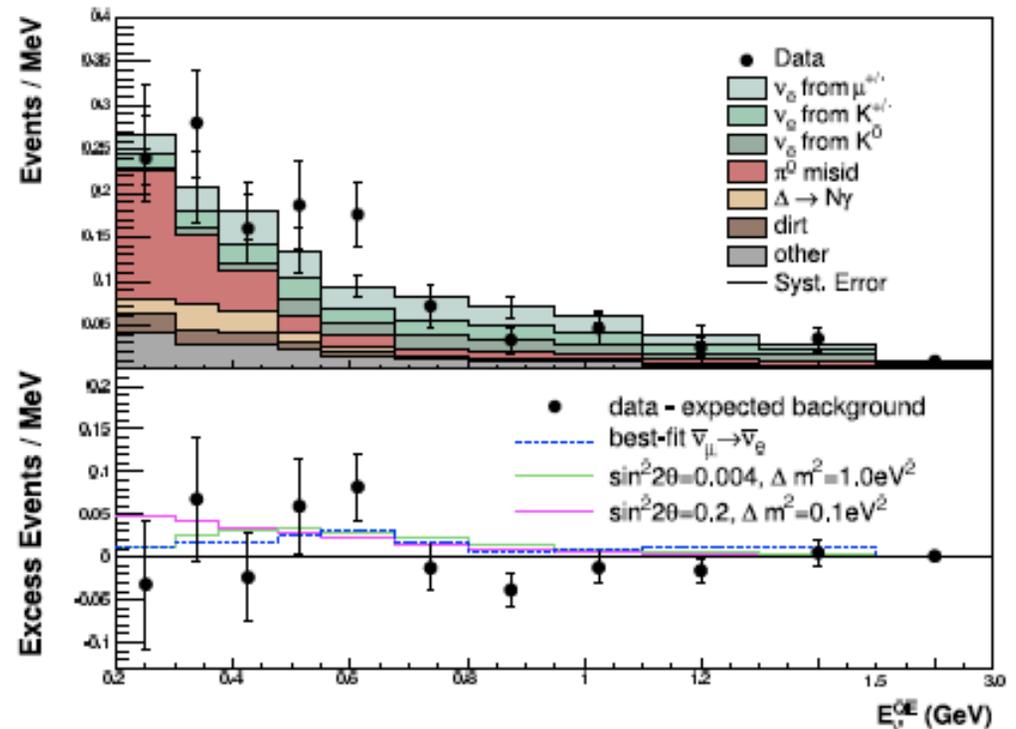
6. Antineutrino oscillation result

Many exotic models have some kind of predictions in antineutrino mode.

Analysis is quite parallel, because MiniBooNE doesn't distinguish e^- and e^+ or μ^- and μ^+ on event-by-event basis.



Bottom line, we don't see the low energy excess.

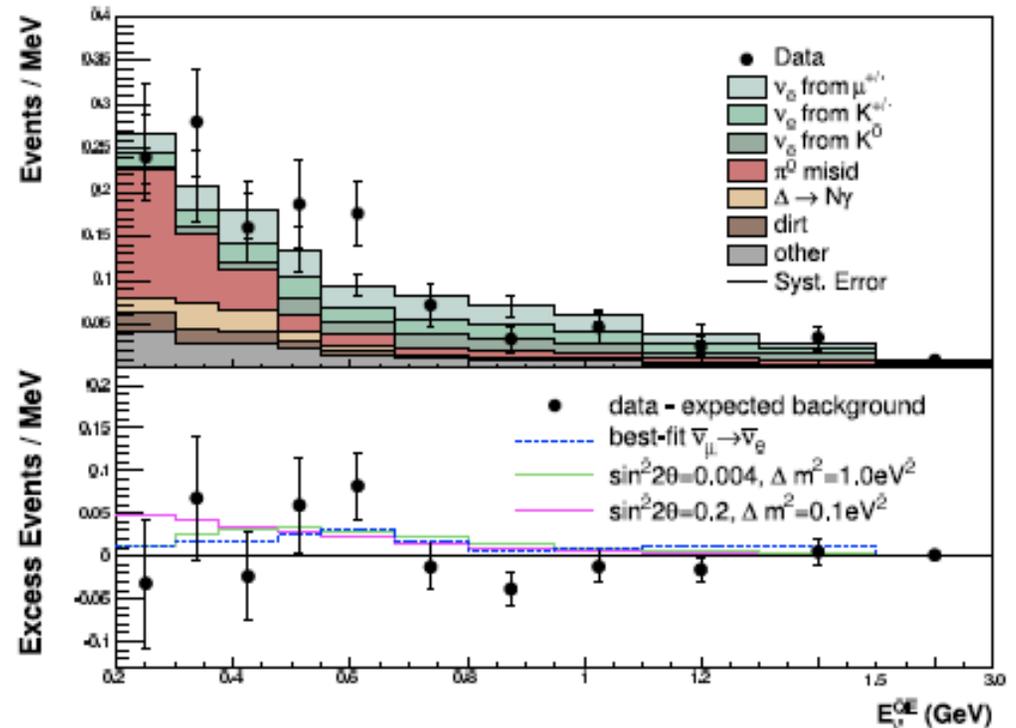


6. Antineutrino oscillation result

Implications

There are many theoretical models to explain low energy excess.

- The models based on same NC cross section for ν and anti- ν are disfavored.
- The models proportioned to POT are disfavored.
- The models which predict all excess only in neutrino mode, but not antineutrino are favored, such as neutrino-only induced excess



Hi theorists! new models are welcome!

6. New antineutrino oscillation result

- Antineutrino mode is the direct test of LSND signal
- Analysis is limited with statistics

New antineutrino oscillation result

	200-475 MeV	475-1250 MeV	200-3000 MeV
anti ν_e candidate	119	120	277

- 70% more data
- low level checks have been done (beam stability, energy scale)
- new dirt event rate measurement (consistent with neutrino mode)
- new $NC\pi^0$ rate measurement (consistent with neutrino mode)
- ν fraction is measured in anti- ν beam

New antineutrino oscillation result
(presented at Neutrino 2010, Athens)



6. New antineutrino oscillation result

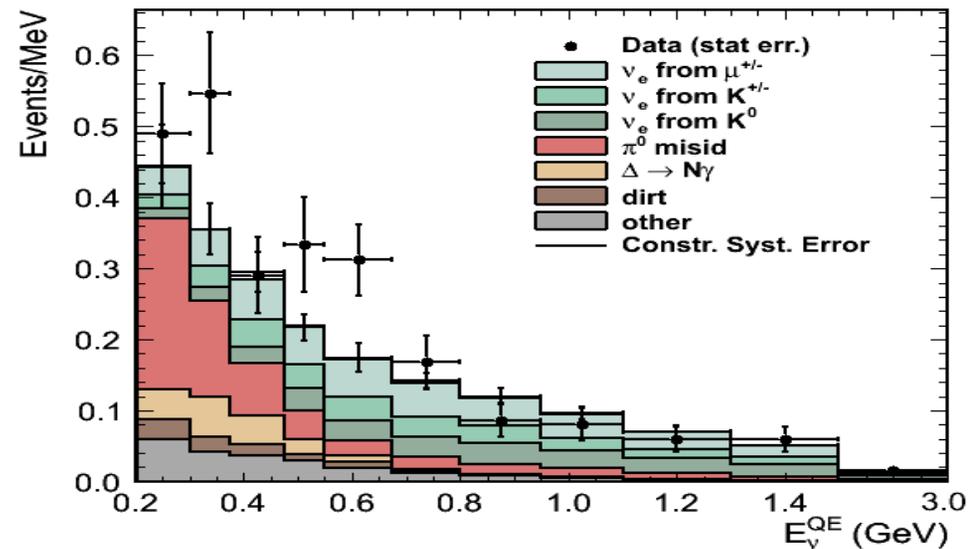
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MiniBooNE now see the excess in LSND-like Δm^2 region!

	200-475 MeV	475-1250 MeV	200-3000 MeV
anti ν_e candidate	119	120	277
MC (stat+sys)	100.5 ± 14.3	99.1 ± 13.9	233.8 ± 22.5
Excess (stat+sys)	$18.5 \pm 14.3 (1.3\sigma)$	$20.9 \pm 13.9 (1.5\sigma)$	$43.2 \pm 22.5 (1.9\sigma)$



6. New antineutrino oscillation result

- Antineutrino mode is the direct test of LSND signal
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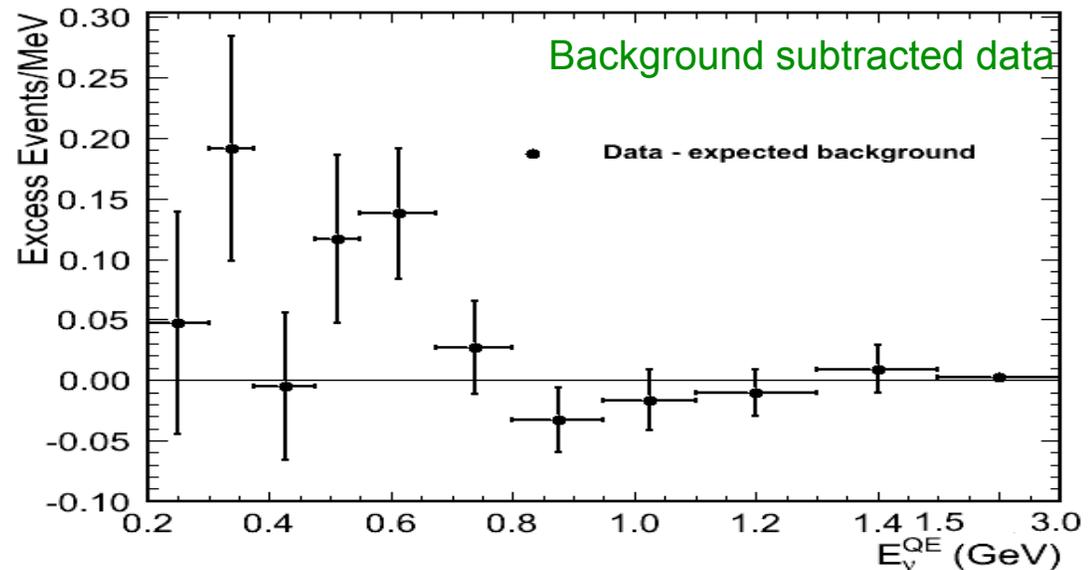
New antineutrino oscillation result

	before fit	
	χ^2/NDF	probability
$475 < E_{\nu}^{\text{QE}} < 1250 \text{ MeV}$	18.5/6	0.5%

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- flatness test (model independent test) shows statistical significance of signal.



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New antineutrino oscillation result

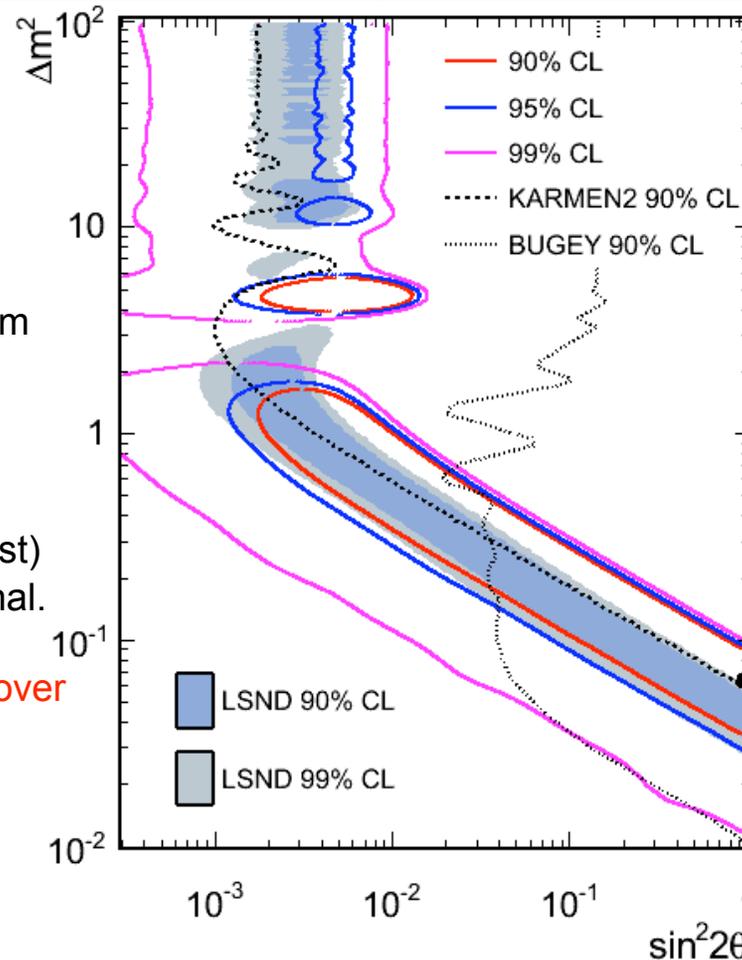
- 70% more data
- low level checks have been done (beam stability, energy scale)
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- new $NC\pi^0$ rate measurement (consistent with neutrino mode)
- $\bar{\nu}$ fraction is measured in anti- ν beam

MiniBooNE now see the excess in LSND-like Δm^2 region!

- flatness test (model independent test) shows statistical significance of signal.

2 massive neutrino model is favored over 99.4% than null hypothesis

	before fit		after fit	
	χ^2/NDF	probability	χ^2/NDF	probability
$475 < E_{\nu}^{\text{QE}} < 1250 \text{ MeV}$	18.5/6	0.5%	8.0/4	8.7%



$E > 475 \text{ MeV}$

Best fit point
 $\Delta m^2 = 0.064 \text{ eV}^2$
 $\sin^2 2\theta = 0.96$

1. Introduction
2. Neutrino beam
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6. Anti-neutrino oscillation result
- 7. Outlook**

7. MiniBooNE oscillation result summary

Neutrino mode analysis

- no excess is observed in the energy region where excess is expected from LSND
- significant excess is observed in low energy region

Antineutrino mode analysis

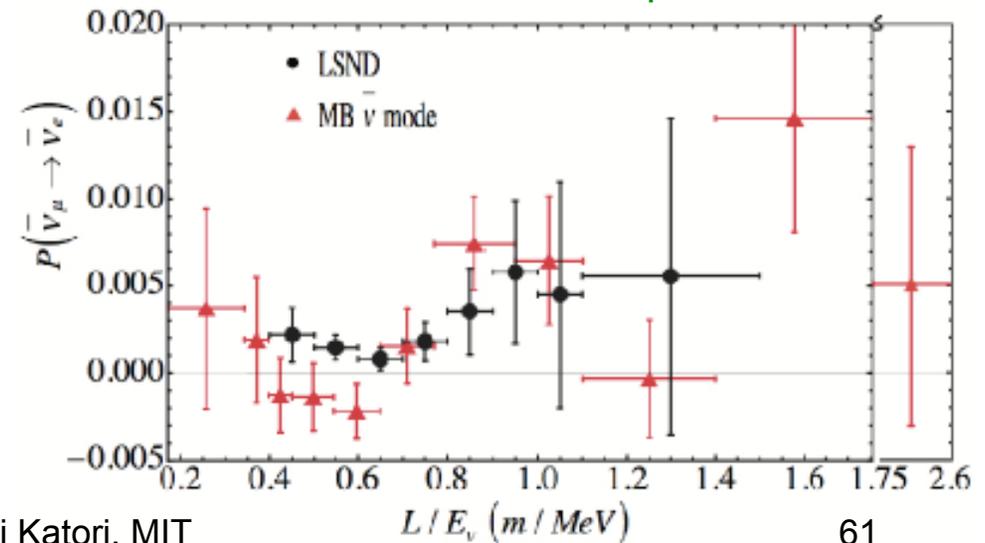
- small excess is observed in low energy region
- LSND consistent excess is observed in the oscillation energy region

These results are not main interest of Neutrino community (this is not θ_{13} nor leptonic CP violation nor Majorana mass measurement).

There is no convincing theoretical model to solve all mysteries.

Is MiniBooNE wrong?

MiniBooNE-LSND comparison in L/E



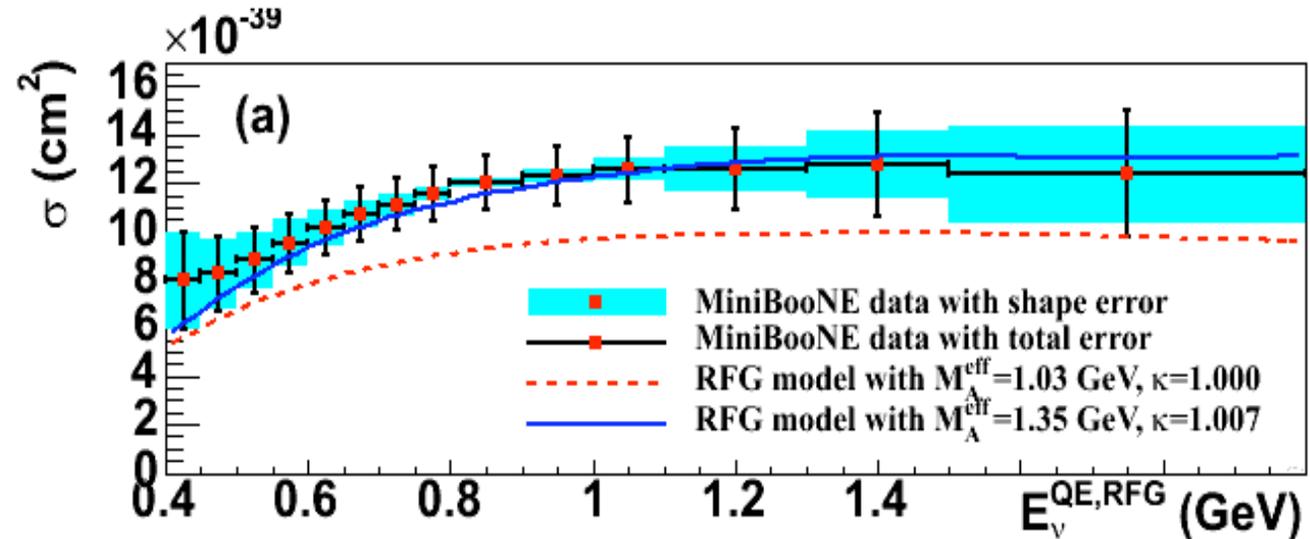
7. MiniBooNE CCQE absolute cross section

CCQE total cross section from MiniBooNE

MiniBooNE observed 30% higher neutrino cross section from RFG model with world averaged nuclear parameter from all past precise bubble chamber experiments.

When we first published this, we got so many criticism. Even a theorist claimed “MiniBooNE overestimate cross section!”

...but there is a turning point...



7. MiniBooNE CCQE absolute cross section



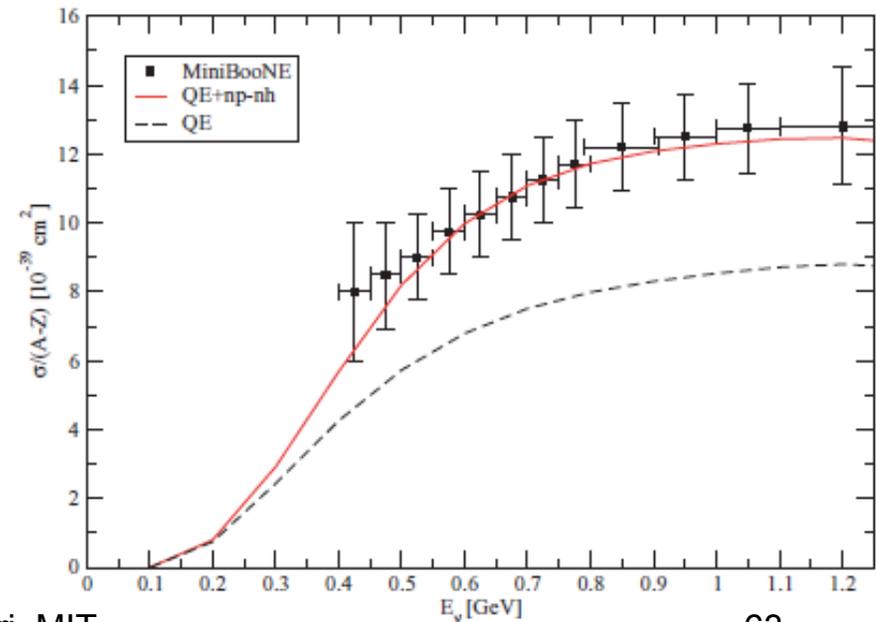
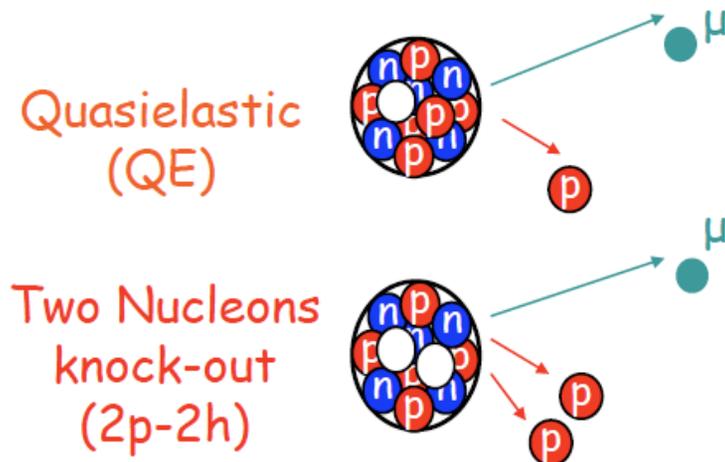
CCQE total cross section from MiniBooNE and RPA model

Martini et al.,
PRC80(2009)065501

Martini et al published their new RPA calculation result. They took into account the detail of nucleon emission channel (np-nh effect) and they explained MiniBooNE data.

Suddenly, many theorists start to appreciate this discovery by MiniBooNE.

So why all past experiments couldn't find this?



7. MiniBooNE CCQE absolute cross section



CCQE total cross section from MiniBooNE and RPA model

Martini et al.,
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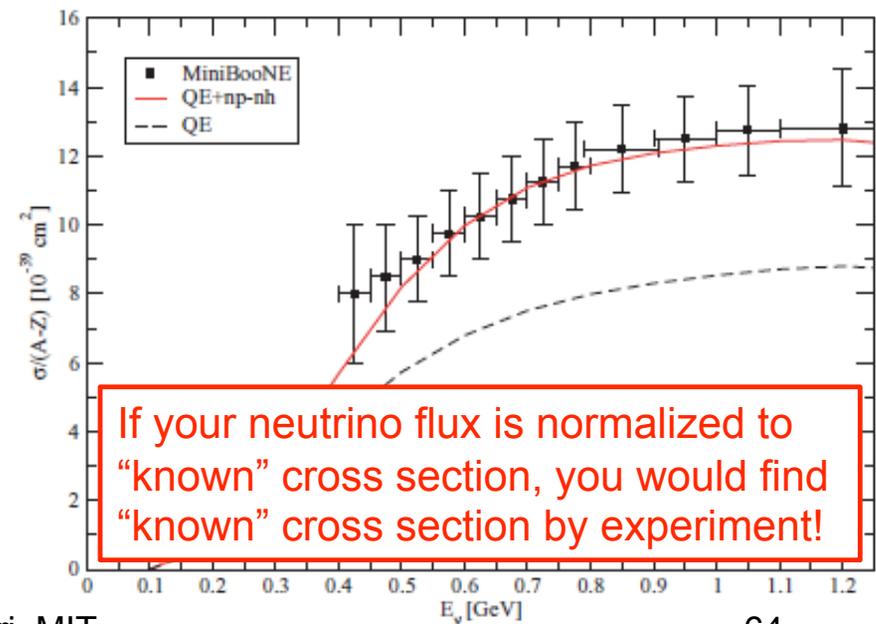
Suddenly, many theorists start to appreciate this discovery by MiniBooNE.

So why all past experiments couldn't find this?

There is a tendency for people to **measure and discover what is predicted**

Phys. Rev. DXX, (19XX)

The distribution of events in neutrino energy for the $3C \nu d \rightarrow \mu^- pp_s$ events is shown in Fig. 4 together with the quasielastic cross section $\sigma(\nu n \rightarrow \mu^- p)$ calculated using the standard $V-A$ theory with $M_A = 1.05 \pm 0.05$ GeV and $M_V = 0.84$ GeV. The absolute cross sections for the CC interactions have been measured using the quasielastic events and its known cross section.⁴



7. MiniBooNE CCQE absolute cross section

We shouldn't do this kind of mistake.

Many of MiniBooNE result are unexpected, and unexplained. But that cannot be a reason to be wrong. Remember, how much our naïve assumptions were correct for what we call now standard neutrino model.

(Neutrino 2006, Murayama)

Solar neutrino oscillation solution is SMA, because it's pretty

-> Wrong, LMA is the right solution

Natural scale of neutrino mass is $\sim 10-100\text{eV}^2$ because it's cosmologically interesting

-> Wrong, much smaller

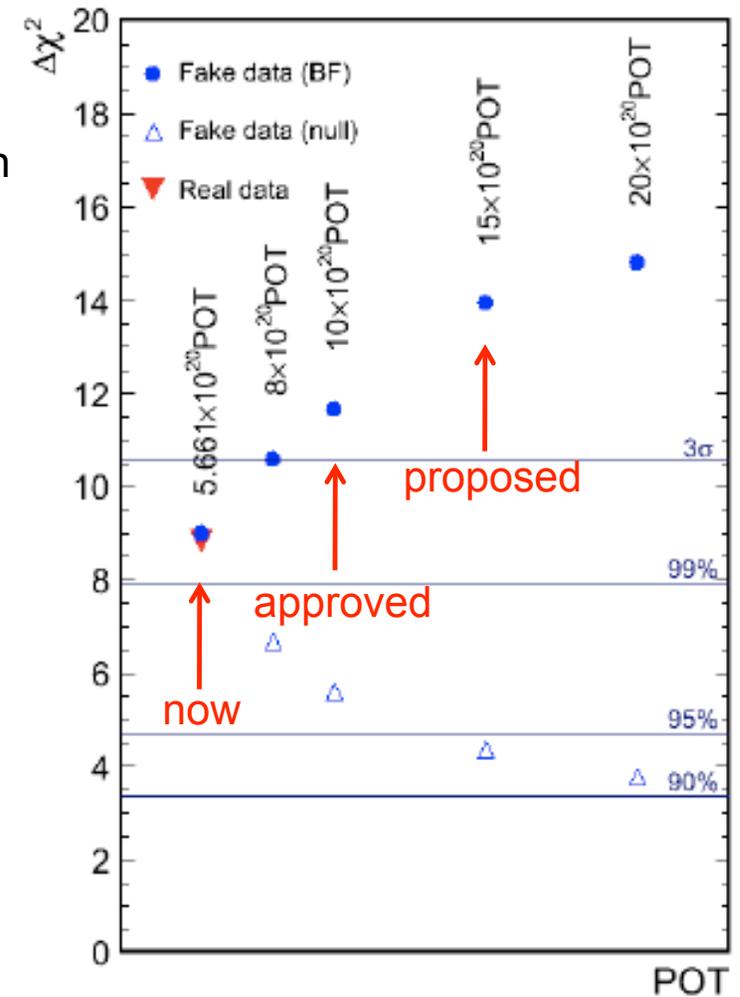
Atmospheric mixing should be small like CKM matrix element $V_{cb} \sim 0.04$, cannot be large

-> Wrong, much larger

Neutrino physics keep surprising us, so does MiniBooNE!

7. MiniBooNE future plan

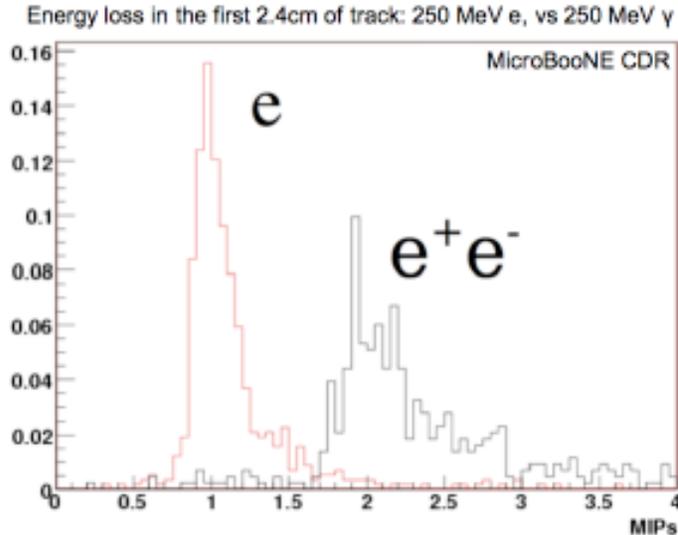
We continue to take data until March 2012 (approved), then we will double the statistics and expect 3σ excess in antineutrino mode. We are putting a proposal for $15E20$ extension.



7. MicroBooNE

Liquid Argon TPC experiment at Fermilab

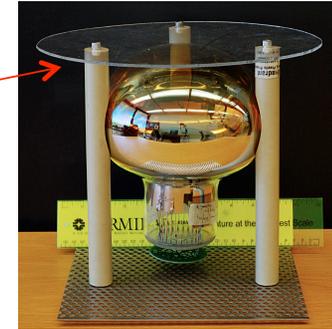
- 70 ton fiducial volume LiqAr TPC
- R&D detector for future large LiqAr TPC for DUSEL
- 3D tracker (modern bubble chamber)
- data taking will start from 2013
- dE/dx can separate single electron from gamma ray (e^+e^- pair)



scintillation from LiqAr

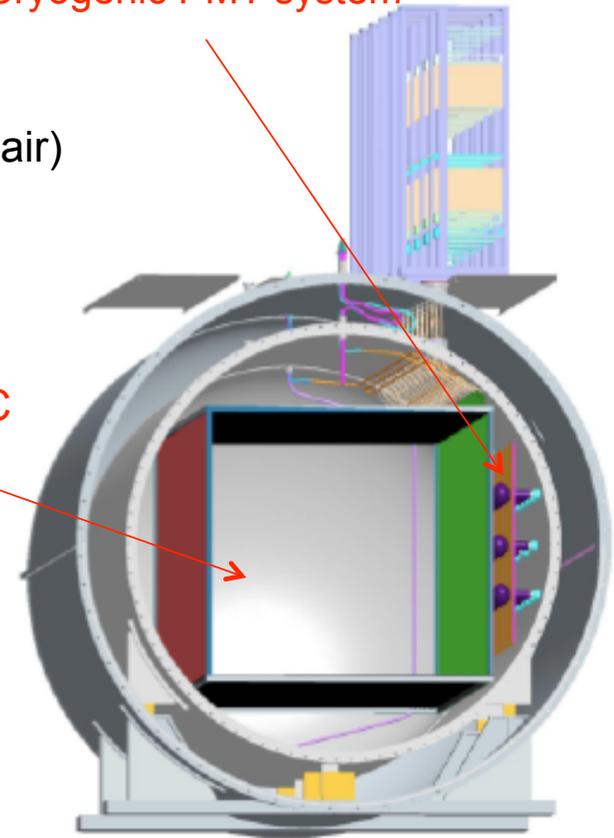
128nm
450nm

TPB (wave length shifter)
coated acrylic plate



Cryogenic PMT system

liquid Argon TPC



MiniBooNE collaboration

University of Alabama
Bucknell University
University of Cincinnati
University of Colorado
Columbia University
Embry Riddle Aeronautical University
Fermi National Accelerator Laboratory
Indiana University
University of Florida

Los Alamos National Laboratory
Louisiana State University
Massachusetts Institute of Technology
University of Michigan
Princeton University
Saint Mary's University of Minnesota
Virginia Polytechnic Institute
Yale University



Thank you for your attention!

Buck up

2. Neutrino beam

HARP experiment (CERN)

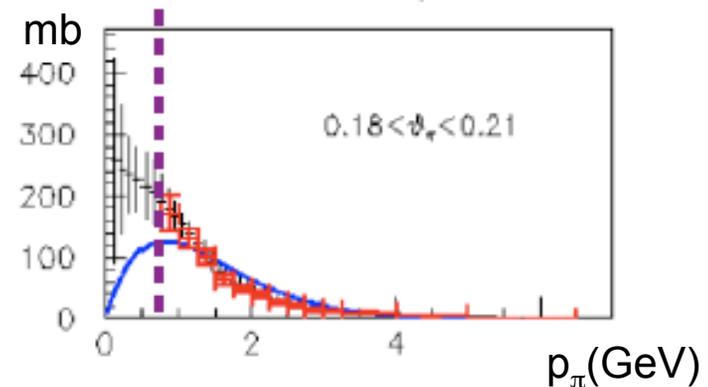
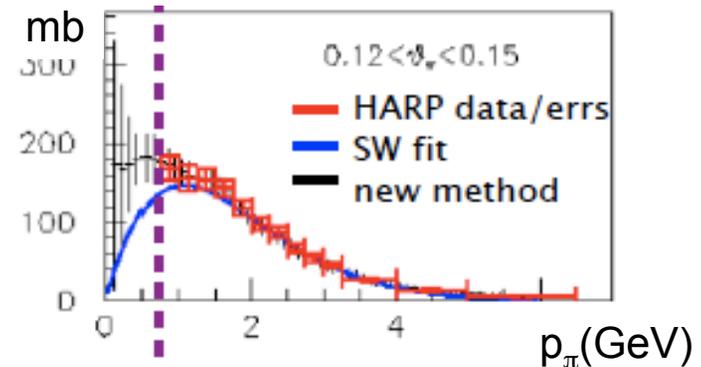


Modeling of meson production is based on the measurement done by HARP collaboration.

- Identical, but 5% λ Beryllium target
- 8.9 GeV/c proton beam momentum

HARP collaboration,
Eur.Phys.J.C52(2007)29

HARP data with 8.9 GeV/c proton beam momentum



Majority of pions create neutrinos in MiniBooNE are directly measured by HARP (>80%)

The error on the HARP data ($\sim 7\%$) directly propagates.
The neutrino flux error is the dominant source of normalization error for an absolute cross section in MiniBooNE, however it doesn't affect oscillation analysis.

decay r

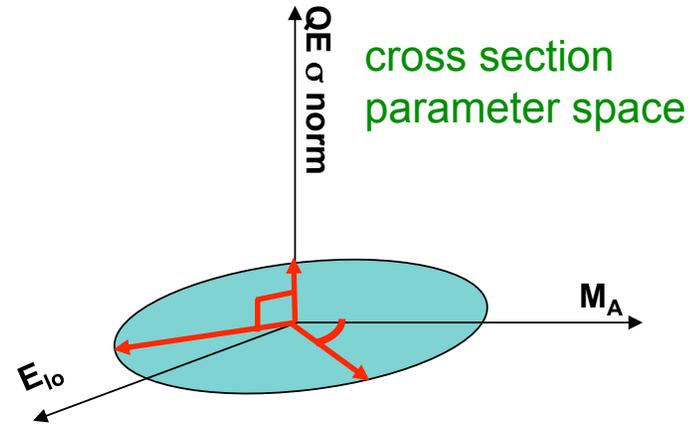
tor

Kator

5. Multisim

ex) cross section uncertainties

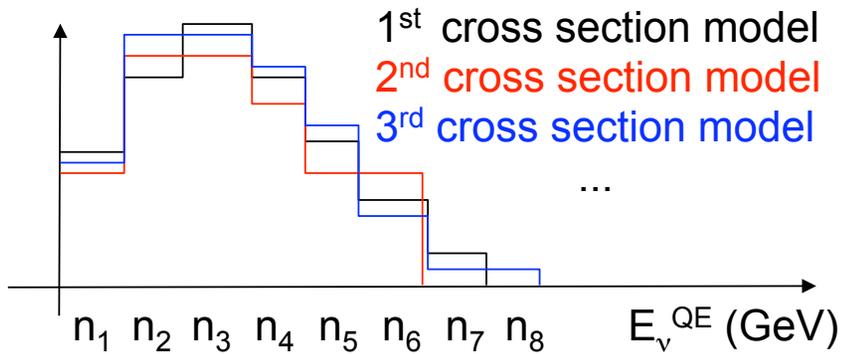
M_A^{QE}	6%	↑ correlated
E_{lo}^{sf}	2%	
QE σ norm	10%	↓ uncorrelated



Input cross section error matrix

$$M_{input}(XS) = \begin{pmatrix} \text{var}(M_A) & \text{cov}(M_A, E_{lo}) & 0 \\ \text{cov}(M_A, E_{lo}) & \text{var}(E_{lo}) & 0 \\ 0 & 0 & \text{var}(\sigma - \text{norm}) \end{pmatrix}$$

cross section error for E_v^{QE}



repeat this exercise many times to create smooth error matrix for E_v^{QE}

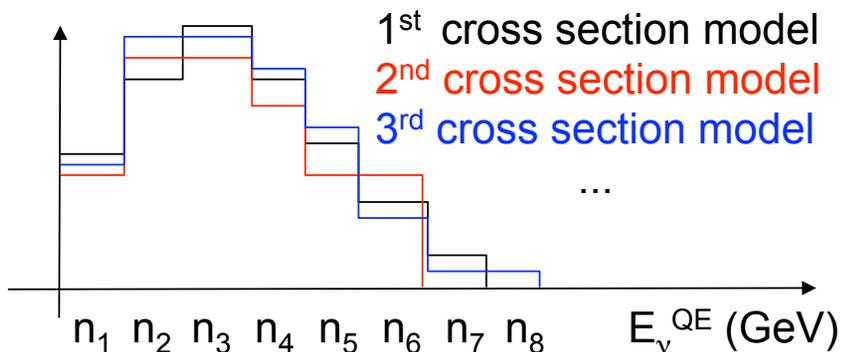
5. Multisim

Output cross section error matrix for E_ν^{QE}

$$[\mathbf{M}_{\text{output}}(\mathbf{xS})]_{ij} \approx \frac{1}{S} \sum_k^S (\mathbf{N}_i^k(\mathbf{xS}) - \mathbf{N}_i^{\text{MC}})(\mathbf{N}_j^k(\mathbf{xS}) - \mathbf{N}_j^{\text{MC}})$$

$$\mathbf{M}_{\text{output}}(\mathbf{xS}) = \begin{pmatrix} \text{var}(n_1) & \text{cov}(n_1, n_2) & \text{cov}(n_1, n_3) & \cdots \\ \text{cov}(n_1, n_2) & \text{var}(n_2) & \text{cov}(n_2, n_3) & \cdots \\ \text{cov}(n_1, n_3) & \text{cov}(n_2, n_3) & \text{var}(n_3) & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

cross section error for E_ν^{QE}



Oscillation analysis use output error matrix for χ^2 fit;

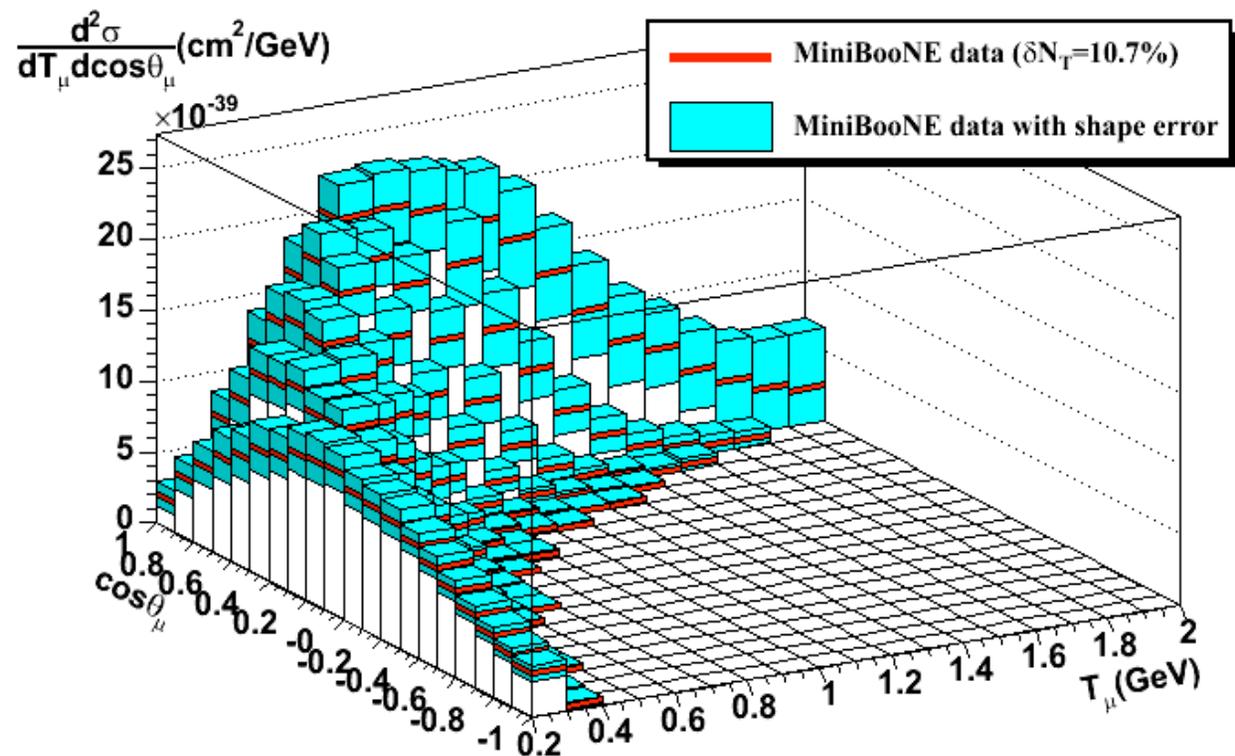
$$\chi^2 = (\text{data} - \text{MC})^T (\mathbf{M}_{\text{output}})^{-1} (\text{data} - \text{MC})$$

6. CCQE double differential cross section

Flux-integrated double differential cross section (T_μ - $\cos\theta$)

This is the most complete information about neutrino cross section based on muon kinematic measurement.

The error shown here is shape error, a total normalization error ($\delta N_T=10.7\%$) is separated.



6. Paradigm shift in neutrino cross section!?

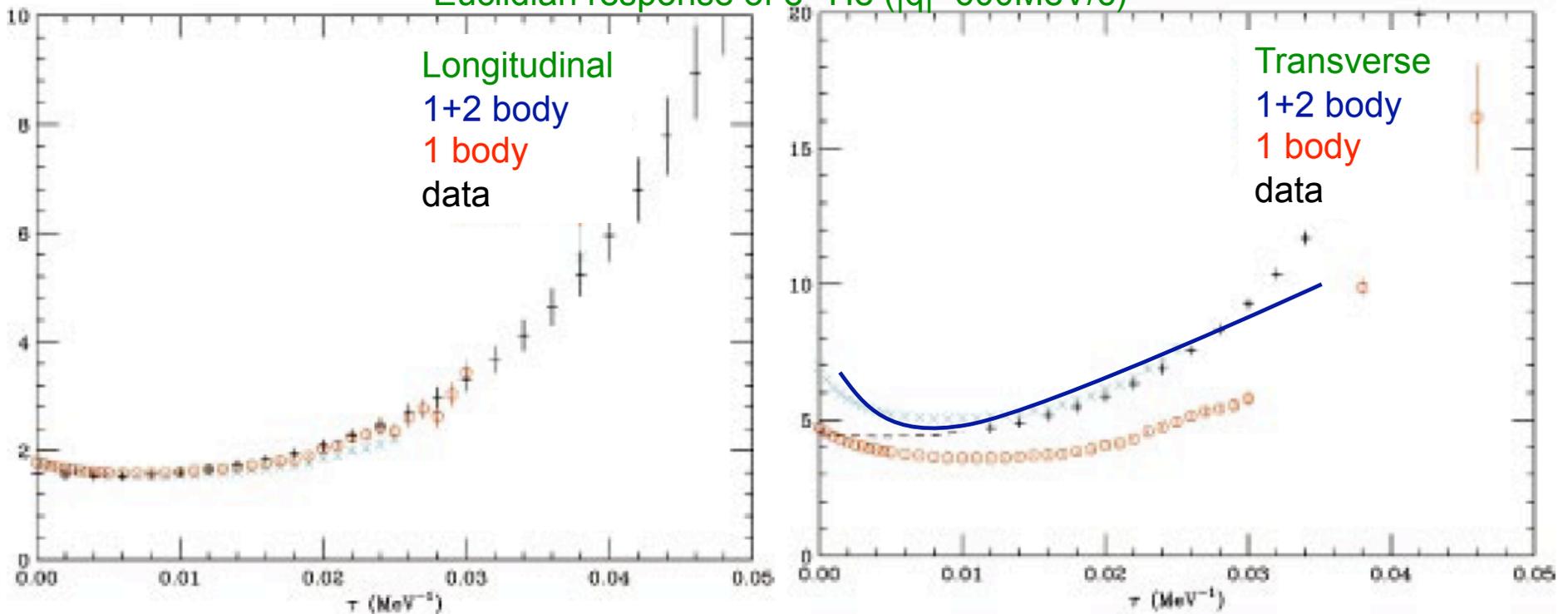
Theoretical approaches for the large cross section and harder Q^2 spectrum

RPA formalism Martini et al., PRC80(2009)065501

SRC+MEC Carlson et al., PRC65(2002)024002

Transverse response is enhanced by presence of short range correlation (SRC) and 2-body current (meson exchange current, MEC).

Euclidian response of $e^{-4}\text{He}$ ($|q|=600\text{MeV}/c$)



Jon Link, Nov. 18, 2005

Fermilab Wine & Cheese seminar

PHYSICAL REVIEW D

VOLUME 23, NUMBER 11

1 JUNE 1981

Quasielastic neutrino scattering: A measurement of the weak nucleon axial-vector form factor

N. J. Baker, A. M. Cnops,* P. L. Connolly, S. A. Kahn, H. G. Kirk, M. J. Murtagh, R. B. Palmer, N. P. Samios, and M. Tanaka

Brookhaven National Laboratory, Upton, New York 11973

(Received 12 February 1981)

The quasielastic reaction $\nu_{\mu} n \rightarrow \mu^{-} p$ was studied in an experiment using the BNL 7-foot deuterium bubble chamber exposed to the wide-band neutrino beam with an average energy of 1.6 GeV. A total of 1138 quasielastic events in the momentum-transfer range $Q^2 = 0.06 - 3.00 \text{ (GeV}/c)^2$ were selected by kinematic fitting and particle identification and were used to extract the axial-vector form factor $F_A(Q^2)$ from the Q^2 distribution. In the framework of the conventional $V - A$ theory, we find that the dipole parametrization is favored over the monopole. The value of the axial-vector mass M_A in the dipole parametrization is $1.07 \pm 0.06 \text{ GeV}$, which is in good agreement with both recent neutrino and electroproduction experiments. In addition, the standard assumptions of conserved vector current and no second-class currents are checked.

We have used a maximum likelihood method to extract M_A from the shape of the Q^2 distribution for each observed neutrino energy. This likelihood function \mathcal{L}^I is independent of the shape of the neutrino spectrum ...

In subsequent cross section analyses the theoretical ("known") quasi-elastic cross section and observed quasi-elastic events were used to determine the flux.

They didn't even try to determine their ν flux from pion production and beam dynamics.

Phys. Rev. D 25, 617 (1982)

The distribution of events in neutrino energy for the $3C \nu d \rightarrow \mu^{-} pp_s$ events is shown in Fig. 4 together with the quasielastic cross section $\sigma(\nu n \rightarrow \mu^{-} p)$ calculated using the standard $V - A$ theory with $M_A = 1.05 \pm 0.05 \text{ GeV}$ and $M_V = 0.84 \text{ GeV}$. The absolute cross sections for the CC interactions have been measured using the quasielastic events and its known cross section.⁴

Jon Link, Nov. 18, 2005

Fermilab Wine & Cheese seminar

VOLUME 49, NUMBER 2

PHYSICAL REVIEW LETTERS

12 JULY 1982

Neutrino Flux and Total Charged-Current Cross Sections in High-Energy Neutrino-Deuterium Interactions

T. Kitagaki, S. Tanaka, H. Yuta, K. Abe, K. Hasegawa, A. Yamaguchi, K. Tamai,
T. Hayashino, Y. Ohtani, and H. Hayano
Tohoku University, Sendai 980, Japan

Fermilab
15ft D₂ Bubble Chamber

To obtain the total cross section from the number of events, the neutrino flux has to be measured on an absolute scale. In this analysis, we determine the neutrino flux using 362 quasielastic events identified in our data¹⁰ and the cross section for reaction (2) derived from the $V-A$ theory.

Again, they use QE events and theoretical cross section to calculate the ν .

When they try to get the flux from meson (π and K) production and decay kinematics they fail miserably for $E_\nu < 30$ GeV.

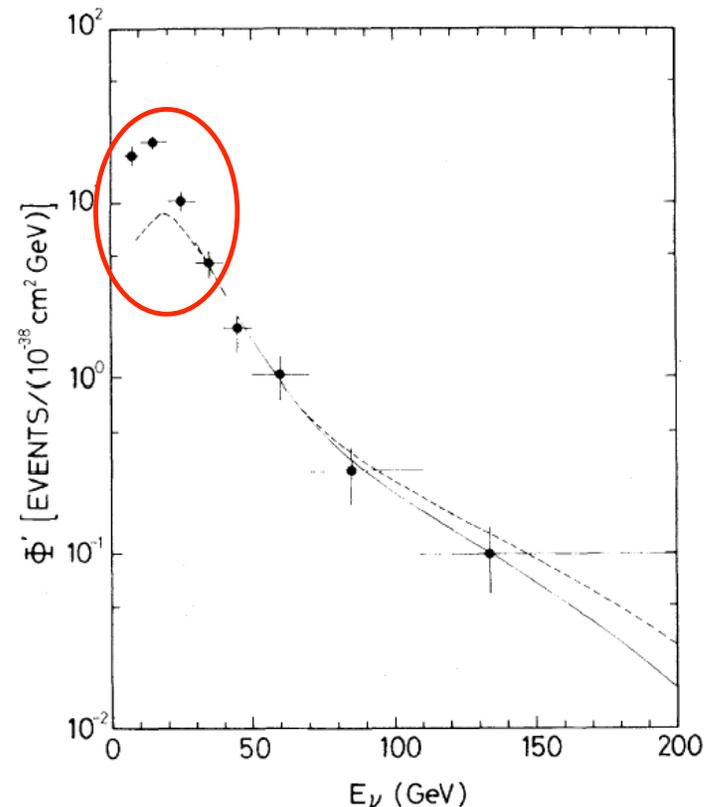


FIG. 2. Neutrino flux distribution obtained from the quasielastic events and the predicted cross section with $M_A = 1.05$ GeV. The solid curve is obtained from the best fit to the flux data for $E_\nu > 30$ GeV. The dashed curve is taken from the Monte Carlo simulation of the flux.

Jon Link, Nov. 18, 2005
Fermilab Wine & Cheese seminar

PHYSICAL REVIEW D

VOLUME 34, NUMBER 1

1 JULY 1986

Determination of the neutrino fluxes in the Brookhaven wide-band beams

L. A. Ahrens, S. H. Aronson, P. L. Connolly,* B. G. Gibbard, M. J. Murtagh, S. J. Murtagh,[†]
S. Terada, and D. H. White

Physics Department, Brookhaven National Laboratory, Upton, New York 11973

Brookhaven
AGS
Liquid Scintillator

The beam calculations described here were based on the Grote, Hagedorn, and Ranft (GHR) (Ref. 11) parametrization; that of Sanford and Wang was used for comparison. An estimate was made of pion production by reinteracting protons guided by the shape of the observed ν_μ spectrum and the observed angular distribution of muons from quasielastic events. The procedure is described¹² in the Appendix.

The Procedure

- Pion production cross sections in some low momentum bins are scaled up by 18 to 79%.
- The K^+ to π^+ ratio is increased by 25%.
- Overall neutrino (anti-neutrino) flux is increased by 10% (30%).

All driven by the neutrino events observed in the detector!

Jon Link, Nov. 18, 2005

Fermilab Wine & Cheese seminar

PHYSICAL REVIEW D

VOLUME 16, NUMBER 11

1 DECEMBER 1977

Study of neutrino interactions in hydrogen and deuterium:

Description of the experiment and study of the reaction $\nu + d \rightarrow \mu^- + p + p_s$

S. J. Barish,* J. Campbell,† G. Charlton,§ Y. Cho, M. Derrick, R. Engelmann,|| L. G. Hyman, D. Jankowski, A. Mann,|| B. Musgrave, P. Schreiner, P. F. Schultz, R. Singer, M. Szczekowski,** T. Wangler, and H. Yuta††

Argonne National Laboratory, Argonne, Illinois 60439

Argonne ZGS
12ft D₂ Bubble Chamber

Flux derived from pion production data. Were able to test assumptions about the form of the cross section using absolute rate and shape information.

TABLE IV. Results of axial-form-factor fits.

Likelihood function	M_A^{Dipole} (GeV)	M_A^{Monopole} (GeV)	M_A^{Tripole} (GeV)
Rate	$0.75^{+0.13}_{-0.11}$	$0.45^{+0.11}_{-0.07}$	$0.96^{+0.17}_{-0.14}$
Shape	1.010 ± 0.09	0.56 ± 0.08	1.32 ± 0.11
Rate and shape	0.95 ± 0.09	0.52 ± 0.08	1.25 ± 0.11
Flux independent	0.95 ± 0.09	0.53 ± 0.08	1.25 ± 0.11

- Pion production measured in ZGS beams were used in this analysis
- A very careful job was done to normalize the beam.
- Yet they have a 25% inconsistency between the axial mass they measure considering only rate information verses considering only spectral information.

Interpretation: Their normalization is wrong.

4. MiniBooNE cross section results

NuInt09, May18-22, 2009, Sitges, Spain

All talks proceedings are available on online (open access),
<http://proceedings.aip.org/proceedings/confproceed/1189.jsp>

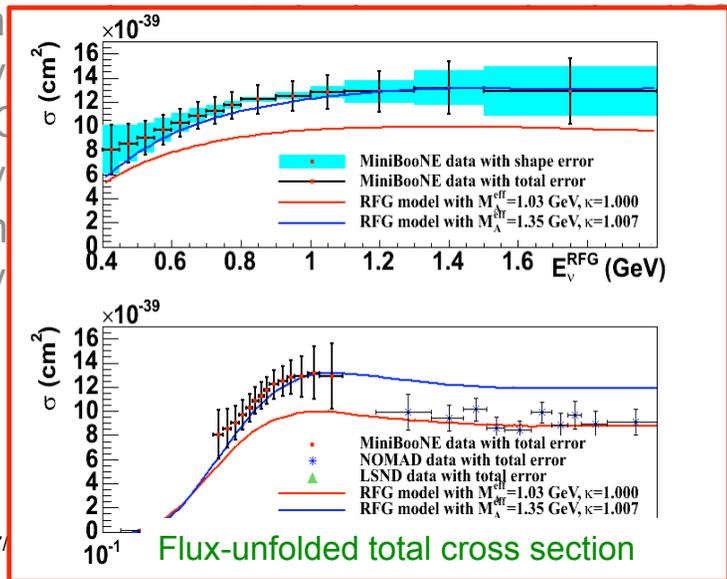
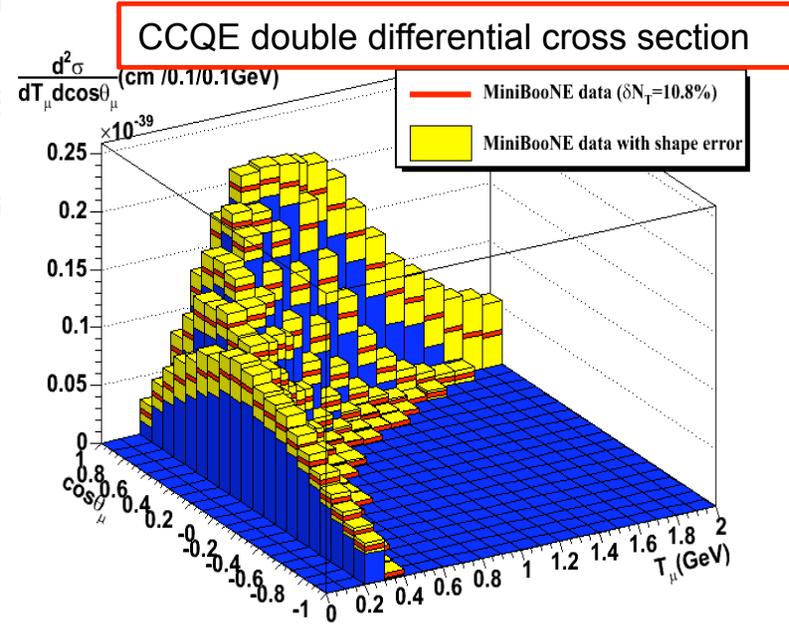


1. charged current quasielastic (CCQE) cross section measurement by Teppei Katori, PRD81(2010)092005
2. neutral current elastic (NCE) cross section measurement by Denis Perevalov, PRD82(2010)092005
3. neutral current π^0 production (NC π^0) cross section measurement (by and anti-)

$$\nu_\mu + n \rightarrow p + \mu^-$$

$$(\nu_\mu + {}^{12}\text{C} \rightarrow X + \mu^-)$$

- first double differential cross section measurement
- observed large absolute cross section



03/17/

pei Katori, MIT

by Denis Perevalov

4. MiniBooNE cross section results

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3. neutral current π^0 production (NC π^0) cross section measurement (ν and anti- ν)
by Colin Anderson, PRD81(2010)013005

4. charged current single pion production (CC π^+) cross section measurement
by Mike Wilking, arXiv:1011.3572

- highest statistics cross section measurement
- new Δ s (strange quark spin) extraction method

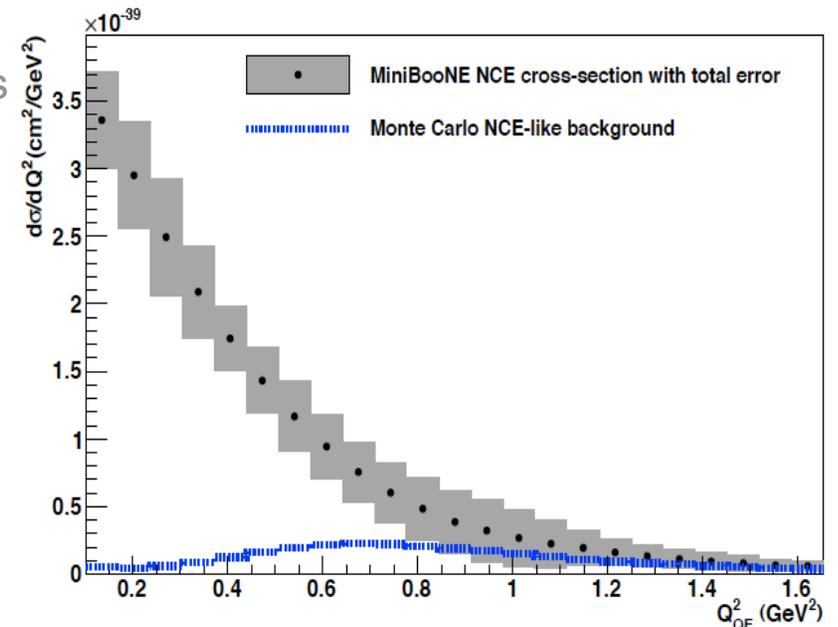
6. CC π^0 /CCQE cross section ratio measurement
by Steve Linden, PRL103(2009)081801

7. anti- ν CCQE measurement
by Joe Grange, arXiv:1102.1964

$$\nu_{\mu} + p \rightarrow \nu_{\mu} + p$$

$$\nu_{\mu} + n \rightarrow \nu_{\mu} + n$$

Flux-averaged NCE p+n differential cross section



by Colin Anderson

4. MiniBooNE cross section results

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by Mike Wilking, arXiv:1011.3572

5. charged current single ρ production (CC ρ) cross section measurement

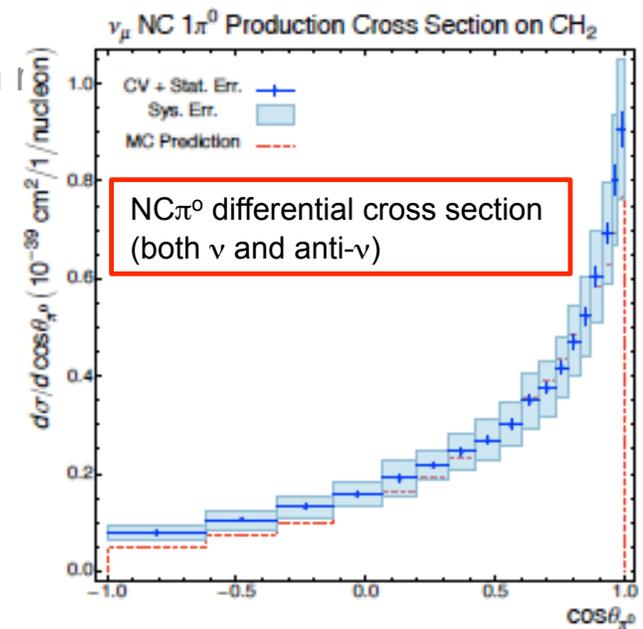
- first differential cross section measurement

6. - observed large absolute cross section

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7. anti- ν CCQE measurement

by Joe Grange, arXiv:1102.1964



by Mike Wilking

4. MiniBooNE cross section results

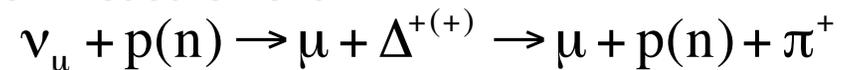
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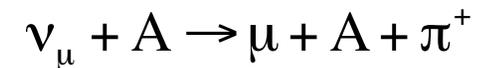
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3. neutral current π^0 production (NC π^0) cross section measurement (ν and anti- ν)

by Colin Anderson, PRD81(2010)013005

4. charged current single pion production (CC π^+) cross section measurement

by Mike Wilking, [arXiv:1011.3572](https://arxiv.org/abs/1011.3572)

- first double differential cross section measurement
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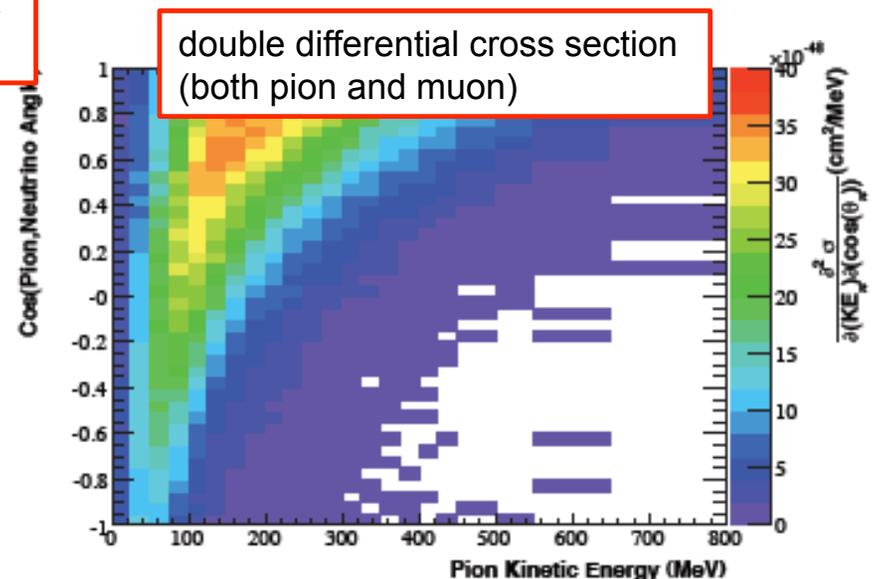
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cross section measurement



by Bob Nelson



4. MiniBooNE cross section results

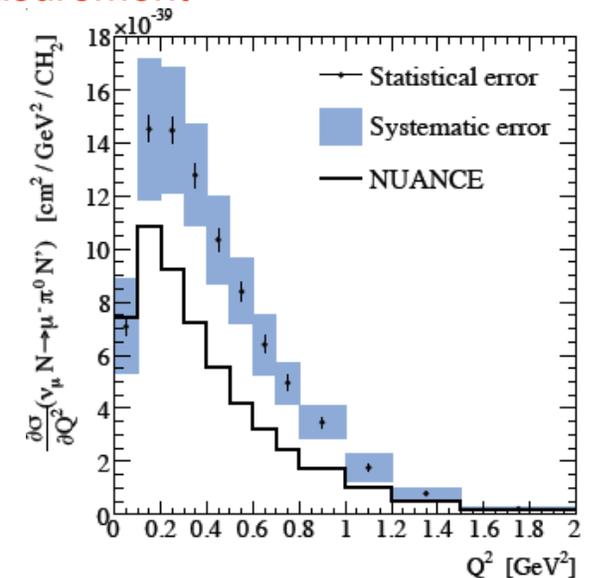
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by Bob Nelson, arXiv:1010.3264
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CC π^0 Q^2 differential cross section



by Steve Linden

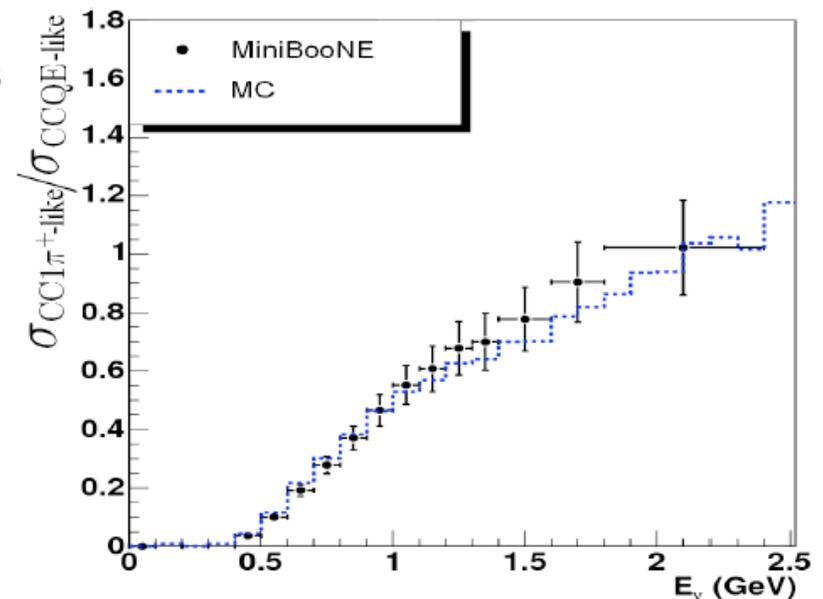


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6. CC π^+ /CCQE cross section ratio measurement
by Steve Linden, PRL103(2009)081801
7. a - data is presented in theorist friendly style
by Bob Strange, arXiv:1102.1504



CC π^+ like/CCQElike cross section ratio

by Joe Grange



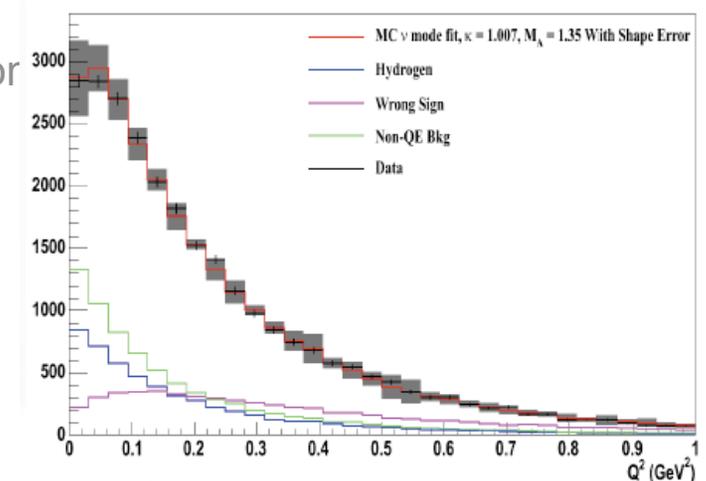
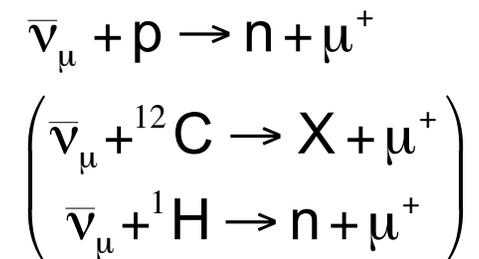
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7. anti- ν CCQE measurement by Joe Grange, [arXiv:1102:1964](https://arxiv.org/abs/1102.1964)

- highest statistics in this channel
- support neutrino mode result
- new method to measure neutrino contamination

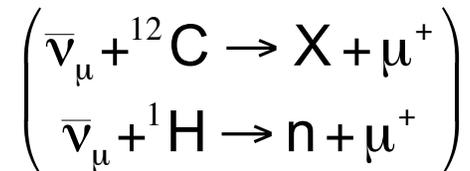
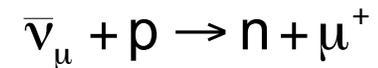


anti- ν CCQE Q^2 distribution

4. MiniBooNE cross section results

NuInt11, Mar. 07-11, 2011, Dehradun, India

In NuInt11, MiniBooNE presented 2 anti-neutrino measurement results.



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5. charged current single π^0 production (CC π^0) cross section measurement
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6.
 - background subtracted absolute rate
 - key to understand nuclear effect
7.
 - key to understand role of axial form factor
8. new anti- ν CCQE measurement
by Joe Grange, arXiv:1102.1504
9. anti- ν NCE measurement
by Ranjan Dharmapalan

4. MiniBooNE cross section results

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by Lee Grange, arXiv:1102.1064

$$\bar{\nu}_\mu + p \rightarrow \bar{\nu}_\mu + p$$

$$\bar{\nu}_\mu + n \rightarrow \bar{\nu}_\mu + n$$

- last channel to measure in non-resonance CC/NC interaction
- understanding of role of strange quark in nucleon static limit

9. anti- ν NCE measurement
by Ranjan Dharmapalan