

Relativity, Mechanics and Quantum Theory.

Solutions to Problem Sheet 6.

(6.1.) a.). Let $A_R \equiv \int_R dt \mathcal{L}(q(t), \dot{q}(t))$.

then under a spacetime transformation:

$$\begin{aligned} \delta A &\equiv A'_R - A_R \\ &= \underbrace{A'_R - A_R}_{\text{Temporal transformation}} + \underbrace{A'_R - A_R}_{\text{Spatial transformation}}. \end{aligned}$$

Temporal transformation. Spatial transformation.

We are concerned, in this problem, with a temporal translation:

$$t' \equiv t + \varepsilon b.$$

The action principle used to derive the equations of motion fixes two points $q(t_1)$ and $q(t_2)$ and extremises the value of the action evaluated over all paths between these points. The path which extremises the action is the dynamical path.

When we consider a temporal translation we do not wish to transform the start and end points of the path, therefore a compensating spatial transformation acts on the coordinates:

$$\begin{array}{ccccc} q(t_1) & \xrightarrow{\quad} & q(t_1 - \varepsilon b) & \xrightarrow{\quad} & q(t_1 + \varepsilon b - \varepsilon b) = q(t_1). \\ \downarrow & & \downarrow & & \\ \text{Compensating} & & \text{Temporal} & & \\ \text{coordinate} & & \text{translation} & & \\ \text{transformation.} & & t \rightarrow t + \varepsilon b & & \end{array}$$

Hence $q(t_1)$ and $q(t_2)$ are not shifted, but we have two transformations:

$$t \rightarrow t + \varepsilon b$$

and

$$q(t) \rightarrow q(t - \varepsilon b) \approx q(t) - \varepsilon b \dot{q}(t) + O(\varepsilon^2).$$

From Noether's Theorem the conserved charge Q if time translation is a symmetry of the action (i.e. if $\delta A = 0$) is:

$$\begin{aligned} Q &\equiv b \mathcal{L} - b \frac{\partial \mathcal{L}}{\partial \dot{q}_i} \dot{q}_i \\ &= b \mathcal{L} - b p_i \dot{q}_i \\ &= -b (p_i \dot{q}_i - \mathcal{L}). \end{aligned}$$

Where $p_i \dot{q}_i - \mathcal{L}$ may be recognised as the energy function.

(6.2).i.)

$$\mathcal{L} = \frac{m}{2} g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu + e A_\mu \dot{x}^\mu.$$

Let $g_{\mu\nu}$ and A_μ be time-independent:

$$\therefore \frac{d}{dt} (g_{\mu\nu}(x)) = \partial_\beta g_{\mu\nu} \dot{x}^\beta = 0.$$

$$\frac{d}{dt} (A_\mu(x)) = \partial_\beta A_\mu \dot{x}^\beta = 0.$$

Lagrange's equation is:

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{x}^\beta} \right) - \frac{\partial \mathcal{L}}{\partial x^\beta} = 0.$$

$$\begin{aligned} \therefore \frac{d}{dt} \left(\frac{m}{2} g_{\mu\nu} \delta^\mu_\beta \dot{x}^\nu + \frac{m}{2} g_{\mu\nu} \dot{x}^\mu \delta^\nu_\beta + e A_\mu \delta^\mu_\beta \right) \\ - \frac{m}{2} (\partial_\beta g_{\mu\nu}) \dot{x}^\mu \dot{x}^\nu - e (\partial_\beta A_\mu) \dot{x}^\mu = 0. \end{aligned}$$

$$\begin{aligned} \therefore \frac{d}{dt} \left(\frac{m}{2} g_{\beta\nu} \dot{x}^\nu + \frac{m}{2} g_{\mu\beta} \dot{x}^\mu + e A_\beta \right) \\ - \frac{m}{2} (\partial_\beta g_{\mu\nu}) \dot{x}^\mu \dot{x}^\nu - e (\partial_\beta A_\mu) \dot{x}^\mu = 0. \end{aligned}$$

$$\therefore m \ddot{x}_\beta + e \cancel{\dot{A}_\beta} - \frac{m}{2} \partial_\beta g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu - e (\partial_\beta A_\mu) \dot{x}^\mu = 0.$$

The Levi-Civita connection* is defined by:

$$\Gamma_{\mu\nu}^{\kappa} \equiv \frac{1}{2} g^{\kappa\beta} (\partial_{\mu} g_{\beta\nu} + \partial_{\nu} g_{\mu\beta} - \partial_{\beta} g_{\mu\nu}).$$

$$\text{So, } \Gamma_{\mu\nu}^{\kappa} \dot{x}^{\mu} \dot{x}^{\nu} = -\frac{1}{2} g^{\kappa\beta} \partial_{\beta} g_{\mu\nu} \dot{x}^{\mu} \dot{x}^{\nu}$$

as $\partial_{\mu} g_{\beta\nu} \dot{x}^{\mu} = 0.$

Premultiplying the equation of motion by $g^{\kappa\beta}$ gives:

$$m \ddot{x}^{\kappa} - \frac{m}{2} g^{\kappa\beta} \partial_{\beta} g_{\mu\nu} \dot{x}^{\mu} \dot{x}^{\nu} - e g^{\kappa\beta} \partial_{\beta} A_{\mu} \dot{x}^{\mu} = 0.$$

$$\therefore m \ddot{x}^{\kappa} + m \Gamma_{\mu\nu}^{\kappa} \dot{x}^{\mu} \dot{x}^{\nu} = e g^{\kappa\beta} \partial_{\beta} A_{\mu} \dot{x}^{\mu} //$$

(ii) The canonical momentum is:

$$p_{\beta} \equiv \frac{\partial \mathcal{L}}{\partial \dot{x}^{\beta}} = \frac{m}{2} g_{\mu\nu} \delta_{\beta}^{\mu} \dot{x}^{\nu} + \frac{m}{2} g_{\mu\nu} \dot{x}^{\mu} \delta_{\beta}^{\nu} + e A_{\mu} \delta_{\beta}^{\mu}.$$

$$p_{\beta} = m g_{\beta\nu} \dot{x}^{\nu} + e A_{\beta}.$$

The Hamiltonian is:

$$p_{\beta} \dot{x}^{\beta} - \mathcal{L} = m g_{\beta\nu} \dot{x}^{\nu} \dot{x}^{\beta} + e A_{\beta} \dot{x}^{\beta} - \mathcal{L}.$$

$$= \frac{m}{2} g_{\mu\nu} \dot{x}^{\mu} \dot{x}^{\nu} //$$

Now

$$p^{\kappa} = m \delta_{\nu}^{\kappa} \dot{x}^{\nu} + e A^{\kappa}$$

$$\therefore \dot{x}^{\kappa} = (p^{\kappa} - e A^{\kappa}) / m.$$

$$\therefore H(x^{\mu}, p_{\nu}) = \frac{m}{2m^2} g_{\mu\nu} (p^{\mu} - e A^{\mu})(p^{\nu} - e A^{\nu}).$$

$$= \frac{1}{2m} g_{\mu\nu} (p^{\mu} - e A^{\mu})(p^{\nu} - e A^{\nu}) //$$

* N.B. the Levi-Civita connection is not part of the course syllabus this year.

(iii.) From question 6.1 we know that the conserved charge if time translation is a symmetry of the motion is 4.

$$p_k \dot{x}^k - \mathcal{L} = \frac{m}{2} g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu.$$

Let us show that this is conserved:

$$\begin{aligned} \frac{d}{dt} \left(\frac{m}{2} g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu \right) &= \frac{m}{2} g_{\mu\nu} \ddot{x}^\mu \dot{x}^\nu + \frac{m}{2} g_{\mu\nu} \dot{x}^\mu \ddot{x}^\nu \\ &= m g_{\mu\nu} \ddot{x}^\mu \dot{x}^\nu \\ &= m g_{\kappa\sigma} \ddot{x}^\kappa \dot{x}^\sigma \\ &= m g_{\kappa\sigma} \dot{x}^\sigma \left(-m \Gamma_{\mu\nu}^\kappa \dot{x}^\mu \dot{x}^\nu + e g^{\kappa\beta} \partial_\beta A_\mu \dot{x}^\mu \right) \\ &= g_{\kappa\sigma} \dot{x}^\sigma \left(+m \dot{x}^\mu \dot{x}^\nu \frac{1}{2} g^{\kappa\beta} \partial_\beta g_{\mu\nu} \right) + e \dot{x}^\sigma g_{\sigma}^\kappa \partial_\beta A_\mu \dot{x}^\mu \\ &= m \dot{x}^\mu \dot{x}^\nu \left(\cancel{\dot{x}^\beta \partial_\beta g_{\mu\nu}} \right) + e \left(\cancel{\dot{x}^\beta \partial_\beta A_\mu} \right) \dot{x}^\mu \\ &\quad \begin{matrix} \swarrow 0. \\ \parallel \\ \left(\frac{d}{dt} (g_{\mu\nu}) \right) \end{matrix} \quad \begin{matrix} \swarrow 0. \\ \parallel \\ \left(\frac{d}{dt} (A_\mu) \right) \end{matrix} \\ &= 0. // \end{aligned}$$

$$(6.3.) \quad \mathcal{L} = \frac{m}{2} (\dot{x}^2 + \dot{y}^2) - \frac{k}{2} (x^2 + y^2).$$

$$(a.) \quad z = x + iy.$$

$$\bar{z} \equiv z^* = x - iy$$

$$\begin{aligned} \therefore z \bar{z} &= (x + iy)(x - iy) \\ &= x^2 + y^2. \end{aligned}$$

$$\dot{z} = \dot{x} + i\dot{y}$$

$$\dot{\bar{z}} = \dot{x} - i\dot{y}$$

$$\therefore \dot{z} \dot{\bar{z}} = \dot{x}^2 + \dot{y}^2.$$

$$\therefore \mathcal{L} = \frac{m}{2} (\dot{z} \dot{\bar{z}}) - \frac{k}{2} z \bar{z}. //$$

(b.) $z \rightarrow z' = e^{i\omega} z.$

$$\therefore z \bar{z} \rightarrow e^{i\omega} z e^{-i\omega} \bar{z} = z \bar{z}.$$

Similarly (as ω is constant):

$$\dot{z} \rightarrow \dot{z}' = e^{i\omega} \dot{z}$$

$$\therefore \dot{z} \dot{\bar{z}} \rightarrow \dot{z}' \dot{\bar{z}}'.$$

$$\begin{aligned} \therefore \mathcal{L} \rightarrow \mathcal{L}' &= \frac{m}{2} (\dot{z}' \dot{\bar{z}}') - \frac{k}{2} (z' \bar{z}') \\ &= \frac{m}{2} (\dot{z} \dot{\bar{z}}) - \frac{k}{2} (z \bar{z}). \\ &= \mathcal{L}. // \end{aligned}$$

(c.) N.B. under this transformation:

$$\begin{aligned} z \rightarrow z' &= e^{i\omega} z = (1 + i\omega + O(\omega^2)) z \\ &= z + i\omega z + O(\omega^2). \\ &= z + \delta z + O(\omega^2). \end{aligned}$$

Noether's charge for a spatial transformation is:

$$Q \equiv X^\mu \frac{\partial \mathcal{L}}{\partial x^\mu}.$$

Here this becomes:

$$\begin{aligned} \hat{Q} &\equiv \delta z \left(\frac{\partial \mathcal{L}}{\partial \dot{z}} \right) + \delta \bar{z} \left(\frac{\partial \mathcal{L}}{\partial \dot{\bar{z}}} \right) \\ &= (i\omega z) \left(\frac{m}{2} \dot{\bar{z}} \right) + (-i\omega \bar{z}) \left(\frac{m}{2} \dot{z} \right). \\ &= \frac{i m \omega}{2} (z \dot{\bar{z}} - \bar{z} \dot{z}). \end{aligned}$$

As ω is a constant ^{infinitesimal parameter} $\uparrow$ we may remove it and

$$Q = \frac{i m}{2} (z \dot{\bar{z}} - \bar{z} \dot{z}) \quad \text{is the Noether charge for the symmetry.}$$

(d.) The equations of motion are:

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{z}} \right) - \frac{\partial \mathcal{L}}{\partial z} = 0$$

$$\Rightarrow \frac{d}{dt} \left(\frac{3}{2} \dot{z} \right) - \frac{k}{2} z = 0.$$

$$\therefore \frac{3}{2} \ddot{z} = \frac{k}{2} z //$$

$$\therefore m \ddot{z} = k z // \quad (*)$$

$$\text{and } m \ddot{z} = -k z // \quad (**)$$

Now,

$$\begin{aligned} \frac{d}{dt} (Q) &= \frac{im}{2} \left(\cancel{\dot{z}\dot{z}} + \cancel{\dot{\bar{z}}\dot{\bar{z}}} - \cancel{\dot{z}\dot{\bar{z}}} - \cancel{\dot{\bar{z}}\dot{z}} \right) \quad \downarrow \text{Using } (*) \text{ and } (**). \\ &= \frac{i}{2} \left(\cancel{z \times \dot{z}} - \cancel{\bar{z} \times \dot{z}} \right) \\ &= 0. // \end{aligned}$$