

Charged Current Interaction measurements in MiniBooNE

hep-ex/XXX

Teppei Katori for the MiniBooNE collaboration
Indiana University
NuInt 07, Fermilab, May., 31, 07

Teppei Katori, Indiana University,
NuInt '07



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outline

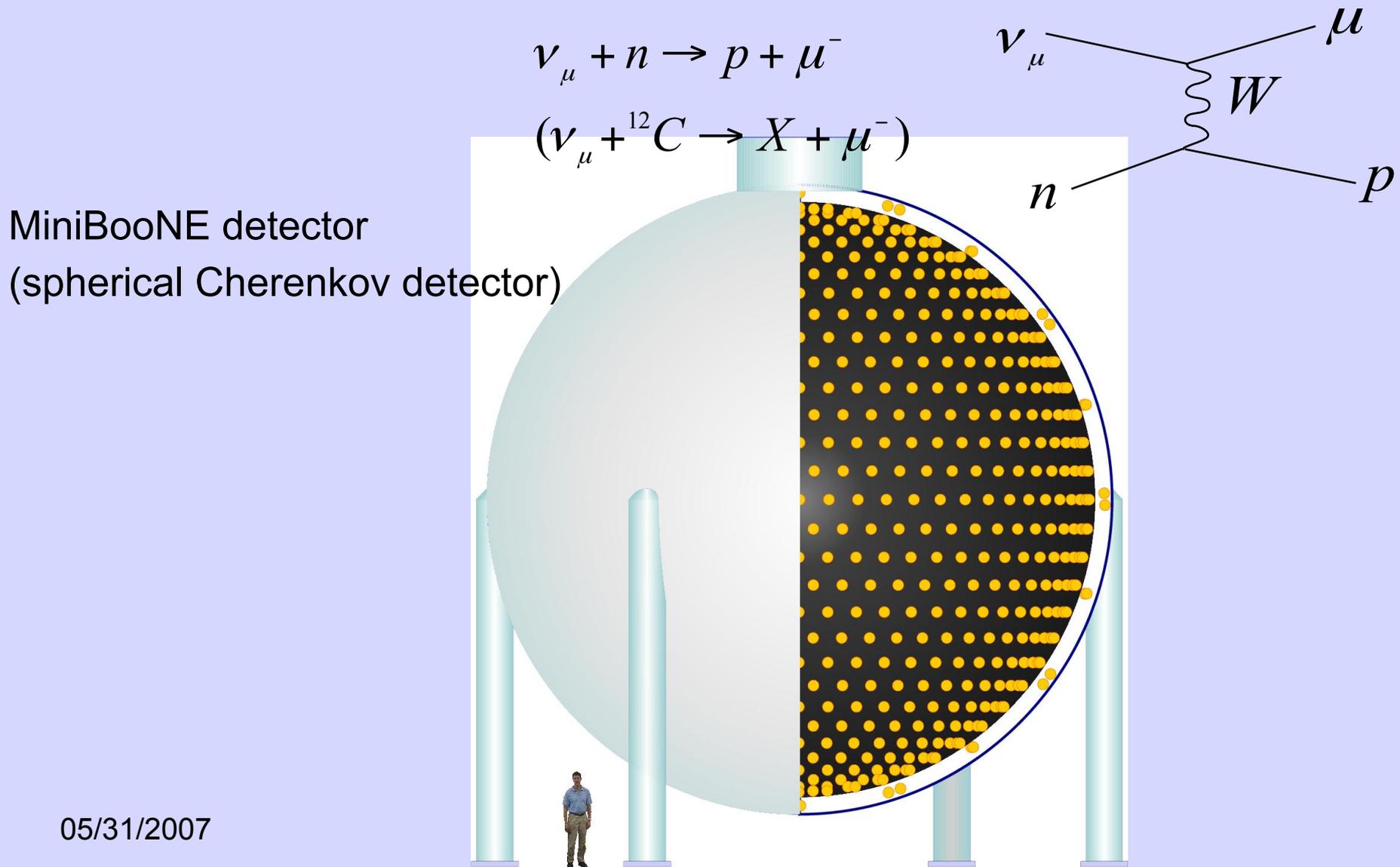
1. CCQE events in MiniBooNE
2. Prediction for CCQE events
3. CCQE data-MC comparison
4. Fit results
5. Anti-neutrino CCQE events
6. Conclusion



1. CCQE events in MiniBooNE

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ν_μ charged current quasi-elastic (ν_μ CCQE) interaction is the most abundant (~40%) and the fundamental interaction in MiniBooNE detector

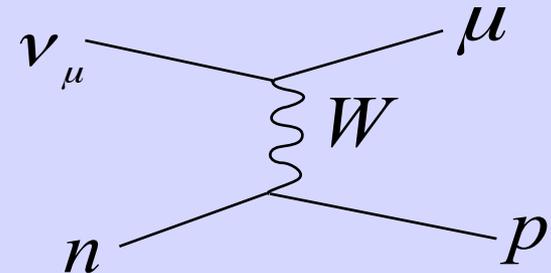


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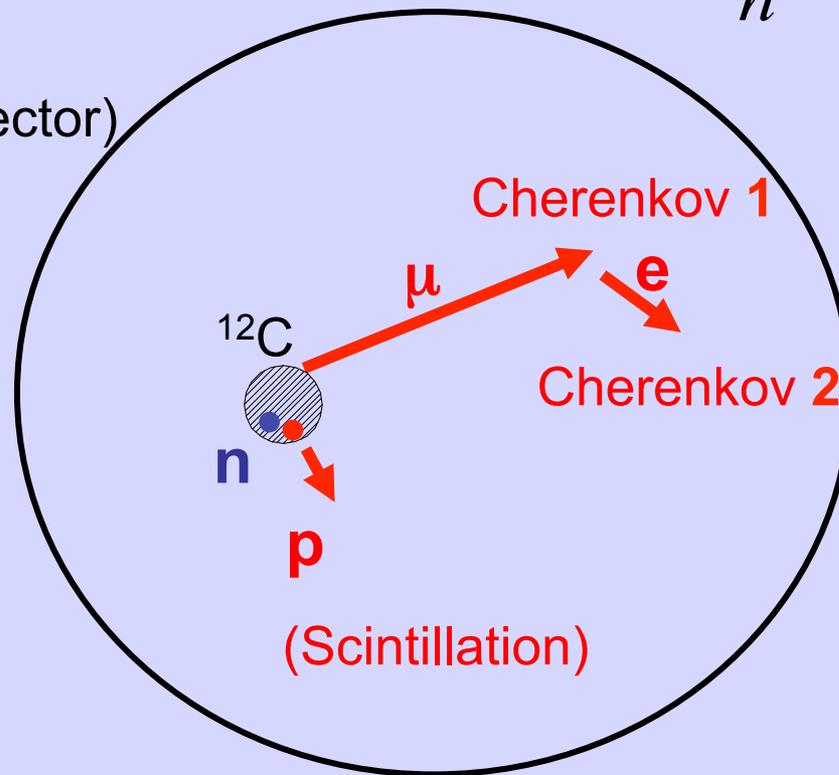
$$\nu_\mu + n \rightarrow p + \mu^-$$

$$(\nu_\mu + {}^{12}\text{C} \rightarrow X + \mu^-)$$



MiniBooNE detector
(spherical Cherenkov detector)

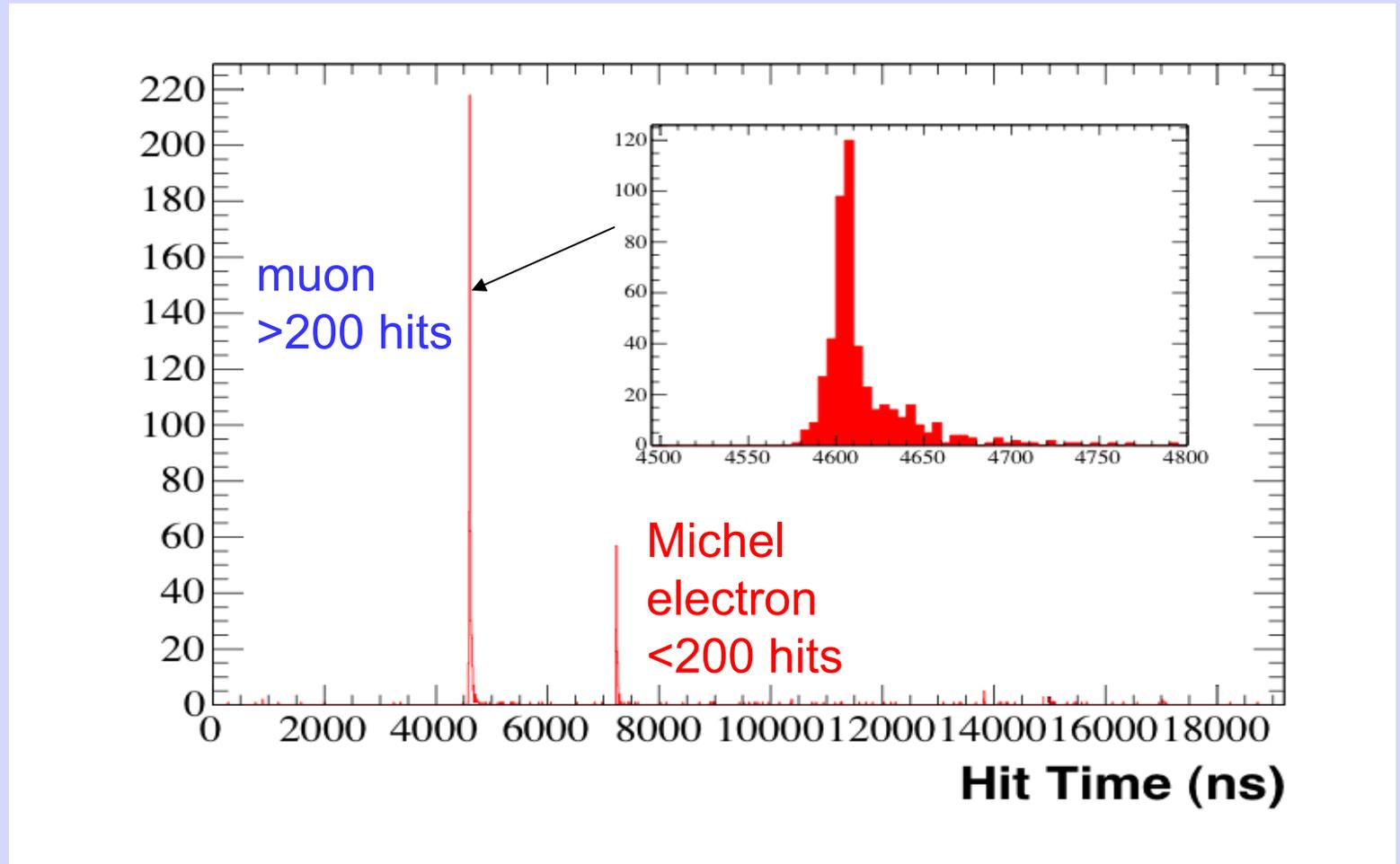
ν -beam



muon like Cherenkov light and subsequent decayed electron (Michel electron) like Cherenkov light are the signal of CCQE event

1. CCQE events in MiniBooNE

ν_μ CCQE interactions ($\nu+n \rightarrow \mu+p$) has characteristic two “subevent” structure from muon decay



35.0% cut efficiency

197,308 events with 5.58E20POT

1. CCQE events in MiniBooNE

Cut and efficiency summary

total 2 subevents	54.2%
muon in beam window ($4400\text{ns} < \text{Time} < 6400\text{ns}$)	52.9%
muon veto hits < 6 and Michel electron veto hits < 6	46.4%
muon tank hits > 200 and Michel electron tank hits < 200	41.6%
fiducial reconstruction for muon	41.3%
muon and electron distance $< 100\text{cm}$	35.0%

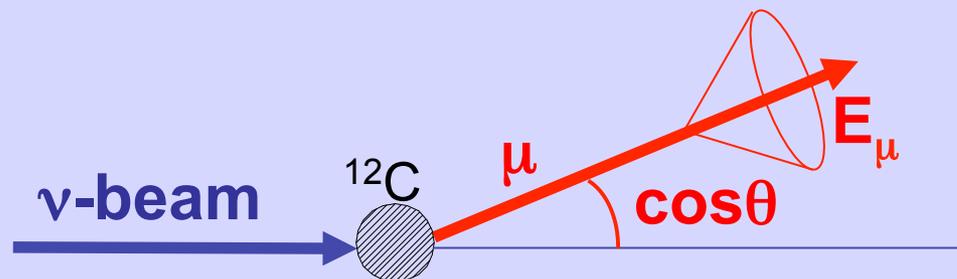
1. CCQE events in MiniBooNE

All kinematics are specified from 2 observables, muon energy E_μ and muon scattering angle θ

Energy of the neutrino E_ν and 4-momentum transfer Q^2 can be reconstructed by these 2 observables

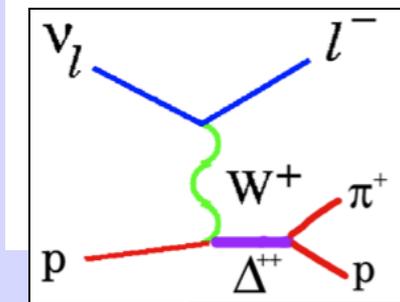
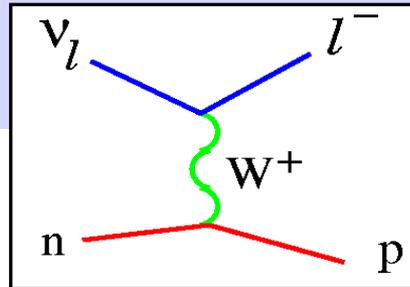
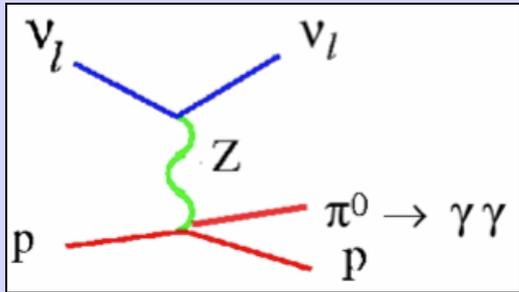
$$E_\nu = \frac{2(M - E_B)E_\mu - (E_B^2 - 2ME_B + m_\mu^2 + \Delta M^2)}{2[(M - E_B) - E_\mu + p_\mu \cos\theta_\mu]}$$

$$Q^2 = -m_\mu^2 + 2E_\nu(E_\mu - p_\mu \cos\theta_\mu)$$

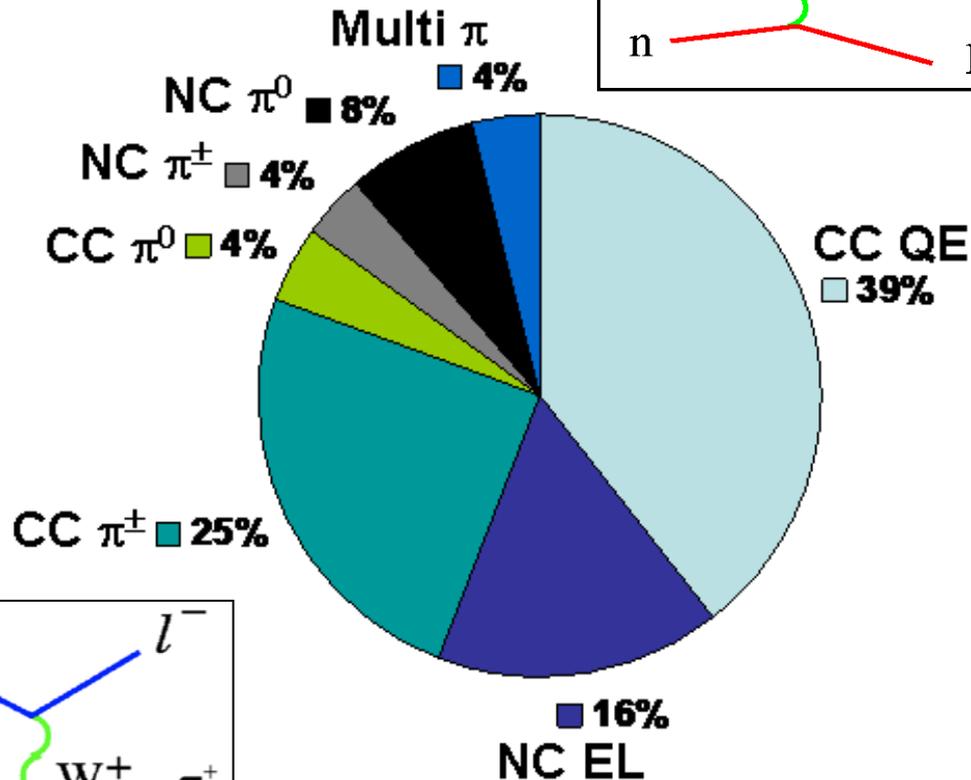


2. Prediction for CCQE events

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Predicted event rates
(NUANCE Monte Carlo)



Casper, Nucl.Phys.Proc.Suppl.
112 (2002) 161

2. Prediction for CCQE events

Smith and Moniz,
Nucl.,Phys.,B43(1972)605

Relativistic Fermi Gas (RFG) Model

Carbon is described by the collection of incoherent Fermi gas particles.

All details come from hadronic tensor.

$$(W_{\mu\nu})_{ab} = \int_{E_{lo}}^{E_{hi}} f(\vec{k}, \vec{q}, w) T_{\mu\nu} dE : \text{hadronic tensor}$$

$f(\vec{k}, \vec{q}, w)$: nucleon phase space density function

$T_{\mu\nu} = T_{\mu\nu}(F_1, F_2, F_A, F_P)$: nucleon tensor

$F_A(Q^2) = g_A / (1 + Q^2/M_A^2)^2$: Axial form factor

E_{hi} : the highest energy state of nucleon = $\sqrt{(\mathbf{p}_F^2 + M^2)}$

E_{lo} : the lowest energy state of nucleon = $\sqrt{(\mathbf{p}_F^2 + M^2)} - w + E_B$

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E_{hi} : the highest energy state of nucleon = $\sqrt{(p_F^2 + M^2)}$

E_{lo} : the lowest energy state of nucleon = $\sqrt{(p_F^2 + M^2)} - w + E_B$

3 parameters are especially important to control nuclear effect of Carbon;

$M_A = 1.03\text{GeV}$: axial mass

$p_F = 220\text{MeV}$: Fermi momentum

$E_B = 34\text{MeV}$: binding energy

3. CCQE data-MC comparison

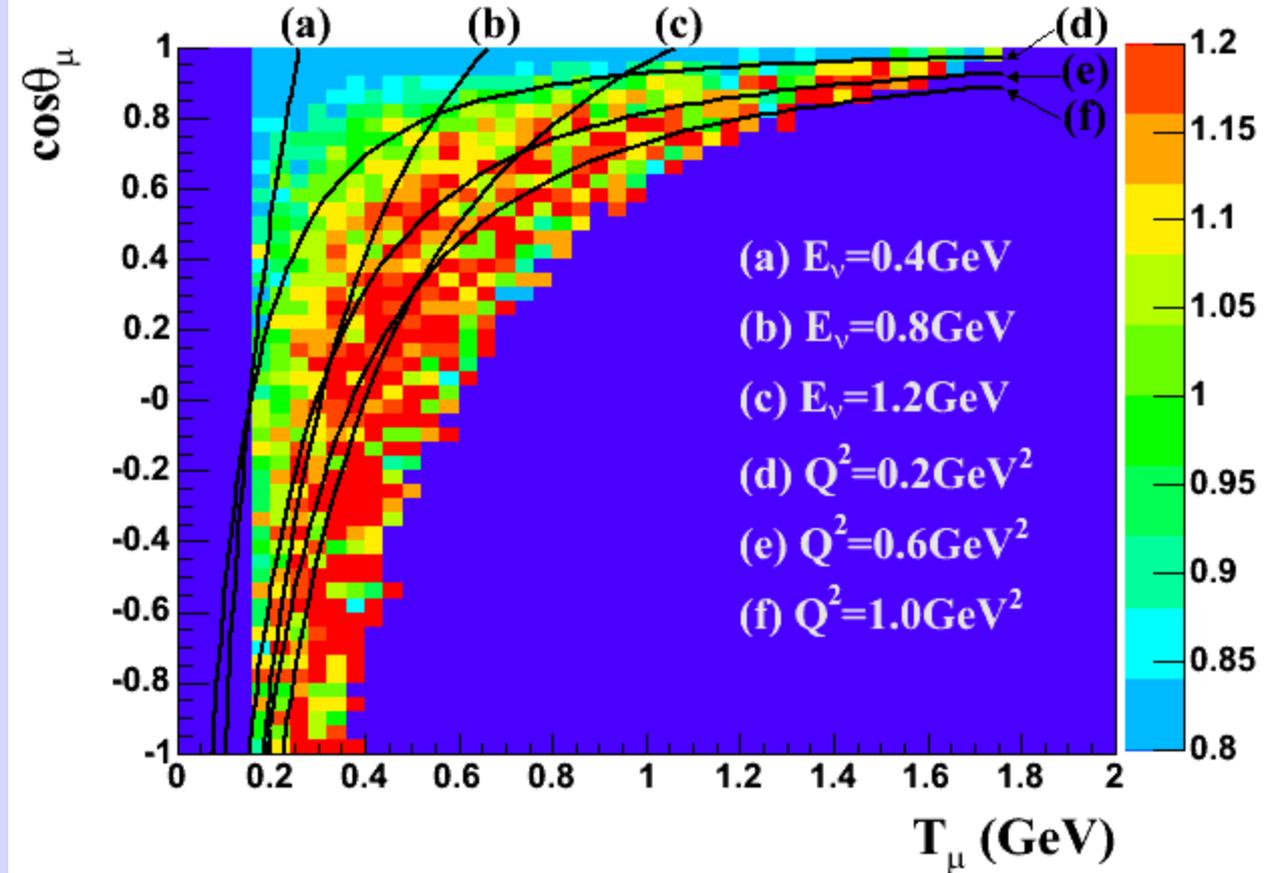
3. CCQE data-MC comparison

CCQE kinematics phase space

The data-MC agreement is not great

Since data-MC disagreements align on the Q^2 lines, not E_ν lines, the source of data-MC disagreement is not the neutrino beam prediction, but the neutrino cross section prediction.

data-MC ratio from RFG model



3. CCQE data-MC comparison

CCQE kinematics phase space

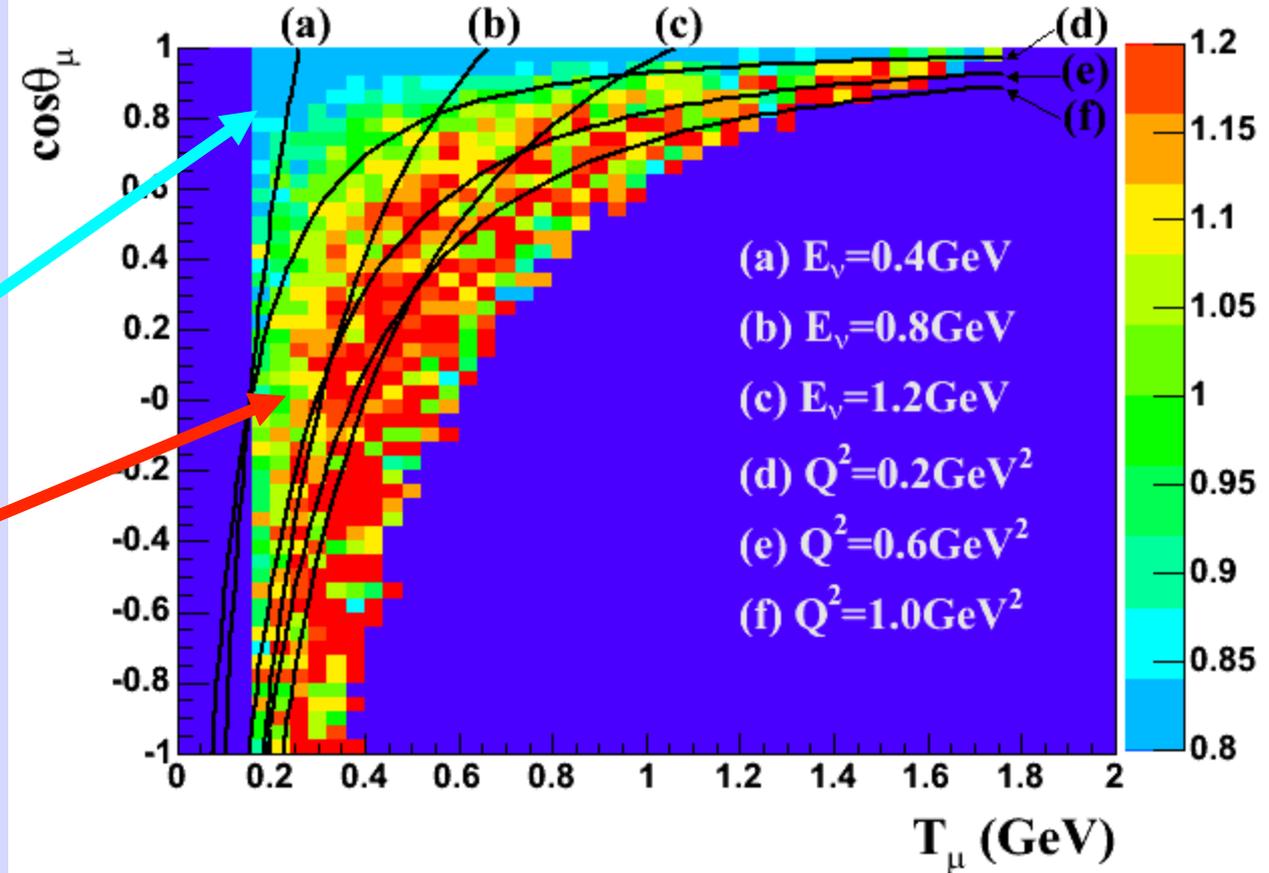
The data-MC agreement is not great

The data-MC disagreement is characterized by 2 features;

(1) data deficit at low Q^2 region

(2) data excess at high Q^2 region

data-MC ratio from RFG model



3. CCQE data-MC comparison

Nuclear model parameters are tuned from electron scattering data, thus the best explanations of observed data-MC disagreements are something one cannot measure from the electron scattering data

(1) data deficit at low Q^2 region
→ Pauli blocking

(2) data excess at high Q^2 region
→ Axial mass M_A

We tune the nuclear parameters in RFG model using Q^2 distribution;

$M_A = \text{tuned}$

$P_F = \text{fixed}$

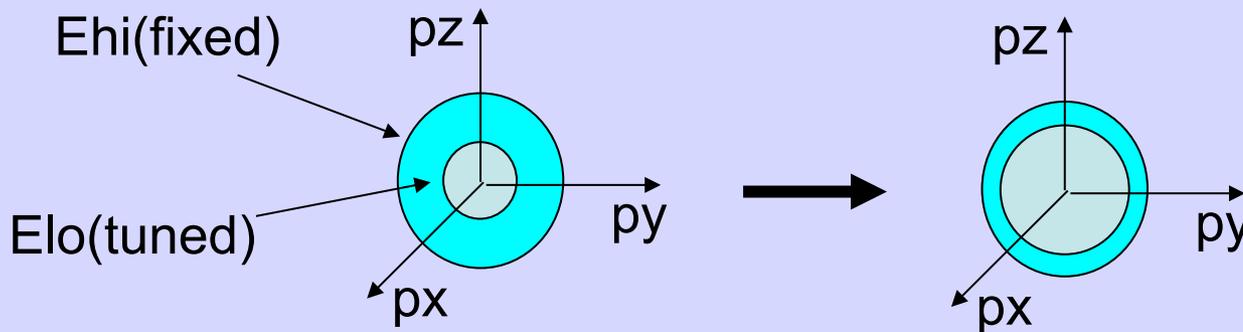
$E_B = \text{fixed}$

3. CCQE data-MC comparison

Pauli blocking parameter "kappa" : κ

To enhance the Pauli blocking at low Q^2 , we introduced a new parameter κ , which is the scale factor of lower bound of nucleon sea and controls the size of nucleon phase space

$$E_{lo} = \kappa \left(\sqrt{(\mathbf{p}_F^2 + M^2)} - w + E_B \right)$$



This modification gives significant effect only at low Q^2 region

We tune the nuclear parameters in RFG model using Q^2 distribution;

M_A = tuned

P_F = fixed

E_B = fixed

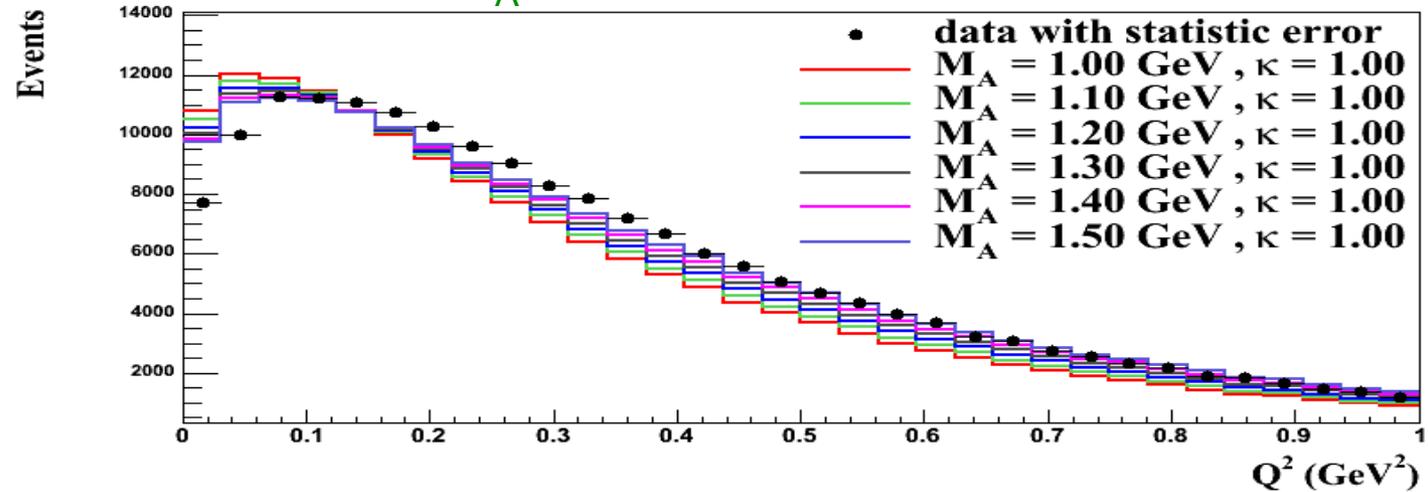
κ = tuned

3. CCQE data-MC comparison

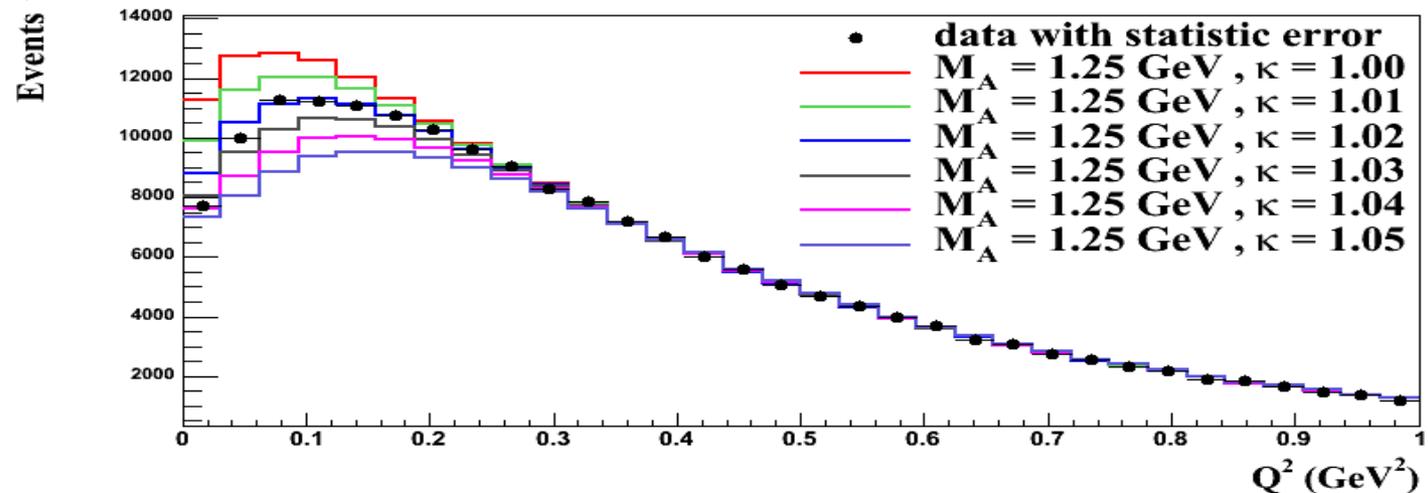
M_A and κ are simultaneously fit to the data

2% change of κ is sufficient to take account the data deficit at low Q^2 region

Q^2 distribution with M_A variation



Q^2 distribution with κ variation



4. Fit results

4. Fit results

Least χ^2 fit for Q^2 distribution

$$\chi^2 = (\text{data} - \text{MC})^T (M_{\text{total}})^{-1} (\text{data} - \text{MC})$$

χ^2 minimum is found by global scan of shape only fit with $0.0 < Q^2 (\text{GeV}^2) < 1.0$

Input error matrices

keep the correlation of systematics

		dependent
		←————→
↑ independent ↓	π^+ production	(8 parameters)
	π^- production	(8 parameters)
	K^+ production	(7 parameters)
	K^0 production	(9 parameters)
	beam model	(8 parameters)
	cross section	(20 parameters)
	detector model	(39 parameters)

The total output error matrix

keep the correlation of Q^2 bins

$$\begin{aligned} M_{\text{total}} = & M(\pi^+ \text{ production}) \\ & + M(\pi^- \text{ production}) \\ & + M(K^+ \text{ production}) \\ & + M(K^0 \text{ production}) \\ & + M(\text{beam model}) \\ & + M(\text{cross section model}) \\ & + M(\text{detector model}) \\ & + M(\text{data statistics}) \end{aligned}$$

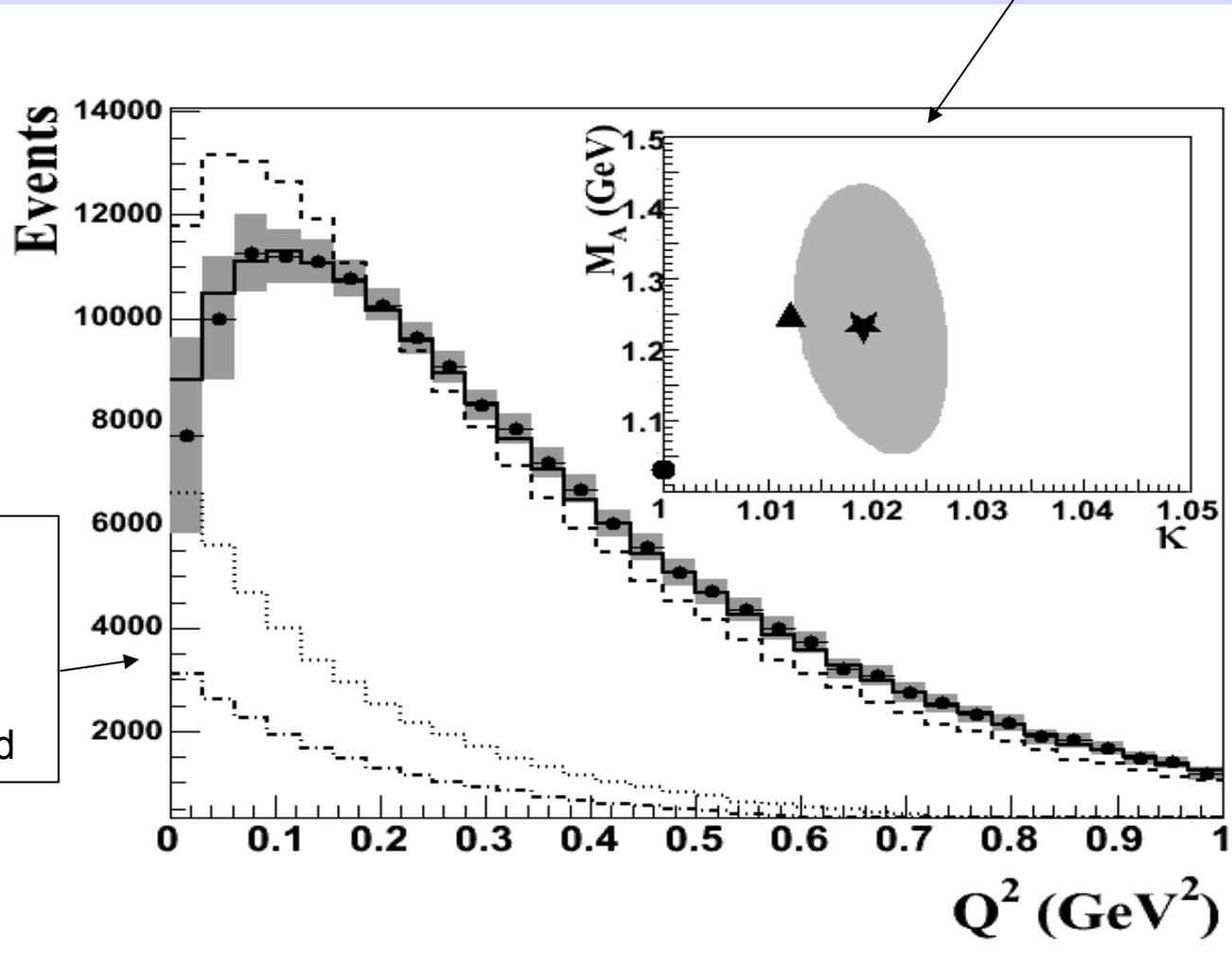
4. Fit results

$M_A - \kappa$ fit result

$$M_A = 1.23 \pm 0.20(\text{stat+sys})$$

$$\kappa = 1.019 \pm 0.011(\text{stat+sys})$$

circle: before fit
star: after fit with 1-sigma contour
triangle: bkgd shape uncertainty



dots : data with error bar
dashed line : before fit
solid line : after fit
dotted line : background
dash-dotted : non-CCQElike bkgd

4. Fit results

Errors

The detector model uncertainty dominates the error in M_A

The error on κ is dominated by Q2 shape uncertainty of background events

	$\delta M_A(\text{GeV})$	$\delta \kappa$
data statistics	0.03	0.003
neutrino flux	0.04	0.003
neutrino cross section	0.06	0.004
detector model	0.10	0.003
CC π^+ background shape	0.02	0.007
total error	0.20	0.011

4. Fit results

M_A - κ fit result

$$M_A = 1.23 \pm 0.20(\text{stat+sys})$$

$$\kappa = 1.019 \pm 0.011(\text{stat+sys})$$

Although fit is done in Q^2 distribution, entire CCQE kinematics is improved

before

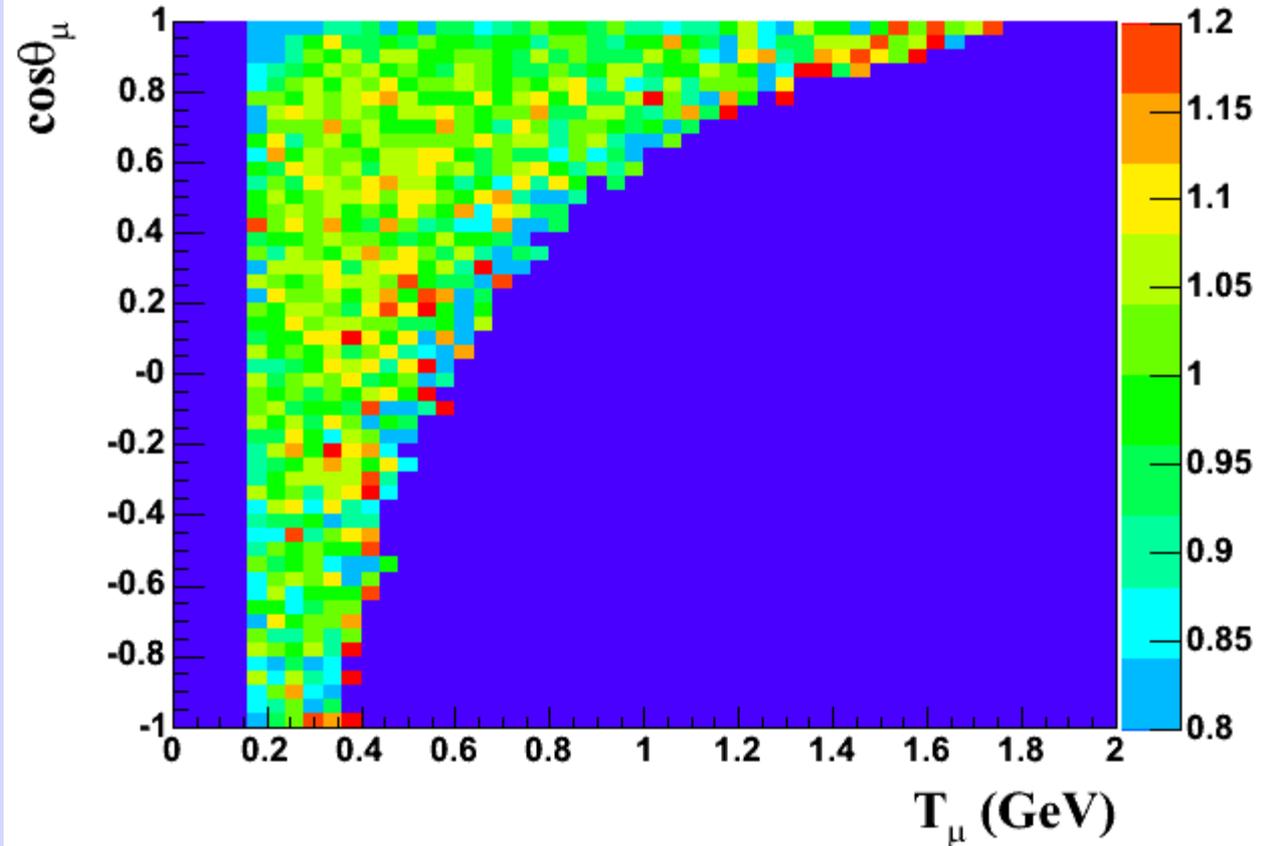
$$\chi^2/\text{dof} = 79.5/53, P(\chi^2) = 1\%$$

↓

after

$$\chi^2/\text{dof} = 45.1/53, P(\chi^2) = 77\%$$

data-MC ratio after the fit



4. Fit results

$M_A - \kappa$ fit result

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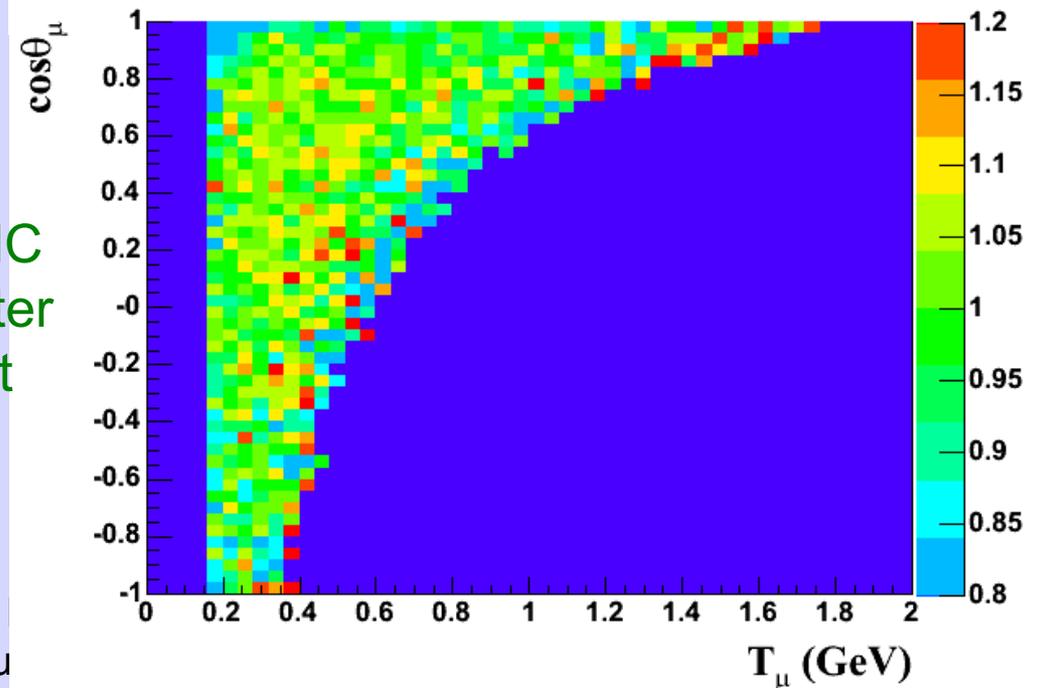
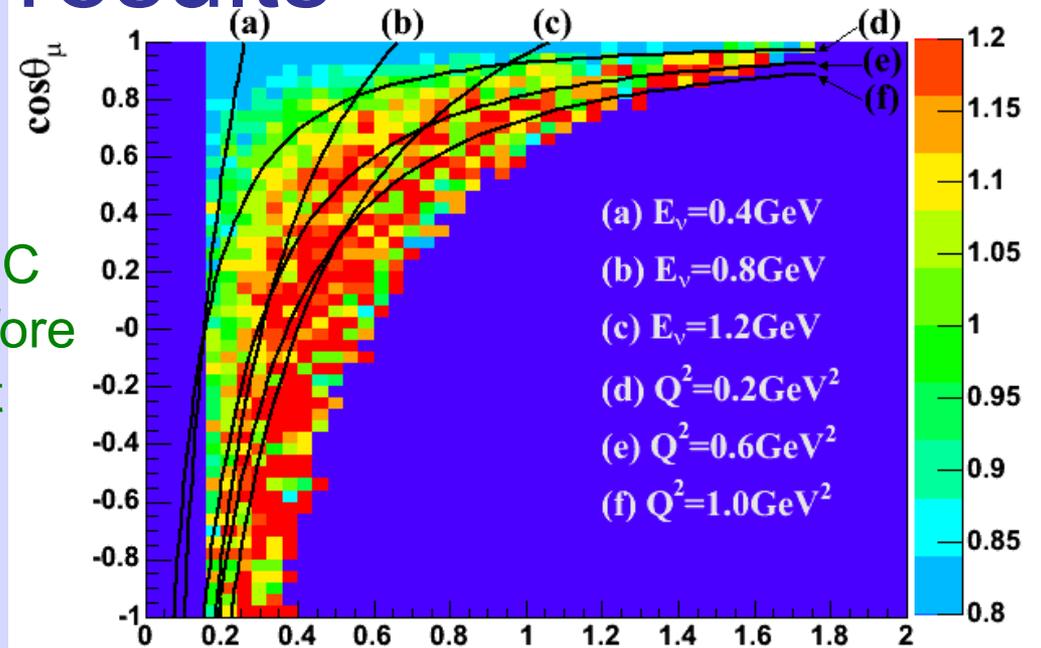


after

$$\chi^2/\text{dof} = 45.1/53, P(\chi^2) = 77\%$$

data-MC
ratio before
the fit

data-MC
ratio after
the fit

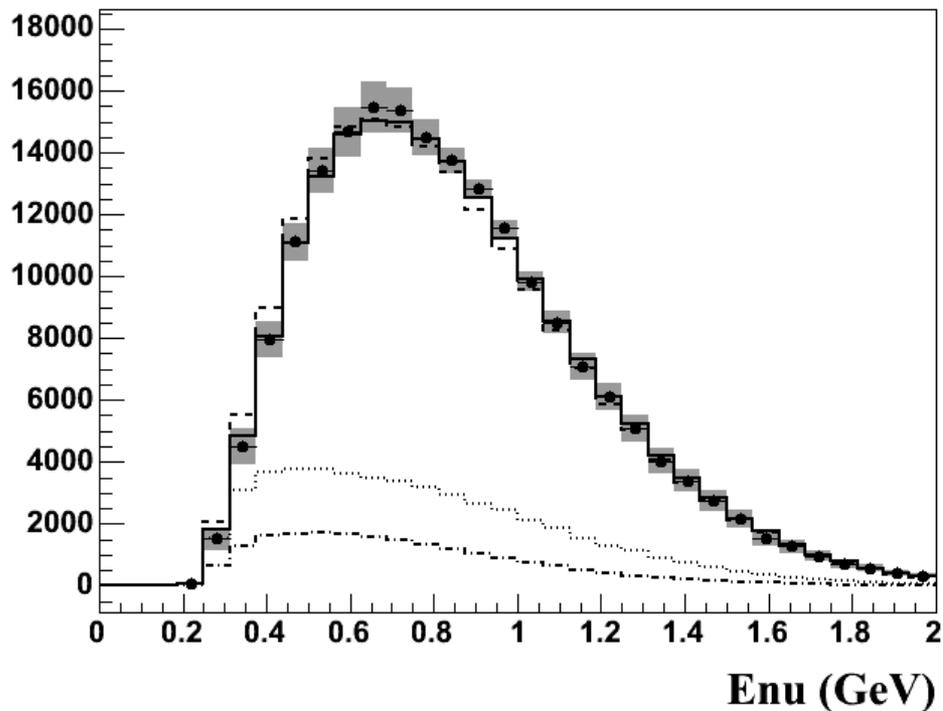


4. Fit results

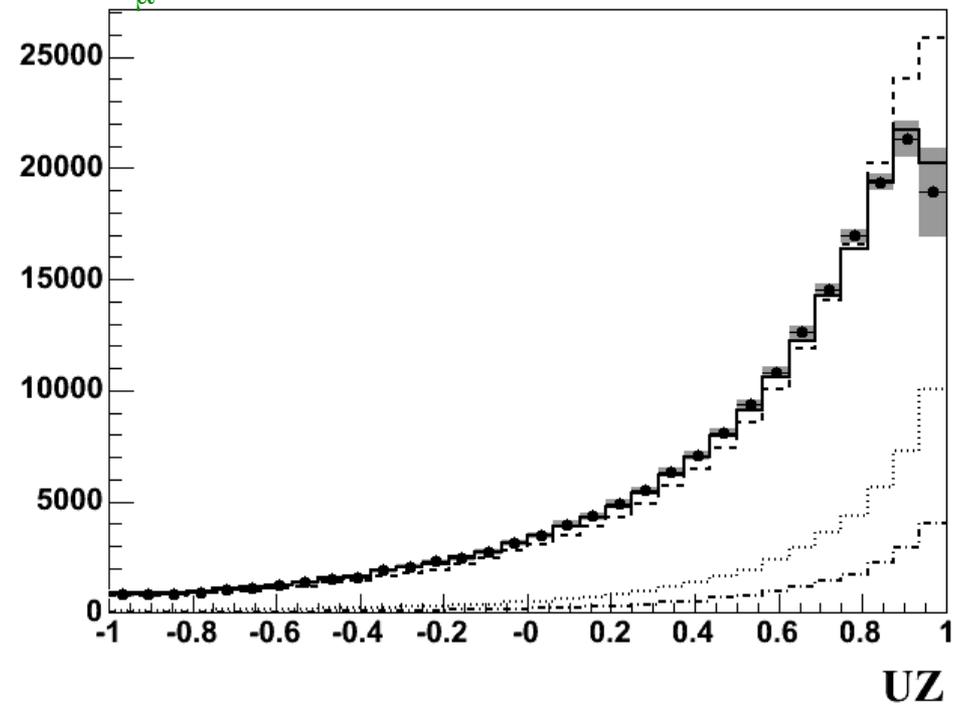
Other kinematics distribution also show very good data-MC agreement
(This is critical for MiniBooNE neutrino oscillation search experiment)

MiniBooNE collaboration,
arXiv:0704.1500 [hep-ex] (2007)

E_ν distribution



$\cos\theta_\mu$ distribution



4. Fit results

M_A only fit result

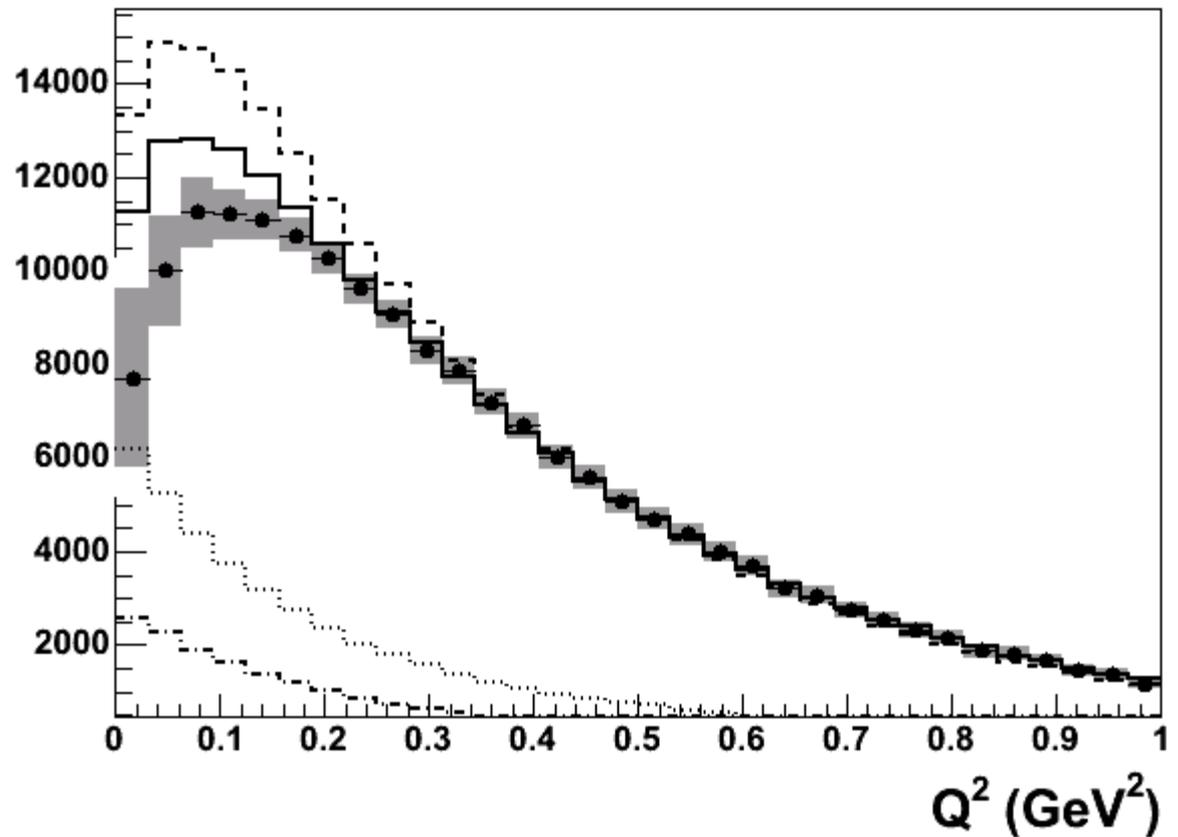
$$M_A = 1.25 \pm 0.12(\text{stat+sys})$$

fit with fixing κ for
 $0.25 < Q^2(\text{GeV}^2) < 1.0$

good agreement above
 0.25GeV^2 but gross
disagreement at low Q^2
region

This fit cannot improve
entire CCQE phase
space

Q^2 distribution



4. Fit results

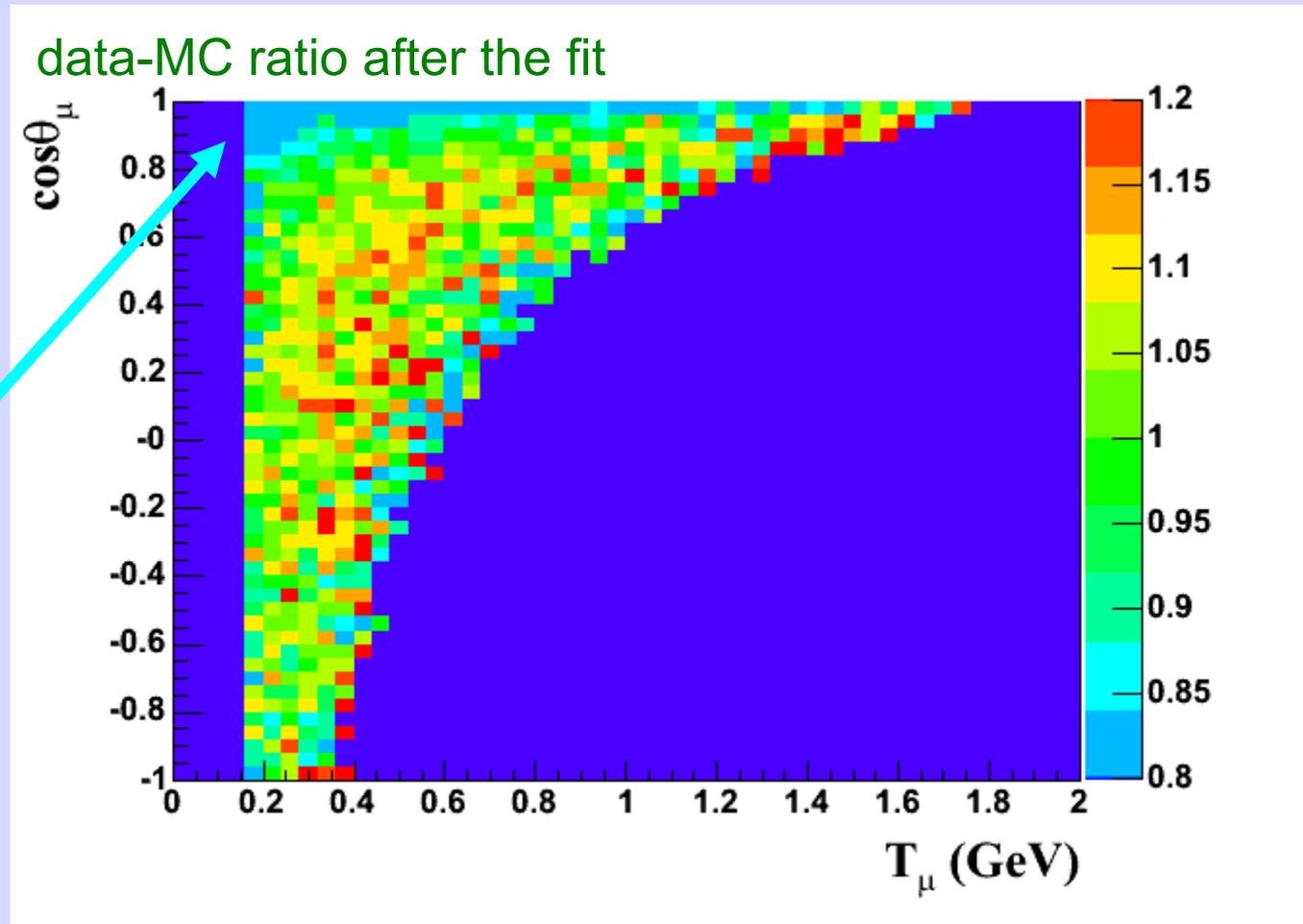
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5. Anti-neutrino CCQE events

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Anti-neutrino Q^2 distribution

MiniBooNE anti-neutrino CCQE

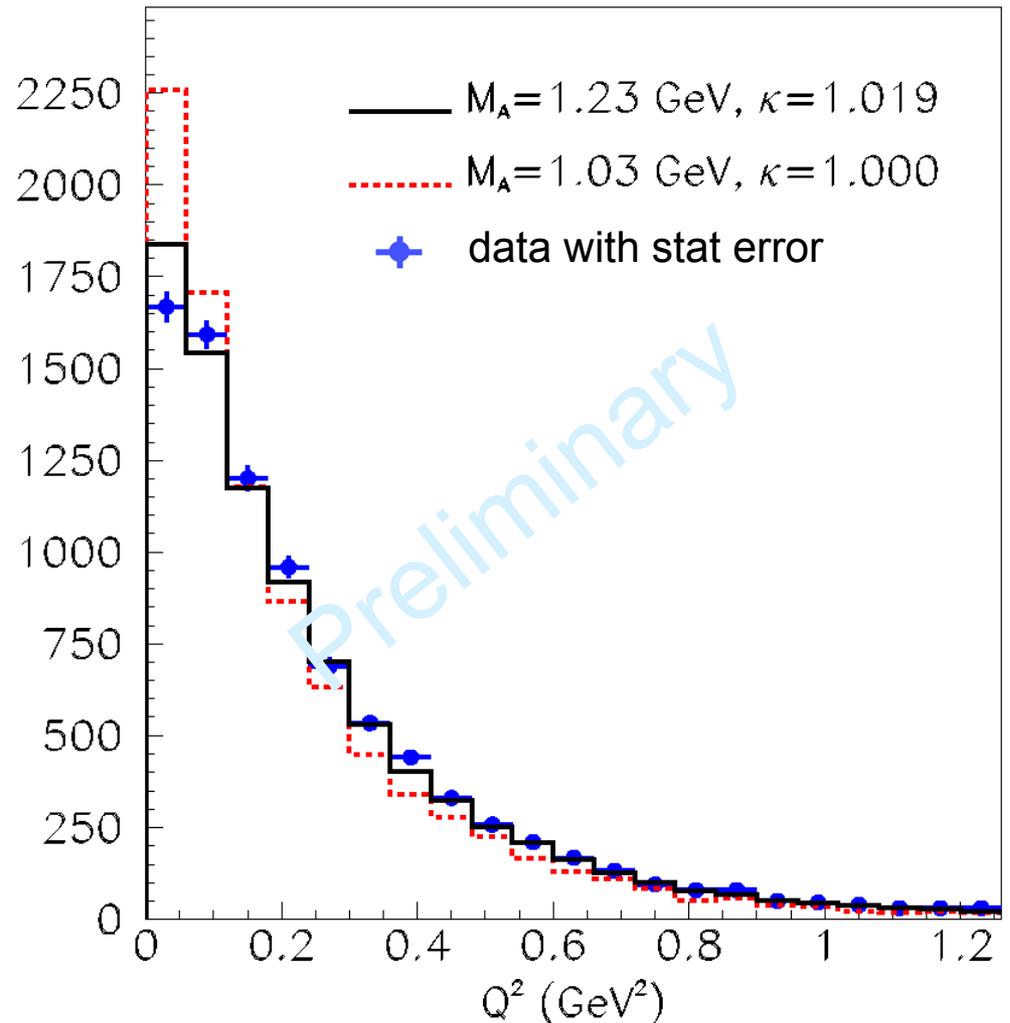
8772 events

(1651 total for pre-MiniBooNE data)

We use same cut with neutrino mode

The values of M_A and κ extracted from neutrino mode are employed to anti-neutrino MC, and they describe data Q^2 distribution well.

Anti-neutrino Q^2 distribution



5. Anti-neutrino CCQE events

Anti-neutrino Q^2 distribution

MiniBooNE anti-neutrino CCQE

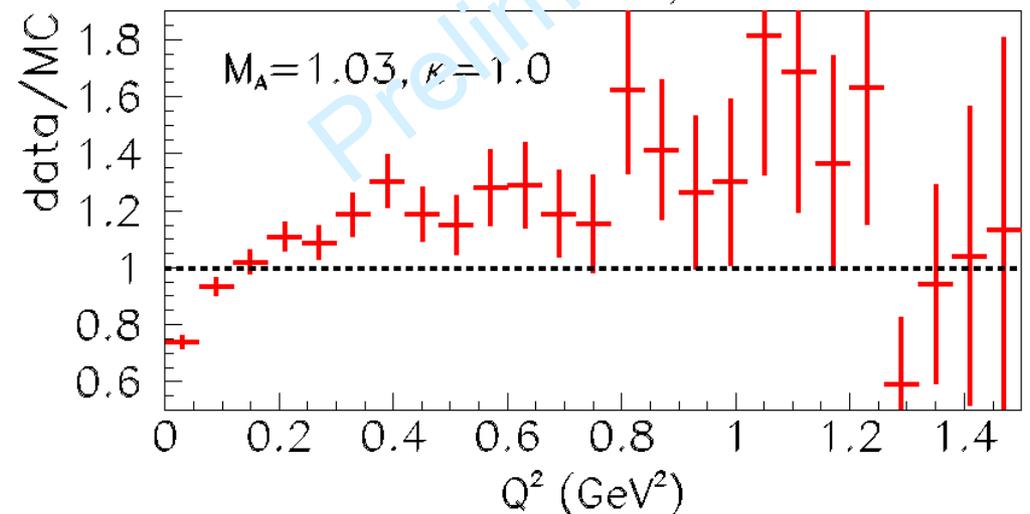
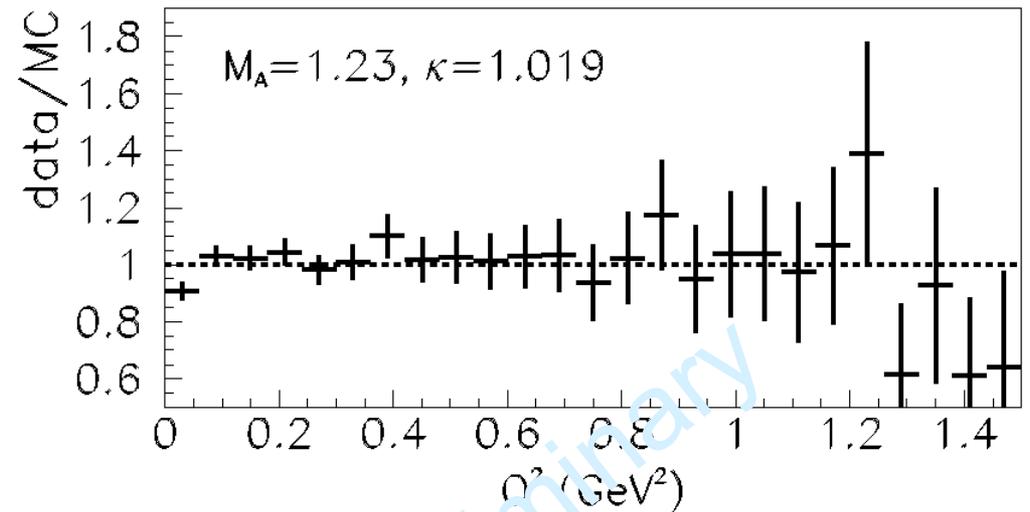
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Anti-neutrino Q^2 distribution data-MC ratio



5. Anti-neutrino CCQE events

Anti-neutrino CCQE kinematics

MiniBooNE anti-neutrino CCQE

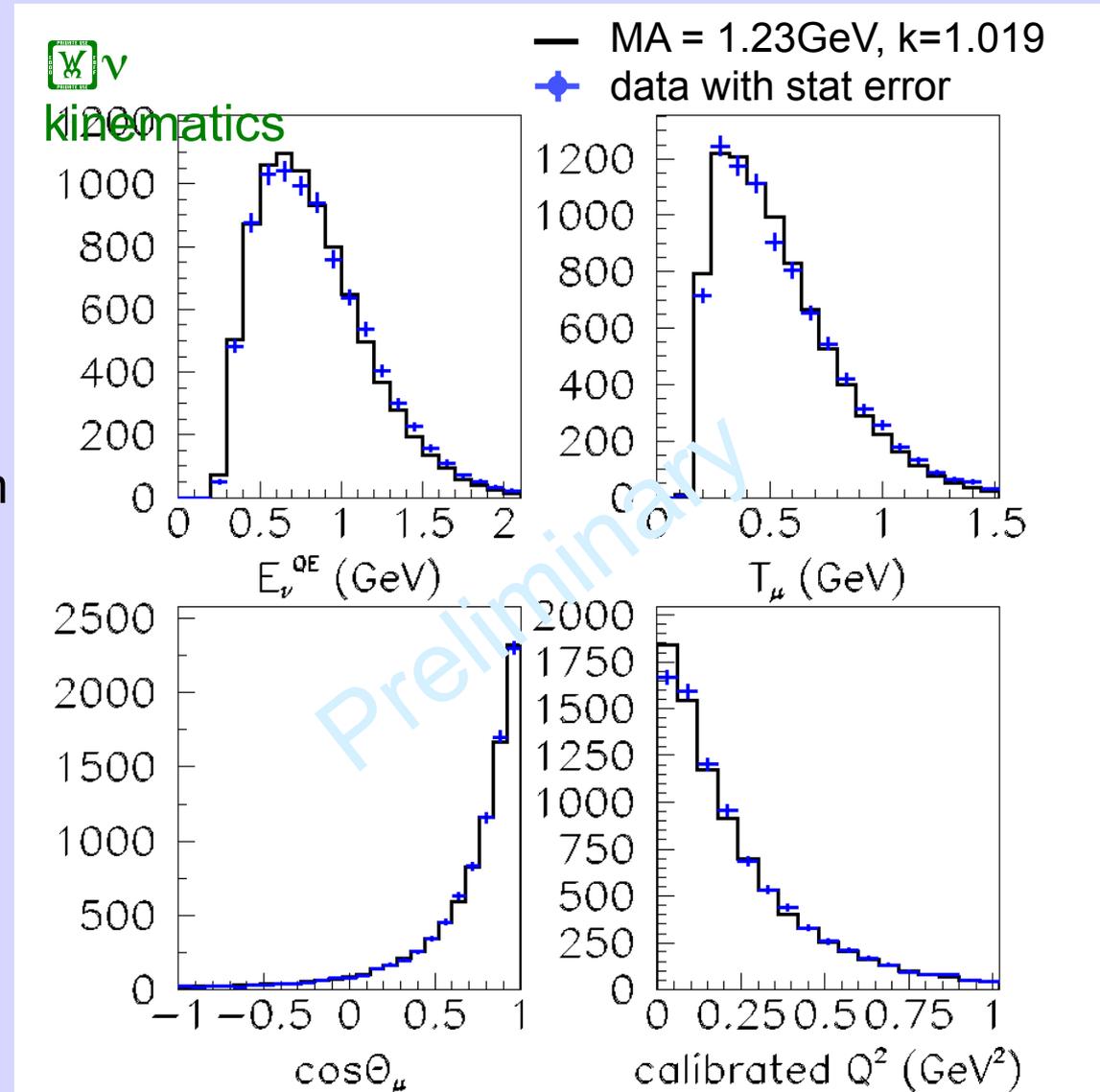
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We use same cut with neutrino mode

The values of M_A and κ extracted from neutrino mode are employed to anti-neutrino MC, and they describe data Q^2 distribution well.

Anti-neutrino CCQE kinematics variables are described by the MC well, too.



6. Conclusion

MiniBooNE has large CCQE data set around 1GeV region

MiniBooNE successfully employed RFG model with appropriate parameter choices for M_A and κ

This new model can describe entire CCQE phase space well

The best fit parameters for MiniBooNE CCQE data are;

$$M_A = 1.23 \pm 0.20(\text{stat+sys})$$
$$\kappa = 1.019 \pm 0.011(\text{stat+sys})$$

Our new model also works well in anti-neutrino data

MiniBooNE is currently taking the data with anti-muon neutrino beam

MiniBooNE collaboration

University of Alabama

Bucknell University

University of Cincinnati

University of Colorado

Columbia University

Embry Riddle University

Fermi National Accelerator Laboratory

Indiana University

Los Alamos National Laboratory

Louisiana State University

University of Michigan

Princeton University

Saint Mary's University of Minnesota

Virginia Polytechnic Institute

Western Illinois University

Yale University



Thank you for your attention!

10. Back up

4. Fit results

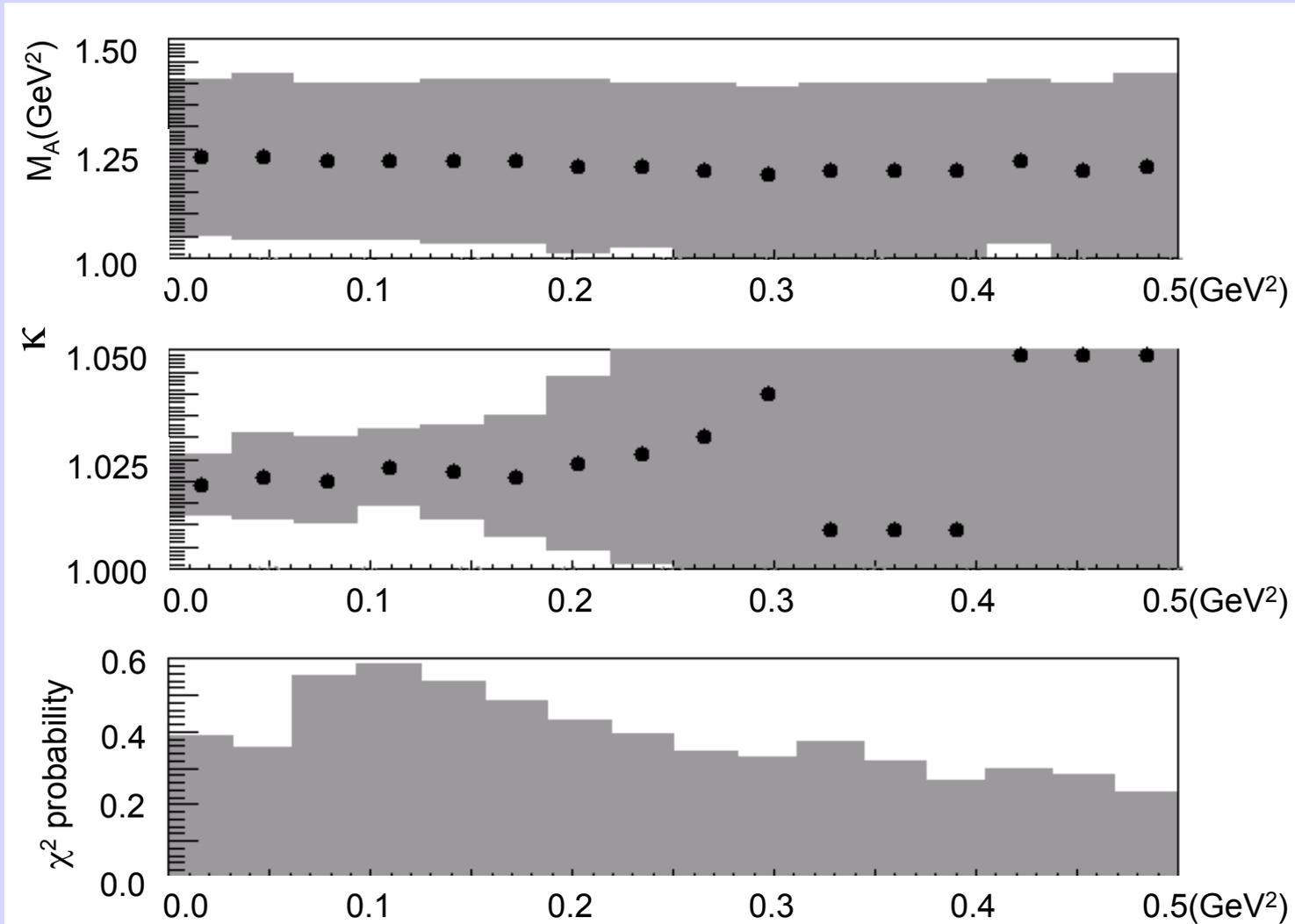
Fit result with varying Q^2_{\min} , $Q^2_{\min} < Q^2 < 1.0\text{GeV}^2$

Fit is repeated with changing the Q^2_{\min}

Fit quality (χ^2 probability)
is good even Q^2_{\min}
 $=0.0\text{GeV}^2$

M_A is stable in wide
range of Q^2_{\min}

Since κ is only important
for low Q^2 region, it has
no power for fit for high
 Q^2



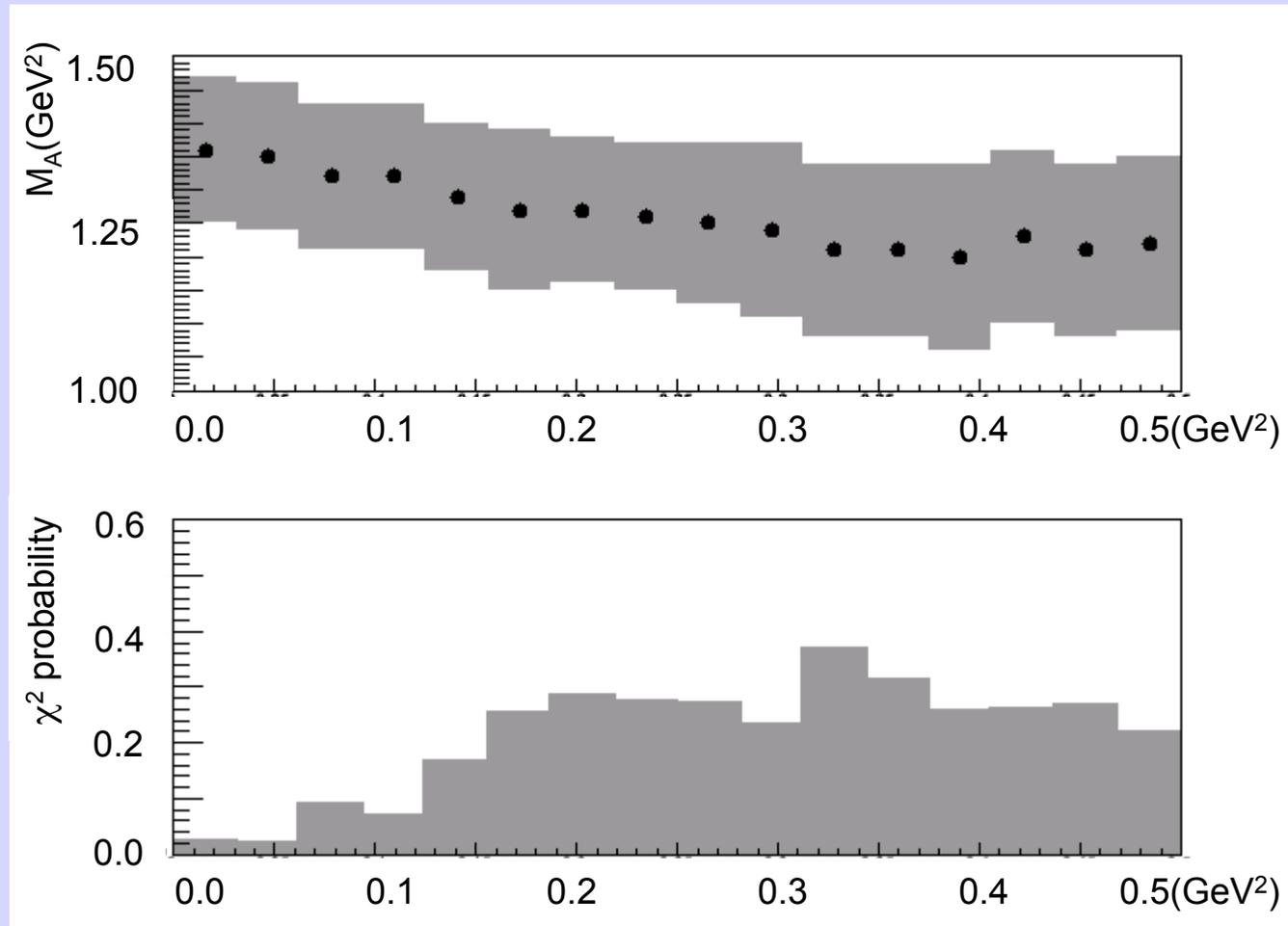
4. Fit results

M_A only fit with varying Q^2_{\min} , $Q^2_{\min} < Q^2 < 1.0\text{GeV}^2$

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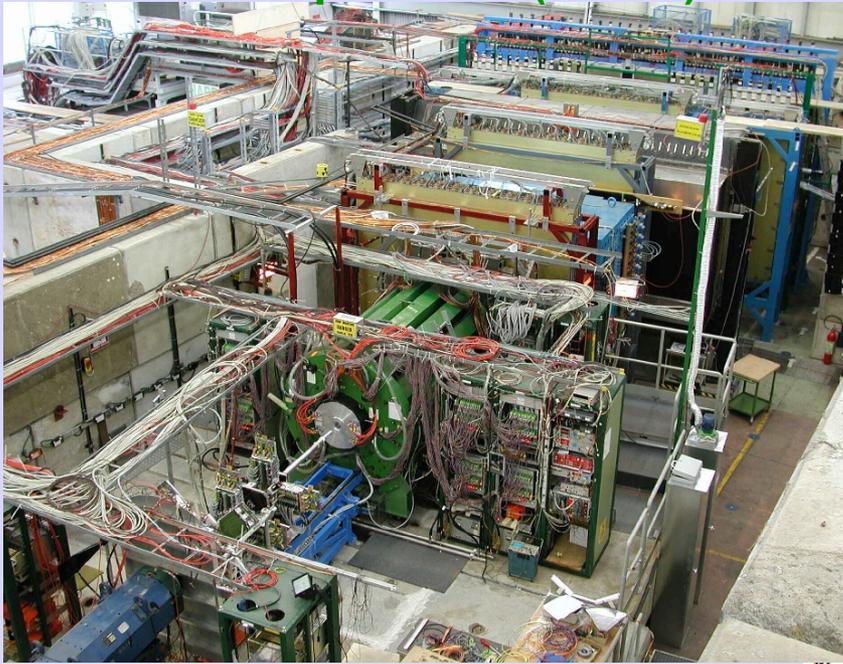
Fit quality (χ^2 probability) is low for $Q^2_{\min} < 0.2\text{GeV}^2$

M_A is stable in wide range of Q^2_{\min} for $Q^2_{\min} > 0.2\text{GeV}^2$



3. Neutrino beam

HARP experiment (CERN)



Modeling Production of Secondary Pions

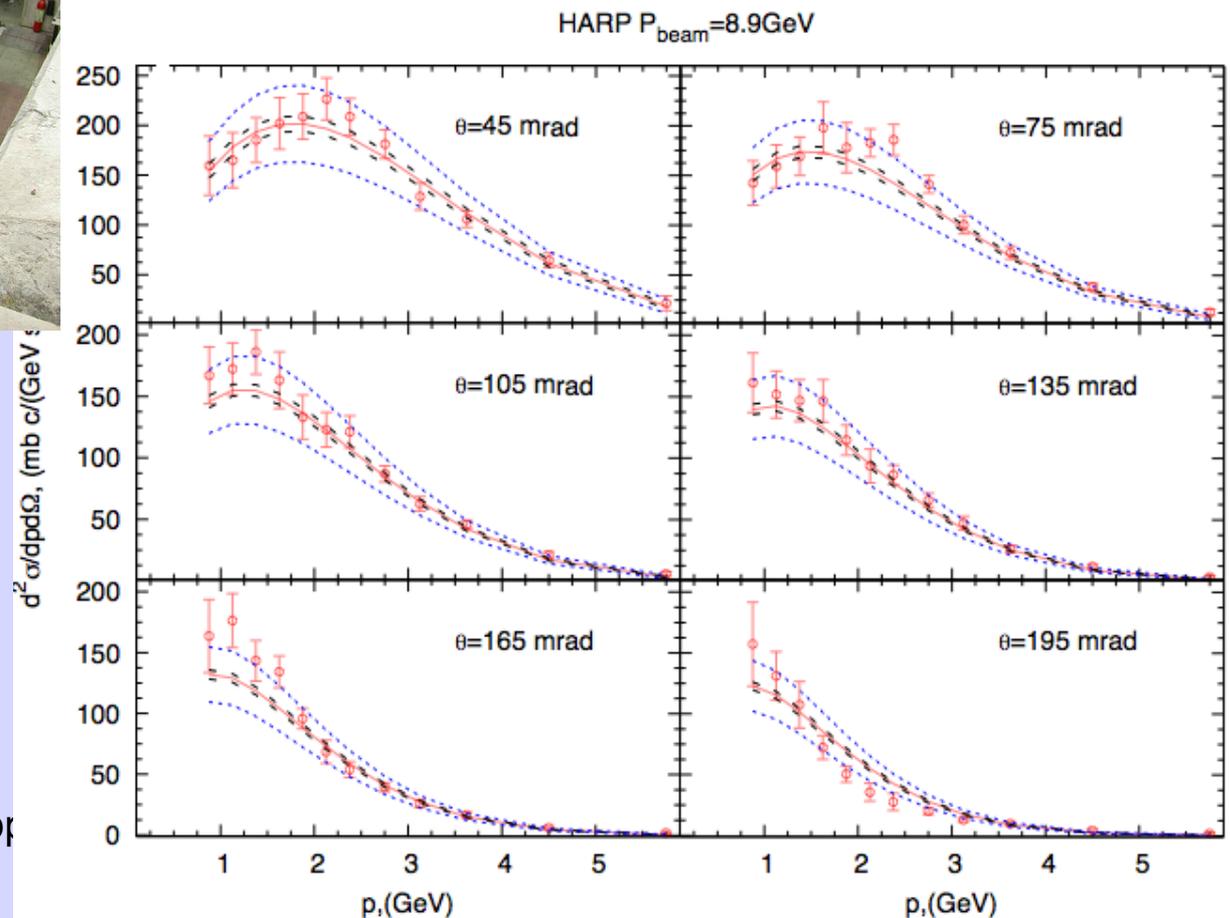
- 5% λ Beryllium target
- 8.9 GeV proton beam momentum

Data are fit to a Sanford-Wang parameterization.

HARP collaboration,
hep-ex/0702024

05/31/2007

Tepp



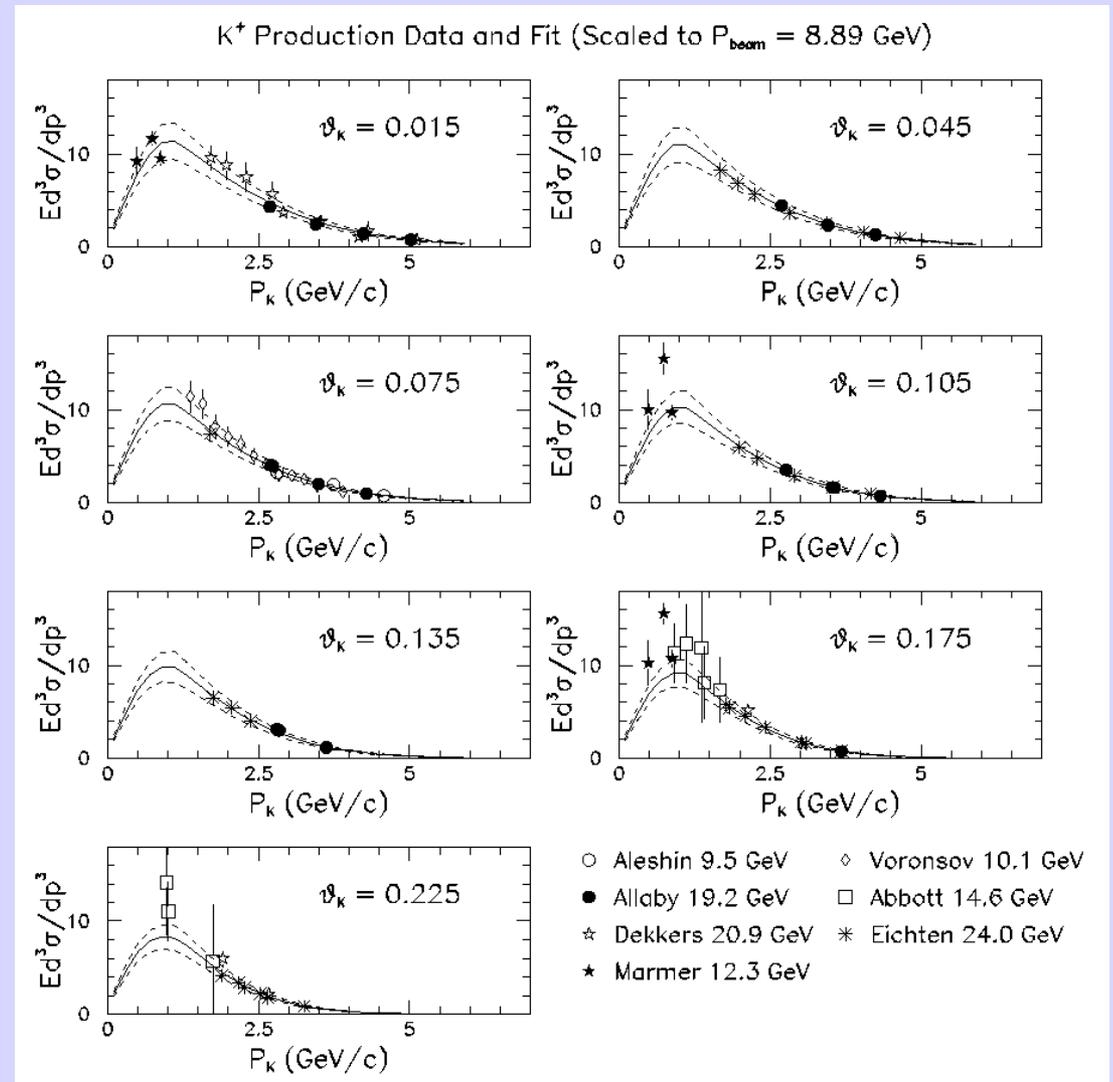
3. Neutrino beam

Modeling Production of Secondary Kaons

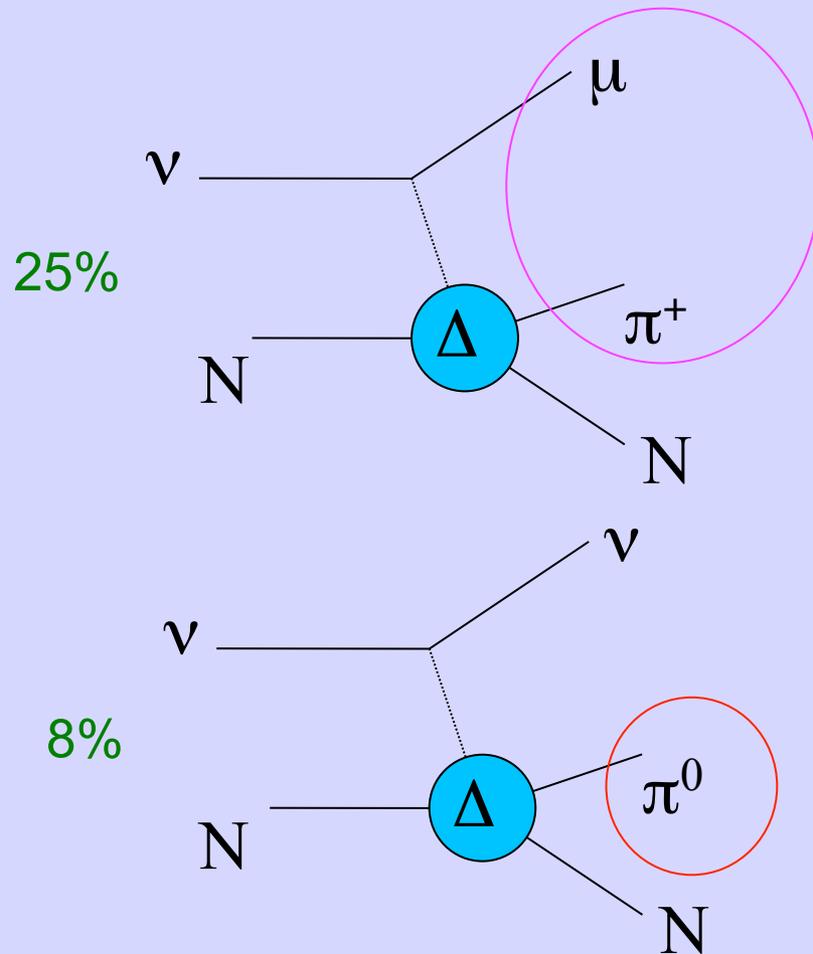
K^+ Data from 10 - 24 GeV.
Uses a Feynman Scaling Parameterization.

K^0 data are also parameterized.

In situ measurement of K^+ from LMC agrees within errors with parameterization



5. Cross section model



(also decays to a single photon
with 0.56% probability)

Events producing pions

$CC\pi^+$

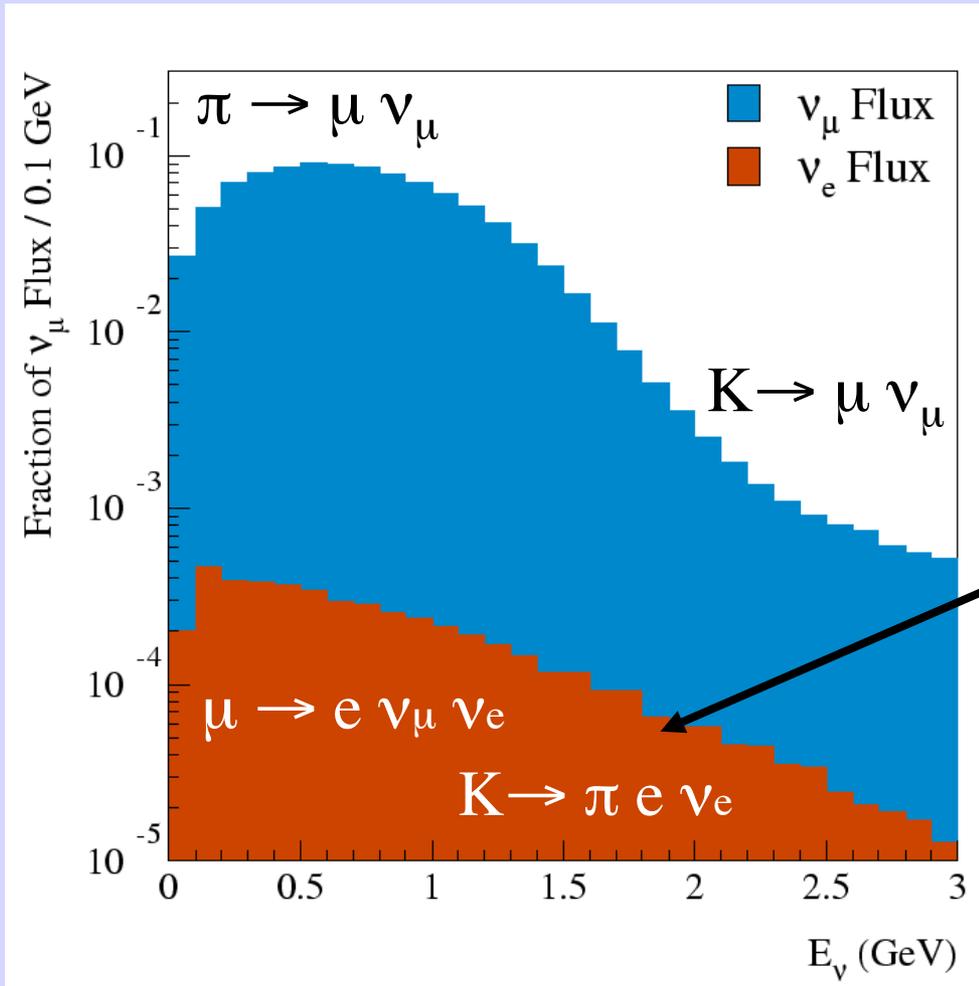
Easy to tag due to 3 subevents.
Not a substantial background to
the oscillation analysis.

$NC\pi^0$

The π^0 decays to 2 photons,
which can look “electron-like”
mimicking the signal...

<1% of π^0 contribute
to background.

6. Blind analysis



“Intrinsic” $\nu_e + \bar{\nu}_e$ sources:

$\mu^+ \rightarrow e^+ \bar{\nu}_\mu \nu_e$ (52%)

$K^+ \rightarrow \pi^0 e^+ \nu_e$ (29%)

$K^0 \rightarrow \pi e \nu_e$ (14%)

Other (5%)

Since MiniBooNE is **blind analysis experiment**, we need to constraint **intrinsic ν_e background** without measuring directly

- (1) μ decay ν_e background
- (2) K decay ν_e background

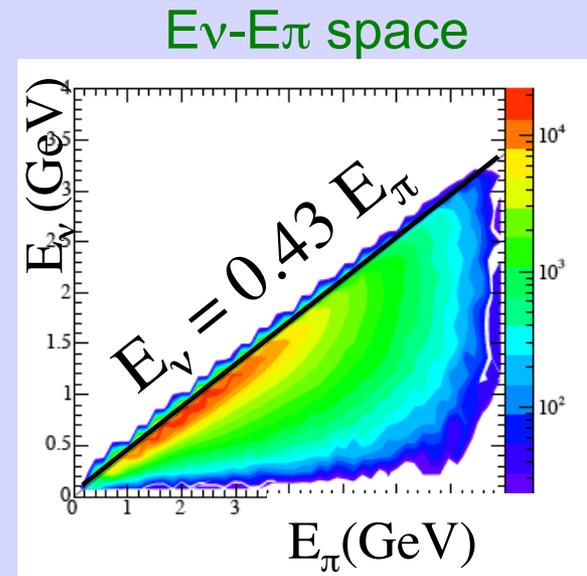
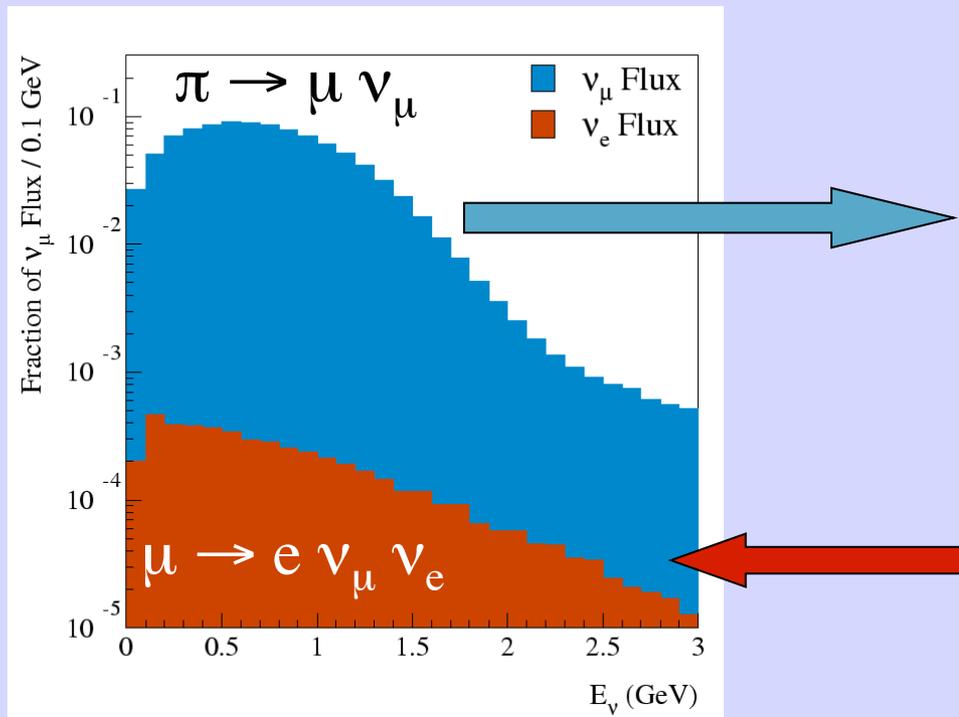
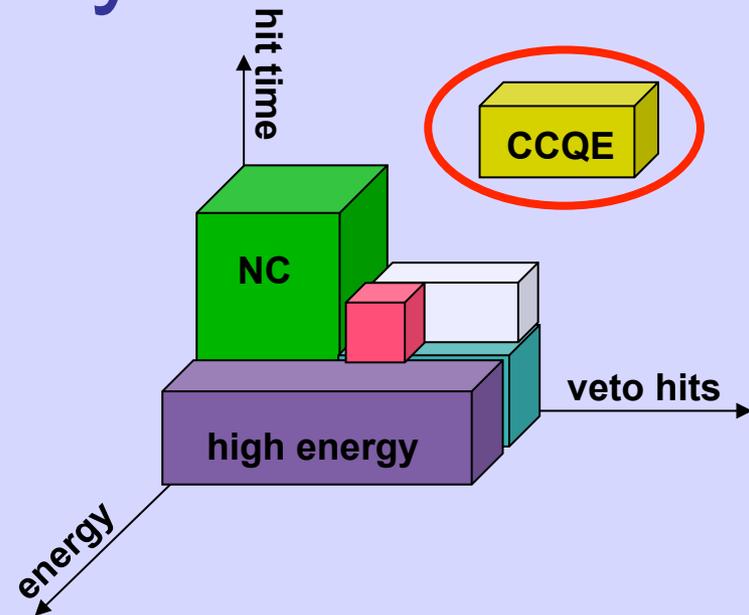
$$\nu_e/\nu_\mu = 0.5\%$$

Antineutrino content: 6%

6. Blind analysis

(1) measure ν_μ flux from ν_μ CCQE event to constraint ν_e background from μ decay

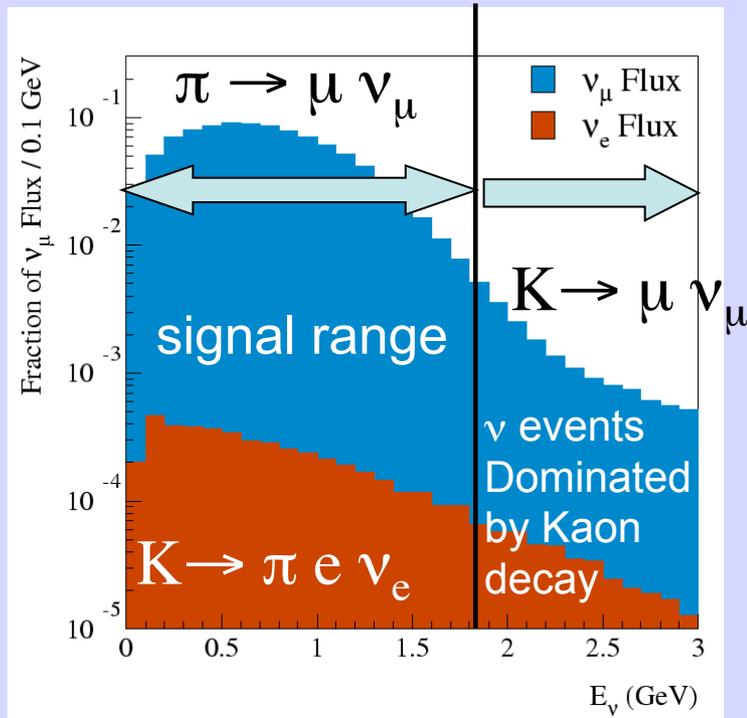
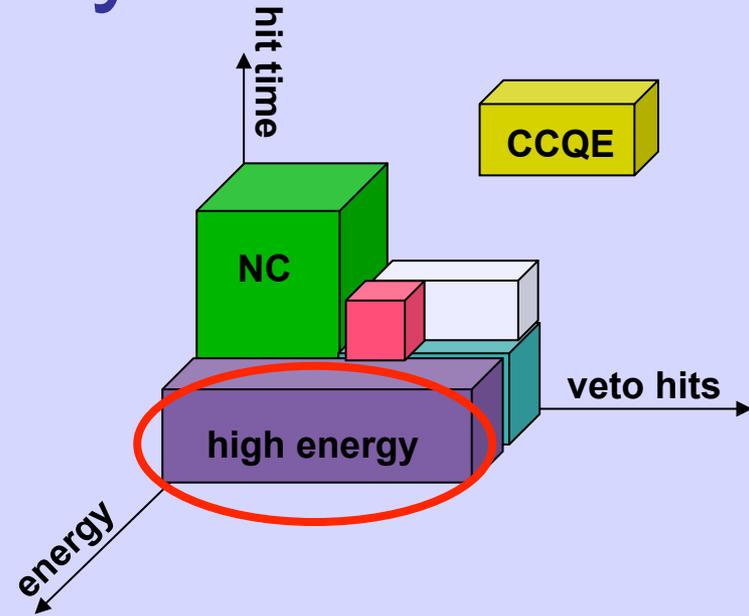
ν_μ CCQE is one of the open boxes. Kinematics allows connection to π flux, hence intrinsic ν_e background from μ decay is constraint.



6. Blind analysis

(2) measure high energy ν_μ events to constraint ν_e background from K decay

At high energies, above “signal range” ν_μ and “ ν_e -like” events are largely due to kaon decay



example of open boxes;

- ν_μ CCQE
 - high energy event
 - $CC\pi^+$
 - NC elastics
 - NC π^0
 - NC electron scattering
 - Michel electron
- etc....

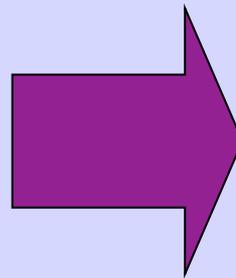
7. Error analysis

Handling uncertainties in the analyses:

What we begin with...

... what we need

For a given source
of uncertainty,
Errors on a wide range
of parameters
in the underlying model



For a given source
of uncertainty,
Errors in bins of
 E_{ν}^{QE}
and information on
the correlations
between bins

Input error matrix
keep the all correlation
of systematics

"multisim"
nonlinear error propagation

Output error matrix
keep the all correlation
of E_{ν}^{QE} bins

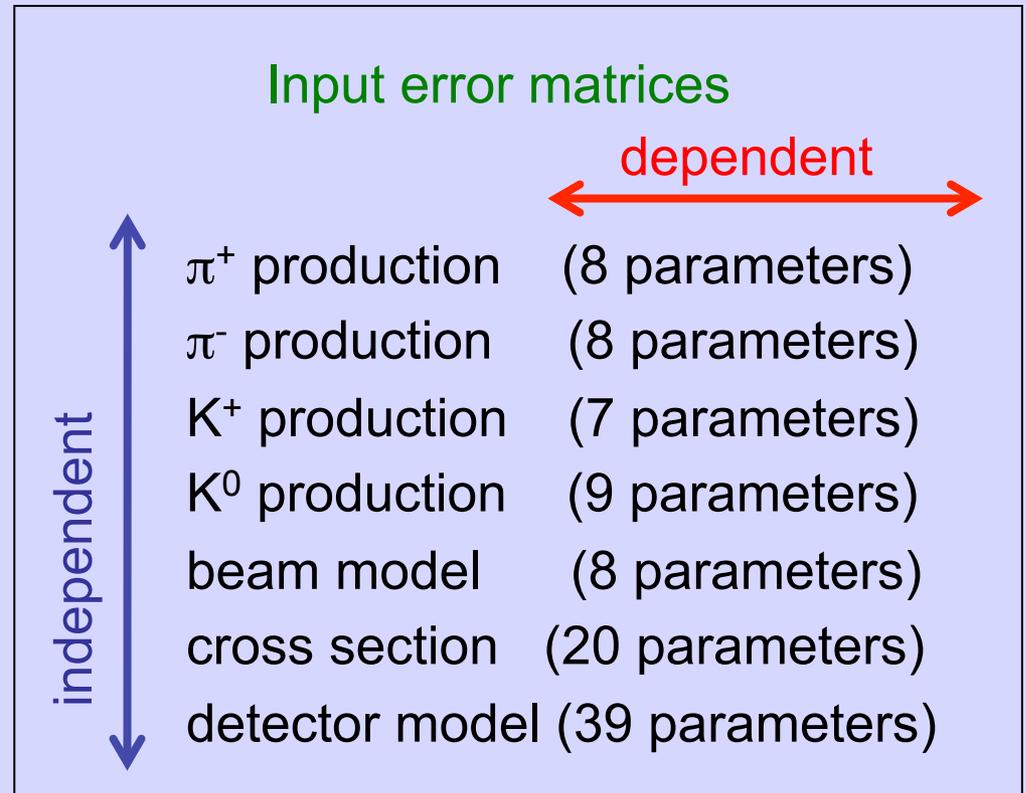
7. Multisim

Multi-simulation (Multisim) method

many fake experiments with different parameter set give the variation of correlated systematic errors for each independent error matrix

total error matrix is the sum of all independent error matrix

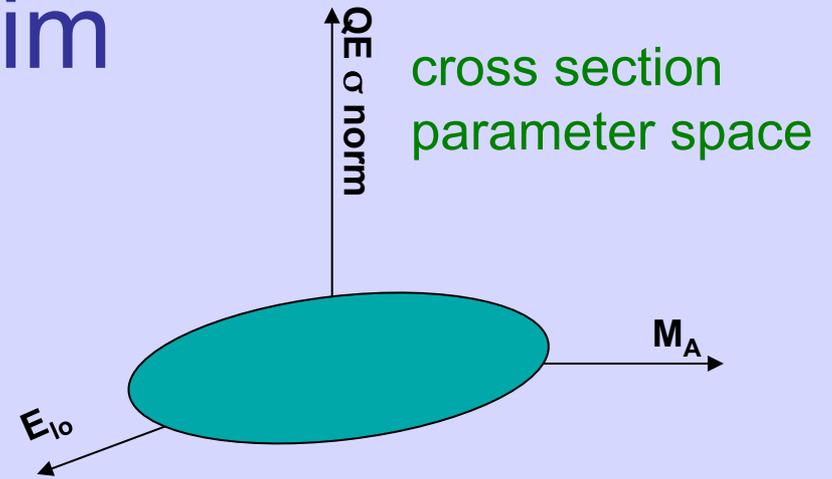
B.P.Roe,
Nucl.,Instrum.,Meth,A570(2007)157



7. Multisim

ex) cross section uncertainties

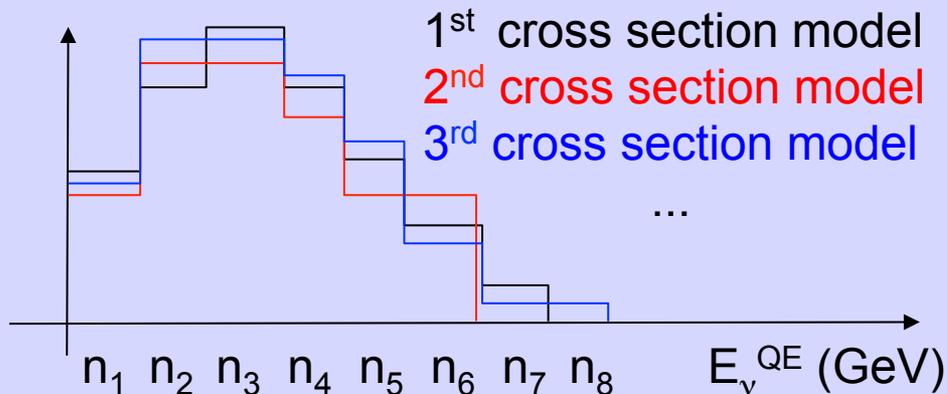
M_A^{QE}	6%	↑ correlated
E_{lo}^{sf}	2%	
QE σ norm	10%	uncorrelated



Input cross section error matrix

$$M_{input}(xs) = \begin{pmatrix} \text{var}(M_A) & \text{cov}(M_A, E_{lo}) & 0 \\ \text{cov}(M_A, E_{lo}) & \text{var}(E_{lo}) & 0 \\ 0 & 0 & \text{var}(\sigma - norm) \end{pmatrix}$$

cross section error for E_v^{QE}

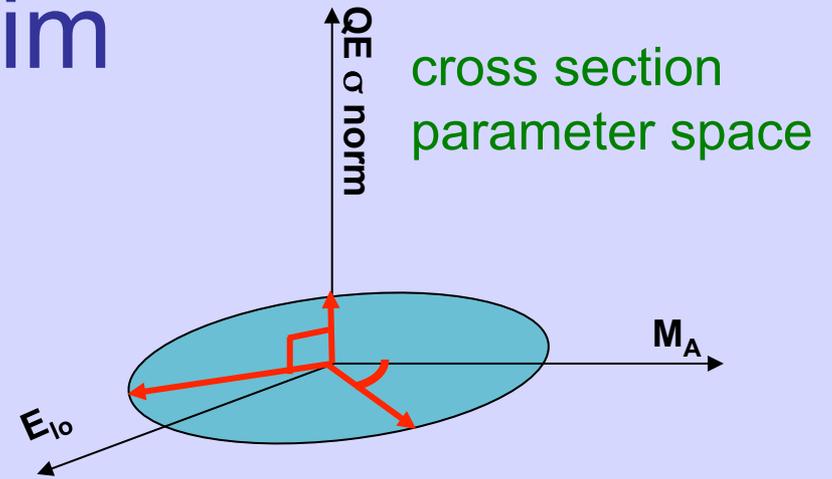


repeat this exercise many times to create smooth error matrix for E_v^{QE}

7. Multisim

ex) cross section uncertainties

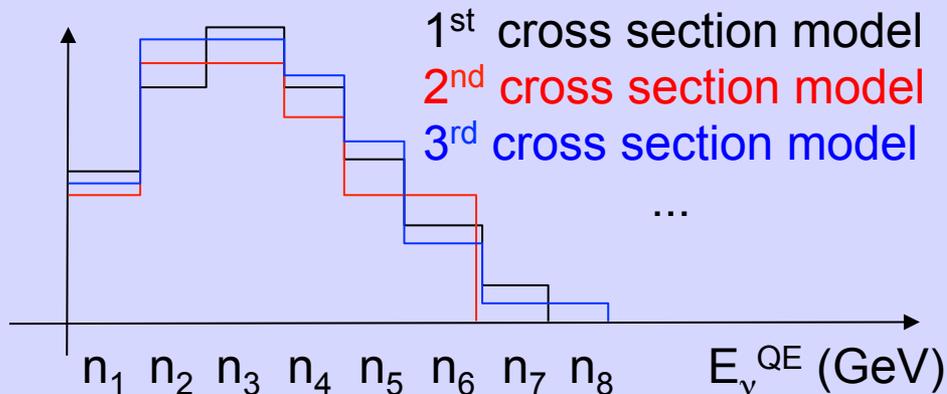
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Input cross section error matrix

$$M_{input}(xs) = \begin{pmatrix} \text{var}(M_A) & \text{cov}(M_A, E_{lo}) & 0 \\ \text{cov}(M_A, E_{lo}) & \text{var}(E_{lo}) & 0 \\ 0 & 0 & \text{var}(\sigma - norm) \end{pmatrix}$$

cross section error for E_v^{QE}



repeat this exercise many times to create smooth error matrix for E_v^{QE}

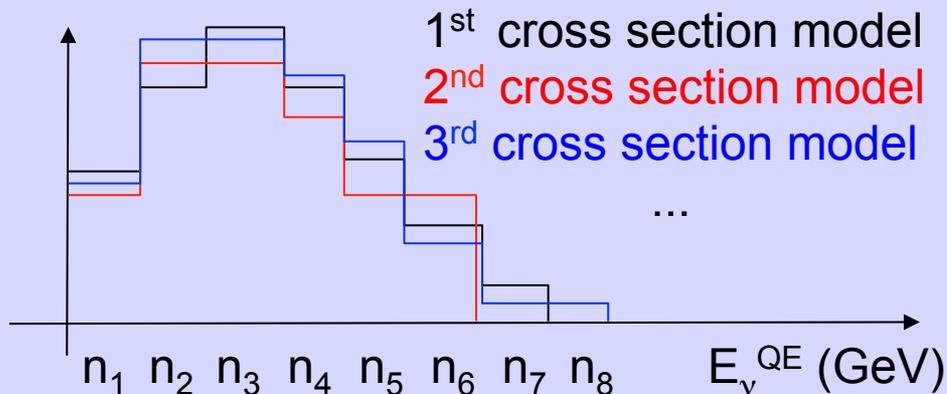
7. Multisim

Output cross section error matrix for E_ν^{QE}

$$[M_{\text{output}}(xS)]_{ij} \approx \frac{1}{S} \sum_k^S (N_i^k(xS) - N_i^{\text{MC}})(N_j^k(xS) - N_j^{\text{MC}})$$

$$M_{\text{output}}(xS) = \begin{pmatrix} \text{var}(n_1) & \text{cov}(n_1, n_2) & \text{cov}(n_1, n_3) & \cdots \\ \text{cov}(n_1, n_2) & \text{var}(n_2) & \text{cov}(n_2, n_3) & \cdots \\ \text{cov}(n_1, n_3) & \text{cov}(n_2, n_3) & \text{var}(n_3) & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

cross section error for E_ν^{QE}



Oscillation analysis use output error matrix for χ^2 fit;

$$\chi^2 = (\text{data} - \text{MC})^T (M_{\text{output}})^{-1} (\text{data} - \text{MC})$$

7. Multisim

ex) cross section uncertainties

M_A^{QE}	6%
$E_{\text{lo}}^{\text{sf}}$	2%
QE σ norm	10%
QE σ shape	function of E_ν
ν_e/ν_μ QE σ	function of E_ν

determined from
MiniBooNE
 ν_μ QE data

NC π^0 rate	function of π^0 mom
$M_A^{\text{coh}}, \text{coh } \sigma$	$\pm 25\%$
$\Delta \rightarrow N\gamma$ rate	function of γ mom + 7% BF

determined from
MiniBooNE
 ν_μ NC π^0 data

E_B, p_F	9 MeV, 30 MeV
Δs	10%
$M_A^{1\pi}$	25%
$M_A^{N\pi}$	40%
DIS σ	25%

determined
from other
experiments

etc...

05/31/2007

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