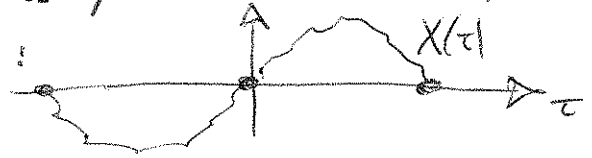


1.2.7 Fourier decomposition

$$\bullet \quad X(\tau) = X_c(\tau) + X(\tau), \quad \left. \begin{array}{l} X_c(0) = X_0, \quad X_c(t) = X_t \\ X(0) = X(t) = 0 \end{array} \right\} \begin{array}{l} (1.2.191) \\ (1.2.192) \\ (1.2.193) \end{array}$$

$X(\tau)$ between 0 and t is expressed as periodic with the period of $2t$ and odd ($X(-\tau) = -X(\tau)$):

~~$$X(\tau) = \sum$$~~



$$X(\tau) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left[a_n \sin \frac{n\pi\tau}{t} + b_n \cos \frac{n\pi\tau}{t} \right]$$

$$b_n = \frac{1}{t} \int_{-t}^t X(\tau) \cos \frac{n\pi\tau}{t} d\tau = 0 \quad \text{as } X(-\tau) = -X(\tau)$$

$$a_n = \frac{1}{t} \int_{-t}^t X(\tau) \sin \frac{n\pi\tau}{t} d\tau = \frac{2}{t} \int_0^t X(\tau) \sin \frac{n\pi\tau}{t} d\tau$$

$$a_0 = \frac{1}{t} \int_{-t}^t X(\tau) d\tau = 0 \quad \text{as } X(\tau) \text{ is odd.}$$

$$\hookrightarrow \boxed{X(\tau) = \sqrt{\frac{2}{t}} \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi\tau}{t}\right)} \quad (1.2.194)$$

$\hookrightarrow$ for convenience

$\bullet \quad \left[\phi_n(\tau) = \sqrt{\frac{2}{t}} \sin\left(\frac{n\pi\tau}{t}\right) \right]$ - forms an orthogonal & complete set of functions: over $\tau \in [0, t]$:

$$\int_0^t \phi_n \phi_m d\tau = \frac{2}{t} \int_0^t \sin\left(\frac{n\pi\tau}{t}\right) \sin\left(\frac{m\pi\tau}{t}\right) d\tau$$

$$\text{If } n \neq m, \text{ then} \quad -\frac{1}{2} \cos\left(\frac{(n+m)\pi\tau}{t}\right) + \frac{1}{2} \cos\left(\frac{(n-m)\pi\tau}{t}\right)$$

$$\hookrightarrow = \frac{2}{t} \left\{ -\frac{1}{2} \frac{t}{(n+m)} \sin\left(\frac{(n+m)\pi\tau}{t}\right) + \frac{1}{2} \frac{t}{(n-m)} \sin\left(\frac{(n-m)\pi\tau}{t}\right) \right\}_0^t = 0$$

$$\text{If } n = m, \text{ then} \quad \int_0^t \phi_n \phi_n d\tau = \frac{2}{t} \int_0^t \left\{ \frac{1}{2} \sin^2\left(\frac{n\pi\tau}{t}\right) - \frac{1}{2} \cos\left(\frac{2n\pi\tau}{t}\right) + \frac{1}{2} \right\} d\tau$$

$$= \frac{2}{t} \cdot \frac{1}{2} \cdot 2t = 1$$

$$\boxed{\int_0^t \psi_n \psi_m d\tau = \delta_{nm}}$$

- So, instead of $X(t)$, coefficients a_n are used as random variables.

Finite difference calculus

$$\boxed{\begin{cases} \partial^{(\varepsilon)} x(t) = \frac{1}{\varepsilon} [x(t+\varepsilon) - x(t)] \rightarrow \partial_t x(t) \\ \bar{\partial}^{(\varepsilon)} x(t) = \frac{1}{\varepsilon} [x(t) - x(t-\varepsilon)] \rightarrow \partial_t x(t) \end{cases}} \quad (1.2.195)$$

$\hookrightarrow \text{as } \varepsilon \rightarrow 0$

$$\partial^{(\varepsilon)} x_j = \frac{1}{\varepsilon} [x_{j+1} - x_j] \quad , \quad 0 \leq j \leq N$$

$$\bar{\partial}^{(\varepsilon)} x_j = \frac{1}{\varepsilon} [x_j - x_{j-1}] \quad , \quad 1 \leq j \leq N+1$$

Integration by parts:

$$\boxed{\varepsilon \sum_{j=1}^{N+1} p_j \bar{\partial}^{(\varepsilon)} x_j = p_j x_j \Big|_0^{N+1} - \varepsilon \sum_{j=0}^N (\partial^{(\varepsilon)} p_j) x_j} \quad (1.2.197)$$

Indeed:

$$\text{LHS} = \varepsilon \sum_{j=1}^{N+1} p_j \frac{1}{\varepsilon} [x_{j+1} - x_j] = \sum_{j=1}^{N+1} (p_j x_{j+1} - p_j x_j)$$

$$\text{RHS} = p_{N+1} x_{N+1} - p_0 x_0 - \varepsilon \sum_{j=0}^N \frac{1}{\varepsilon} (p_{j+1} - p_j) x_j$$

$$= p_{N+1} x_{N+1} - p_0 x_0 - \sum_{j=0}^N (p_{j+1} x_j - p_j x_j)$$

$$= \sum_{j=1}^N p_j x_j - \sum_{j=0}^N p_{j+1} x_j + (p_{N+1} x_{N+1} - p_0 x_0) + p_0 x_0 =$$

$j \rightarrow j+1$

$$= \sum_{j=1}^N p_j x_j - \sum_{j=1}^{N+1} p_j x_{j-1} + p_{N+1} x_{N+1} \equiv \sum_{j=1}^{N+1} p_j x_j - \sum_{j=1}^{N+1} p_j x_{j-1} \equiv \text{RHS}$$

If $\underline{X_{N+1} = X_0 = 0}$

$$\sum_{j=1}^{N+1} p_j \bar{\partial}^{(\varepsilon)} x_j = - \sum_{j=0}^N (\partial^{(\varepsilon)} p_j) x_j \quad (1.2.198)$$

If $p_0 = p_{N+1}$, $x_0 = x_{N+1}$, then

$$\sum_{j=1}^{N+1} p_j \bar{\partial}^{(\varepsilon)} x_j = - \sum_{j=0}^{N+1} (\partial^{(\varepsilon)} p_j) x_j - \cancel{(\partial^{(\varepsilon)} p_0) x_0} \quad (1.2.199)$$

since $(\partial^{(\varepsilon)} p_0) x_0 = \frac{1}{\varepsilon} [p_1 - p_0] x_0 = \frac{1}{\varepsilon} [p_1 - p_{N+1}] x_{N+1} = \frac{1}{\varepsilon} [p_{N+1} - p_{N+1}] x_{N+1} = (\partial^{(\varepsilon)} p_{N+1}) x_{N+1}$ simply rearrangement of terms

• for functions with vanishing endpoints $x_{N+1} = x_0 = 0$:

$$\begin{aligned} \sum_{j=1}^{N+1} (\bar{\partial}^{(\varepsilon)} x_j)^2 &= \sum_{j=1}^{N+1} \frac{1}{\varepsilon^2} (x_j - x_{j-1})^2 = \sum_{j=1}^{N+1} \frac{1}{\varepsilon^2} (x_j^2 + x_{j-1}^2 - 2x_j x_{j-1}) \\ &= - \sum_{k,j=0}^N x_j \left[\partial^{(\varepsilon)} \bar{\partial}^{(\varepsilon)} \right]_{jk} x_k \quad \left[\equiv \sum_{j=0}^N x_j \partial^{(\varepsilon)} \bar{\partial}^{(\varepsilon)} x_j \right] \end{aligned}$$

$$\left[\partial^{(\varepsilon)} \bar{\partial}^{(\varepsilon)} \right]_{jk} = \frac{1}{\varepsilon^2} \begin{pmatrix} 2 & -1 & & \\ -1 & 2 & -1 & \\ & -1 & 2 & \\ & & & \ddots & \\ & & & & 2 & -1 \\ & & & & -1 & 2 \end{pmatrix}$$

$(N+1) \times (N+1)$
matrix

(1.2.202)

• $i \partial^{(\varepsilon)} e^{-i\omega t} = i \frac{1}{\varepsilon} [e^{-i\omega(t+\varepsilon)} - e^{-i\omega t}] = \frac{i}{\varepsilon} (e^{-i\omega\varepsilon} - 1) e^{-i\omega t} \equiv \Omega e^{-i\omega t}$ (1.2.206)
 $i \bar{\partial}^{(\varepsilon)} e^{-i\omega t} = \bar{\Omega} e^{-i\omega t}$ with $\bar{\Omega} = \frac{i}{\varepsilon} (1 - e^{i\omega\varepsilon}) \equiv \Omega^*$ (1.2.207)

In the continuum limit $\varepsilon \rightarrow 0$:

$$\Omega \rightarrow \frac{i}{1} (e^{-i\omega\varepsilon} (-i\omega)) \Rightarrow \omega, \quad \bar{\Omega} \rightarrow \omega \text{ as well}$$

The Wiener measure in terms of Fourier coefficients

$$X(t) \Rightarrow X(\tau_j), \quad j = 0, \dots, N+1$$

Then the sum (1.2.197) will be carried out only up to N :

$$X(\tau_j) = \sum_{n=1}^N \sqrt{\frac{2}{(N+1)\varepsilon}} a_n \sin(\nu_n \tau_j) \equiv X_j \quad (1.2.209)$$

$$\nu_n = \frac{j\pi n}{t} \rightarrow \frac{j\pi n}{(N+1)\varepsilon}$$

$$\tau_j = 2j$$

The expansion functions are orthogonal:

$$\psi_n(\tau_0) = 0$$

$$\psi_n(\tau_{N+1}) = 0$$

$$\psi_n(\tau_j) = \sqrt{\frac{2}{N+1}} \sin(\nu_n \tau_j)$$

$$\sum_{j=1}^N \psi_n(\tau_j) \psi_m(\tau_j) = \frac{2}{N+1} \sum_{j=1}^N \sin(\nu_n \tau_j) \sin(\nu_m \tau_j)$$

$$= \frac{2}{N+1} \sum_{j=1}^N \frac{1}{2i} [e^{i\nu_n \tau_j} - e^{-i\nu_n \tau_j}] \frac{1}{2i} [e^{i\nu_m \tau_j} - e^{-i\nu_m \tau_j}]$$

$$= \frac{2}{N+1} \left(-\frac{1}{4}\right) \sum_{j=1}^N \left\{ e^{i(\nu_n + \nu_m)\tau_j} - e^{-i(\nu_n + \nu_m)\tau_j} - e^{i(\nu_n - \nu_m)\tau_j} + e^{-i(\nu_n - \nu_m)\tau_j} \right\}$$

~~$$= -\frac{1}{2(N+1)} \cdot 2 \operatorname{Re} \cdot \sum_{j=1}^N \left\{ e^{i\nu_{n+m}\tau_j} - e^{-i\nu_{n-m}\tau_j} \right\}$$

$$\tau_j = 2j, \quad \sum_{j=1}^N e^{i\nu \tau_j} = \sum_{j=1}^N (e^{i\nu \varepsilon})^j = \frac{e^{i\nu \varepsilon} - e^{i\nu \varepsilon(N+1)}}{1 - e^{i\nu \varepsilon}}$$

$$\hookrightarrow -\frac{1}{2(N+1)} \operatorname{Re} \left[\frac{e^{i\nu_{n+m}\varepsilon} - e^{i\nu_{n+m}\varepsilon(N+1)}}{1 - e^{i\nu_{n+m}\varepsilon}} - \frac{e^{i\nu_{n-m}\varepsilon} - e^{i\nu_{n-m}\varepsilon(N+1)}}{1 - e^{i\nu_{n-m}\varepsilon}} \right]$$

$$= -\frac{1}{N+1} \left\{ \frac{(1 - e^{-i\nu_{n+m}\varepsilon}) (e^{i\nu_{n+m}\varepsilon})}{2(1 - \cos \nu_{n+m}\varepsilon)} - \frac{(1 - e^{-i\nu_{n-m}\varepsilon}) (e^{i\nu_{n-m}\varepsilon})}{2(1 - \cos \nu_{n-m}\varepsilon)} \right\}$$~~

~~$\frac{1}{2(N+1)}$~~ But $\sum_{j=1}^N e^{i\nu_k \tau_j} = \sum_{j=1}^N e^{i\nu_k \delta j} = \sum_{j=1}^N (e^{i\nu_k \delta})^j$

$$= \sum_{j=0}^N \dots - 1 = \frac{1 - e^{i\nu_k \delta(N+1)}}{1 - e^{i\nu_k \delta}} - 1$$

Here

$$\nu_k \delta = \frac{\pi k}{(N+1)\delta} \cdot \delta = \frac{\pi k}{(N+1)\delta}$$

$$e^{i\nu_k \delta(N+1)} = e^{i\pi k} = (-1)^k$$

$$\text{and } \sum_{j=1}^N = \begin{cases} 0 - 1 = -1, & \text{if } k = \text{even} \\ \frac{2}{1 - e^{i\nu_k \delta}} - 1, & \text{if } k = \text{odd} \end{cases}$$

~~if~~ ~~when~~ $n \neq m$, then. We have:

$$\sum_{j=1}^N \psi_n(\tau_j) \psi_m(\tau_j) = -\frac{1}{2(N+1)} \left\{ \left[\frac{1 - (-1)^{n+m}}{1 - e^{i\nu_{n+m} \delta}} + \frac{1 - (-1)^{n+m}}{1 - e^{-i\nu_{n+m} \delta}} \right] \right.$$

$$\left. - \left[\frac{1 - (-1)^{n-m}}{1 - e^{-i\nu_{n-m} \delta}} + \frac{1 - (-1)^{n-m}}{1 - e^{i\nu_{n-m} \delta}} \right] \right\}$$

If $n-m = \text{even}$, then, $n+m = (n-m) + 2m = \text{even}$ as well. $\rightarrow 0$

If $n \neq m$. If $n-m = \text{odd}$, then $n+m = (n-m) + 2m = \text{odd}$ as well.

$$\left[\dots \right]_{(k=n+m)} = 2 \left[\frac{1}{1 - e^{i\nu_k \delta}} + \frac{1}{1 - e^{-i\nu_k \delta}} \right] = 2 \frac{1 - e^{-i\nu_k \delta} + 1 - e^{i\nu_k \delta}}{2 - 2 \cos \nu_k \delta}$$

$$= 2 \frac{2(1 - \cos \nu_k \delta)}{1 - \cos \nu_k \delta} = 2$$

The same for $k=n-m$ term $\rightarrow 0$ again.

If, however, $n=m$, then

$$\sum_{j=1}^N \psi_n(\tau_j)^2 = -\frac{1}{2(N+1)} \sum_{j=1}^N \left\{ \left(e^{i\nu_{2n} \tau_j} - e^{-i\nu_{2n} \tau_j} \right) + 2 \right\}$$

$$= -\frac{1}{2(N+1)} \left\{ \cancel{-2N} - 2N + \left[\underbrace{\frac{1 - e^{i\nu_{2N}\epsilon(N+1)}}{1 - e^{i\nu_{2N}\epsilon}}}_{=0} - 1 + \frac{1 - e^{-i\nu_{2N}\epsilon(N+1)}}{1 - e^{-i\nu_{2N}\epsilon}} - 1 \right] \right\}$$

$$= -\frac{1}{2(N+1)} \{-2N - 2\} = 1$$

$$\boxed{\sum_{j=1}^N \psi_n(\tau_j) \psi_m(\tau_j) = \delta_{nm} \equiv \frac{2}{N+1} \sum_{j=1}^N \sin(\nu_n \tau_j) \sin(\nu_m \tau_j)} \quad (1.2.2)$$

So, they are orthogonal. Completeness:

$$\begin{aligned} \sum_{n=1}^N \psi_n(\tau_j) \psi_n(\tau_{j'}) &= \frac{2}{N+1} \sum_{n=1}^N \sin(\nu_n \tau_j) \sin(\nu_n \tau_{j'}) \\ &= \frac{2}{N+1} \left(-\frac{1}{4} \right) \sum_{n=1}^N \left\{ e^{i\nu_n \epsilon(j+j')} + e^{-i\nu_n \epsilon(j+j')} - e^{-i\nu_n \epsilon(j-j')} - e^{i\nu_n \epsilon(j-j')} \right\} \end{aligned}$$

If $j=j'$, then

$$\hookrightarrow -\frac{1}{2(N+1)} \sum_{n=1}^N \left\{ e^{i\nu_n \epsilon 2j} + e^{-i\nu_n \epsilon 2j} - 2 \right\},$$

$$\sum_{n=1}^N e^{i\nu_n \epsilon 2j} = \sum_{n=1}^N \left(e^{i\frac{\pi k}{N+1}} \right)^n = \sum_{n=0}^N \left(e^{i\frac{\pi k}{N+1}} \right)^n - 1$$

$$= \frac{1 - e^{i\frac{\pi k}{N+1}(N+1)}}{1 - e^{i\frac{\pi k}{N+1}}} - 1 = \frac{1 - e^{i\pi k}}{1 - e^{i\pi k/(N+1)}} - 1 = \frac{1 - (-1)^k}{1 - e^{i\frac{\pi k}{N+1}}} - 1$$

So, if $k=2j \rightarrow$ we obtain -1 from each sum,

$$\hookrightarrow -\frac{1}{2(N+1)} \{-2 - 2N\} = 1$$

If $j \neq j'$, then

$$\hookrightarrow -\frac{1}{2(N+1)} \left\{ \left[\frac{1 - (-1)^{j+j'}}{1 - e^{i\pi(j+j')/(N+1)}} + \frac{1 - (-1)^{j-j'}}{1 - e^{-i\pi(j-j')/(N+1)}} \right] - \right.$$

-95a-

$$- \left[\frac{1 - (-1)^{j-j'}}{1 - e^{\frac{i\pi(j-j')}{N+1}}} + \frac{1 - (-1)^{j+j'}}{1 - e^{\frac{i\pi(j+j')}{N+1}}} \right]$$

It is 0 if $j-j'$ = even. If $j-j'$ = odd, then

$$[\dots] = 2 \frac{2}{1 - e^{\frac{i\pi k}{N+1}}} + \frac{2}{1 - e^{-\frac{i\pi k}{N+1}}}$$

$$= 2 \frac{2(1 - \cos \frac{\pi k}{N+1})}{2(1 - \cos \frac{\pi k}{N+1})} = 2, \quad k = j+j' \text{ or } j-j'$$

and both brackets cancel out. We hence obtain:

$$\sum_{k=1}^N \psi_k(\tau_j) \psi_k(\tau_{j'}) = \frac{2}{N+1} \sum_{k=1}^N \sin(\nu_k \tau_j) \sin(\nu_k \tau_{j'}) = \delta_{j,j'} \quad (1.2.21)$$

• Jacobian of transformation $X(\tau_j) \leftrightarrow a_m$:

$$\frac{\partial X(\tau_j)}{\partial a_m} = \sqrt{\frac{2}{(N+1)\epsilon}} \sin(\nu_m \tau_j) = \frac{1}{\sqrt{\epsilon}} \psi_m(\tau_j) = \epsilon^{-1/2} \psi_m(\tau_j)$$

and $J = \left| \frac{\partial X(\tau_j)}{\partial a_m} \right| = \epsilon^{-N/2} |\psi_m(\tau_j)| \equiv \epsilon^{-N/2} \quad (1.2.212)$
 as ψ_m are orthogonal

and we obtain in new variables:

$$d_w X(\tau) = \left(\prod_{j=1}^N \frac{dX_j}{\sqrt{4\pi\epsilon}} \right) \frac{1}{\sqrt{4\pi\epsilon}} \exp \left\{ - \frac{1}{4\epsilon} \sum_{j=1}^{N+1} \frac{(X_j - X_{j-1})^2}{\epsilon} \right\}$$

• Consider the exp. term. It is diagonalised in the new variables:

$$\sum_{j=1}^{N+1} \frac{(X_j - X_{j-1})^2}{\epsilon} = \sum_{j=1}^N \frac{(\sum_{k=1}^N \frac{1}{\sqrt{2}} \psi_k(\tau_j) a_k)^2}{\epsilon} = \sum_{j=1}^N \sum_{k=1}^N \frac{1}{2} \psi_k(\tau_j)^2 a_k^2$$

NOT FINISHED

-968-

$$\sum_{j=0}^N \frac{1}{\varepsilon} (X_{j+1} - X_j)^2 = \frac{1}{\varepsilon} \sum_{j=0}^N (X_{j+1} - X_j)^2,$$

$$X_{j+1} - X_j = \sum_n a_n [\psi_n(\tau_{j+1}) - \psi_n(\tau_j)]$$

$$= \sum_{n=1}^N a_n \sqrt{\frac{2}{N+1}} \left[\sin\left(\frac{\nu_n \tau_{j+1}}{\nu_n \tau_j + \nu_n \varepsilon}\right) - \sin(\nu_n \tau_j) \right]$$

$$\equiv \sqrt{\frac{2}{N+1}} \sum_{n=1}^N a_n 2 \cos\left(\nu_n \tau_j + \frac{\nu_n \varepsilon}{2}\right) \sin\left(\frac{1}{2} \nu_n \varepsilon\right)$$

$$= \sqrt{\frac{2}{N+1}} \sum_{n=1}^N a_n \frac{1}{2} \left[e^{i \nu_n \tau_j} e^{i \nu_n \varepsilon / 2} + e^{-i \nu_n \tau_j} e^{-i \nu_n \varepsilon / 2} \right] \sin\left(\frac{1}{2} \nu_n \varepsilon\right)$$

Denote $\lambda_n = \exp(i \nu_n \varepsilon) = \exp\left(\frac{i \pi n}{N+1}\right)$, then

$$X_{j+1} - X_j = \sqrt{\frac{2}{N+1}} \sum_{n=1}^N a_n \sin\left(\frac{1}{2} \nu_n \varepsilon\right) \left[\sqrt{\lambda_n} e^{i \nu_n \tau_j} + \frac{1}{\sqrt{\lambda_n}} e^{-i \nu_n \tau_j} \right]$$

and hence

$$\sum_{j=0}^N \frac{1}{\varepsilon} (X_{j+1} - X_j)^2 = \left(\frac{2}{N+1}\right) \frac{1}{\varepsilon} \sum_{n,m=1}^N a_n a_m \sin\left(\frac{1}{2} \nu_n \varepsilon\right) \sin\left(\frac{1}{2} \nu_m \varepsilon\right) \times$$

$$\times \left\{ \sum_{j=0}^N \left[\sqrt{\lambda_n} e^{i \nu_n \tau_j} + \frac{1}{\sqrt{\lambda_n}} e^{-i \nu_n \tau_j} \right] \left[\sqrt{\lambda_m} e^{i \nu_m \tau_j} + \frac{1}{\sqrt{\lambda_m}} e^{-i \nu_m \tau_j} \right] \right\} \quad (*)$$

Consider $\{.. \}$ for $n \neq m$:

$$\{.. \} = \sqrt{\lambda_n \lambda_m} \left(\sum_j e^{i \nu_{n+m} \tau_j} \right) + \sqrt{\frac{\lambda_m}{\lambda_n}} \left(\sum_j e^{-i \nu_{n-m} \tau_j} \right) + \sqrt{\frac{\lambda_n}{\lambda_m}} \left(\sum_j e^{i \nu_{n-m} \tau_j} \right)$$

$$+ \frac{1}{\sqrt{\lambda_n \lambda_m}} \left(\sum_j e^{-i \nu_{n+m} \tau_j} \right)$$

Here $\sum_{j=0}^N e^{i \nu_k \tau_j} = \sum_{j=0}^N e^{\frac{i \pi k}{(N+1) \varepsilon} \varepsilon j} = \sum_{j=0}^N \left(e^{\frac{i \pi k}{N+1}} \right)^j = \frac{1 - e^{\frac{i \pi k (N+1)}{N+1}}}{1 - e^{\frac{i \pi k}{N+1}}}$

-97e-

$$= \frac{1 - e^{i\pi k}}{1 - \lambda_k} = \frac{1 - (-1)^k}{1 - \lambda_k}$$

Therefore, $\delta B = 0$ if $k = \text{even}$.

$$\begin{aligned} \{c_n\} &= \sqrt{\lambda_n \lambda_m} \frac{1 - (-1)^{n+m}}{1 - \lambda_n \lambda_m} + \frac{1}{\sqrt{\lambda_n \lambda_m}} \frac{1 - (-1)^{n+m}}{1 - \frac{1}{\lambda_n \lambda_m}} + \\ &+ \sqrt{\frac{\lambda_m}{\lambda_n}} \frac{1 - (-1)^{n+m}}{1 - \frac{\lambda_m}{\lambda_n}} + \sqrt{\frac{\lambda_n}{\lambda_m}} \frac{1 - (-1)^{n+m}}{1 - \frac{\lambda_n}{\lambda_m}} \\ &= \cancel{\lambda_n} [1 - (-1)^{n+m}] \left\{ \cancel{\sqrt{\lambda_n \lambda_m}} \frac{1}{1 - \lambda_n \lambda_m} + \cancel{\sqrt{\lambda_n \lambda_m}} \frac{1}{\lambda_n \lambda_m - 1} \right. \\ &\left. + \cancel{\sqrt{\frac{\lambda_m}{\lambda_n}}} \frac{\lambda_n}{\lambda_m - \lambda_n} + \cancel{\sqrt{\frac{\lambda_n}{\lambda_m}}} \frac{\lambda_m}{\lambda_m - \lambda_n} \right\} = 0 \end{aligned}$$

Therefore, terms $\ln(X)$ with $n \neq m$ vanish.

Consider now $n = m$; the $\{c_n\}$ brackets from (x):

$$\begin{aligned} \{c_n\} &= \lambda_n \sum_{j=0}^N e^{i\varphi_{2n} \tau_j} + 2(N+1) + \frac{1}{\lambda_n} \sum_{j=0}^N e^{-i\varphi_{2n} \tau_j} \\ &= 2(N+1) + \lambda_n \underbrace{\frac{1 - (-1)^{2n}}{\dots}}_0 + \frac{1}{\lambda_n} \underbrace{\frac{1 - (-1)^{2n}}{\dots}}_0 = 2(N+1) \end{aligned}$$

Therefore,

$$\begin{aligned} \sum_{j=0}^N \frac{1}{\varepsilon} (X_{j+1} - X_j)^2 &= \frac{2}{\varepsilon(N+1)} \sum_{n=1}^N a_n^2 \sin^2\left(\frac{1}{2} \varphi_n \varepsilon\right) 2(N+1) \\ &= \frac{4}{\varepsilon} \sum_{n=1}^N \sin^2\left(\frac{1}{2} \varphi_n \varepsilon\right) a_n^2 \equiv \frac{4}{\varepsilon} \sum_{n=1}^N \frac{1}{2} (1 - \cos(\varphi_n \varepsilon)) a_n^2 \\ &= \frac{1}{\varepsilon} \sum_{n=1}^N \underbrace{[2(1 - \cos(\varphi_n \varepsilon))]}_{\varepsilon^2 \Omega_n \bar{\Omega}_n} a_n^2 = \varepsilon \sum_{n=1}^N \underbrace{\Omega_n \bar{\Omega}_n}_{|\Omega_n|^2} a_n^2 \end{aligned}$$

Now

$$d_w X(\tau) = \left(\prod_{j=1}^N \frac{dX_j}{\sqrt{4\pi\Delta\tau}} \right) \frac{1}{\sqrt{4\pi\Delta\tau}} \cdot \exp \left\{ -\frac{1}{4\Delta\tau} \sum_{n=1}^N \Omega_n \bar{\Omega}_n q_n^2 \right\} \quad (1.2.213)$$

Show the diagonalisation via the discrete differential calculus.
(as in the book).

$$\begin{aligned} \sum_{j=0}^N \frac{1}{\varepsilon} (X_{j+1} - X_j)^2 &= \frac{1}{\varepsilon} \sum_{j=0}^N (X_{j+1}^2 + X_j^2 - 2X_j X_{j+1}) \\ &= \frac{1}{\varepsilon} \left[\underbrace{X_1^2}_{\sim} + \underbrace{(X_2^2 + X_1^2 - 2X_1 X_2)}_{\sim} + \underbrace{(X_3^2 + X_2^2 - 2X_2 X_3)}_{\sim} + \dots \right. \\ &\quad \left. + \underbrace{(X_N^2 + X_{N-1}^2 - 2X_{N-1} X_N)}_{\sim} + \underbrace{X_N^2}_{\sim} \right] \end{aligned}$$

if $X_0 = X_{N+1} \equiv 0$.

This can also be written as:

$$\begin{aligned} -\varepsilon \sum_{j=0}^N \frac{1}{\varepsilon} \partial^{(\varepsilon)} X_j \bar{\partial}^{(\varepsilon)} X_j &= -\frac{1}{\varepsilon} \sum_{j=0}^N X_j \partial^{(\varepsilon)} \frac{1}{\varepsilon} [X_j - X_{j-1}] \\ &= -\sum_{j=0}^N X_j (\partial^{(\varepsilon)} X_j - \partial^{(\varepsilon)} X_{j-1}) = -\frac{1}{\varepsilon} \sum_{j=0}^N X_j [X_{j+1} - X_j - (X_j - X_{j-1})] \\ &= -\frac{1}{\varepsilon} \sum_{j=0}^N X_j (X_{j+1} - 2X_j + X_{j-1}) = -\frac{1}{\varepsilon} \sum_{j=0}^N (-2X_j^2 + X_j X_{j+1} + X_{j-1} X_j) \\ &= \frac{1}{\varepsilon} \sum_{j=1}^N (2X_j^2 - X_j X_{j+1} - X_{j-1} X_j) \\ &= \frac{1}{\varepsilon} \left[(2X_1^2 - X_1 X_2) + (2X_2^2 - X_2 X_3 - X_1 X_2) + (2X_3^2 - X_3 X_4 - X_2 X_3) + \dots \right. \\ &\quad \left. + (2X_{N-1}^2 - X_{N-1} X_N - X_{N-2} X_{N-1}) + (2X_N^2 - X_{N-1} X_N) \right] \end{aligned}$$

which is the same thing. Hence, the term in the exponential

$$\boxed{\sum_{j=0}^N \frac{1}{\varepsilon} (X_{j+1} - X_j)^2 = -\varepsilon \sum_{j=0}^N X_j \partial^{(\varepsilon)} \bar{\partial}^{(\varepsilon)} X_j} \quad (1.2.200)$$

Now,

$$X_j = \sum_{n=1}^N a_n \sqrt{\frac{2}{\varepsilon(N+1)}} \sin(\nu_n \tau_j)$$

so that

$$\sum_{j=0}^N \frac{1}{\varepsilon} (X_{j+1} - X_j)^2 = -\varepsilon \frac{2}{\varepsilon(N+1)} \sum_{n,m=1}^N a_n a_m \sum_{j=0}^N \sin(\nu_n \tau_j) \left[\partial^{(\varepsilon)} \bar{\partial}^{(\varepsilon)} \sin(\nu_m \tau_j) \right]$$

$$\equiv -\frac{2}{N+1} \sum_{n,m} a_n a_m \left(\frac{1}{2i} \right) \sin \quad \text{Consider}$$

$$\partial^{(\varepsilon)} \bar{\partial}^{(\varepsilon)} \sin \nu_m \tau_j = \frac{1}{2i} \partial^{(\varepsilon)} \bar{\partial}^{(\varepsilon)} \left[e^{i\nu_m \tau_j} - e^{-i\nu_m \tau_j} \right]$$

$$= -\frac{1}{2i} \partial^{(\varepsilon)} \left\{ \bar{\Omega}_m(-\nu_m) e^{i\nu_m \tau_j} - \bar{\Omega}_m(+\nu_m) e^{-i\nu_m \tau_j} \right\}$$

$$= -\frac{1}{2i} \left\{ \bar{\Omega}_m(-\nu_m) \Omega_m(-\nu_m) e^{i\nu_m \tau_j} - \bar{\Omega}_m(\nu_m) \Omega_m(\nu_m) e^{-i\nu_m \tau_j} \right\}$$

where

$$\bar{\Omega}_m(-\nu_m) \Omega_m(-\nu_m) = \frac{i}{\varepsilon} (1 - e^{-i\nu_m \varepsilon}) \frac{i}{\varepsilon} (e^{i\nu_m \varepsilon} - 1)$$

$$= \frac{-1}{\varepsilon^2} (e^{i\nu_m \varepsilon} - 1 - 1 + e^{-i\nu_m \varepsilon}) = \frac{-1}{\varepsilon^2} (-2 + 2 \cos \nu_m \varepsilon)$$

$$= \frac{2}{\varepsilon^2} (1 - \cos \nu_m \varepsilon)$$

and also

$$\bar{\Omega}_m(\nu_m) \Omega_m(\nu_m) = \frac{i}{\varepsilon} (1 - e^{+i\nu_m \varepsilon}) \frac{i}{\varepsilon} (e^{-i\nu_m \varepsilon} - 1)$$

$$= -\frac{1}{\varepsilon^2} (-1 - 1 + e^{-i\nu_m \varepsilon} + e^{i\nu_m \varepsilon}) = \frac{2}{\varepsilon^2} (1 - \cos \nu_m \varepsilon) \equiv \Omega_m \bar{\Omega}_m,$$

i.e. the same. Therefore,

$$\partial^{(\varepsilon)} \bar{\partial}^{(\varepsilon)} \sin \nu_m \tau_j = -\frac{1}{2i} \bar{\Omega}_m \Omega_m (e^{i\nu_m \tau_j} - e^{-i\nu_m \tau_j})$$

$$= -\bar{\Omega}_m \Omega_m \sin(\nu_m \tau_j)$$

which gives

$$\sum_{j=0}^N \frac{1}{\varepsilon} (X_{j+1} - X_j)^2 = -\frac{2}{N+1} \sum_{n,m=1}^N a_n a_m \sum_{j=0}^N (-\bar{\Omega}_m \Omega_m) \sin \nu_n \tau_j \sin \nu_m \tau_j$$

-100 -

10000

$$= -\frac{2}{N+1} \sum_{n,m} (-\Omega_m \bar{\Omega}_m) a_n a_m \underbrace{\left(\sum_{j=0}^N \sin \nu_n \bar{\nu}_j \sin \nu_m \bar{\nu}_j \right)}_{\delta_{nm} \frac{N+1}{2}}$$

$$= + \sum_{n,m} \Omega_m \bar{\Omega}_m a_n^2 = \sum_{n,m} |\Omega_m|^2 a_n^2,$$

ie the same result.

Q: How can we show that the energy is conserved?

A: We can show that the energy is conserved by showing that the time derivative of the energy is zero.

Let's consider the energy of the system:

$$E = \frac{1}{2} \sum_{n=1}^N \left(\dot{q}_n^2 + \omega_n^2 q_n^2 \right)$$

where q_n is the displacement of the mass m_n .

The time derivative of the energy is:

$$\dot{E} = \sum_{n=1}^N \dot{q}_n \ddot{q}_n + \omega_n^2 q_n \dot{q}_n$$