

large c behaviour given by contribution from states which do not quate c in commutation relations.

only $L_{-1}^n |h\rangle \Rightarrow$ leading c behaviour is

$$\begin{array}{c} 2 \quad 3 \\ | \quad | \\ \hline 1 \quad h \quad 4 \end{array} = z^{h-h_3-h_4} \underbrace{F(h-h_{12}, h+h_{34}, 2h, z)}_{\text{global block}} + \dots$$

We know

$$\begin{array}{c} 2 \quad 3 \\ | \quad | \\ \hline 1 \quad h \quad 4 \end{array} = z^{h-h_3-h_4} \left\{ 1 + \frac{(h-h_{12})(h+h_{34})}{2h} z + \dots \right.$$

$$+ z^2 \frac{(\quad) \cdot M_2(\quad)}{\det M_2} \quad (\text{from } \S 5) + \dots$$

(ii) What is remainder after taking off global block contribution

$$\frac{(h-h_{12})_2 (h+h_{34})_2}{2! (2h)_2} z^2 \quad ?$$

Answer: remarkably simple.

$$z^2 \cdot \left(h_1^2 - 3h_1^2 + h_2 + 6h_1 h_2 - 3h_2^2 - h + 2h_1 h + 2h_2 h + h^2 \right)$$

$$\left(\begin{array}{l} h_1 \rightarrow h_4, h_2 \rightarrow h_3 \end{array} \right)$$

$$2 (2h+1)^2 (C - C_{12})$$

These factors ensure the residue of the pole vanishes when the null state means the term doesn't actually appear.

Recall: If $h = h_{12}$, the state space is reduced. This factor is not present.

Check: Factor vanishes if $(C=1/2)$ $h=1/2$, $h_1=h_2=1/4$

Result:

$F(h-h_{12}, h+h_{34}, h, z)$ gets correction

$$+ \sum_{\substack{m \geq 1 \\ n \geq 2}} \frac{R_{mn}}{c - c_{mn}} F(c_{mn}, \{h_i\}, \underbrace{h+h_{mn}}_{\text{weight of vanishing intermediate state}}, z)$$

\uparrow value at the pole $\underbrace{\hspace{2cm}}$

$$R_{mn} = P_{mn} \times A_{mn}$$

\uparrow Factor that ensure residue vanishes \uparrow overall h -dependent factor.

$\underbrace{\hspace{2cm}}$ calculated
 $\underbrace{\hspace{2cm}}$ guessed.

Formula works.

Can formally give a closed form expression for the recursion (Perlmutter).

2nd recursion: repeat for large h / poles in h

leading behaviour: can read off from series expansion:

$$1 + \frac{h}{2}x + \frac{h^2 x^2}{8} + \frac{h^3 x^3}{48} + \frac{h^4 x^4}{384} \dots = e^{xh/2}$$

Actually, much more complicated

$$e_{h_1} \frac{e_{h_2} e_{h_3}}{h} e_{h_4} \sim (16q)^h \left(\frac{z(1-z)}{16q} \right)^{\frac{c-1}{24}} \theta_3(q)^{\frac{c-1}{2} - \sum h_i}$$

($z^{-h_3-h_4}$ missing...)

Here $\theta_3(q) = \sum_{n=-\infty}^{\infty} q^{n^2}$ $q = e^{-\pi \frac{K(1-z)}{K(z)}}$

$$K(z) = \frac{\pi}{2} F\left(\frac{1}{2}, \frac{1}{2}; 1; z\right) = \frac{1}{2} \int_0^1 \frac{dt}{\sqrt{t(1-t)(1-zt)}}$$

$$= \frac{\pi}{2} \theta_3^2(q)$$

Recursion: $e_{h_1} \frac{e_{h_2} e_{h_3}}{h} e_{h_4} = \left[\dots \right] \cdot H(c, \{h_i\}, h, q)$

$$H(c, \{h_i\}, h, q) = 1 + \sum_{\substack{m, n \\ \geq 1}} \frac{(16q)^{mn} R'_{mn}}{h - h_{mn}(c)} \cdot H(c, \{h_i\}, h_{mn} + mn, q)$$

(Again, can be resummed)

These are

- Very effective numerically
- Somewhat tricky in cases when poles arise.
- Very effective at large c / large h .

Next: Exact formula.