

MARTINGALE PROPERTY OF GENERALIZED STOCHASTIC EXPONENTIALS

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ABSTRACT. For a real Borel measurable function b , which satisfies certain integrability conditions, it is possible to define a stochastic integral of the process $b(Y)$ with respect to a Brownian motion W , where Y is a diffusion driven by W . It is well known that the stochastic exponential of this stochastic integral is a local martingale. In this paper we consider the case of an arbitrary Borel measurable function b where it may not be possible to define the stochastic integral of $b(Y)$ directly. However the notion of the stochastic exponential can be generalized. We define a non-negative process Z , called *generalized stochastic exponential*, which is not necessarily a local martingale. Our main result gives deterministic necessary and sufficient conditions for Z to be a local, true or uniformly integrable martingale.

1. INTRODUCTION

A *stochastic exponential* of X is a process $\mathcal{E}(X)$ defined by

$$\mathcal{E}(X)_t = \exp \left\{ X_t - X_0 - \frac{1}{2} \langle X \rangle_t \right\}$$

for some continuous local martingale X , where $\langle X \rangle$ denotes a quadratic variation of X . It is well known that the process $\mathcal{E}(X)$ is also a continuous local martingale. Sufficient conditions for the martingale property of $\mathcal{E}(X)$ have been studied extensively in the literature because this question appears naturally in many situations. Novikov's and Kazamaki's sufficient conditions (see [9] and [6]) for $\mathcal{E}(X)$ to be a martingale are particularly well-known. Novikov's condition consists in $\mathbb{E} \exp\{(1/2)\langle X \rangle_t\} < \infty$, $t \in [0, \infty)$. Kazamaki's condition consists in that $\exp\{(1/2)X\}$ should be a submartingale. Novikov's criterion is of narrower scope than Kazamaki's one but often easier to apply. In this respect let us note that none of conditions $\mathbb{E} \exp\{(1/2 - \varepsilon)\langle X \rangle_t\} < \infty$ (with $\varepsilon > 0$) is sufficient for $\mathcal{E}(X)$ to be a martingale (see Liptser and Shiryaev [7, Ch. 6]). For a further literature review see the bibliographical notes in the monographs Karatzas and Shreve [5, Ch. 3], Liptser and Shiryaev [7, Ch. 6], Protter [10, Ch. III], and Revuz and Yor [11, Ch. VIII].

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In the case of one dimensional processes, necessary and sufficient conditions for the process $\mathcal{E}(X)$ to be a martingale were recently studied by Engelbert and Senf in [3], Blei and Engelbert in [1] and Mijatović and Urusov in [8]. In [3] X is a general continuous local martingale and the characterisation is given in terms of the Dambis-Dubins-Schwartz time-change that turns X into a Brownian motion. In [1] X is a strong Markov continuous local martingale and the condition is deterministic, expressed in terms of the speed measure of X .

In [8] the local martingale X is of the form $X_t = \int_0^t b(Y_u) dW_u$ for some measurable function b and a one-dimensional diffusion Y with drift μ and volatility σ driven by a Brownian motion W . In order to define the stochastic integral X , an assumption that the function $\frac{b^2}{\sigma^2}$ is locally integrable on the entire state space of the process Y is required. Under this restriction the characterization of the martingale property of $\mathcal{E}(X)$ is studied in [8], where the necessary and sufficient conditions are deterministic and are expressed in terms of functions μ, σ and b only.

In the present paper we consider an arbitrary Borel measurable function b . In this case the stochastic integral X can only be defined on some subset of the probability space. However, it is possible to define a non-negative possibly discontinuous process Z , known as a generalized stochastic exponential, on the entire probability space. It is a consequence of the definition that, if the function b satisfies the required local integrability condition, the process Z coincides with $\mathcal{E}(X)$. We show that the process Z is not necessarily a local martingale. In fact Z is a local martingale if and only if it is continuous. We find a deterministic necessary and sufficient condition for Z to be a local martingale, which is expressed in terms of local integrability of the quotient $\frac{b^2}{\sigma^2}$ multiplied by a linear function. We also characterize the processes Z that are true martingales and/or uniformly integrable martingales. All the necessary and sufficient conditions are deterministic and are given in terms of functions μ, σ and b .

The paper is structured as follows. In Section 2 we define the notion of generalized stochastic exponential and study its basic properties. The main results are stated in Section 3, where we give a necessary and sufficient condition for the process Z defined by (8) and (12) to be a local martingale, a true martingale or a uniformly integrable martingale. Finally, in Section 4 we prove Theorem 3.4 that is central in obtaining the deterministic characterisation of the martingale property of the process Z . Appendix A contains an auxiliary fact that is used in Section 2.

2. DEFINITION OF GENERALIZED STOCHASTIC EXPONENTIAL

Let $J = (l, r)$ be our state space, where $-\infty \leq l < r \leq \infty$. Let us define a J -valued diffusion Y on a probability space $(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \in [0, \infty)}, \mathbb{P})$ driven by a stochastic differential equation

$$dY_t = \mu(Y_t) dt + \sigma(Y_t) dW_t, \quad Y_0 = x_0 \in J,$$

where W is a (\mathcal{F}_t) -Brownian motion, and μ and σ real, Borel measurable functions defined on J that satisfy the Engelbert–Schmidt conditions

$$(1) \quad \sigma(x) \neq 0 \quad \forall x \in J,$$

$$(2) \quad \frac{1}{\sigma^2}, \frac{\mu}{\sigma^2} \in L^1_{\text{loc}}(J).$$

With $L^1_{\text{loc}}(J)$ we denote the class of locally integrable functions, i.e. real Borel measurable functions defined on J that are integrable on every compact subset of J . Engelbert–Schmidt conditions guarantee existence of a weak solution that might exit the interval J and is unique in law (see [5, Chapter 5]). Denote by ζ the exit time of Y . In addition, we assume that the boundary points are absorbing, i.e. the solution Y stays at the boundary point at which it exits on the set $\{\zeta < \infty\}$. Let us note that we assume that (\mathcal{F}_t) is generated neither by Y nor by W .

We would like to define a process X as a stochastic integral of a process $b(Y)$ with respect to Brownian motion W , where $b : J \rightarrow \mathbb{R}$ is an arbitrary Borel measurable function. Before further discussion, we should establish if the stochastic integral can be defined.

Define a set

$$A = \{x \in J; \frac{b^2}{\sigma^2} \notin L^1_{\text{loc}}(x)\},$$

where $L^1_{\text{loc}}(x)$ denotes a space of real, Borel measurable functions f with $\int_{x-\varepsilon}^{x+\varepsilon} |f(y)| dy < \infty$ for some $\varepsilon > 0$. Then A is closed and its complement is a union of open intervals. Furthermore we can define maps α and β on $J \setminus A$ so that

$$(3) \quad \alpha(x), \beta(x) \in A \cup \{l, r\} \quad \text{and} \quad x \in (\alpha(x), \beta(x)) \subseteq J \setminus A.$$

In other words $\alpha(x)$ is the point in $A \cup \{l, r\}$ that is closest to x from the left side and $\beta(x)$ is the closest point in $A \cup \{l, r\}$ from the right side. Let us also note that the equality

$$(4) \quad L^1_{\text{loc}}(I) = \bigcap_{x \in I} L^1_{\text{loc}}(x)$$

holds for any interval I by a simple compactness argument.

For any $x, y \in J$ we define the stopping times

$$(5) \quad \tau_x = \inf\{t \geq 0; Y_t = x\},$$

$$(6) \quad \tau_{x,y} = \tau_x \wedge \tau_y,$$

with the convention $\inf \emptyset = \infty$, where $c \wedge d := \min\{c, d\}$. Define the stopping time $\zeta^A = \zeta \wedge \tau_A$, where

$$(7) \quad \tau_A = \inf\{t \geq 0; Y_t \in A\}.$$

Then for all $t \geq 0$ we have

$$\int_0^t b^2(Y_u) du < \infty \quad \mathbb{P}\text{-a.s. on } \{t < \zeta^A\}.$$

This follows from Proposition A.1 and the fact that a continuous process Y on $\{t < \zeta^A\}$ reaches only values in an open interval that is a component of the complement of A , where $\frac{b^2}{\sigma^2}$ is locally integrable (note that (4) is applied here).

Let us define $A_n = \{x \in J; \rho(x, A \cup \{l, r\}) \leq \frac{1}{n}\}$, where $\rho(x, y) = |\arctan x - \arctan y|$, $x, y \in \bar{J}$, and set $\zeta_n^A = \inf\{t \geq 0; Y_t \in A_n\}$. Since $\zeta_n^A < \zeta^A$ on the set $\{\zeta^A < \infty\}$, we have $\int_0^{t \wedge \zeta_n^A} b^2(Y_u) du < \infty$ \mathbb{P} -a.s. Thus, we can define the stochastic integral $\int_0^{t \wedge \zeta_n^A} b(Y_u) dW_u$ for every n . Since the integrals $\int_0^{t \wedge \zeta_n^A} b(Y_u) dW_u$ and $\int_0^{t \wedge \zeta_{n+1}^A} b(Y_u) dW_u$ coincide on $\{t < \zeta_n^A\}$ and $\zeta_n^A \nearrow \zeta^A$, we can define $\int_0^{t \wedge \zeta^A} b(Y_u) dW_u$ as a \mathbb{P} -a.s limit of integrals

$$\int_0^{t \wedge \zeta^A} b(Y_u) dW_u = \lim_{n \rightarrow \infty} \int_0^{t \wedge \zeta_n^A} b(Y_u) dW_u \text{ on } \{t < \zeta^A\} \cup \left\{ \int_0^{\zeta^A} b^2(Y_u) du < \infty \right\}.$$

In the case where A is not empty or Y exits the interval J , the stochastic exponential cannot be defined. However, we can define a *generalized stochastic exponential* Z in the following way for every $t \in [0, \infty)$

$$(8) \quad Z_t = \begin{cases} \exp\{\int_0^t b(Y_u) dW_u - \frac{1}{2} \int_0^t b^2(Y_u) du\} & , t < \zeta^A \\ \exp\{\int_0^\zeta b(Y_u) dW_u - \frac{1}{2} \int_0^\zeta b^2(Y_u) du\} & , t \geq \zeta^A = \zeta, \int_0^\zeta b^2(Y_u) du < \infty \\ 0 & , t \geq \zeta^A = \tau_A \text{ or} \\ & t \geq \zeta^A = \zeta, \int_0^\zeta b^2(Y_u) du = \infty \end{cases}$$

The different behaviour of Z on $\{t \geq \zeta^A = \zeta\}$ and $\{t \geq \zeta^A = \tau_A\}$ follows from the fact that after the exit time ζ the process Y is stopped, while this does not happen after τ_A . The definition of the set A and Proposition A.1 imply that the integral $\int_0^t b^2(Y_u) du$ is infinite on the event $\{t > \tau_A\}$ \mathbb{P} -a.s. Therefore, we set $Z = 0$ on the set $\{t \geq \zeta^A = \tau_A\}$.

Let us define the processes

$$(9) \quad \bar{Z}_t = \exp \left\{ \int_0^{t \wedge \zeta^A} b(Y_u) dW_u - \frac{1}{2} \int_0^{t \wedge \zeta^A} b^2(Y_u) du \right\},$$

where we set $\bar{Z}_t = 0$ for $t \geq \zeta^A$ on $\{\zeta^A < \infty, \int_0^{\zeta^A} b^2(Y_u) du = \infty\}$. The process \bar{Z} has continuous trajectories: continuity at time ζ^A on the set $\{\zeta^A < \infty, \int_0^{\zeta^A} b^2(Y_u) du = \infty\}$ follows from the Dambis–Dubins–Schwarz theorem on stochastic intervals; see Th. 1.6 and Ex. 1.18 [11, Ch. V]. Note further that \bar{Z}_t is strictly positive on the event $\{t < \zeta^A\} \cup \left\{ \int_0^{\zeta^A} b^2(Y_u) du < \infty \right\}$ \mathbb{P} -a.s. and is equal to zero on its complement.

Lemma 2.1. *The process $\bar{Z} = (\bar{Z}_t)_{t \geq 0}$ defined in (9) is a continuous local martingale.*

Remark. Fatou’s lemma and Lemma 2.1 imply that the process \bar{Z} is a continuous non-negative supermartingale.

Proof. If the starting point x_0 of the diffusion Y is in A , then $\tau_A = 0$ \mathbb{P} -a.s and \bar{Z} is constant and hence a local martingale. If $x_0 \in J \setminus A$, then pick a decreasing (resp. increasing) sequence $(\alpha_n)_{n \in \mathbb{N}}$ (resp. $(\beta_n)_{n \in \mathbb{N}}$) in the open interval $(\alpha(x_0), \beta(x_0))$ such that $\alpha_n \searrow \alpha(x_0)$ (resp. $\beta_n \nearrow \beta(x_0)$), where $\alpha(x_0), \beta(x_0)$ are defined in (3). Assume also that $\alpha_n < x_0 < \beta_n$ for all $n \in \mathbb{N}$. Note that $\tau_{\alpha_n, \beta_n} \nearrow \zeta^A$ \mathbb{P} -a.s.

The process $M^n = (M_t^n)_{t \in [0, \infty)}$, defined by

$$M_t^n = \int_0^{t \wedge \tau_{\alpha_n, \beta_n}} b(Y_u) dW_u \quad \text{for all } t \in [0, \infty),$$

is a local martingale. Therefore its stochastic exponential $\mathcal{E}(M^n)$ is also a local martingale and, since $\tau_{\alpha_n, \beta_n} < \zeta^A$ \mathbb{P} -a.s., the equality $\bar{Z}_{t \wedge \tau_{\alpha_n, \beta_n}} = \mathcal{E}(M^n)_t$ holds \mathbb{P} -a.s. for all $t \geq 0$. For any $m \in \mathbb{N}$ define the stopping time $\eta_m = \inf\{t \geq 0; \bar{Z}_t \geq m\}$. The stopped process $\bar{Z}^{\eta_m \wedge \tau_{\alpha_n, \beta_n}}$ is a bounded local martingale and hence a martingale. Furthermore note that $\eta_m \nearrow \infty$ \mathbb{P} -a.s. as m tends to infinity. To prove that \bar{Z}^{η_m} is a martingale for each $m \in \mathbb{N}$, note that the process \bar{Z} is stopped at ζ^A , which implies the almost sure limit $\lim_{n \rightarrow \infty} \bar{Z}_t^{\eta_m \wedge \tau_{\alpha_n, \beta_n}} = \bar{Z}_t^{\eta_m}$ for every $t \in [0, \infty)$. Since $\bar{Z}_t^{\eta_m \wedge \tau_{\alpha_n, \beta_n}} \leq m$ \mathbb{P} -a.s., the conditional dominated convergence theorem implies the martingale property of \bar{Z}^{η_m} for every $m \in \mathbb{N}$. This concludes the proof. \square

Define the process $S = (S_t)_{t \in [0, \infty)}$ by

$$(10) \quad S_t = \exp \left\{ \int_0^{\tau_A} b(Y_u) dW_u - \frac{1}{2} \int_0^{\tau_A} b^2(Y_u) du \right\} \mathbb{1}_{\{t \geq \tau_A, \int_0^{\tau_A} b^2(Y_u) du < \infty\}}.$$

Note that for any $t \in [0, \infty)$, \mathbb{P} -a.s. on the set $\{\tau_A \leq t\}$ we have $\tau_A < \zeta$ and hence the integrals in (10) are well-defined. We can express Z as

$$(11) \quad Z = \bar{Z} - S.$$

It is clear from this representation that Z is not necessarily a continuous process. Moreover, since that paths of S are non-decreasing, we have

$$\mathbb{E}[Z_t | \mathcal{F}_s] \leq \bar{Z}_s - \mathbb{E}[S_t | \mathcal{F}_s] \leq \bar{Z}_s - S_s = Z_s.$$

It follows that Z is a non-negative supermartingale and we can define

$$(12) \quad Z_\infty = \lim_{t \rightarrow \infty} Z_t.$$

Note further that if $x_0 \in A$, we have $Z \equiv 0$.

A path of the process Z defined by (8) and (12) is equal to a path of a stochastic exponential if $\zeta^A = \infty$. Otherwise, if $\zeta^A < \infty$, it has one of the following forms:

- (i) $\tau_A < \zeta$ and $\int_0^{\tau_A} b^2(Y_t) dt < \infty$ (see Figure 1);
- (ii) $\zeta^A < \infty$ and $\int_0^{\zeta^A} b^2(Y_t) dt = \infty$ (see Figure 2);
- (iii) $\zeta < \tau_A$ and $\int_0^\zeta b^2(Y_t) dt < \infty$ (see Figure 3).

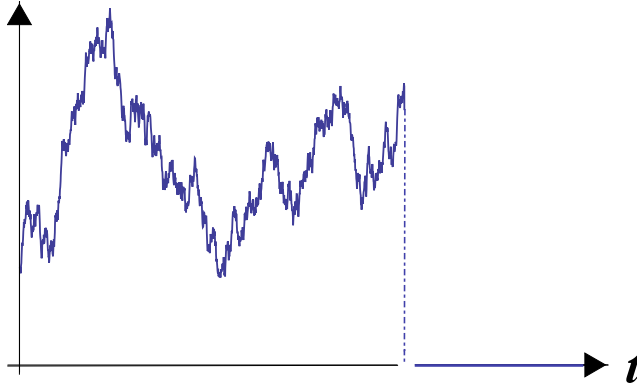


FIGURE 1. If $\tau_A < \zeta$, then the process Z is positive up to time τ_A and is equal to zero afterwards. If the integral $\int_0^{\tau_A} b^2(Y_t) dt$ is finite, then Z_t approaches a positive value as t approaches τ_A . Therefore, there is a jump at $t = \tau_A$.

3. MAIN RESULTS

The case $A = \emptyset$ was studied by Mijatović and Urusov in [8]. We generalize their result for the case where $A \neq \emptyset$.

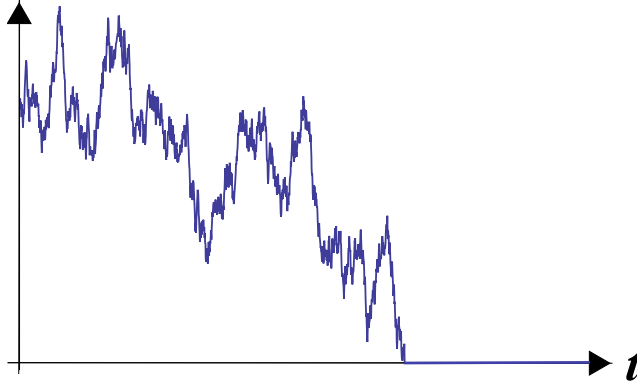


FIGURE 2. If $\zeta^A < \infty$ and $\int_0^{\zeta^A} b^2(Y_t)dt = \infty$, then the process Z is zero after the time ζ^A . Since the limit of Z_t is zero as t approaches ζ^A , there is no jump.

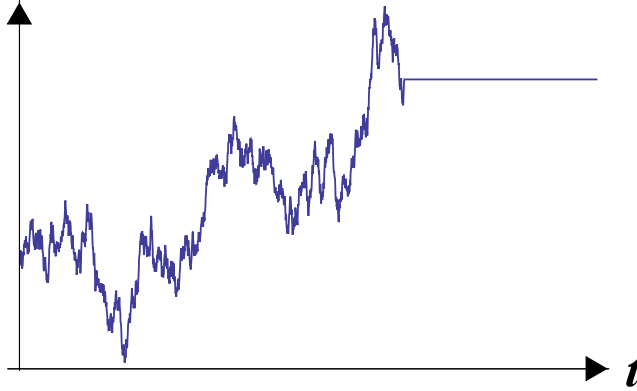


FIGURE 3. If $\zeta < \tau_A$, the process Z is stopped after the exit time. Since $\int_0^\zeta b^2(Y_t)dt$ is finite, Z_t is equal to a positive constant for $t \geq \zeta$.

3.1. **The Case $A = \emptyset$.** In this case we have

$$(13) \quad \frac{b^2}{\sigma^2} \in L^1_{\text{loc}}(J).$$

The generalized stochastic exponential Z defined by (8) and (12) can now be written as

$$Z_t = \exp \left\{ \int_0^{t \wedge \zeta} b(Y_u) dW_u - \frac{1}{2} \int_0^{t \wedge \zeta} b^2(Y_u) du \right\},$$

where we set $Z_t = 0$ for $t \geq \zeta$ on $\{\zeta < \infty, \int_0^\zeta b^2(Y_u) du = \infty\}$. Note that in this case Z is a local martingale by Lemma 2.1, since in this case $\zeta^A = \zeta$ and hence $Z = \bar{Z}$.

Let us now define an auxiliary J -valued diffusion \tilde{Y} governed by the SDE

$$d\tilde{Y}_t = (\mu + b\sigma) \left(\tilde{Y}_t \right) dt + \sigma \left(\tilde{Y}_t \right) d\tilde{W}_t, \quad \tilde{Y}_0 = x_0,$$

on some probability space $(\tilde{\Omega}, \tilde{\mathcal{F}}, (\tilde{\mathcal{F}}_t)_{t \in [0, \infty)}, \tilde{\mathbb{P}})$. The coefficients $\mu + b\sigma$ and σ satisfy Engelbert–Schmidt conditions since $\frac{b}{\sigma} \in L^1_{\text{loc}}(J)$ (this follows from (13)). Hence the SDE has a weak solution, unique in law and possibly explosive. As with diffusion Y , we denote by $\tilde{\zeta}$ the exit time of \tilde{Y} and assume that the boundary points are absorbing.

For an arbitrary $c \in J$ we define the scale functions s, \tilde{s} and their derivatives $\rho, \tilde{\rho}$:

$$(14) \quad \begin{aligned} \rho(x) &= \exp \left\{ - \int_c^x \frac{2\mu(y)}{\sigma^2(y)} dy \right\}, \quad x \in J, \\ \tilde{\rho}(x) &= \rho(x) \exp \left\{ - \int_c^x \frac{2b(y)}{\sigma(y)} dy \right\}, \quad x \in J, \\ s(x) &= \int_c^x \rho(y) dy, \quad x \in \bar{J}, \\ \tilde{s}(x) &= \int_c^x \tilde{\rho}(y) dy, \quad x \in \bar{J}. \end{aligned}$$

For any $a \in (l, r]$, $L^1_{\text{loc}}(a-)$ denotes the set of all Borel measurable functions $f : J \rightarrow \mathbb{R}$ such that there exists $b < a$, $b \in J$, with $\int_b^a |f(x)| dx < \infty$. Similarly, we define the set $L^1_{\text{loc}}(a+)$ for any $a \in [l, r)$.

We say that the endpoint r is *good* if

$$s(r) < \infty \quad \text{and} \quad \frac{(s(r) - s)b^2}{\rho\sigma^2} \in L^1_{\text{loc}}(r-).$$

It is equivalent to show that

$$\tilde{s}(r) < \infty \quad \text{and} \quad \frac{(\tilde{s}(r) - \tilde{s})b^2}{\tilde{\rho}\sigma^2} \in L^1_{\text{loc}}(r-).$$

The endpoint l is *good* if

$$s(l) > -\infty \quad \text{and} \quad \frac{(s - s(l))b^2}{\rho\sigma^2} \in L^1_{\text{loc}}(l+),$$

or equivalently

$$\tilde{s}(l) > -\infty \quad \text{and} \quad \frac{(\tilde{s} - \tilde{s}(l))b^2}{\tilde{\rho}\sigma^2} \in L^1_{\text{loc}}(l+).$$

If an endpoint is not good, we say it is *bad*. The good and bad endpoints were introduced in [8], where one can also find the proof of equivalences above.

We will use the following terminology:

$$\begin{aligned} \tilde{Y} \text{ exits at } r &\text{ means } \tilde{\mathbb{P}}(\tilde{\zeta} < \infty, \lim_{t \nearrow \tilde{\zeta}} \tilde{Y}_t = r) > 0; \\ \tilde{Y} \text{ exits at } l &\text{ means } \tilde{\mathbb{P}}(\tilde{\zeta} < \infty, \lim_{t \nearrow \tilde{\zeta}} \tilde{Y}_t = l) > 0. \end{aligned}$$

Define

$$(15) \quad \tilde{v}(x) = \int_c^x \frac{\tilde{s}(x) - \tilde{s}(y)}{\tilde{\rho}(y)\sigma^2(y)} dy, \quad x \in J,$$

and

$$(16) \quad \tilde{v}(r) = \lim_{x \nearrow r} \tilde{v}(x), \quad \tilde{v}(l) = \lim_{x \searrow l} \tilde{v}(x)$$

(note that \tilde{v} is decreasing on $(l, c]$ and increasing on $[c, r)$).

Feller's test for explosions (see [5, Chapter 5, Theorem 5.29]) tells us that:

(i) \tilde{Y} exits at the boundary point r if and only if

$$\tilde{v}(r) < \infty.$$

It is equivalent to check (see [2, Chapter 4.1])

$$\tilde{s}(r) < \infty \quad \text{and} \quad \frac{\tilde{s}(r) - \tilde{s}}{\tilde{\rho}\sigma^2} \in L_{\text{loc}}^1(r-);$$

(ii) \tilde{Y} exits at the boundary point l if and only if

$$\tilde{v}(l) < \infty,$$

which is equivalent to

$$\tilde{s}(l) > -\infty \quad \text{and} \quad \frac{\tilde{s} - \tilde{s}(l)}{\tilde{\rho}\sigma^2} \in L_{\text{loc}}^1(l+).$$

Remark. The endpoint r (resp. l) is bad whenever one of the processes Y and \tilde{Y} exits at r (resp. l) and the other does not.

The following theorems are reformulations of Theorems 2.1 and 2.3 in [8].

Theorem 3.1. *Let the functions μ, σ and b satisfy conditions (1), (2) and (13). Then the process Z is a martingale if and only if \tilde{Y} does not exit at the bad endpoints.*

Theorem 3.2. *Let the functions μ, σ and b satisfy conditions (1), (2) and (13). Then Z is a uniformly integrable martingale if and only if one of the conditions (a) – (d) below is satisfied:*

- (a) $b = 0$ a.e. on J with respect to the Lebesgue measure;
- (b) r is good and $\tilde{s}(l) = -\infty$;
- (c) l is good and $\tilde{s}(r) = \infty$;
- (d) l and r are good.

3.2. **The Case $A \neq \emptyset$.** In the rest of the paper we assume that

$$x_0 \notin A,$$

since $x_0 \in A$ implies that $Z \equiv 0$. The following example shows that even when A is not empty we can get a martingale or a uniformly integrable martingale defined by (8) and (12).

Example 3.3. (i) Let us consider the state space $J = \mathbb{R}$, coefficients of the SDE $\mu = 0, \sigma = 1$, starting point of the diffusion $x_0 > 0$ and function $b(x) = \frac{1}{x}$ for $x \in \mathbb{R} \setminus \{0\}$ and $b(0) = 0$. Then $A = \{0\}$ and $Y_t = W_t, W_0 = x_0$. Using Itô's formula and the fact that Brownian motion does not exit at infinity, we get for $t < \tau_0$

$$\begin{aligned} Z_t &= \exp \left\{ \int_0^t \frac{1}{W_u} dW_u - \frac{1}{2} \int_0^t \frac{1}{W_u^2} du \right\} \\ &= \frac{1}{x_0} W_t \end{aligned}$$

and $Z_t = 0$ for $t \geq \tau_0$. Hence, $Z_t = \frac{1}{x_0} W_{t \wedge \tau_0}$ that is a martingale.

(ii) Using the same functions μ, σ and b as above on a state space $J = (-\infty, x_0 + 1)$ we get

$$Z_t = \frac{1}{x_0} W_{t \wedge \tau_{0, x_0+1}},$$

which is a uniformly integrable martingale.

Let for any $x \in J \setminus A$ the points $\alpha(x), \beta(x) \in A \cup \{l, r\}$ be as in (3). Then $\frac{b^2}{\sigma^2} \in L_{\text{loc}}^1(\alpha(x), \beta(x))$. Therefore, on $(\alpha(x), \beta(x))$ functions μ, σ and b satisfy the same conditions as in the previous subsection.

For any starting point $x_0 \in J \setminus A$ we can define an auxiliary diffusion \tilde{Y} with state space $(\alpha(x_0), \beta(x_0))$ driven by the SDE

$$d\tilde{Y}_t = (\mu + b\sigma) \left(\tilde{Y}_t \right) dt + \sigma \left(\tilde{Y}_t \right) d\tilde{W}_t, \quad \tilde{Y}_0 = x_0,$$

on some probability space $(\tilde{\Omega}, \tilde{\mathcal{F}}, (\tilde{\mathcal{F}}_t)_{t \in [0, \infty)}, \tilde{\mathbb{P}})$. There exists a unique weak solution of this equation since coefficients satisfy the Engelbert–Schmidt conditions.

As in the previous subsection we can define good and bad endpoints. Let functions $\rho, \tilde{\rho}, s, \tilde{s}$ and \tilde{v} be defined by (14), (15) and (16) with $c \in (\alpha(x_0), \beta(x_0))$. We say that the endpoint $\beta(x_0)$ is *good* if

$$s(\beta(x_0)) < \infty \quad \text{and} \quad \frac{(s(\beta(x_0)) - s)b^2}{\rho\sigma^2} \in L_{\text{loc}}^1(\beta(x_0)-).$$

It is equivalent to show and sometimes easier to check that

$$\tilde{s}(\beta(x_0)) < \infty \quad \text{and} \quad \frac{(\tilde{s}(\beta(x_0)) - \tilde{s})b^2}{\tilde{\rho}\sigma^2} \in L_{\text{loc}}^1(\beta(x_0)-).$$

The endpoint $\alpha(x_0)$ is *good* if

$$s(\alpha(x_0)) > -\infty \quad \text{and} \quad \frac{(s - s(\alpha(x_0)))b^2}{\rho\sigma^2} \in L_{\text{loc}}^1(\alpha(x_0)+),$$

or equivalently

$$\tilde{s}(\alpha(x_0)) > -\infty \quad \text{and} \quad \frac{(\tilde{s} - \tilde{s}(\alpha(x_0)))b^2}{\tilde{\rho}\sigma^2} \in L_{\text{loc}}^1(\alpha(x_0)+).$$

If an endpoint is not good, we say it is *bad*.

The following theorem plays a key role in all that follows.

Theorem 3.4. *Let $\alpha : J \setminus A \rightarrow A \cup \{l, r\}$ be the function defined by (3) and assume that the starting point $x_0 \in J \setminus A$ of diffusion Y satisfies $\alpha(x_0) > l$. Denote $\alpha_0 = \alpha(x_0) \in A$. Then:*

- (a) $(x - \alpha_0)\frac{b^2}{\sigma^2}(x) \in L_{\text{loc}}^1(\alpha_0+) \iff \int_0^{\tau_{\alpha_0}} b^2(Y_t) dt < \infty$ \mathbb{P} -a.s. on $\{\tau_{\alpha_0} = \tau_A < \infty\}$;
- (b) $(x - \alpha_0)\frac{b^2}{\sigma^2}(x) \notin L_{\text{loc}}^1(\alpha_0+) \iff \int_0^{\tau_{\alpha_0}} b^2(Y_t) dt = \infty$ \mathbb{P} -a.s. on $\{\tau_{\alpha_0} = \tau_A < \infty\}$.

Remarks. (i) The assumption $\alpha(x_0) > l$ means that Theorem 3.4 deals with the situation where the set A contains points to the left of the starting point x_0 .

(ii) Clearly, Theorem 3.4 has its analogue for $\beta_0 = \beta(x_0) < r$, i.e. the case when A contains points to the right of x_0 . The deterministic criterion in this case takes the form

$$(17) \quad (\beta_0 - x)\frac{b^2}{\sigma^2}(x) \in L_{\text{loc}}^1(\beta_0-) \iff \int_0^{\tau_{\beta_0}} b^2(Y_t) dt < \infty \text{ } \mathbb{P}\text{-a.s. on } \{\tau_{\beta_0} = \tau_A < \infty\},$$

$$(18) \quad (\beta_0 - x)\frac{b^2}{\sigma^2}(x) \notin L_{\text{loc}}^1(\beta_0-) \iff \int_0^{\tau_{\beta_0}} b^2(Y_t) dt = \infty \text{ } \mathbb{P}\text{-a.s. on } \{\tau_{\beta_0} = \tau_A < \infty\}.$$

(iii) Note that $\mathbb{P}(\tau_{\alpha_0} = \tau_A < \infty) > 0$. Indeed, if $\beta_0 = r$, then $\{\tau_{\alpha_0} = \tau_A < \infty\} = \{\tau_{\alpha_0} < \infty\}$; if $\beta_0 < r$, then $\{\tau_{\alpha_0} = \tau_A < \infty\} = \{\tau_{\alpha_0} < \tau_{\beta_0}\}$. In both cases $\mathbb{P}(\tau_{\alpha_0} = \tau_A < \infty) > 0$ by [2, Theorem 2.11].

(iv) Since the diffusion Y starts at x_0 , \mathbb{P} -a.s. we have the following implications

$$(19) \quad \alpha_0, \beta_0 \in A \implies \tau_A = \tau_{\alpha_0, \beta_0},$$

$$(20) \quad \alpha_0 \in A, \beta_0 = r \implies \tau_A = \tau_{\alpha_0},$$

$$(21) \quad \alpha_0 = l, \beta_0 \in A \implies \tau_A = \tau_{\beta_0}.$$

The case $\alpha_0 = l, \beta_0 = r$ cannot occur since $A \neq \emptyset$.

(v) Note that right-hand sides of the equivalences in (a) and (b) in Theorem 3.4 are not

negations of each other. If it does not hold that the integral $\int_0^{\tau_{\alpha_0}} b^2(Y_t) dt$ is finite \mathbb{P} -a.s. on the set $\{\tau_{\alpha_0} = \tau_A < \infty\}$, then it must be infinite on some subset of $\{\tau_{\alpha_0} = \tau_A < \infty\}$ of positive probability, which may be strictly smaller than $\mathbb{P}(\tau_{\alpha_0} = \tau_A < \infty)$.

For any starting point $x_0 \in J \setminus A$ of diffusion Y , define the set

$$B(x_0) = \left\{ x \in J \cap \{\alpha_0, \beta_0\}; \int_0^{\tau_x} b^2(Y_t) dt = \infty \text{ } \mathbb{P}\text{-a.s. on } \{\tau_x = \tau_A < \infty\} \right\}.$$

Note that $B(x_0)$ is contained in A and that it contains at most two points. Theorem 3.4 implies that $\alpha_0 \in B(x_0)$ if and only if the deterministic condition in (b) is satisfied. Similarly $\beta_0 \in B(x_0)$ is equivalent to the deterministic condition in (18). Therefore Theorem 3.4 yields a deterministic description of the set $B(x_0)$.

We can now give a deterministic characterisation for a generalized stochastic exponential Z to be a local martingale and a true martingale.

Theorem 3.5. (i) *The generalized stochastic exponential Z is a local martingale if and only if $\alpha(x_0), \beta(x_0) \in B(x_0) \cup \{l, r\}$.*

(ii) *The generalized stochastic exponential Z is a martingale if and only if Z is a local martingale and at least one of the conditions (a)-(b) below is satisfied and at least one of the conditions (c)-(d) below is satisfied:*

(a) \tilde{Y} does not exit at $\beta(x_0)$, i.e. $\tilde{v}(\beta(x_0)) = \infty$ or, equivalently, we have

$$\tilde{s}(\beta(x_0)) = \infty \quad \text{or} \quad \left(\tilde{s}(\beta(x_0)) < \infty \text{ and } \frac{\tilde{s}(\beta(x_0)) - \tilde{s}}{\tilde{\rho}\sigma^2} \notin L_{\text{loc}}^1(\beta(x_0)-) \right);$$

(b) $\beta(x_0)$ is good,

(c) \tilde{Y} does not exit at $\alpha(x_0)$, i.e. $\tilde{v}(\alpha(x_0)) = -\infty$ or, equivalently, we have

$$\tilde{s}(\alpha(x_0)) = -\infty \quad \text{or} \quad \left(\tilde{s}(\alpha(x_0)) > -\infty \text{ and } \frac{\tilde{s} - \tilde{s}(\alpha(x_0))}{\tilde{\rho}\sigma^2} \notin L_{\text{loc}}^1(\alpha(x_0)+) \right);$$

(d) $\alpha(x_0)$ is good.

Remark. Part (ii) of Theorem 3.5 says that Z is a martingale if and only if the $(\alpha(x_0), \beta(x_0))$ -valued process \tilde{Y} can exit only at the good endpoints.

Proof. (i) We can write $Z = \bar{Z} - S$ as in (11). The process \bar{Z} is a continuous local martingale by Lemma 2.1. Suppose that Z is a local martingale. Then S can be written as a sum of two local martingales and therefore, it is also a local martingale. It follows that S is a supermartingale (since it is non-negative). Since $\zeta^A > 0$ (we are assuming that $x_0 \in J \setminus A$) and $S_0 = 0$, S should be almost surely equal to 0. By definition (10) of S

this happens if and only if $\mathbb{P}(\tau_A < \infty, \int_0^{\tau_A} b^2(Y_u) du < \infty) = 0$, which is by the definition of the set $B(x_0)$ and (19)–(21) equivalent to $\alpha(x_0), \beta(x_0) \in B(x_0) \cup \{l, r\}$.

(ii) To get at least a local martingale S needs to be zero \mathbb{P} -a.s. Then $Z = \bar{Z}$. Since the values of Y on $[0, \zeta^A)$ do not exit the interval $(\alpha(x_0), \beta(x_0))$, the conditions of Theorem 3.1 are satisfied and the result follows. \square

Similarly, we can characterize uniformly integrable martingale. We can use characterization in Theorem 3.2 for the process \bar{Z} defined by (9). As above, for $\alpha(x_0), \beta(x_0) \in B(x_0) \cup \{l, r\}$ the process Z defined by (8) and (12) coincides with \bar{Z} . Otherwise, Z is not even a local martingale.

Theorem 3.6. *The process Z is a uniformly integrable martingale if and only if Z is a local martingale and at least one of the conditions (a) – (d) below is satisfied:*

- (a) $b = 0$ a.e. on $(\alpha(x_0), \beta(x_0))$ with respect to the Lebesgue measure;
- (b) $\alpha(x_0)$ is good and $\tilde{s}(\beta(x_0)) = \infty$;
- (c) $\beta(x_0)$ is good and $\tilde{s}(\alpha(x_0)) = -\infty$;
- (d) $\alpha(x_0)$ and $\beta(x_0)$ are good.

The following remark simplifies the application of Theorems 3.5 and 3.6 in specific situations.

Remark. If $\alpha(x_0) \in B(x_0)$, then $\alpha(x_0)$ is not a good endpoint. Indeed, if $s(\alpha(x_0)) > -\infty$, then we can write

$$\frac{(s(x) - s(\alpha(x_0)))b^2(x)}{\rho(x)\sigma^2(x)} = \frac{(s(x) - s(\alpha(x_0)))}{(x - \alpha(x_0))\rho(x)} \frac{b^2}{\sigma^2}(x).$$

The first fraction is bounded away from zero, since it is continuous for $x > \alpha(x_0)$ and has a limit equal to 1 as x approaches $\alpha(x_0)$. Since $\alpha(x_0) \in B(x_0)$, (b) of Theorem 3.4 implies $(x - \alpha(x_0))\frac{b^2}{\sigma^2}(x) \notin L_{\text{loc}}^1(\alpha(x_0)+)$. Therefore $\frac{(s - s(\alpha(x_0)))b^2}{\rho\sigma^2} \notin L_{\text{loc}}^1(\alpha(x_0)+)$ and the conclusion follows. Similarly, $\beta(x_0) \in B(x_0)$ implies that $\beta(x_0)$ is not a good endpoint.

4. PROOF OF THEOREM 3.4

For the proof of Theorem 3.4 we first consider the case of a Brownian motion. Let W be a Brownian motion with $W_0 = x_0$. Denote by $L_t^y(W)$ the local time of W at time t and level y . Let us consider a Borel function $b : \mathbb{R} \rightarrow \mathbb{R}$ and set

$$A = \{x \in \mathbb{R} : b^2 \notin L_{\text{loc}}^1(x)\}.$$

We assume that $x_0 \notin A$ and define α_0, β_0 ($\alpha_0 < \beta_0$) so that

$$\alpha_0, \beta_0 \in A \cup \{-\infty, \infty\} \quad \text{and} \quad x_0 \in (\alpha_0, \beta_0) \subseteq \mathbb{R} \setminus A.$$

We additionally assume that $\alpha_0 > -\infty$. Below we use the notations $\tau_x^W, \tau_{x,y}^W$, and τ_A^W for the stopping times defined by (5), (6) and (7) respectively with Y replaced by W .

Lemma 4.1. *If $(x - \alpha_0)b^2(x) \in L_{\text{loc}}^1(\alpha_0+)$, then*

$$\int_0^{\tau_{\alpha_0}^W} b^2(W_t) dt < \infty \quad \mathbb{P}\text{-a.s. on } \{\tau_{\alpha_0}^W = \tau_A^W\}.$$

Remark. In the setting of Lemma 4.1 we have $\mathbb{P}(\tau_{\alpha_0}^W < \infty) = 1$ since Brownian motion reaches every level in finite time almost surely. Therefore the events $\{\tau_{\alpha_0}^W = \tau_A^W\}$ and $\{\tau_{\alpha_0}^W = \tau_A^W < \infty\}$ are equal. Cf. with the formulation of Theorem 3.4.

Proof. Let $(\beta_n)_{n \in \mathbb{N}}$ be an increasing sequence such that $x_0 < \beta_n < \beta_0$ and $\beta_n \nearrow \beta_0$. By [11, Chapter VII, Corollary 3.8] we get

$$\mathbb{E} \left[\int_0^{\tau_{\alpha_0}^W \wedge \tau_{\beta_n}^W} b^2(W_t) dt \right] = \frac{\beta_n - x_0}{\beta_n - \alpha_0} \int_{\alpha_0}^{x_0} (y - \alpha_0) b^2(y) dy + \frac{x_0 - \alpha_0}{\beta_n - \alpha_0} \int_{x_0}^{\beta_n} (\beta_n - y) b^2(y) dy$$

for every β_n . Both integrals are finite since $b^2 \in L_{\text{loc}}^1(\alpha_0, \beta_0)$ and $(x - \alpha_0)b^2(x) \in L_{\text{loc}}^1(\alpha_0+)$. Thus, we have $\mathbb{E}[\int_0^{\tau_{\alpha_0}^W \wedge \tau_{\beta_n}^W} b^2(W_t) dt] < \infty$ and therefore $\int_0^{\tau_{\alpha_0}^W \wedge \tau_{\beta_n}^W} b^2(W_t) dt < \infty$ almost surely for every n . It remains to note that \mathbb{P} -a.s. on $\{\tau_{\alpha_0}^W = \tau_A^W\}$ we have $\tau_{\alpha_0}^W < \tau_{\beta_n}^W$ for sufficiently large n . This concludes the proof. \square

Lemma 4.2. *If $\int_0^{\tau_{\alpha_0}^W} b^2(W_t) dt < \infty$ on a set U with $\mathbb{P}(U) > 0$, then $(x - \alpha_0)b^2(x) \in L_{\text{loc}}^1(\alpha_0+)$.*

Proof. The idea of the proof comes from [4]. Using the occupation times formula we can write

$$\int_0^{\tau_{\alpha_0}^W} b^2(W_t) dt = \int_{\alpha_0}^{\infty} b^2(y) L_{\tau_{\alpha_0}^W}^y(W) dy \geq \int_{\alpha_0}^{x_0} b^2(y) L_{\tau_{\alpha_0}^W}^y(W) dy.$$

Let us define a process $R_y = \frac{1}{y - \alpha_0} L_{\tau_{\alpha_0}^W}^y(W)$. Then R is positive and we have

$$(22) \quad \int_0^{\tau_{\alpha_0}^W} b^2(W_t) dt \geq \int_{\alpha_0}^{x_0} R_y (y - \alpha_0) b^2(y) dy.$$

By [11, Chapter VI, Proposition 4.6], Laplace transform of R_y is

$$\mathbb{E}[\exp\{-\lambda R_y\}] = \frac{1}{1 + 2\lambda} \quad \text{for every } y.$$

Hence, every random variable R_y has exponential distribution with $\mathbb{E}[R_y] = 2$.

Denote by L an indicator function of a measurable set. We can write

$$\mathbb{E}[LR_y] = \mathbb{E} \left[L \int_0^\infty \mathbb{1}_{\{R_y > u\}} du \right] = \int_0^\infty \mathbb{E}[L \mathbb{1}_{\{R_y > u\}}] du.$$

By Jensen's inequality we get a lower bound for the integrand

$$\begin{aligned} \mathbb{E}[L \mathbb{1}_{\{R_y > u\}}] &= \mathbb{E}[(L - \mathbb{1}_{\{R_y \leq u\}})^+] \\ &\geq (\mathbb{E}[L] - \mathbb{P}[R_y \leq u])^+ \\ &= (\mathbb{E}[L] + e^{-\frac{u}{2}} - 1)^+. \end{aligned}$$

Hence,

$$(23) \quad \mathbb{E}[LR_y] \geq \int_0^\infty (\mathbb{E}[L] + e^{-\frac{u}{2}} - 1)^+ du = C,$$

where C is a strictly positive constant if $\mathbb{E}[L]$ is strictly positive.

Let us suppose that we can choose L , so that $\mathbb{E}[L \int_0^{\tau_{\alpha_0}^W} b^2(W_t) dt]$ is finite. Using Fubini's Theorem and inequalities (22) and (23), we get

$$\mathbb{E} \left[L \int_0^{\tau_{\alpha_0}^W} b^2(W_t) dt \right] \geq \int_{\alpha_0}^{x_0} \mathbb{E}[LR_y] (y - \alpha_0) b^2(y) dy \geq C \int_{\alpha_0}^{x_0} (y - \alpha_0) b^2(y) dy.$$

Therefore, $(y - \alpha_0) b^2(y) \in L_{\text{loc}}^1(\alpha_0 +)$ if we can find an indicator function L such that $\mathbb{E}[L]$ is strictly positive and $\mathbb{E}[L \int_0^{\tau_{\alpha_0}^W} b^2(W_t) dt]$ is finite.

Since $\int_0^{\tau_{\alpha_0}^W} b^2(W_t) dt < \infty$ on a set with positive measure, such L exists. Indeed, denote by L_n an indicator function of the set $U_n = \{\int_0^{\tau_{\alpha_0}^W} b^2(W_t) dt \leq n\}$. Then, for every integer n , we have $\mathbb{E}[L_n \int_0^{\tau_{\alpha_0}^W} b^2(W_t) dt] < \infty$. Since the sequence $(U_n)_{n \in \mathbb{N}}$ is increasing, $U \subseteq \bigcup_{n \in \mathbb{N}} U_n$ and $\mathbb{P}(U) > 0$, there exists an integer N such that $\mathbb{P}(U_N) > 0$ and therefore $\mathbb{E}[L_N] > 0$. \square

Now we return to the setting of Section 2.

Proof of Theorem 3.4. Suppose that $\mu \equiv 0$ and σ is arbitrary. Since Y_t is a continuous local martingale, by the Dambis–Dubins–Schwarz theorem we have $Y_t = B_{\langle Y \rangle_t}$ for a Brownian motion B with $B_0 = x_0$, defined possibly on an enlargement of the initial probability space. Using the substitution $u = \langle Y \rangle_t$, we get

$$(24) \quad \int_0^{\tau_{\alpha_0}^Y} b^2(Y_t) dt = \int_0^{\tau_{\alpha_0}^Y} \frac{b^2}{\sigma^2}(Y_t) d\langle Y \rangle_t = \int_0^{\langle Y \rangle_{\tau_{\alpha_0}^Y}} \frac{b^2}{\sigma^2}(B_u) du \quad \mathbb{P}\text{-a.s.}$$

Furthermore, it is easy to see from $Y_t = B_{\langle Y \rangle_t}$ that

$$(25) \quad \langle Y \rangle_{\tau_{\alpha_0}^Y} = \tau_{\alpha_0}^B \quad \mathbb{P}\text{-a.s. on } \{\tau_{\alpha_0}^Y < \infty\}$$

and

$$(26) \quad \{\tau_{\alpha_0}^Y = \tau_A^Y < \infty\} \subseteq \{\tau_{\alpha_0}^B = \tau_A^B\} \quad \mathbb{P}\text{-a.s.}$$

Now (24)–(26) imply the theorem in the case $\mu \equiv 0$ and σ arbitrary.

It only remains to prove the general case when both μ and σ are arbitrary. Let $\tilde{Y}_t = s(Y_t)$, where s is the scale function of Y . Then \tilde{Y} satisfies SDE

$$d\tilde{Y}_t = \tilde{\sigma}(\tilde{Y}_t) dW_t,$$

where $\tilde{\sigma}(x) = s'(q(x))\sigma(q(x))$ and q is the inverse of s .

Define $\tilde{b} = b \circ q$. Since s is strictly increasing, $s(\alpha_0) > -\infty$ by the assumption $\alpha_0 > l$, and $\tilde{Y}_{\tau_{\alpha_0}} = s(Y_{\tau_{\alpha_0}}) = s(\alpha_0)$, it follows that the equality $\tau_{\alpha_0}^Y = \tau_{s(\alpha_0)}^{\tilde{Y}}$ holds \mathbb{P} -a.s. Then we have

$$\int_0^{\tau_{\alpha_0}^Y} b^2(Y_t) dt = \int_0^{\tau_{s(\alpha_0)}^{\tilde{Y}}} \tilde{b}^2(\tilde{Y}_t) dt.$$

Besides, for some small positive ε we have

$$\int_{s(\alpha_0)}^{s(\alpha_0+\varepsilon)} \frac{\tilde{b}^2(x)}{\tilde{\sigma}^2(x)} (x - s(\alpha_0)) dx = \int_{\alpha_0}^{\alpha_0+\varepsilon} \frac{b^2(y)}{\sigma^2(y)} \frac{s(y) - s(\alpha_0)}{s'(y)} dy.$$

The fraction $\frac{s(y)-s(\alpha_0)}{s'(y)(y-\alpha_0)}$ is continuous for $y > \alpha_0$ and tends to 1 as $y \searrow \alpha_0$. Hence it is bounded and bounded away from zero on $(\alpha_0, \alpha_0 + \varepsilon]$. It follows that $(x - \alpha_0) \frac{b^2}{\sigma^2}(x) \in L_{\text{loc}}^1(\alpha_0+)$ if and only if $(x - s(\alpha_0)) \frac{\tilde{b}^2}{\tilde{\sigma}^2}(x) \in L_{\text{loc}}^1(s(\alpha_0)+)$. Then the result follows from the first part of the proof. \square

Remark. It is interesting to note that, in fact, both sets in (26) are \mathbb{P} -a.s. equal. One can prove the reverse inclusion using the Engelbert–Schmidt construction of solutions of SDEs and the fact that $\alpha_0 > l$.

APPENDIX A.

Let Y be a J -valued diffusion starting from x_0 with a drift μ and volatility σ that satisfy the Engelbert–Schmidt conditions. Let $b : J \rightarrow \mathbb{R}$ be a Borel measurable function and let $(c, d) \subseteq J$, $c < x_0 < d$. Recall that, for any $x, y \in J$, the stopping times τ_x and $\tau_{x,y}$ are defined in (5) and (6). We now extend the definition in (6) by setting $\tau_{c,d} := \zeta$ if $c = l, d = r$; $\tau_{c,d} := \tau_c \wedge \zeta$ if $c > l, d = r$; $\tau_{c,d} := \zeta \wedge \tau_d$ if $c = l, d < r$.

Proposition A.1. (i) *The condition*

$$\frac{b^2}{\sigma^2} \in L_{\text{loc}}^1(c, d)$$

implies that for all $t \in [0, \infty)$,

$$\int_0^t b^2(Y_u) du < \infty \text{ } \mathbb{P}\text{-a.s. on } \{t < \tau_{c,d}\}.$$

(ii) For any $\alpha \in J$ such that $\frac{b^2}{\sigma^2} \notin L_{\text{loc}}^1(\alpha)$ we have

$$\int_0^t b^2(Y_u) du = \infty \text{ } \mathbb{P}\text{-a.s. on } \{\tau_\alpha < t < \zeta\}.$$

Remark. Let us note that $\mathbb{P}(\tau_\alpha < \infty) > 0$ by [2, Theorem 2.11]. Clearly, $\{\tau_\alpha < \infty\} = \{\tau_\alpha < \zeta\}$. Hence there exists $t \in [0, \infty)$ such that $\mathbb{P}(\tau_\alpha < t < \zeta) > 0$.

Proof. (i) Using the occupation times formula, \mathbb{P} -a.s. we get

$$(27) \quad \int_0^t b^2(Y_u) du = \int_0^t \frac{b^2}{\sigma^2}(Y_u) d\langle Y \rangle_u = \int_J \frac{b^2}{\sigma^2}(y) L_t^y(Y) dy, \quad t \in [0, \zeta).$$

\mathbb{P} -a.s. on the set $\{t < \tau_{c,d}\}$ the function $y \mapsto L_t^y(Y)$ is càdlàg (see [11, Chapter VI, Theorem 1.7] with a compact support in the interval (c, d)). Now the first statement follows from (27).

(ii) We have

$$\int_{\alpha-\varepsilon}^{\alpha+\varepsilon} \frac{b^2}{\sigma^2}(y) dy = \infty \text{ for all } \varepsilon > 0.$$

By [2, Theorem 2.7], we have for any $t \geq 0$

$$L_t^\alpha(Y) > 0 \text{ and } L_t^{\alpha-}(Y) > 0 \text{ } \mathbb{P}\text{-a.s. on } \{\tau_\alpha < t < \zeta\}.$$

Then, \mathbb{P} -a.s. on $\{\tau_\alpha < t < \zeta\}$, there exists $\varepsilon > 0$ such that the function $y \mapsto L_t^y(Y)$ is bounded away from zero on the interval $(\alpha - \varepsilon, \alpha + \varepsilon)$. It follows from (27) that $\int_0^t b^2(Y_u) du = \infty$ \mathbb{P} -a.s on $\{\tau_\alpha < t < \zeta\}$. This concludes the proof. \square

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